

Neutrino history:

from the beginning of β -decay
to modern discoveries of neutrino
oscillations.

Prof. Yu. Gaponov, RRC "KI"

(Theses of the report)

(Without of references!)

2.

Development of neutrino conception in particle physics of XX century

Yu.V. Gaponov
R.R.C. "Kurchatov Institute", Moscow.

1. Careful study of β -spectra (1899 - 20's years)
2. Neutrino conception occurs (1929 - 37)
3. Experimental observations of nuclei
and particles β -decays (30's - 50's years)
4. Theoretical discoveries of 50's,
Standard Model (SM) delivery (1955 - 67)
5. Experimental statement of the SM (50's - 70's y.)
6. Theoretical generalization of the SM - Grand
Unification Theory and return to Majorana
neutrino idea (80's years)
7. In search for an exit beyond SM -
experiments (80's - 90's years)
8. Conclusion: Neutrino physics on the
threshold of the new century

3.

Riddles of β -ray spectra (1899 - end of 20's y.)

β -ray discovery, β -spectrometer

Bequerel, 1900
Paschen, 1904

Problem of absorption

Exponential law
new β -sources
Linear law

Schmidt, 06
Hahn, Meitner, 08
Thomson, Bragg,
Wilson, 09

Problem of β -lines

β -lines spectrum (photoplates)
(magnetic meth.)

nuclear origin of β
internal conversion γ into β
experiments of β -lines of Ra β +C

Danysz J., 11
Hahn, Meitner, 12
Bohr, 13
Rutherford, 14
Rutherford, Robinson
Rawlinson, 14

Continuous spectrum (Geiger count)

Fine β -line spectra
 β - α analogy
 γ -spectra of nuclei

Chadwick, 14
Ellis, 20
Meitner, 22
Smekal 22

Secondary electron effects,
 β - γ time order of emission

Radiationless processes
 γ -ray absence in RaE
Primary of β -rays

Roseland, 23
Meitner, 23
Ellis-Skinner, 23
Ellis-Wooster, 25

Continuous spectrum riddle

Heat measurements
Upper limit in β -spectrum

Ellis-Wooster, 27
Meitner, 29
Ellis, 29-32
Sargent, 32

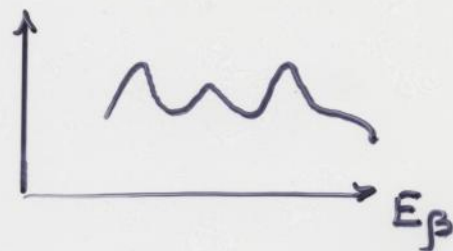
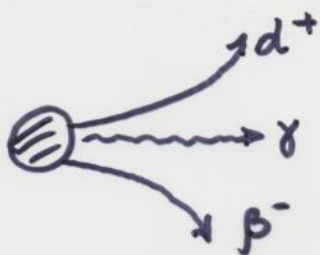
Energy conservation in
ThC - ThD β -chain cycle

Ellis, Nevil, 33

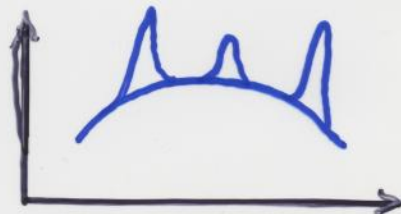
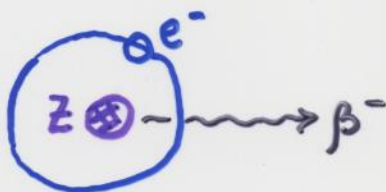
Riddles of β -rays

β -rays charge v , e/m , absorption	1900 - 08	magnetic properties of β -rays	fast e^-
β -spectrometers photoplates, Geiger's counters	1910 - 14	diff. spectra + monolines contin. spectrum	Atom. model of Bohr-Reserf- outer electrons
external source of β -lines excitations	1914	connect. of convers. lines with atomic electr.	photoeffect on outer e^-
β -line spectra, coordination with atomic ones	1920 - 23	systematics of conv. lines γ -sp. from nuclei	coordin. with K, L, M ... atomic spectra, nuclear γ -rays
Z estimations from convers. lines β - γ order	1923 - 25	priority of β absence of monoenerg. nucl. electrons	e-p nucl. model β -d. analogy of L. Meitner
Investigation of RaE continuous spectrum, heat effect experim.	1927 - 29	Problem of continuous spectrum $E_\beta \sim E_0/2$	Inner nucl. electrons? E conserv. law ?
Upper limit of β -sp. d- β chains in Th	1929 - 33	β -sp. is nonstati- stic? Energy conserv. law for $E_{0\beta}$	Quantum mech. ? Paradox ? Where is ? energy ?

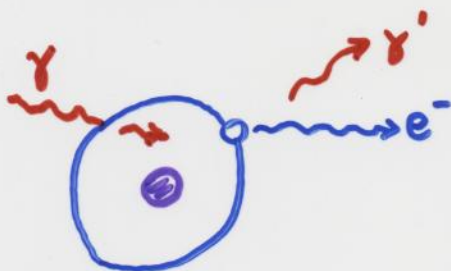
1900-08



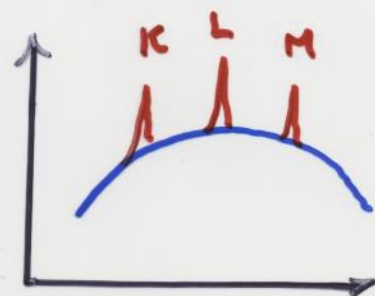
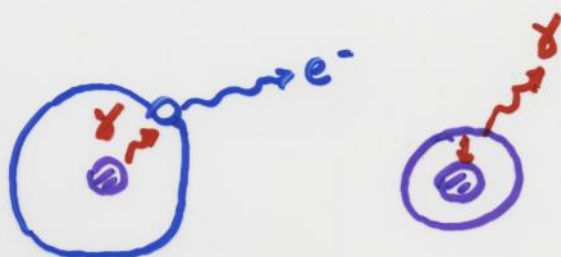
1910-14



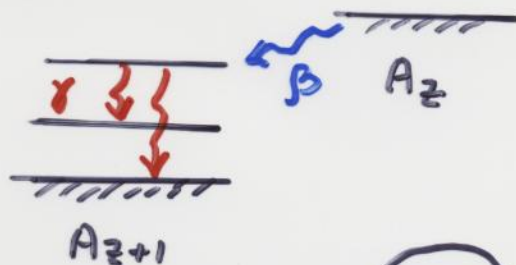
1914



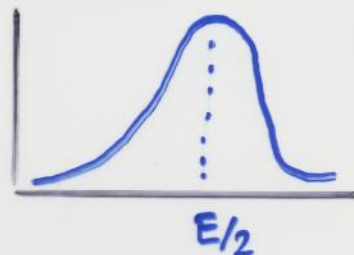
1920-23



1923-25

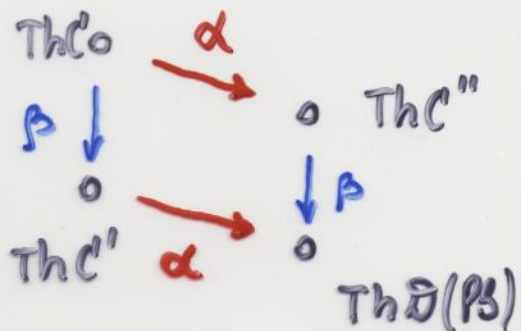
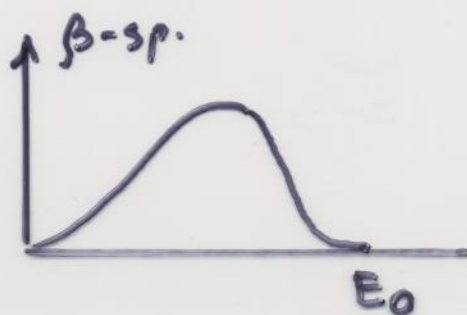


1927-29



RaE

1929-33



First experimental period (1899-20s) ^{6.}

(Radioact. families heavier Pb)

- β -particles are high-vel. electrons from nuclei
(N. Bohr)
- Exper. β -spectra consist of β -lines + contin. β
- β -lines are result of internal conversion of
nucl. γ -rays excited from atomic orbits (Rutherford)
- nuclei have excited states decayed through γ -rays
 β -rays preceded nuclear γ .
- Contin. β -spectrum is of primary origin and
 γ -rays can be absent (RaE). It has the β -sp.
upper limit.
- heat effect of β -rays shows its mean energy
that is in contradiction with energy cons. law.
However it can be restored if one uses in
 β - α chains the upper limits (ThC - ThD)
- nuclear structure is explained on p-e model base.

Thus it was formulated an experimental paradox

- observed violation of energy conservation law.

3.

Neutrino conception Birth (1929 - 37)

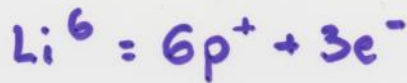
Dirac equation	Dirac, 28
Weyl equation	Weyl, 29
Contradictions in nuclear p-e model (odd-odd nuclei)	Kronig, 28
Nonconservation of energy	Bohr, 29
"neutron" Pauli hypothesis	Pauli, 30
Discovery of n	Chadwick, 32
Discovery of e^+	Anderson, 32
Proton-neutron nuclear model	Heisenberg, 32 Ivanenko, Majorana
β -decay theory with E-nonconserv.	Beck, Sitte, 33
Neutrino hypothesis	Pauli, 33
Fermi theory (in analogy with electrodynamics)	Fermi, 33
Proposal on ν mass estimation	Fermi, Perren, 33
Artificial radioactivity	Joliot, Curie, 34
Estimation of ν interaction $g \sim 4 \cdot 10^{-50}$ erg cm ³	Bethe, Peierls, 34
Double β -decay	Heppert-Mayer, 35
Recoil effect measurm. in β -decay	Leipunsky, 35
Generalization of Fermi theory	Gamow-Teller, 36
Majorana neutrino	Majorana, 37
Konopinsky-Uhlenbeck scheme	Konopinskiy, Uhlenbeck, 35
Conclusions of 30's beginning	

$$H\psi = i\hbar \frac{\partial}{\partial t} \psi$$

$$\gamma_\mu \partial_\mu \phi + m\phi = 0$$

Dirac

$m=0$ Weyl
P-nonconservation?



β -decay: $A(N, Z) \rightarrow A(N-1, Z+1) + e^- + \nu$ Pauli

$H_{ee} = e^2 (\bar{p} \gamma_\mu p) (\bar{e} \gamma_\mu e) \frac{1}{q^2}$
quant. electrodynamics

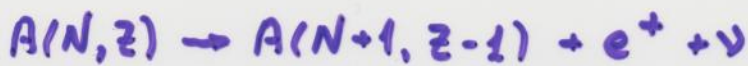
$e^2 = 1/137$

$H_\beta = G_F (\bar{p} \gamma_\mu n) (\bar{e} \gamma_\mu \nu)$

Fermi theory

$G_F m_e^2 \sim 10^{-10}$

nondetectable neutrino!



β^+ -decay



K'-capture



2β -decay

$H_\beta = G_F (\bar{p} O_i n) (\bar{e} O_i \nu)$

$O_i = 1 (S) \quad \gamma_\mu (V) \quad \sigma_{\mu\nu} (T) \quad i\gamma_\mu \gamma_5 (A) \quad \gamma_5 (PS)$

1937 . Beginning of Majorana Physics.

Ettore Majorana

"Symmetric Theory of Electron and Positron."

Nuovo Cim. v. 14, 1937, p. 171 (Phys. P. & N. v. 34, 2003 p. 124)
(in Russian p. 239)

main items:

$$\left(\frac{\hat{H}}{c} + (\vec{\alpha}, \hat{\vec{p}}) + \beta mc\right) \psi(\vec{x}, t) = 0$$

$$\alpha_x = \rho_1 \sigma_x, \quad \alpha_y = \rho_1 \sigma_y, \quad \alpha_z = \rho_1 \sigma_z, \quad \beta = \rho_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$



$$\alpha_x$$



$$\alpha_y^m = \rho_3$$



$$\alpha_z$$



$$\beta^m = -\rho_1 \sigma_y = \begin{pmatrix} 0 & (\sigma_y) \\ (\sigma_y) & 0 \end{pmatrix}$$

major.

$$\left(\frac{1}{c} \frac{\partial}{\partial t} - (\vec{\alpha}^m, \vec{\nabla}) - i \beta^m \mu\right) \psi(x, t) = 0$$

real real

$$\psi(\vec{x}, t) = U(\vec{x}, t) + i V(\vec{x}, t) = \text{Re} \psi + i \text{Im} \psi$$

$$\left(\frac{1}{c} \frac{\partial}{\partial t} - (\vec{\alpha}^m, \vec{\nabla}) + \beta^m \mu\right) U(\vec{x}, t) = 0, \quad \left(\frac{1}{c} \frac{\partial}{\partial t} - (\vec{\alpha}^m, \vec{\nabla}) + \beta^m \mu\right) V(\vec{x}, t) = 0$$

Two independent equations for new particles (truly neutral)

C-transformation: $\psi^{(c)}(\vec{x}, t) = \rho_3 \psi^*(\vec{x}, t) = \pm \psi^*(\vec{x}, t)$

$U(\vec{x}, t) :$ $\psi^*(\vec{x}, t) = \psi(\vec{x}, t) = \psi^{(c)}(\vec{x}, t)$

$i V(\vec{x}, t) :$ $\psi^*(\vec{x}, t) = -\psi(\vec{x}, t) = \psi^{(c)}(\vec{x}, t)$

majorana conditions: $\psi^{(c)}(\vec{x}, t) = \pm \psi(\vec{x}, t)$

for neutral particles: $\nu, n \text{ ?}$

Here is the beginning of Nonstandard Neutrino Properties!

