

SPIN EFFECTS IN THE STRONGLY CORRELATED QUARK MODEL

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DSPIN-07

Introduction

Where does the Proton Spin come from?

Spin "Crisis":

DIS experiments: $\Delta\Sigma = \Delta u + \Delta d + \Delta s \approx 1$

SU(6) $\rightarrow 1$

Sum rule for the nucleon spin:

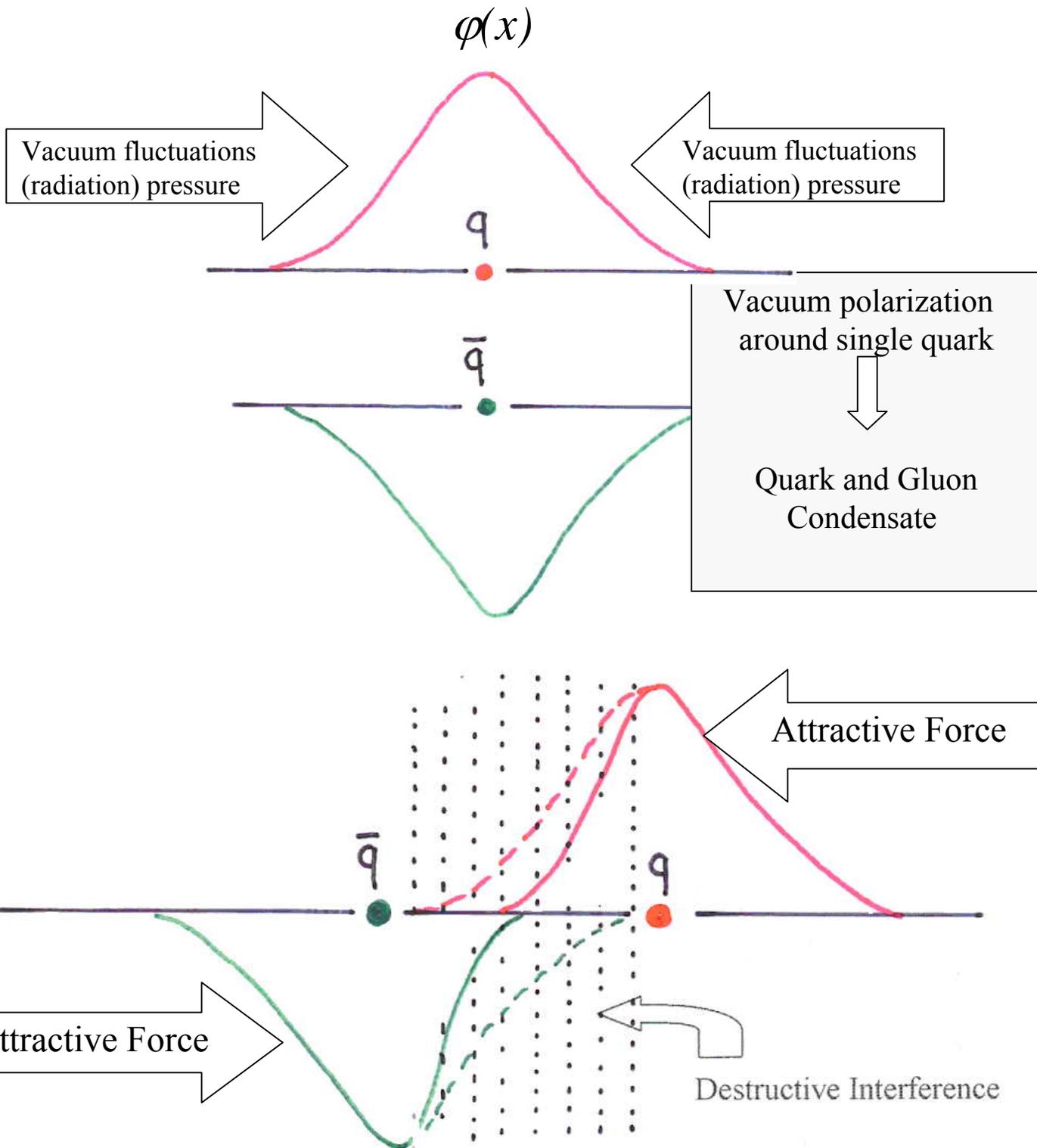
$$1/2 = (1/2)\Delta\Sigma(Q^2) + \Delta g(Q^2) + L(Q^2)_{q+g}$$

SCQM:

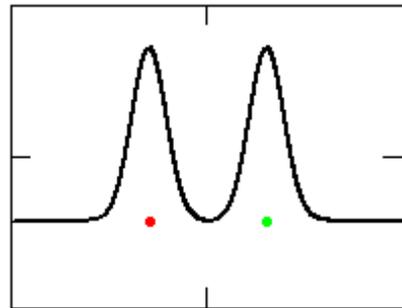
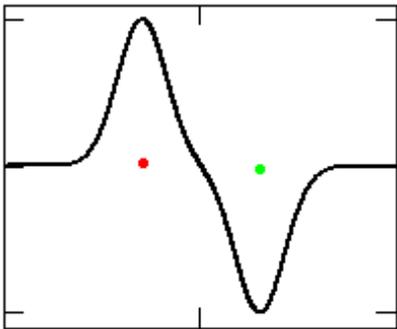
Total nucleon spin comes from circulating around each of three valence quarks gluon and quark-antiquark condensate.

$$s = L_{\bar{q}q+g}$$

Strongly Correlated Quark Model (SCQM)



Quark – antiQuark Oscillations



Constituent Quarks – Solitons

Sine- Gordon (SG) equation

$$\partial_{\mu} \partial^{\mu} \phi(x, t) + \sin \phi(x, t) = 0$$

Breather – oscillating soliton-antisoliton pair, the periodic solution of SG:

$$\phi(x, t)_{s-as} = 4 \tan^{-1} \left[\frac{\sinh \left(ut / \sqrt{1-u^2} \right)}{u \cosh \left(x / \sqrt{1-u^2} \right)} \right]$$

The density profile of the breather:

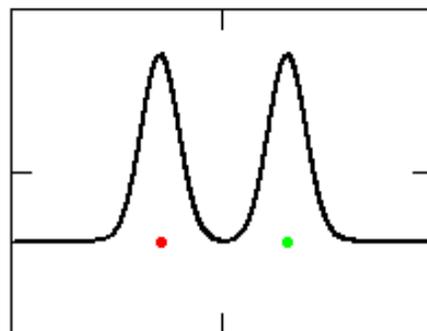
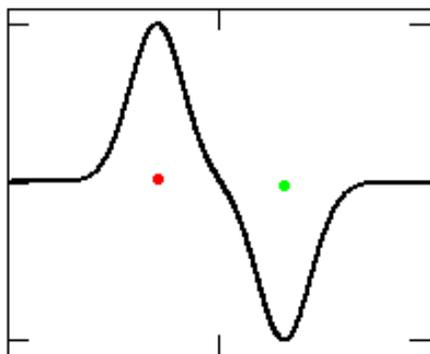
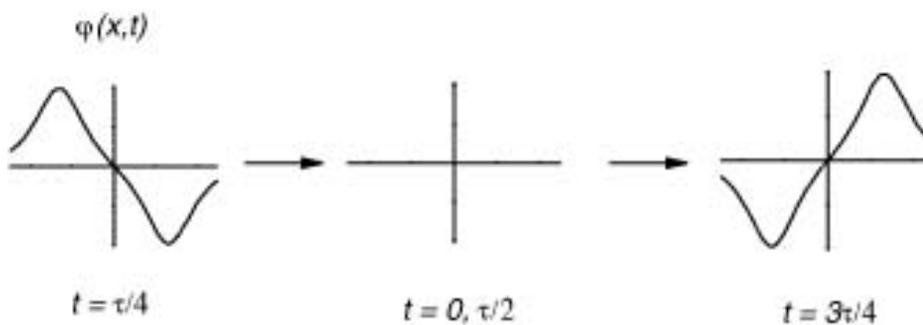
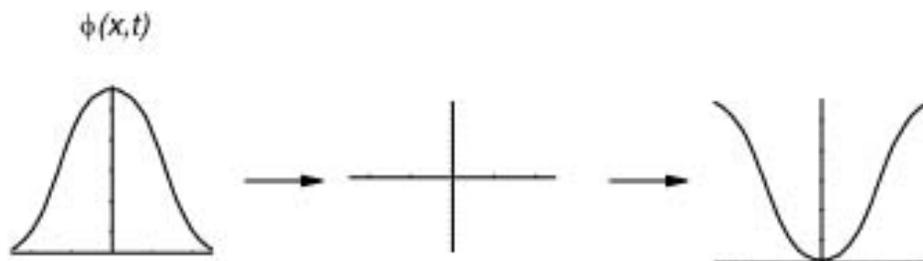
$$\varphi(x, t)_{s-as} = \frac{\partial \phi(x, t)_{s-as}}{\partial x}$$

Breather solution of SG is Lorenz – invariant.