First theoretical period (1929-37)

(on quant. - mech. picture base)

- The fundament of quant. - mech. picture for atomic phenomena (Schrodinger) was described

However:

- There was no nucl. phen. picture
 - There was no relativistic phen. picture
Dirac equation but no e^+
 - There was no β -decay picture
 - There was Weyl equation but with P -nonconserv.
- Two exits: Bohr's hypothesis on E -nonconserv on nuclear level
Pauli's hypothesis on new, neutral particle

Experim. discoveries of e^+ , n

Consequences: p - n nucl. model

relativ. quantum electrodynamics (QED)

neutrino Pauli hypoth.

Fermi β -decay theory in analogy with QED

Estimations of m_ν , G_F

Generalization of Fermi sch. to 5 variants

" " " by Konopinsky-Uhlenbeck

But Majorana hypoth. about neutr. particles $Q=0$

Thus, new particles were discovered and theory of weak processes ^{was} constructed, but

- what variant of theory?

- what is Weyl Eq. and Major. hyp.?

neutrino in decay processes (30's - 50's years)^{41.}

Nuclear β -decay:

β^+ -decay
K-capture
 2β in Majorana sch.
neutron star collapse
W-boson idea
solar neutrino (pp cycle)

Curie, Jolio, 34
Alvarez, 37
Furry, 39
Landau, 38
Klein, 39
Bethe, 39

μ -meson decay:

meson theory of nucl. forces
 μ -discovery
 μ -radioactivity

(Tamm) Yukawa, 35
Anderson, Nedemeyer, 38
Rasetti, Rossi, 41

μ -decay mechanism:

No $\mu \rightarrow e\gamma$
3-particle decay
 μ -e analogy
 μ -decay parameters

Hincks, Pontecorvo, Sard
Aethaus, Piccioni, 48
Anderson, Steinberg, Hincks, Pontecorvo, Zhdanov
Pontecorvo, Klein 48
michel' 50

Reactor neutrinos

First reactors
Reactor for neutrino search
n decay
 ν mass limit from T-decay
 ν registration
CE-AR experiment ($\nu \neq \bar{\nu}$)

Fermi '42 Kurchatov '46
Pontecorvo, Alvarez 46
Snell, Robson, Spivak 48
Curran, Hanna, Pontecorvo 49
Reines, Cowan 53
Devis 56-59

Neutrino from accelerators:

Pion flux
 $\pi \rightarrow \mu\nu$ decay
Strange particles in cosmic rays
 K^0 -duality, oscillations
 K_2^0 -discovery

Gardner, Latter '48
Latter, Orsini, Powell '47
Rochester, Butler '47
Gell-Mann, Pais, Piccioni, 53
Lande 56

Summaries to middle 50's:

Neutrino of 40's years

① $m_\nu < 500 \text{ eV (T)}$, $(\delta_{\mu\nu} p_\mu + m) \phi_\nu = 0$
Dirac ν

β^- $A(N, Z) \rightarrow A(N-1, Z+1) + e^- + \nu$

β^+ $A(N, Z) \rightarrow A(N+1, Z-1) + e^+ + \nu$

κ $A(N, Z) + e^- \rightarrow A(N+1, Z-1) + \nu$

$\pi \rightarrow \mu + \nu$, $\mu \rightarrow e + \nu + \nu$

$H_\beta = \underline{G_F} [\bar{p} \gamma_\mu n] (\bar{e} \gamma_\mu \nu) + \text{h.c.}$ (Fermi)

$H_\beta = \underline{G_F} \sum_i [\bar{p} O_i n] (\bar{e} O_i \nu) \cdot \lambda_i + \text{h.c.}$ (General form)

$O_i = 1 (S)$, $\gamma_\mu (V)$, $\sigma_{\mu\nu} (T)$, $\gamma_\mu \gamma_5 (A)$, $\gamma_5 (PS)$

Experim. data : S + T

$(ft)_\beta$, $e\nu$ -correl. , $\beta\gamma$ -correl.

Allowed and forbid. β -transitions.

Problems : ν was not observed ?

ν , $\tilde{\nu}$; or $\nu = \tilde{\nu}$?

Majorana ν

what is it ?

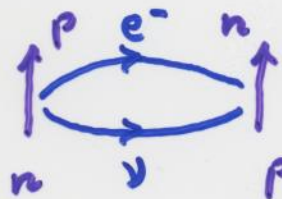
Weyl equation

Road to π -meson!

N-N forces (1932)



Tamm's hypothesis (1935)



Yukawa's hypothesis (1935)



$m_{\pi} \sim 100 \text{ MeV}$

Cosmic rays (1938) $\rightarrow \mu^-$ -showers

μ -radioactivity (1941) $\mu^- \rightarrow e^- + \text{neutr. part.}$

No $\mu \rightarrow e\gamma$ (1948) $\mu^- \not\rightarrow e^- + \gamma$

3-decay of μ (1948) $\mu^- \rightarrow e^- + 2\nu$

μ -e-analogy

π -fluxes: $p + p \rightarrow p + n + \pi^+$ (1948)

π -decay: $\pi^{\pm} \rightarrow \mu^{\pm} + \nu$ (1947)

$\pi^{\pm} \rightarrow e^{\pm} + \nu$

1946 Inverse β Process. β . Pontecorvo.

NRC of Canada, DoAE, Chalk River
Report PD-205, 1946

... Direct proof of the existence of neutrino... must be based on experiments, the interpretation of which does not require the law of conservation of energy, i.e. on experiments in which some characteristic process produced by **free neutrinos** (a process produced by neutrinos after they have been emitted in a disintegration) is observed.

$$\beta^- + Z \rightarrow \nu + (Z-1)$$

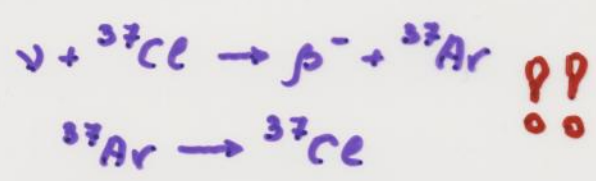
$$\nu + Z \rightarrow \beta^- + (Z-1) \quad , \quad \nu + Z \rightarrow \beta^+ + (Z-1)$$

$$\nu + Z + \beta^-(K') \rightarrow Z-1$$

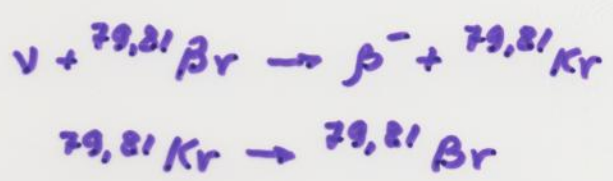
$$\nu + Z \rightarrow \gamma + \beta^- + (Z+1) \quad , \quad \nu + Z \rightarrow \gamma + \beta^+ + (Z-1)$$

Consequently, the radioactivity of the produced nucleus may be looked for as a proof of the inverse process.

An Example:



(34 days, K' -capture)



(34h; emission of positrons of 0.4 MeV)

Cross sections

$$\sigma_{inv} \leq \lambda^2 \cdot \frac{\lambda}{c} \cdot \frac{1}{Z}$$

$$(\sim 10^{-42} \text{ cm}^2 \quad E_\nu \approx 5 \text{ MeV})$$

Sources.

"The neutrino flux from the sun is of order of 10^{16} neutrinos $\text{cm}^{-2}\text{s}^{-2}$. The neutrinos emitted by the sun, however, are not very energetic. The use of high intensity piles permits two possible strong neutrino sources:

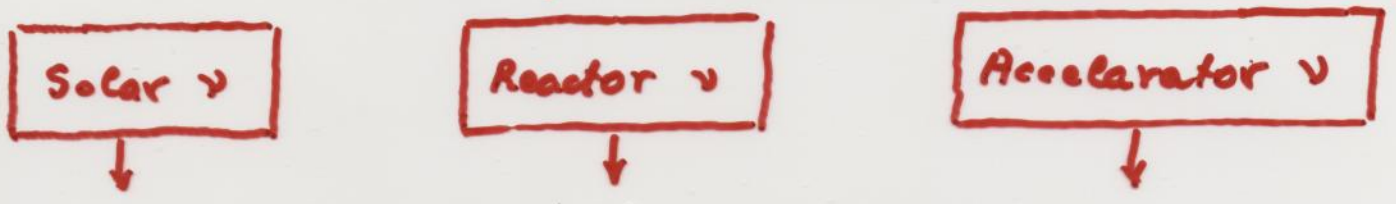
1. The neutrino source is the pile itself, **during operation...**
The advantage of such an arrangement is the possibility of using high energy neutrinos emitted by all the very short period fission fragments. Probably this is the most convenient neutrino source.
2. The neutrino source is the "hot" uranium metal extracted from a pile, or the fission fragment concentrate from "hot" uranium metal...

In the case of the investigation of inverse β processes produced by electrons of γ rays of high energy, the best source is a cyclotron or a synchrotron..."

Chalk River Laboratory
Chalk River, Canada

November 13, 1946

It is the first Neutrino Programme.



1947 Nuclear Capture of Mesons and the Meson Decay

B. Pontecorvo, NRC of Canada, Chalk River, Ontario, Canada
Phys. Rev., 1947, v. 72, p. 246

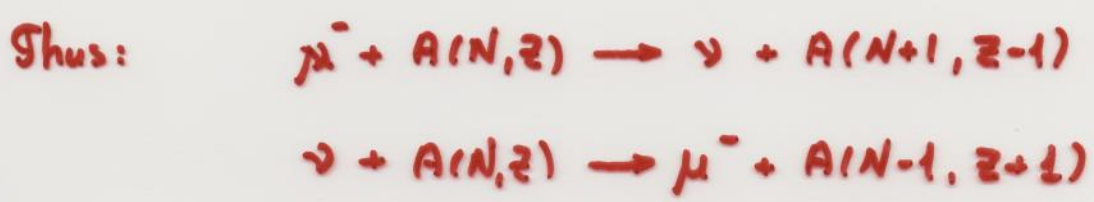
The experiment of Conversi, Pancini and Piccioni [1] indicates that the probability of capture of a meson by nuclei is much smaller than would be expected on the basis of Yukawa theory...

! We notice that the probability (about 10^6 sec^{-1}) of capture of a bound negative meson is of the order of the probability of ordinary K-capture processes, when allowance is made for the difference in the disintegration energy and the difference in the volumes of the K shell and of the meson orbit. We assume that this is significant and wish to discuss the possibility of a fundamental analogy between β processes and processes of emission or absorption of charge mesons... ?

In the hypothesis ... the process of nuclear absorption or production of a single meson would be accompanied by the emission of a neutrino.

June 21, 1947

[1] Conversi M., Pancini E., Piccioni O. // Phys. Rev., 1947, vol. 71, p. 209; see also Sigurgeirson T., Tamakawa A // Phys. Rev., 1947, vol. 71, p. 319



μ -e-decay

"The absorption of charged particles from the 2.2-microsecond meson decay." (with E.P. Hincks)

Phys. Rev., 1948, vol. 74, p. 697

"Search for gamma radiation in the 2.2-microsecond meson decay process" (with E.P. Hincks)

Phys. Rev., 1948, vol. 73, p. 257

"On the absence of photons among the decay products of the 2.2-microsecond meson" (with E.P. Hincks)

Phys. Rev., 1950, vol. 28 A, p. 29

"On the disintegration products of the 2.2-microsecond meson." (with E.P. Hincks)

Phys. Rev., 1950, vol. 77 p. 102

$$\mu \rightarrow e + \nu + \bar{\nu} \quad ; \quad \mu \rightarrow e + \gamma \quad (\text{No?})$$

 β -decay

"The neutrino and the recoil of nuclei in beta disintegration" Reports on Progress in Physics
London: The Physical Society, 1948, vol. 11, p. 32

"The β spectrum of ^3H " (with G.C. Hanna)

Phys. Rev., 1949, vol. 75, p. 983

$$m_\nu < 500 \text{ eV} ?$$

"Recent developments in proportional counter technique"

Helv. Phys. Acta, 1950, vol. 23, Suppl. 3 p. 97-118

18

Second experimental period (30's - 50's)

(wide investigations of decay processes)

- Nucl. physics was developed, new classes of β -processes were observed: β^+ , K -capture
 2β (th.), neutron β -decay
- New particles were discovered - mesons:
 π , K , K^0 , W -exch. meson hypoth., K^0 -oscill.
- μ -meson and its radioactivity was observed
3-part. decay with 2 neutrino participation
- New astrophysical concepts:
star neutron collapse
 p - p cycle in Sun
- Reactors as ν -sources $\nu + p \rightarrow n + e^+$
- ν experiments became reality:
 m_ν estimation from T -decay
Registration ν from reactor
 ν flux from accelerators
 $\sigma_\nu \sim E_\nu^2$

Thus neutrino became a real object
for physical investigations!

Neutrino theory and Standard Model Birth. (1953 - 1967)

19.

Lepton charge	Marx, Zeldovich, Konopinski-Mahmoud 53
Yang-Mills fields	Tang, Mills, 54
CPT-theorem	Lüders, Pauli, 54
CVC-hypothesis	Gershtein, Zeldovich, 55
P-nonconservation in weak processes	Lee, Tang 56 (Shapiro 55)
in β -decay	Wu, 57
in μ, π -decays	Garvin, Lederman, Weinitch Fridman, Telegdi, 57
CP-conservation	Landau, Lee, Yang, 57
2-component neutrino	Landau, Salam, Sakurai, 57
Universal V-A inter.	Gell-Mann, Feynman Sudarshan, Marshak, 57
ν -oscillation hypothesis	Pontecorvo, 57
PCAC hypothesis	Goldberger, Treiman 58
SU(3) symmetry	Gell-Mann, Levi, Okun' Kobzarev 58-63
Cabibo angle	Cabibo 63
CP-violation	Christenson, Cronin, Fitch, 64 Turley
Forming of Standard Model	
Electroweak interaction	Schwinger, 57
Calibration theory = gauge theory	Glashow, 61
Higgs mechanism	Higgs, 64
Nonintegral quark charges	Gell-Mann, Zweig, 64
Standard Model	Glashow, Salam, Weinberg t' Hooft, 71 61-67
SM renormalizability	
Weak interaction in atoms	Zeldovich, Khriplovich 59
Weak NN interaction	Blin-Stoep, Shapiro, Lovashev, 66 Abov, 64

1957 Pontecorvo returns to neutrino physics

(after T.A. Lee - C.N. Yang, 1957)

"Some remarks on slow processes of transformation of elementary particles" (with L.β. Okun)

Zh. Eksp. Teor. Fiz., 1957, vol. 32, p. 1527

Discussion on nonlepton weak processes and estimations of $K^0 \rightarrow \tilde{K}^0$ process by using $|ΔS| = 2$ mechanism

"Mesonium and antimesonium"

Zh. Eksp. Teor. Fiz., 1957, vol. 33, p. 549

"Gell-Mann and Pais [1] first pointed out the interesting consequence of the fact that K^0 and \tilde{K}^0 particles are not identical [2]. The possibility of $K^0 \rightarrow \tilde{K}^0$ transitions induced by weak interactions makes it necessary to consider neutral K meson as a mixture of particles of different (combined) [3] parity (K_1^0 and K_2^0).

We discuss here the problem as to whether there exist other "mixed" neutral particles (not necessary "elementary" ones) which are not identical to the corresponding antiparticles and for which **particle \rightleftharpoons antiparticle transitions are not strictly forbidden.** ... by conservation laws for the number of baryons and light fermions (... nuclear [4] and neutrino [5] charges).

$$(\mu^+e^-) \longrightarrow (\nu + \tilde{\nu}) \longrightarrow (\mu^-e^+)$$

... The law of conservation of neutrino charge is not yet definitely established: it is only established experimentally that neutrino and antineutrino are not identical particles [8]. **If the theory of two component neutrino [9] was not valid ... and if the conservation law for neutrino charge took no place, neutrino \rightarrow antineutrino oscillation in vacuum would be in principle possible ...**

1957

B. Pontecorvo's hypothesis on neutrino oscillation.

"Inverse β -processes and non-conservation of lepton charge."

JINR Preprint P-95, Dubna, 1957

"Not long ago the question was raised [1] as to whether there exist neutral particle mixtures, other than K^0 mesons [2], ... for which the transition **particle \rightarrow antiparticle** is not strictly forbidden...

It was noted that neutrino may be such a particle mixture and consequently that there is a possibility of **real transitions neutrino \rightarrow antineutrino in vacuum**, provided that the lepton (neutrino) charge [3] is not conserved...

Recently there came to our attention a paper by Davis [4], who investigated the production of ^{37}Ar from ^{37}Cl ... Below it is assumed that a) the neutrino (ν) and antineutrino ($\bar{\nu}$) emitted in the processes



are not identical particles; b) the neutrino charge is not strictly conserved ...



are possible, although ... they are less probable ...

It follows from a) and b) that neutrinos in vacuum can transform themselves into antineutrino and vice versa. This means that neutrino and antineutrino are particle mixtures, i.e. symmetrical and antisymmetrical combination of two truly neutral Majorana particles ν_1 and ν_2 having different combined parity [5] ...

... So, for example, a beam of neutral leptons from a reactor which at first consist mainly of antineutrinos will change in composition and at a certain distance R from the reactor will be composed of neutrino and antineutrino in equal quantities ...

... The upper limit of R which can give observable effects in the experiment of Cowan and Rines [6] is of order of a meter, which corresponds to a time for the transformation $\text{neutrino} \rightarrow \text{antineutrino}$ $T \leq 10^{-8} \text{ sec}$. If one takes into account that the neutrino energy ... is always larger by several orders of magnitude than $m_\nu c^2$... and that ... in laboratory system there is a considerable relativistic increase of the transformation time then the question arises as to whether the condition $T \leq 10^{-8} \text{ sec}$ is plausible ... The time T is [9] about $10^{-10} \times \frac{\text{neutrino energy}}{m_\nu c^2} \text{ sec}$ which is considerably greater than 10^{-8} sec ...

(P) ... In conclusion it is interesting to underline that, independently of the plausibility of the concrete effects which were discussed above, non-conservation of neutrino charge under the condition that neutrino and antineutrino are distinguishable entities (or which is the same, the existence of two Majorana neutrinos with different combined particles) inevitably leads to effects of Gell-mann - Pais - Piccioni type [2]. Under the above assumptions, effects of transformation of neutrino into antineutrino and vice versa may be unobservable in the laboratory because of large values of R , but will certainly occur, at least, on an astronomical scale... (P)

1957 .

B. Pontecorvo .

It is interesting to compare Pontecorvo's estimation with modern data:

$$L \sim 10^{-10} \frac{E_\nu}{m_\nu c^2} \text{ sec} \cdot 10^{10} \frac{\text{cm}}{\text{sec}} \sim \frac{E_\nu (\text{MeV}) \cdot 10^6}{10^{-2} \text{eV}} \text{ cm} \sim 10^8 \text{ cm} \sim 10^3 \text{ km}$$

for $E_\nu \sim 1 \text{ MeV}$

Modern Experiment $L \sim 10^2 \text{ km}$

1962 Neutrino mixing in Sakata models

1. "Possible unified models of elementary particles with two neutrinos"
T. Katayama, K. Matumoto, S. Tanaka, E. Yamada

Prog. Theor. Phys. v. 28, 1962, p. 675

2. "Remarks on the unified model of elementary particles."

Z. Maki, M. Nakayama and S. Sakata

Prog. Theor. Phys. v. 28, 1962, p. 870

Sakata - Nagoya model : p, n, Λ SU(3) hadron symmetry
(Sakata, 1956)

Gamba - Marshak - Okubo : ν, e, μ lepton-hadron symmetry
(1959)

↓ β^+ -matter

$$p = (\beta^+ \nu) \quad , \quad n = (\beta^+ e^-) \quad , \quad \Lambda = (\beta^+ \mu^-)$$

Cabibbo mixing: $J_h(x) = (\bar{n}p) + c(\bar{\Lambda}p) \rightarrow \cos\theta_c(\bar{n}p) + \sin\theta_c(\bar{\Lambda}p)$
 $G_\mu^2 = G_n^2 + G_\Lambda^2 \quad \quad n' = \cos\theta_c n + \sin\theta_c \Lambda$

Discovery of two neutrinos - ν_μ, ν_e (1962)

modification: $\nu_1 = \nu_e \cos\delta + \nu_\mu \sin\delta \quad , \quad \nu_2 = -\nu_e \sin\delta + \nu_\mu \cos\delta$

$$p = (\beta^+ \nu_1) \quad , \quad n = (\beta^+ e^-) \quad , \quad \Lambda = (\beta^+ \mu^-) \quad , \quad (\beta^+ \nu) - No!$$

$$\nu_e = \nu_1 \cos\delta - \nu_2 \sin\delta \quad , \quad \nu_\mu = \nu_1 \sin\delta + \nu_2 \cos\delta$$

"However, weak neutrinos are not stable due to the occurrence of virtual transformation $\nu_e \leftrightarrow \nu_\mu \dots$ "

In 1. two majorana neutrino case was also mentioned!

Thus: Hadron - Lepton Sym. Cabibbo angle hypothesis



Second theoretical period (1953-67) :
discovery of parity nonconservation and
Standard Model construction.

- There are divergent problem in QED
- Concepts of discrete operations (P, T, C) and lepton charge in opposite to Q are developed
- In K-decays θ - τ problem appeared that lead to nonconserv. parity hypothesis
- Experimental discov. of P-violation:
 - in nuclear β -processes
 - in particle decays: $\mu, \pi, K \dots$
 - hypothesis of weak inter. in nuclei (NN) atoms (Ne) appeared

Theory returns to Weyl and Fermi:

- 2 component neutrino
- CP-conservation
- V-A univ. interaction
- CVC, PCAC hypothesis

CP-violation experim. discovery

is
SM formulated: EW interaction, calibr. theory = gauge
Higgs' mech, SU(3) quark scheme, gluons, renormalis!

However, Pontecorvo neutrino oscil. hypothesis
in analogy with K^0 -osc. appeared
(then 2-flavour osc. version)

25.

Beginning of experimental neutrino physics, Statement of Standard Model (60's-70's)

ν in nuclei (ν_e)

ν polarization
 $\nu_e \neq \bar{\nu}_e$
Prop. on neutr. curr. (D, Li)
Geochemical exp. on 2β
 m_ν from T-decay ($< 35\text{eV}$)
 $\nu_e e$ in reactor ν flux
 $\nu_e \bar{D}$ in reactor ν flux

Goldhaber '58
Davis '59
Gaponov, Gershtein '62
64
Tretyakov et al. '76
Reines et al. '76
Reines et al. '79

ν from accelerator (ν_μ)

$\nu_\mu \neq \bar{\nu}_\mu$
magnetic furnace for ν fluxes
 ν in bubble chamber
 $\nu_\mu \neq \bar{\nu}_\mu$
Parton model
 $\nu_\mu e$
 $\nu_\mu N$ (neutr. curr.)
 J/ψ (c-quark)
 ψ (b-quark)
 τ -meson
W, Z - bosons discovery

Danby et al. '62
Vander Meer '63
Bloch '63
Bernardini '64
Bjorken, Paschos, Feynman '69
Masert et al. '73
Masert et al. '73
Ting, Richter et al. '74
Herb et al. '77
Perl et al. '75
CERN '84

ν in astrophysics

Hot Universe
Solar neutrino
Underground experim.
 ν sea in Universe
First undergr. detector ($\sim 8700\text{m w.p.}$)
 ν processes in stars
in n collapse
Relict γ radiation
Hyp. relict ν
Theory sol. ν experim
 ν Baksan telescope
Solar ν registration (10 years)

Gamow '46
Fowler '58
Markov, Greison '60
Pontecorvo, Smorodinsky '61
Reines '63
Fauler, Koye '65
Zeldovich '65
Penzias, Wilson '65
Dick, Zeldovich, Novikov
Bahcall '72
Chudakov '77
Davis '78

Summary up to 80's : ν fluxes and ν detectors

Neutrino of 60's

$\nu_e, \tilde{\nu}_e; \nu_\mu, \tilde{\nu}_\mu$ - two types of ν .

$(\gamma_\mu p_\mu) \phi_\nu = 0$ Weyl equation

β^\pm -decay, K -capture, 2β -decay (estimations)

π, K -decays, $\mu^- \rightarrow e^- + \tilde{\nu}_e + \nu_\mu$

Universal weak interaction?

$H_W = \frac{G_F}{\sqrt{2}} J_\mu^+ J_\mu^-$ $J_\mu^h = (\bar{p} \gamma_\mu (1 - \lambda \gamma_5) n)$

$J_\mu^e = (\bar{e} \gamma_\mu (1 + \gamma_5) \nu)$

$G_\beta, G_S, G_\mu^2 = G_\beta^2 + G_S^2$ - empirical rule.

" $G_\mu \cos \theta_c, G_\mu \sin \theta_c$

Experim. data: $V-A, \lambda \approx 1.25$

$(ft)_{0-0}, (ft)_{n,GT}, (ev), (\bar{J}e),$

$(\bar{J}\nu), \bar{J}[\bar{e} \times \bar{\nu}], \beta \gamma_{circ.} - corr., \pi \rightarrow \mu \nu_\mu$

ν -fluxes: $\tilde{\nu}_e, \nu_\mu$

Neutral curr: $\nu + e \rightarrow \nu' + e', \nu + N \rightarrow \nu' + N' ?$
 NN -weak int., Ne -weak inter.

Problems: Neutral currents?
Majorana ν ?

Weak Interaction Hamiltonian

(Modern data) Yu. Gaponov, Ya.F. (2000)

General Theory:

$$H_{\beta} = (G_{F\beta}/\sqrt{2}) \sum_i (\bar{\Psi}_p O_i \Psi_n) (\bar{\Psi}_e O^i (C_i + C'_i \gamma_5) \Psi_{\nu e})$$

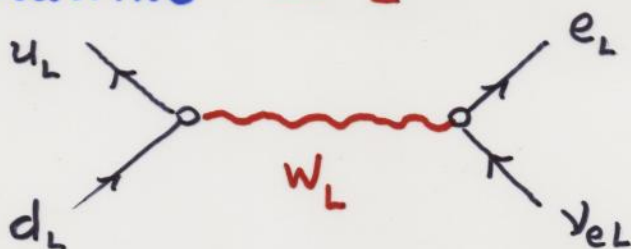
It includes 15 parameters (8 real ones) due to

$$O_i = 1 (S), \gamma_{\mu} (V), \sigma_{\mu\nu} = 1/(2i) (\gamma_{\mu} \gamma_{\nu} - \gamma_{\nu} \gamma_{\mu}) (T), i\gamma_{\mu} \gamma_5 (A)$$

contributions

Theor. Hypotheses:

1. CP - conservation $\rightarrow 8$
2. V, A - currents (W_L - boson) $\rightarrow 4$
3. L - neutrino $\rightarrow 2$



Standard Model:

$$H_{SM} = (G_{\nu\beta}/\sqrt{2}) (\bar{\Psi}_p \gamma_{\mu} (1 - \lambda \gamma_5) \Psi_n) (\bar{\Psi}_e \gamma^{\mu} (1 + \gamma_5) \Psi_{\nu e})$$

$$G_{\nu\beta} = 1.4183(18) \cdot 10^{-62} \text{ J m}^3$$

$$\lambda = -1.2673(10)$$

Neutron experiments

n-decay probab: $dW(E_e, z, \Omega_e, \Omega_\nu) =$

$$G_{F\beta}^2 / (2\pi)^5 F(\pm z, E_e) (E_e - E_{e0})^2 E_e p_e dE_e d\Omega_e d\Omega_\nu$$

$$(1 + (B m_e / E_e) \sqrt{1 - \alpha^2 z^2} + \alpha (\bar{p}_e \bar{p}_\nu) / E_e E_\nu + A (\bar{J} \bar{p}_e) / E_e + \beta (\bar{J} \bar{p}_\nu) / E_\nu$$

$$+ \mathcal{D} (\bar{J} [\bar{p}_e \times \bar{p}_\nu]) / E_e E_\nu + \underbrace{E (\bar{J} [\bar{\sigma}_e \times \bar{p}_\nu]) / E_\nu}_{\text{R-param.}}$$

$b = 0$ Fierz terms absence

Realized Experiments:

$(f\tau)_n$	n life-time
a	$e\nu$ -correlation
A	eJ_n -correlation
B	νJ_n -correlation
\mathcal{D}	triple $e\nu J$ -correlation

Non realized Experiments:

$\vec{\sigma}_e$ e-polar. measur. (or \vec{P}_γ - γ -polar.)

$\vec{\sigma}_p$ p-polar. measur.

$E(R)$ triple correl.

$0^+ - 0^+$ nucl. transitions:

$$(ft)_{0-0} (1 + \delta_R) (1 - \delta_C) = (Ft)_{0-0} = K \ln 2 / 2 G_{V0-0}^2$$

$$G_{V0-0} = 1.4173(11) \cdot 10^{-62} \text{ Jm}^3$$

n decay in Standard model

$$f_n \tau_n = K' / (G_{\nu\beta}^2 (1 + 3\lambda^2))$$

$$f_n = 1.71465(14)$$

$$89(2)$$

$$A = -2(\lambda + \lambda^2) / (1 + 3\lambda^2)$$

$$\beta = -2(\lambda - \lambda^2) / (1 + 3\lambda^2)$$

$$a = (1 - \lambda^2) / (1 + 3\lambda^2)$$

(τ_n)

$$\tau_n = 885.7(10) \text{ s}$$

From 0-0 ; τ_n data

$$\lambda = -1.2675(10)$$

(A)

$$\bar{A} = -0.1161(7)$$

From A data

$$\lambda = -1.2664(19)$$

Exp. discr.

(B)

Old: $\beta = +0.9820(40)$

New: $(A - \beta) / (A + \beta) = -1.2686(49) = \lambda$

$$\beta = +0.9876(4) \text{ in SM}$$

(a)

$$a = -0.1017(51)$$

From a

data

$$\lambda = -1.2591(168)$$

$$\text{Mean value } \lambda = -1.2673(9)$$

Third experimental period (60's - 70's)

↳ physics development and SM statement.

3 main directions in ν physics:

Low energy neutrino (β -decay and reactor ν)

- ν properties - m_ν , P_ν
- improvement of $\nu_e \neq \tilde{\nu}_e$
- investigation of $\nu_e \bar{\nu}_e$, $\nu_e e$ processes

Extremely low cross sections!

High energy neutrino (accelerators)

- new methods of ν registration (Bubble ch.)
- improvement of $\nu_e \neq \nu_\mu$, $\nu_\mu \neq \tilde{\nu}_\mu$, ν_τ from τ -decay

3 fam. of ν

- investigations of $\nu_\mu N$, $\nu_\mu e$ processes
- new quarks discoveries - c , b , t

SM experim. improvement

- neutral processes
- W , Z - bosons

Neutrino astrophysics:

↳ - Solar, Star, Collapse, ν -sea

Solar ν experiments - first plans and Experim. underground Labs, ν -telescopes

Thus, our knowledges on ν properties widen up to ν -scattering, ν was investigated in SM frames and is now a new instrument. ν fluxes, ν detectors appeared. Special role ν in Universe was discovered.

Standard model (GSW-theory)

1. Leptons: Massless \rightarrow Massive

$$\begin{matrix} \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L & \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}_L & \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}_L \\ (e)_R & (\mu)_R & (\tau)_R \end{matrix} \quad m_\nu - ?$$

2. Quarks: Massless \rightarrow Massive

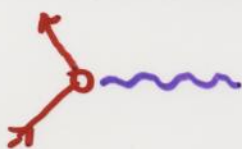
$$\begin{matrix} \begin{pmatrix} u \\ d \end{pmatrix}_L & \begin{pmatrix} c \\ s \end{pmatrix}_L & \begin{pmatrix} t \\ b \end{pmatrix}_L \\ (u)_R & (c)_R & (t)_R \\ (d)_R & (s)_R & (b)_R \end{matrix} \quad m_q - ?$$

3. Quark mixing:

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} \text{Kobayashi} \\ \text{Masakawa} \\ \text{matrix} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix} = \begin{pmatrix} d \cos \theta_c + s \sin \theta_c + \dots \\ -d \sin \theta_c + s \cos \theta_c + \dots \\ \dots \end{pmatrix}$$

Neutrino mixing - ?

4. Gauge fields:



$$\gamma, Z^0, W^\pm; G_{dp} \text{ (8 colour pairs)}$$

Weak + EM int $\rightarrow \gamma, Z^0$, Inter. constants
 $\nu \gamma$ -int. = 0 $g, g'; \sin \theta_w$

5. Higgs boson + spont. broken vacuum
 Constr. Mass mechanism ?

6. Fundamental symmetry:

$$SU(3) \times \underbrace{SU(2) \times U(1)}_{\text{electroweak}}$$

In search for an exit beyond Standard model Theory (80's years)

Modern conception of the Standard Model

Theories of Grand Unification

Unification of interactions: EM, Weak, Strong, Gravity...

Left - Right Handed schemes

Lepto - Quark schemes

Supersymmetrical theories

Superstrings

nonstandard neutrino properties

majorana models

$\nu_e \leftrightarrow \nu_\mu \leftrightarrow \nu_\tau$ oscillations

Neutrino oscillation in matter (MSW-mech.)

neutrino anomalous magnetic moment

neutrino precession (VVO-mech.)

Neutrino flavor precession (LMA-mech.)

Astrophysics and cosmology

Solar neutrino

neutrino from collapsed stars

Origin of elements in star bursts

Dark matter in the Universe

MSW - Mikheev, Smirnov, Wolfenstein

VVO - Voloshin, Vysotsky, Okun

LMA - Livo, Marsiano, Akhmedov

Experiment in search for exit
beyond the Standard Model
(80-90's years)

Supernova SN-1987A

Reactor ν

- ν_p, ν_e, ν_D - experiments
- ν magnetic moment
- ν_e - oscillations
- ν - diagnostics

Solar ν

Ce-Ar, Kamiokande

ν_\odot - deficit ?

SAGE - GALLEX

Super-Kamiokande

SNO

Borexino etc.

? $\nu_e + \nu_{\mu, \tau}$ 2001

Acceleration ν

High E_ν

LSND ?

Measurement on great distance

ν from Earth

Atmospheric ν

$\nu_\mu \leftrightarrow \nu_\tau$ oscillations (98-99) ?

β -decay

Right-handed currents

Lepto-quark bosons

CP-violation

ν mass from T

2β decay

$2\beta (0\nu)$ - decay from gram to tons

Search for dark matter

Solar neutrino physics

Standard solar model

Radius	696 000 km
Surface T°	5773 °K
Central T°	$15.6 \cdot 10^6$ °K
Contents of H	34.1 %
Contents of He	63.9 %
Rest elements	1.96 %
Central density	148 g/cm ³

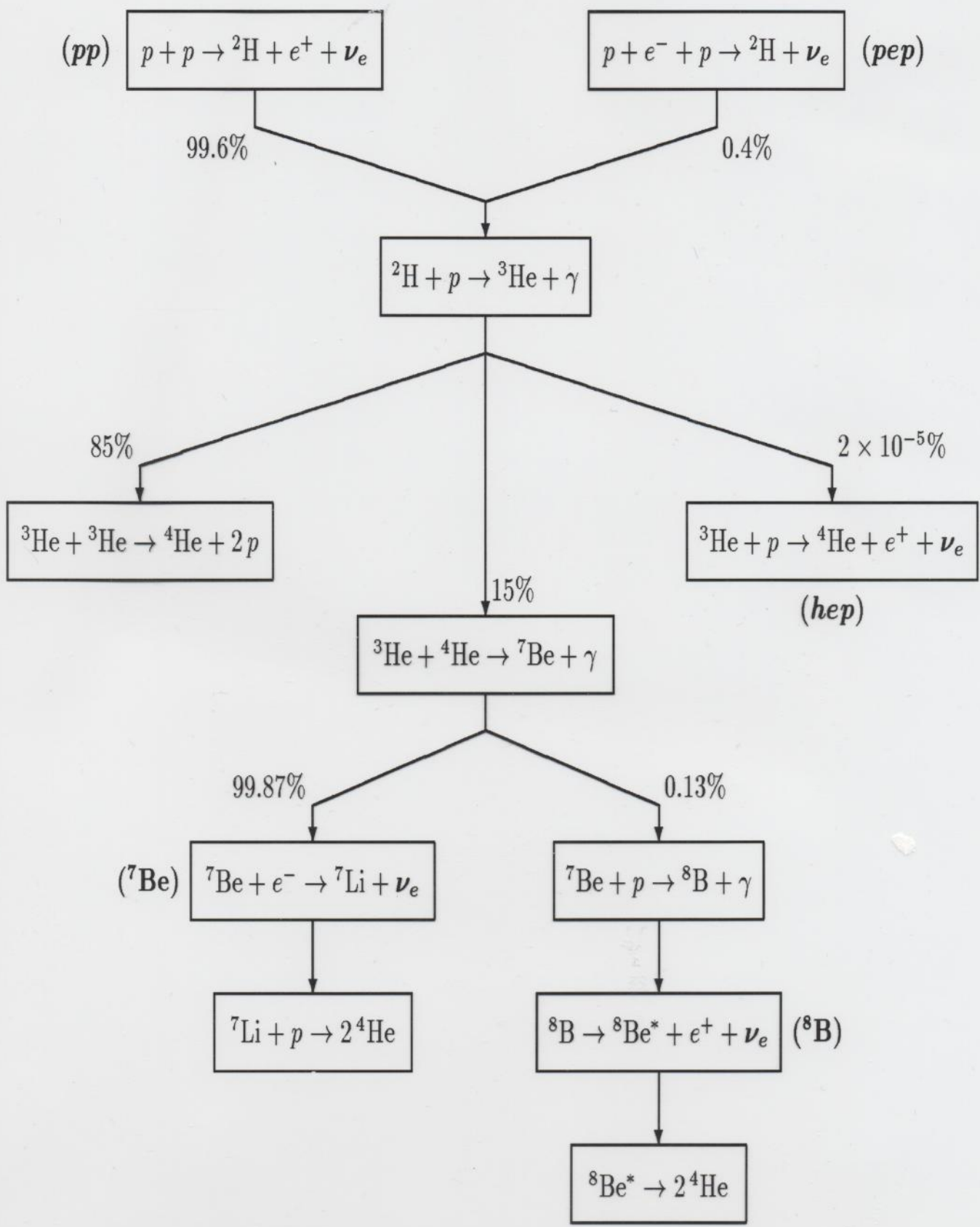
Neutrino fluxes on the Earth

Source	Reaction	E (MeV)	Φ (cm ⁻² s ⁻¹)
pp	$pp \rightarrow D e^+ \gamma_e$	0.27 ÷ 0.42	$5.94 \cdot 10^{10}$
pep	$pep \rightarrow D \gamma_e$	1.445	$1.4 \cdot 10^8$
${}^7\text{Be}$	$e {}^7\text{Be} \rightarrow {}^7\text{Li} \nu_e$	0.385	$4.8 \cdot 10^9$
${}^8\text{B}$	${}^8\text{B} \rightarrow {}^8\text{Be}^* e^+ \nu_e$	0.865 6.73 → 15	$5.15 \cdot 10^6$

ν -oscill. mechanism hypothesis

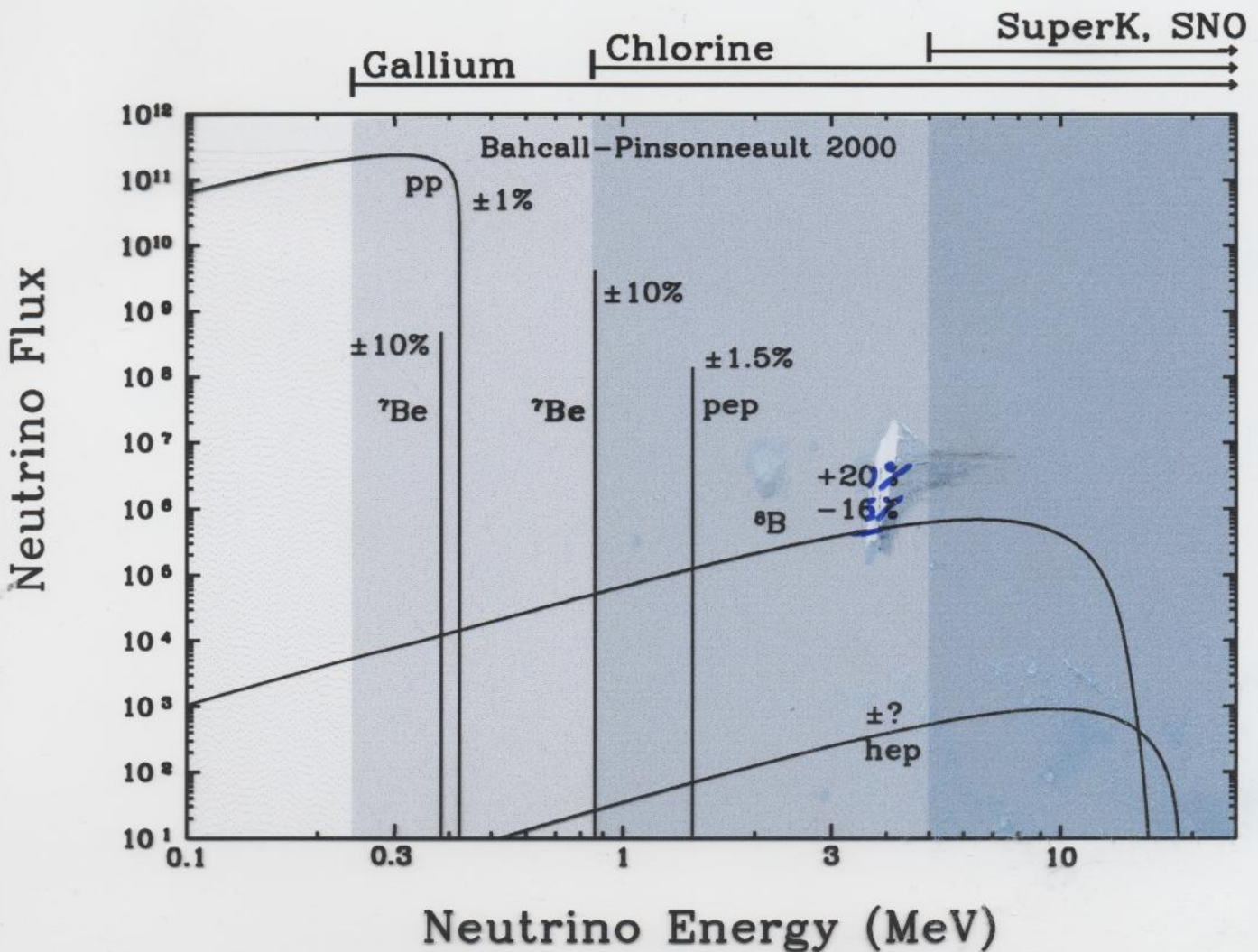
Resonance	MSW
Precession of μ_ν	VVO
Spin-flavor precession	LMA

SSM - pp cycle



SSM - neutrinos

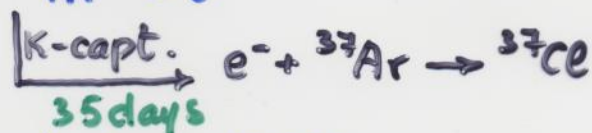
Source	Reaction	$E_{\nu}^{Av.} (MeV)$	$E_{\nu}^{Max.} (MeV)$
pp	$p + p \rightarrow d + e^{+} + \nu_e$	0.2668	0.423
pep	$p + e^{-} + p \rightarrow d + \nu_e$	1.445	1.445
${}^7\text{Be}$	$e^{-} + {}^7\text{Be} \rightarrow {}^7\text{Li} + \nu_e$	0.3855	0.3855
${}^8\text{B}$	${}^8\text{B} \rightarrow {}^8\text{Be}^{*} + e^{+} + \nu_e$	0.8631	0.8631
${}^8\text{B}$	${}^8\text{B} \rightarrow {}^8\text{Be}^{*} + e^{+} + \nu_e$	6.735	~ 15
hep	${}^3\text{He} + p \rightarrow {}^4\text{He} + e^{+} + \nu_e$	9.628	18.778
${}^{13}\text{N}$	${}^{13}\text{N} \rightarrow {}^{13}\text{C} + e^{+} + \nu_e$	0.7063	1.1982
${}^{15}\text{O}$	${}^{15}\text{O} \rightarrow {}^{15}\text{N} + e^{+} + \nu_e$	0.9964	1.7317
${}^{17}\text{F}$	${}^{17}\text{F} \rightarrow {}^{17}\text{O} + e^{+} + \nu_e$	0.9977	1.7364



Solar neutrino detectors of first generation.

$$1 \text{ SNU} = \frac{1}{10^{36}} \text{ event/s}$$

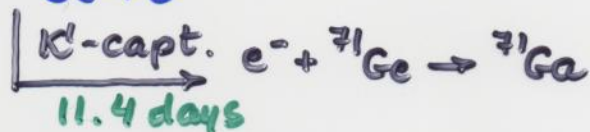
1. Chlorine Argon method (B. Pontecorvo, 1946)



Homestake, R. Davis et al. 4100 m.w.e.

610 t C_2Cl_4 ${}^{37}\text{Cl}$ (24%)

2. Gallium Germanium method (V. Kuzmin, 1966)



SAGE, G. Zatsepin, V. Gavrin et al. 4700 m.w.e.

60 t natur. Ga ${}^{71}\text{Ga}$ (39.9%)

Effect. 0.95(12)

${}^{51}\text{Cr}$ (513gr)

0.5 μCi

GALEX Italy-Germ. Coll. Gran-Sasso 3400 m.w.e.

↓ 30 t GaCl_3 (120) Effect. 0.92(8) ${}^{51}\text{Cr}$ (1.7 μCi)

GNO Gran-Sasso

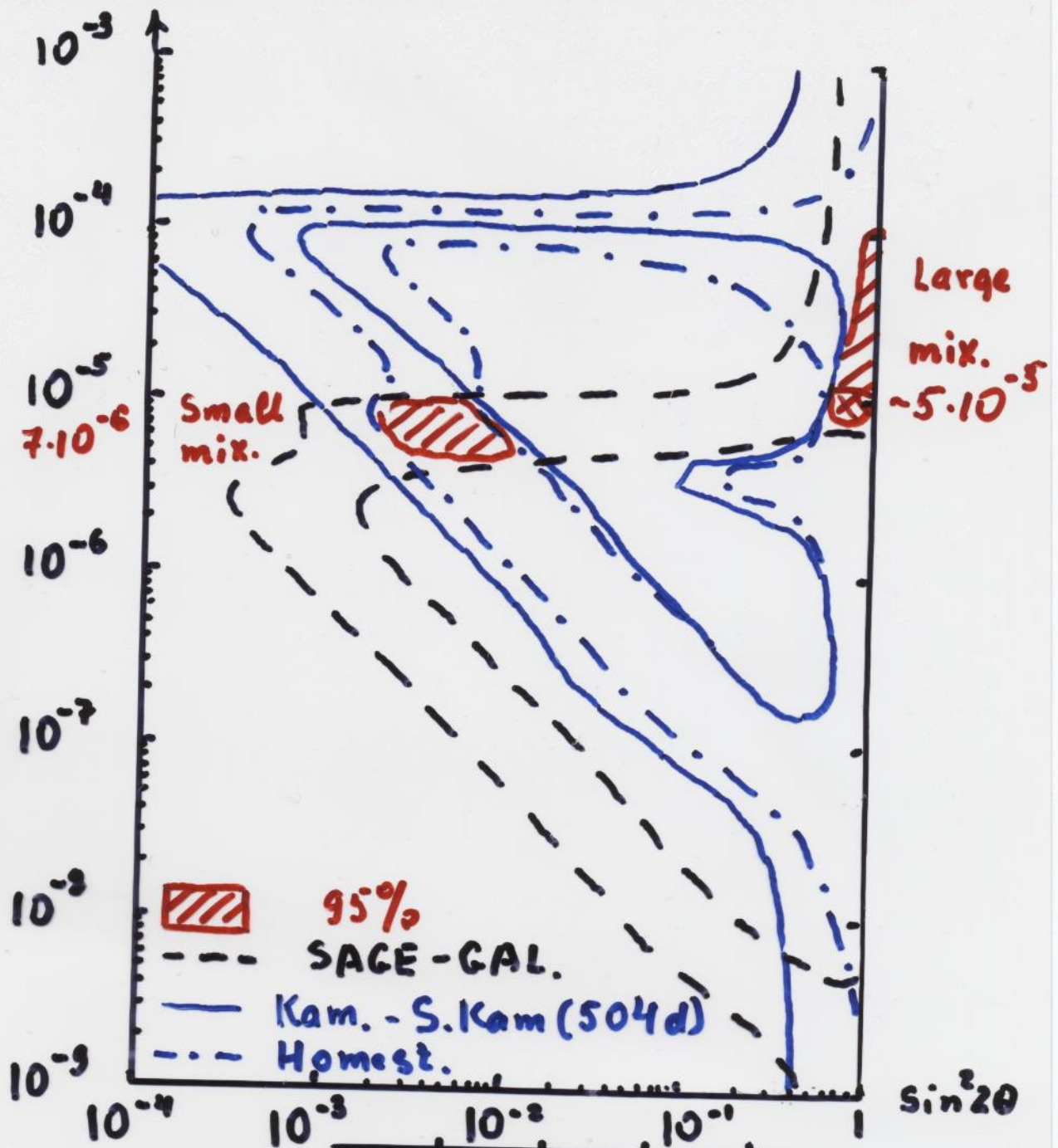
3. Kamiokande, Super Kamiokande



SK 50000 t H_2O (32000 t) 11000 Phm

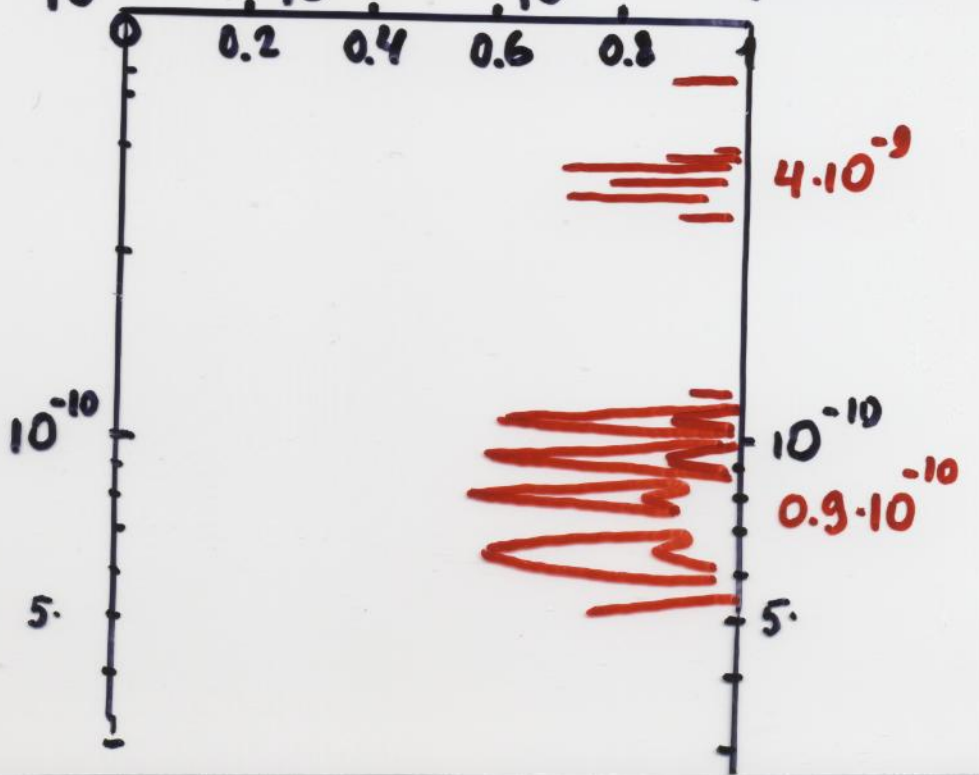
Solar γ -osc.

$\Delta m^2 (eV^2)$



MSW

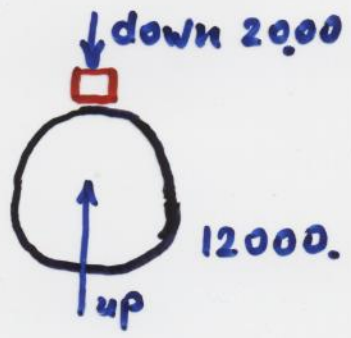
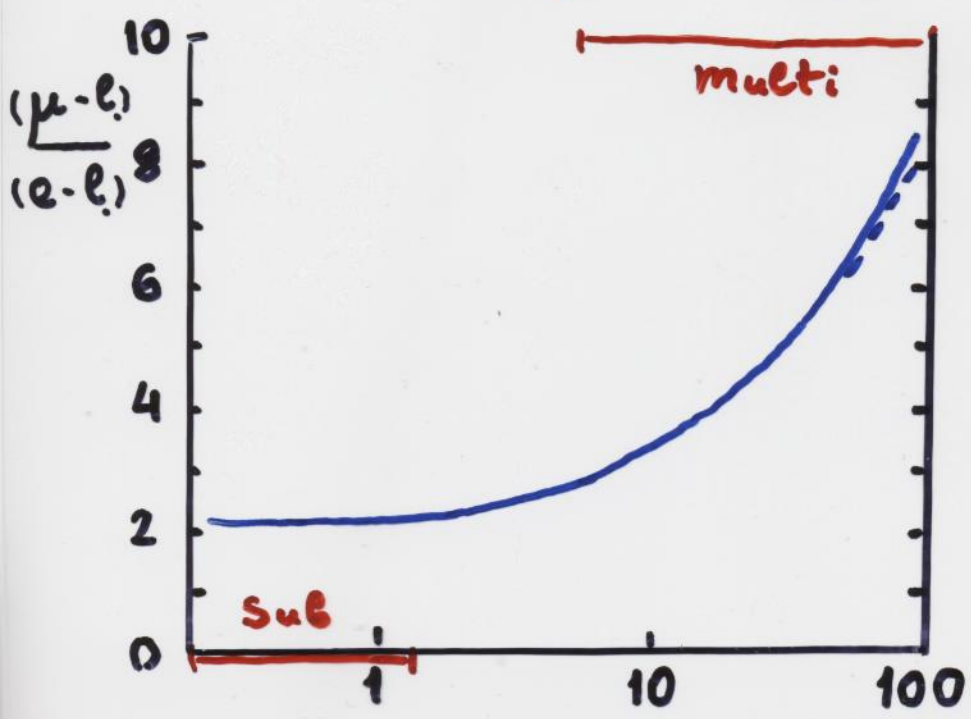
H.-Lang.



Atmospheric ν

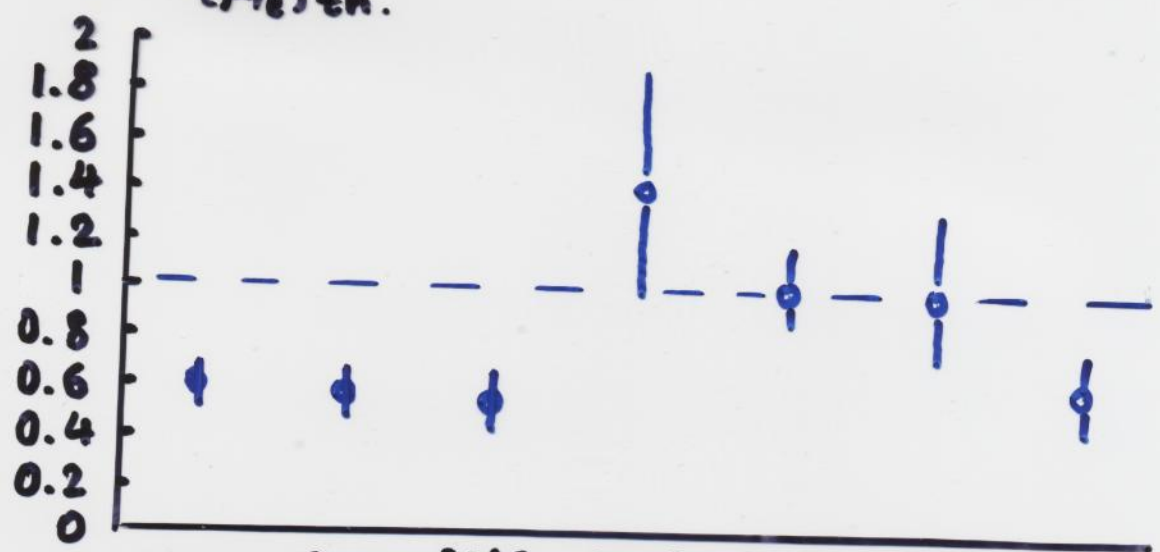


exp.
 $\frac{(\mu\text{-like})/(\nu\text{-like})}{(\mu\text{-like})/(\nu\text{-like})} \neq 1$
 theor.



UP/Down $\neq 1$
 E_{ν} (GeV)

$\frac{(\mu/e)_{exp.}}{(\mu/e)_{th.}}$



Kam	Kam	IMP	IMP	Frej.	Neusey	SoudanII
0.60	0.57	0.54	1.40	1.00	0.99	0.61
± 0.06	± 0.08	± 0.05	± 0.35	± 0.15	± 0.30	± 0.15
± 0.05	± 0.07	± 0.12	± 0.3	± 0.08	-	± 0.05
sub	mult	sub	mult			

Super-Kamiokande (535 days)

		Monte-Carlo	
		Data	Theory
Sub GEV	1 Ring e-like	1231	1049
	μ-like	1158	1574
	Mul. Ring	911	981
Multi: GEV	1 Ring e-like	290	237
	μ-like	230	298
	Mul. Ring	533	560

$$\frac{(N/e)_D}{(N/e)_{Th.}}$$

st sys
↓ +th.

$$0.63 \pm 0.026 \pm 0.05$$

$$0.65 \pm 0.05 \pm 0.08$$

Up/Down

$$0.93 \pm 0.13$$

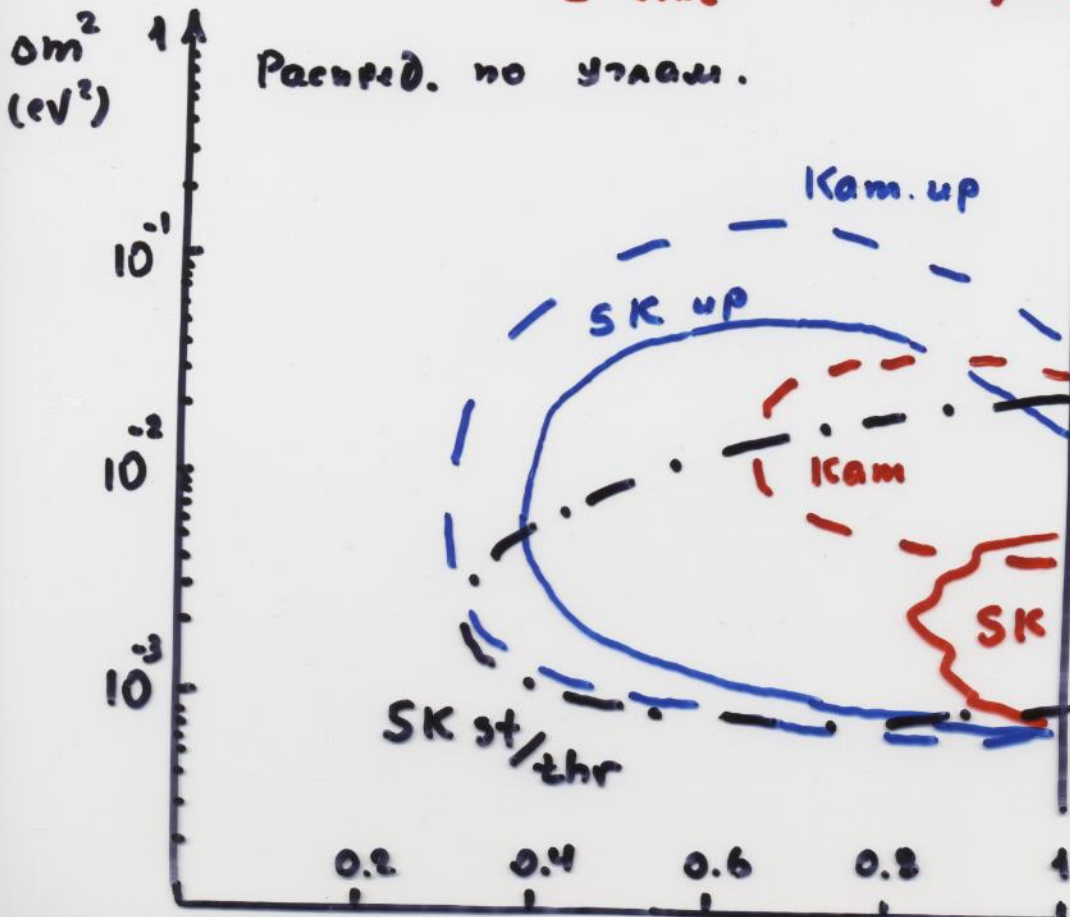
e-like

$$0.54 \pm 0.06$$

μ-like

Δm^2
(eV^2)

Рассред. по углам.



$$\sin^2 2\theta > 0.8$$

$$\Delta m^2 \sim 10^{-3} \div \sim 10^{-2} eV^2$$

$$\nu_\mu \rightarrow \nu_\tau$$

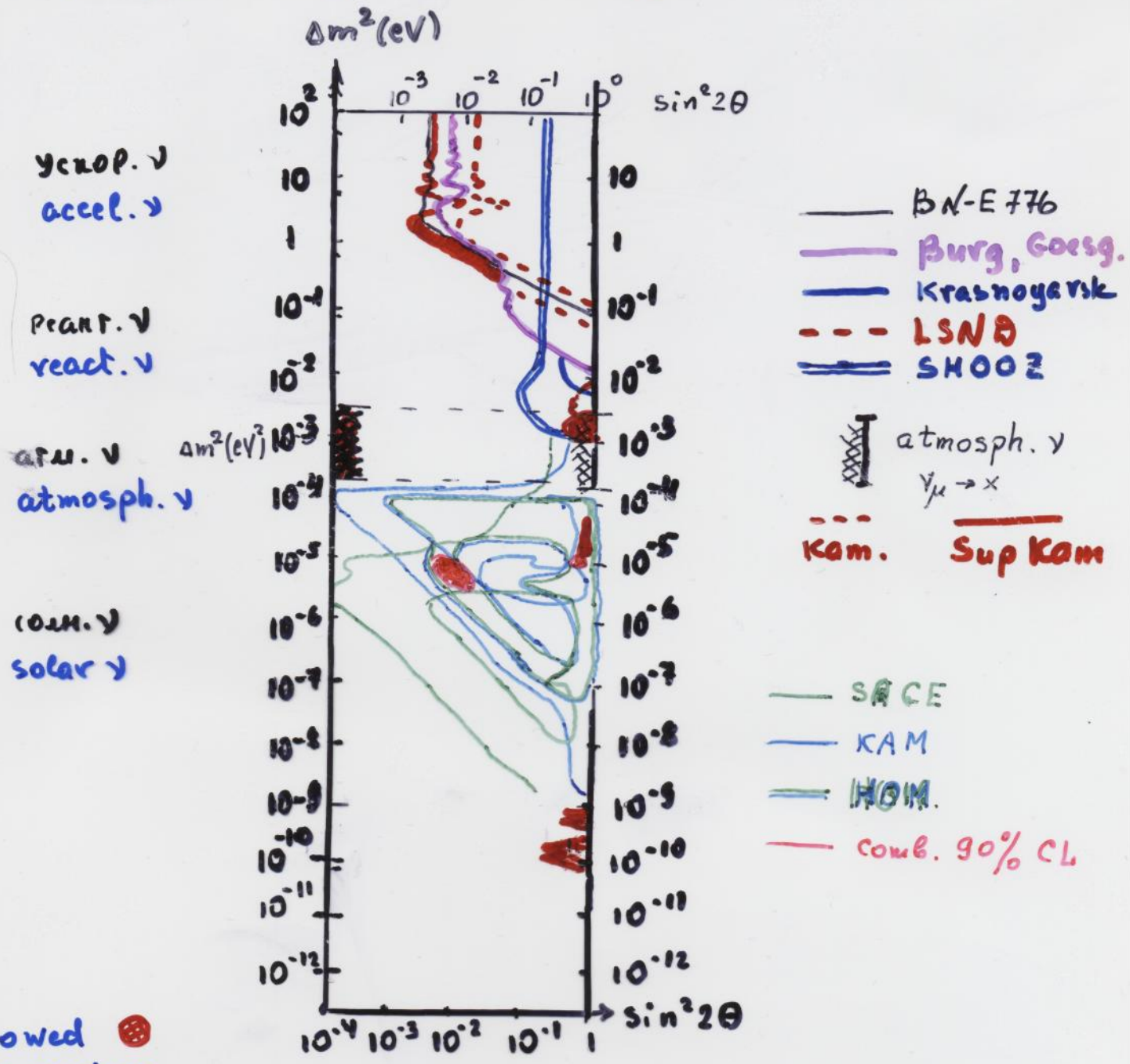
or

$$\nu_\mu \rightarrow \nu_{st.}$$

$\sin^2 2\theta$

Experimental results of Ga-Ge det. (sol. ν)
 γ-accel. experiments and reactor ν-experiments.
 Atmospheric ν-anomaly.

$$P_{\nu_1 \rightarrow \nu_2}(r) = \frac{1}{2} \sin^2(2\theta) \sin^2\left(\frac{r}{E}\right) |m_1^2 - m_2^2|$$



allowed osc. regions

Однородные области ν-осцил.

LSND
 SKamioke.
 MSW -
 Sun "Just so"

$$5 \div 0.5 \quad 10^{-2}$$

$$10^{-2} - 10^{-3} \quad > 0.8$$

$$5 \cdot 10^{-5} \quad 0.85 \quad 7 \cdot 10^{-6} \quad 10^{-2}$$

$$(99\% 4) 10^{-10} \quad > 0.5$$

- sin² 2θ

New Projects (II gener.)

+ SuperKamiokande: 50 000 T H₂O 11 000
(32 000 T)

Aims: Solar ν , SN explosions

+ Kam Land: 1000 T liq. sc. 1300

Aims: react. ν ~ 200 km , solar ν
atm. ν , SN expl. , $\bar{\nu}_e$ from Earth
2 β -decay (Xe)

+ SNO (Sudbury) 1000 T D₂O 9500

Aims: ν_e , $\nu_{\bar{\mu}}$, solar ν ,
atm. ν , SN expl. , $n\bar{n}$

Borexino (Italy) 300 T liq. sc. 2000

Pilot 5 T " " 100

Aims: Sol. ν ^7Be , atm. ν , $\bar{\nu}_e$ from Earth
2 β -dec. Gol, TB, Se

Long dist. ν experim.

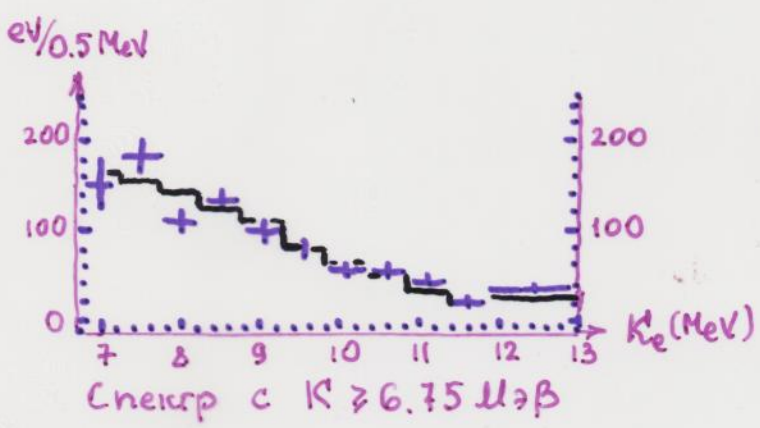
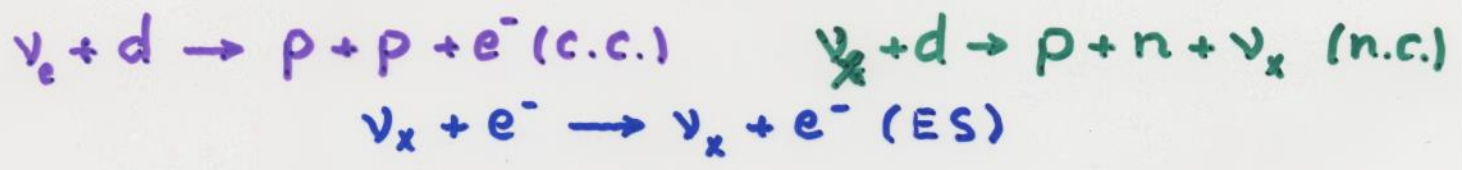
Artif. ν sources

ν magn. moment

^{51}Cr , Pm

$\$NO$ (Canada), 2001

Первые результаты (First results of collab.)



ch. channel

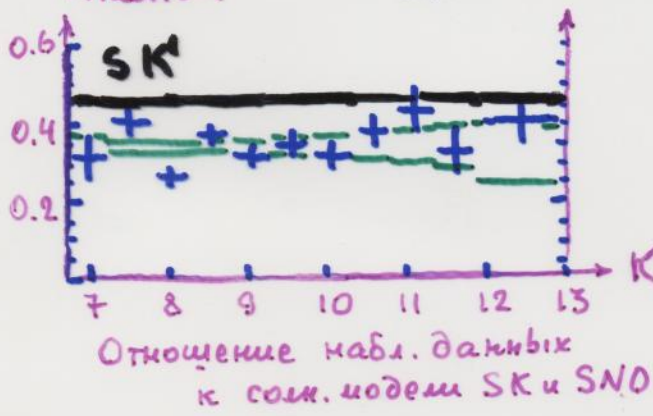
$$\Phi^{CC}(\nu_e) = 1.75 \pm 0.07 \pm 0.12 \pm 0.05 \times 10^6 \text{ cm}^{-2} \text{ сек}^{-1}$$

st. sys th.

$$\Phi^{ES}(\nu_x) = 2.39 \pm 0.34 \pm 0.16 \pm 0.14 \times 10^6 \text{ cm}^{-2} \text{ сек}^{-1}$$

$$\Phi_{SK}^{ES}(\nu_x) = 2.32 \pm 0.03 \pm 0.08 \pm 0.07 \times 10^6 \text{ cm}^{-2} \text{ сек}^{-1}$$

Data/theor. SM



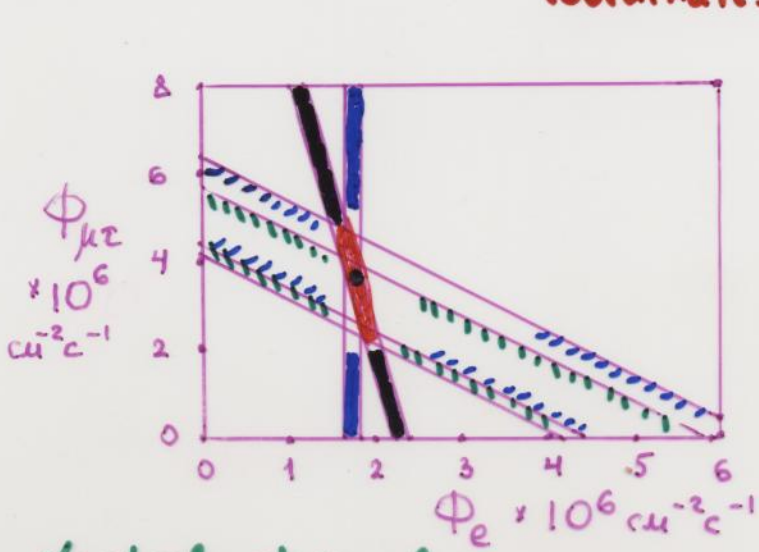
νe - Neutr. channel

$$\Phi(\nu_{\mu\tau}) = 3.69 \pm 1.13 \cdot 10^6 \text{ cm}^{-2} \text{ сек}^{-1}$$

$$\Phi(\nu_e) = 1.75 \pm 0.14 \cdot 10^6$$

$$\Phi(\nu_x) = 5.44 \pm 0.99 \cdot 10^6 \text{ cm}^{-2} \text{ сек}^{-1}$$

согласует с SM coordinates with SM



..... $\Phi^{CC}(\nu_e)$

..... $\Phi_{SK}^{ES}(\nu_x)$

..... SM (ν)

..... $\Phi(\nu_x)$

..... Paspen. (10) paion allowed region

νd - Neutral channel

Нейтрон. канал: $\Phi^{NC}(\nu_x) = 5.09 \pm 0.4 \pm 0.46 \pm 0.43 \times 10^6 \text{ cm}^{-2} \text{ сек}^{-1}$

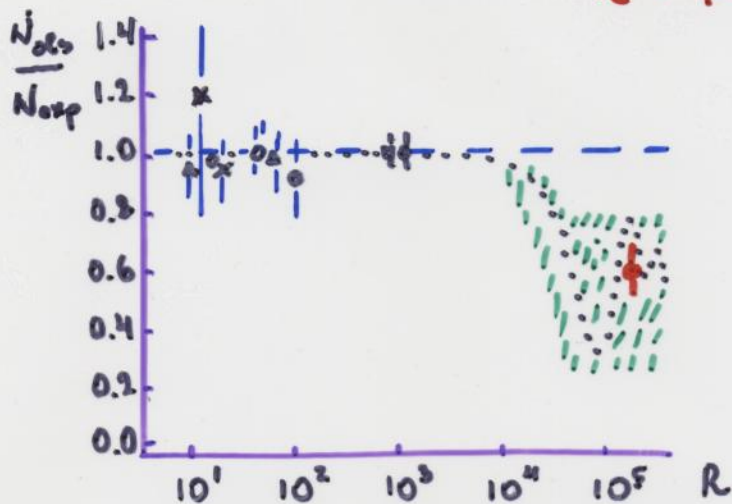
$\nu_x d$

$$4.90 \pm 0.24 \pm 0.29 \times 10^6 \text{ cm}^{-2} \text{ сек}^{-1}$$

KamLAND Data (2003)



162 t

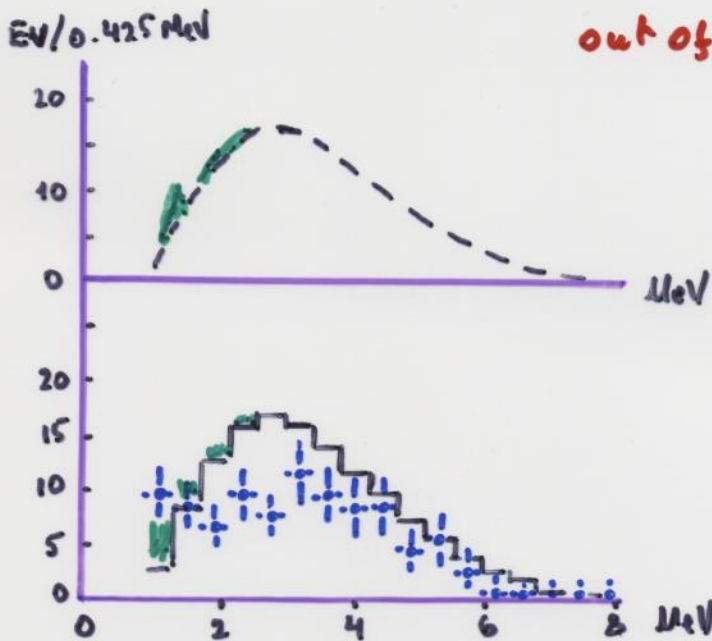


ILL, Savannah River, Bugey, Revno
 Goesgen, Krasnoyarsk, Paolo Verde,
 Chooz, **KamLAND**

$$\frac{N_{obs} - N_{exp}}{N_{exp}} = 0.611 \pm 0.085 \pm 0.041$$

Neutrino
stat
syst

out of SM (99.95%)



$$\sin^2 2\theta = 1.0$$

$$\Delta m^2 = 6.9 \cdot 10^{-5} \text{ eV}^2 \text{ (R)}$$

$$\sin^2 2\theta = 0.91$$

$$\Delta m^2 = 6.9 \cdot 10^{-5} \text{ eV}^2 \text{ (spectrum)}$$

Summary of Neutrino-2004

Oscillation parameters:

$$\Delta m_{\text{atm}}^2 = 2.5 \cdot 10^{-3} \text{ eV}^2 (\pm 0.5) \quad \text{SK} + \text{K2K}$$

$$\Delta m_{\text{sol}}^2 = 8 \cdot 10^{-5} \text{ eV}^2 \begin{pmatrix} +0.6 \\ -0.5 \end{pmatrix} \quad \text{KamI} + \text{SNO(salt)}$$

$$(\Delta m_{\text{LSNB}}^2 \sim 1 \text{ eV}^2 \text{ ?})$$

$$\theta_{23} (\text{atm}) \quad \text{large} \quad \sin^2 2\theta_{23} = 1.00 \quad \text{K2K}$$

$$\theta_{12} (\text{sol}) \quad \text{large} \quad \tan^2 \theta_{12} = 0.40 \begin{matrix} +0.09 \\ -0.07 \end{matrix} \quad \text{KamI} + \text{SNO(salt)}$$

$$\theta_{13} (\text{CHOOZ}) \quad \text{small} \quad \sin^2 \theta_{13} < 0.061$$

$$0.8 \cdot 10^{-3} \leq m_\nu \leq 0.5 \text{ (eV)}$$

Direct mass measurements

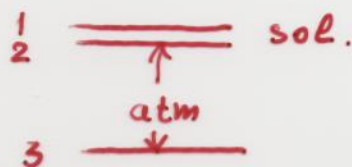
$$m_{\nu e} < 2.2 \text{ eV}$$

$$m_{\nu \mu} < 170 \text{ keV}$$

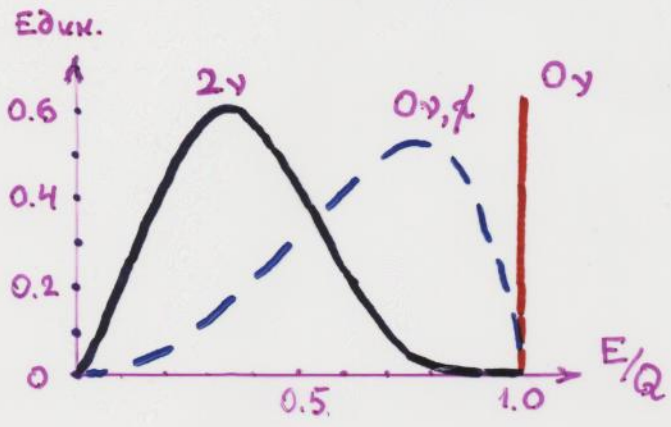
$$m_{\nu \tau} < 18.2 \text{ MeV}$$

$$\text{Cosmology} \quad m_\nu < 0.23 \div 0.7 \text{ eV}$$

$$U = \begin{pmatrix} 0.84 & 0.54 & 0.1 \\ -0.44 & 0.56 & 0.71 \\ 0.32 & -0.63 & 0.71 \end{pmatrix}$$

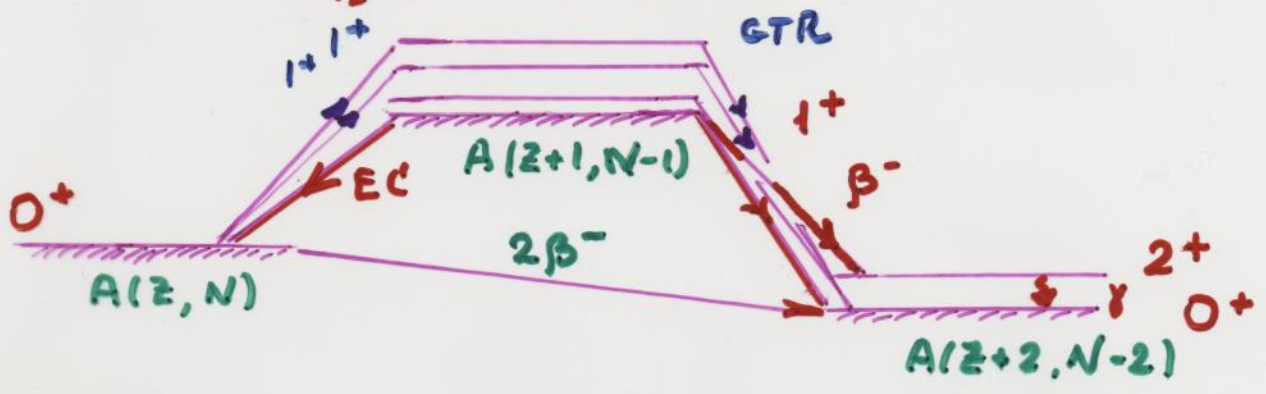


2β - 2ν - распад. - decay



$$\begin{aligned}
 A(z, N) &\rightarrow A(z+2, N-2) + \\
 &+ 2e^- + 2\nu \quad (2\nu) \\
 &+ 2e^- + \gamma_M \quad (0\nu, \gamma) \\
 &+ 2e^- \quad (0\nu)
 \end{aligned}$$

$$(T_{1/2}^{2\nu})^{-1} = F^{2\nu} |M^{2\nu}|^2$$



$$T_{1/2}^{2\nu} = \frac{G_F^4 g_A^4 m_e^9}{96 \ln 2} |M^I M^F|^2 \int d\varepsilon_1 d\varepsilon_2 d\omega F_1 F_2 P_1 \varepsilon_1 P_2 \varepsilon_2 \omega_1^2 \omega_2^2 \mathcal{D}(K, L)$$

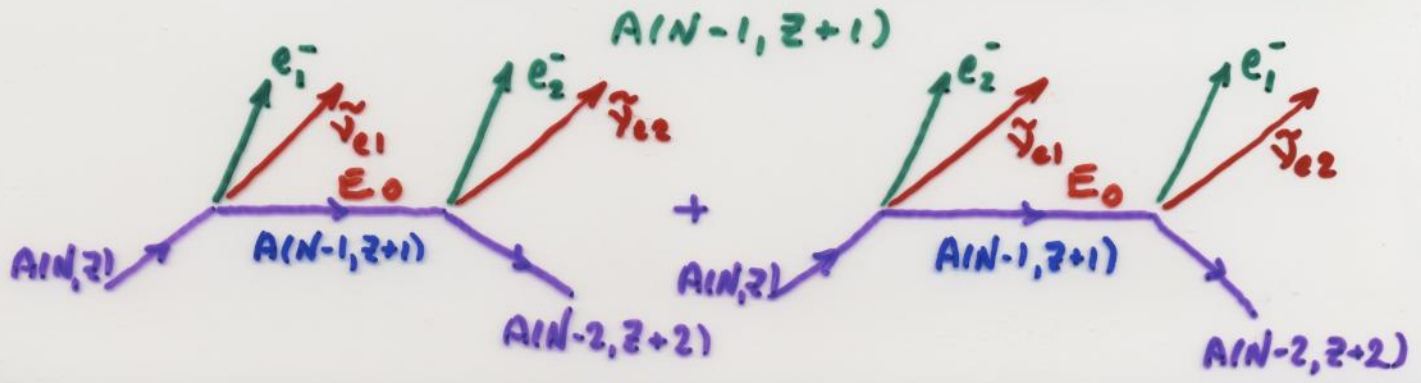
$$\begin{aligned}
 \mathcal{D}(K, L) &= \frac{K^2 + L^2 + KL}{(K-L)^2} \begin{matrix} (00) \\ (02) \end{matrix} \quad K = [\mu_1 + (\varepsilon_1 + \omega_1 - \varepsilon_2 - \omega_2)/2]^{-1} \\
 &\quad + [\mu_1 - (\varepsilon_1 + \omega_1 - \varepsilon_2 - \omega_2)/2]^{-1}, \\
 \mu_1 &= E - \frac{M_i + M_f}{2} \quad L (\omega_1 \leftrightarrow \omega_2).
 \end{aligned}$$

SSD - approximation?

Проблематика: 1+ сост. A(N-1, Z+1) - основное!

$$T_{1/2}^{2\nu} \approx 3 \cdot 10^{14} \text{ yr (ft)}_{EC} \text{ (ft)}_{\beta^-} / H(Q_{\beta\beta}, J_f)$$

для 1+ очн. состояния (1+ ground state)



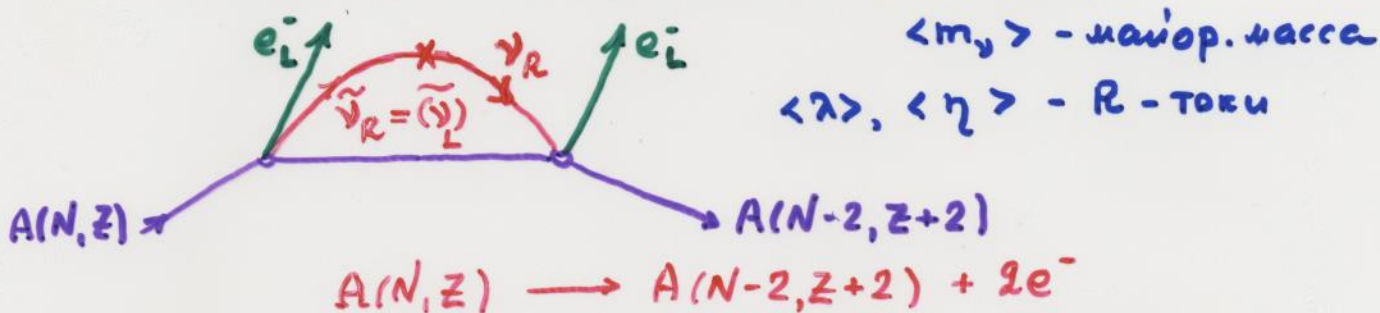
2β - 0ν - распад - decay

$$(T_{1/2}^{0\nu})^{-1} = \langle m_\nu \rangle^2 / m_e^2 F^{0\nu} |M^{0\nu}|^2$$

$$\lim T_{1/2}^{0\nu} \approx (4.18 \cdot 10^{24} \text{ ke}^{-1}) \frac{a}{f} \sqrt{\frac{M \epsilon \tau}{\beta \Delta E}}$$

a - обозначение, M - масса яд

β - фон, ΔE - разрешение, ε - эффект.



$$(T_{1/2}^{0\nu})^{-1} \sim C_{mm} \langle m_\nu \rangle^2 + C_{\eta\eta} \langle \eta \rangle^2 + C_{\lambda\lambda} \langle \lambda \rangle^2 + C_{m\eta} \langle m_\nu \rangle \langle \eta \rangle + C_{m\lambda} \langle m_\nu \rangle \langle \lambda \rangle + C_{\eta\lambda} \langle \eta \rangle \langle \lambda \rangle$$

A(N, Z)	$T_{1/2}^{0\nu}$ лет	$\langle m_\nu \rangle, \text{эВ}$
⁴⁸ Ca	$> 1.5 \cdot 10^{21}$	< 0
⁷⁶ Ge	$> 1.6 \cdot 10^{25}$	$< 0.4 \div 1.07$
⁸² Se	$> 2.7 \cdot 10^{22}$	$< 4.7 \div 9.4$
⁹⁶ Zr	$> 1 \cdot 10^{21}$	< 23
¹⁰⁰ Mo	$> 5.2 \cdot 10^{22}$	$< 2.2 \div 4.9$
¹¹⁶ Cd	$> 3.2 \cdot 10^{22}$	< 4.1
¹²⁸ Te	$> 7.7 \cdot 10^{24}$	< 1.2
¹³⁰ Te	$> 5.6 \cdot 10^{23}$	< 2.9
¹³⁶ Xe	$> 4.4 \cdot 10^{23}$	$< 1.8 \div 5.2$
¹⁵⁰ Nd	$> 1.2 \cdot 10^{21}$	$< 5.3 \div 20$

⁷⁶Ge, данные коллаб.
 $> 1.9 \cdot 10^{25}$ л. $< 0.2 \div \beta$
 Klapdor, 2002
 $0.2 \leq \langle m_\nu \rangle \leq 0.5$

$\langle \lambda \rangle < 10^{-9} \text{ eV}$
 $\langle \eta \rangle < 10^{-7} \text{ eV}$



Нет данных.
 no date!

Resume of the report.

End of XIX

Discovery of radioactivity
 α - β - γ - rays in
radioactive families

1912-13

Rutherford model of atom
Bohr's atom + nucleus

Conversion
 β - lines

Effects of external
electron shell

Continuous spectrum
(RaE)

Nuclear electrons
p-e nuclear model

β -spectrometers
measurements of γ -rays

Excited states of
nuclei, nuclear γ -spectra

heat effect of
 β -rays

Problem of Energy
conservation?
 β -spectrum and quantum
mechanics?

upper limit of
 β -spectrum

β -decay paradoxes

1929-33

Nonconservation of energy
or neutral particle ?

30-32

Neutral particle hyp.
n discovery

Quant. electrodyn. ↔ Fermi β -th.

Experim. n-p nuclear model Theory

nuclear β -decays
↓
elem. particles
↓
New experim. methods:
nuclear decays, accelerat.
cosmic rays, reactors

Phenomenol. β -theory
Isospin
↓
Majorana hypoth.
↓
C, P, T - symmetries
lepton charge
CVC - hypoth.

τ - θ problem

Nonconservation of P
in β, π, μ -decays
Weak interaction phys.

Theory of weak inter, correl.
and polariz. experiments

Particles in
high energy physics
↓
Nucl. ν -physics
↓
Astrophys. objects
↓
New experim. methods
↓
 ν sources and ν detectors
 ν fluxes artif.

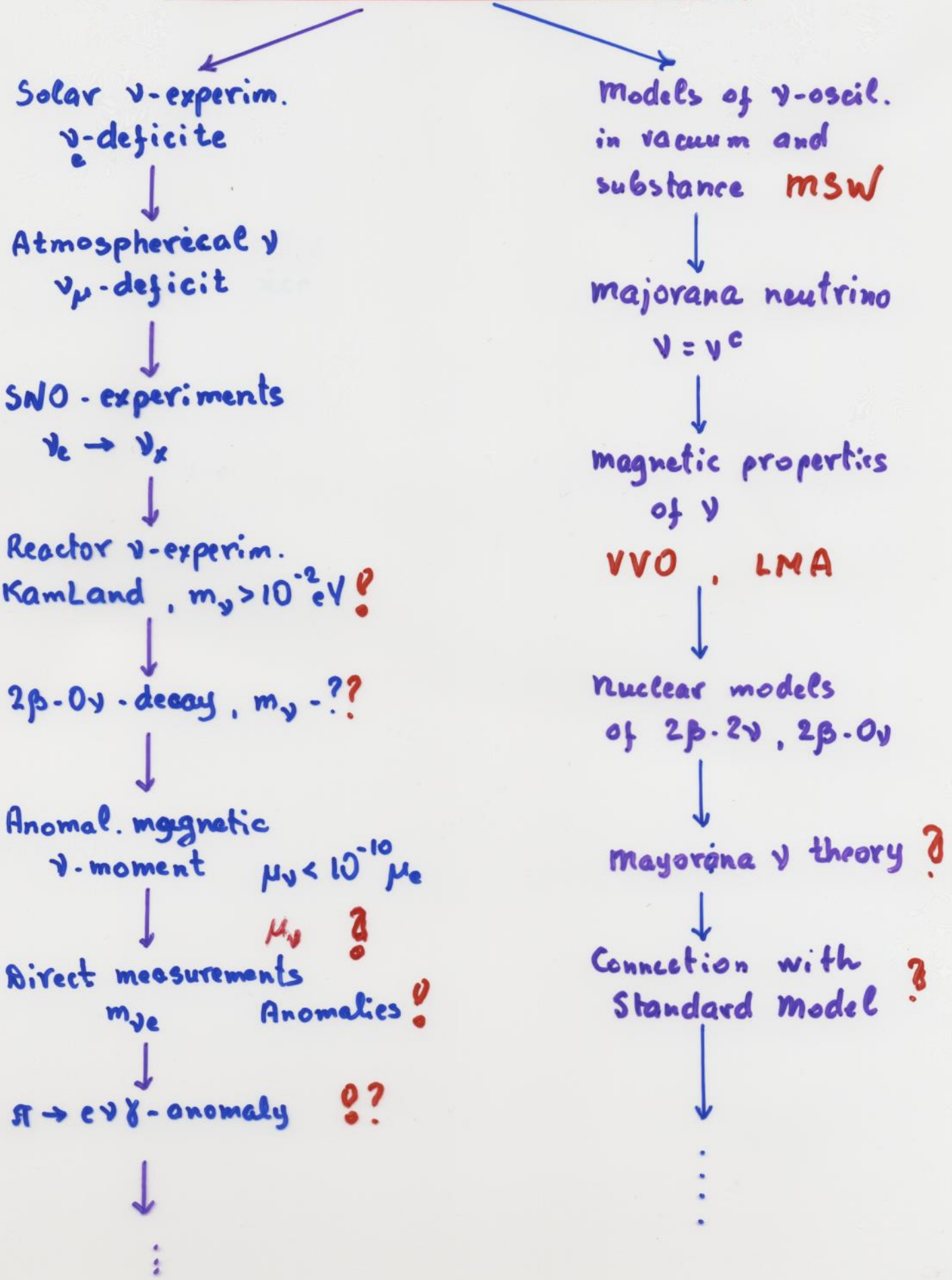
Structure of particles
↓
Nuclear theory
↓
Standard model
↓
New quarks
↓
Nonstandard theories

56-58

80.
-90

??

Beyond Standard Model Phenomena



Acknowledgments :

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