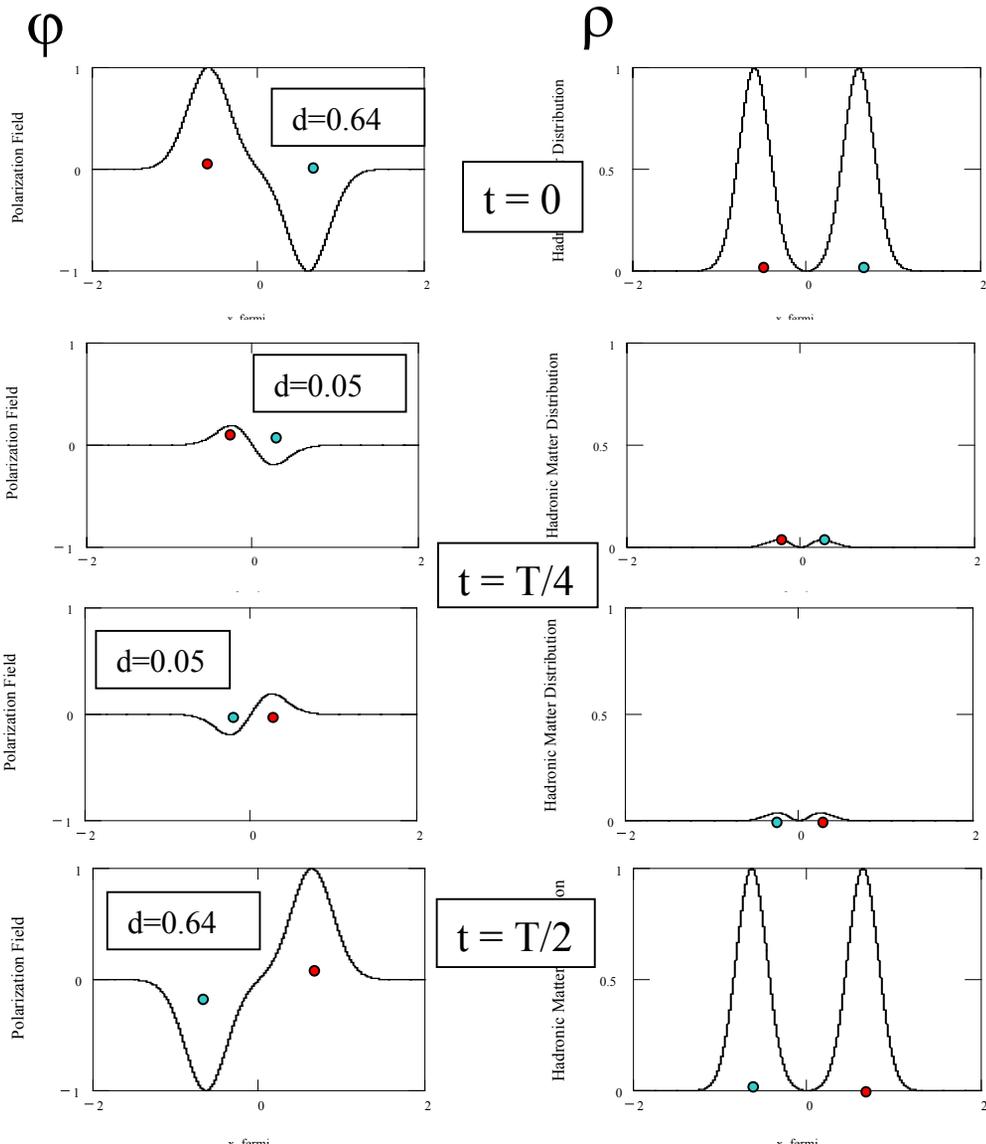
Effective soliton – antisoliton potential

$$U(x) = 2M \tanh^2(mx)$$

Breather (soliton –antisoliton) solution of SG equation



Interplay Between Current and Constituent Quarks \equiv Chiral Symmetry Breaking and Restoration \equiv Dynamical Constituent Mass Generation



Hamiltonian of the Quark – AntiQuark System

$$H = \frac{m_q^-}{(1 - \beta_q^{-2})^{1/2}} + \frac{m_q}{(1 - \beta_q^2)^{1/2}} + V_{qq}^-(2x)$$

m_q^- , m_q are the current masses of quarks,
 $\beta = \beta(x)$ – the velocity of the quark (antiquark),
 V_{qq}^- is the quark–antiquark potential.

$$H = \left[\frac{m_q^-}{(1 - \beta_q^{-2})^{1/2}} + U(x) \right] + \left[\frac{m_q}{(1 - \beta_q^2)^{1/2}} + U(x) \right]$$

$U(x) = \frac{1}{2} V_{qq}^-(2x)$ is the potential energy of the quark.

Conjecture:

$$2U(x) = \int_{-\infty}^{\infty} dz' \int_{-\infty}^{\infty} dy' \int_{-\infty}^{\infty} dx' \rho(x, \vec{r}') \approx 2M_Q(x),$$

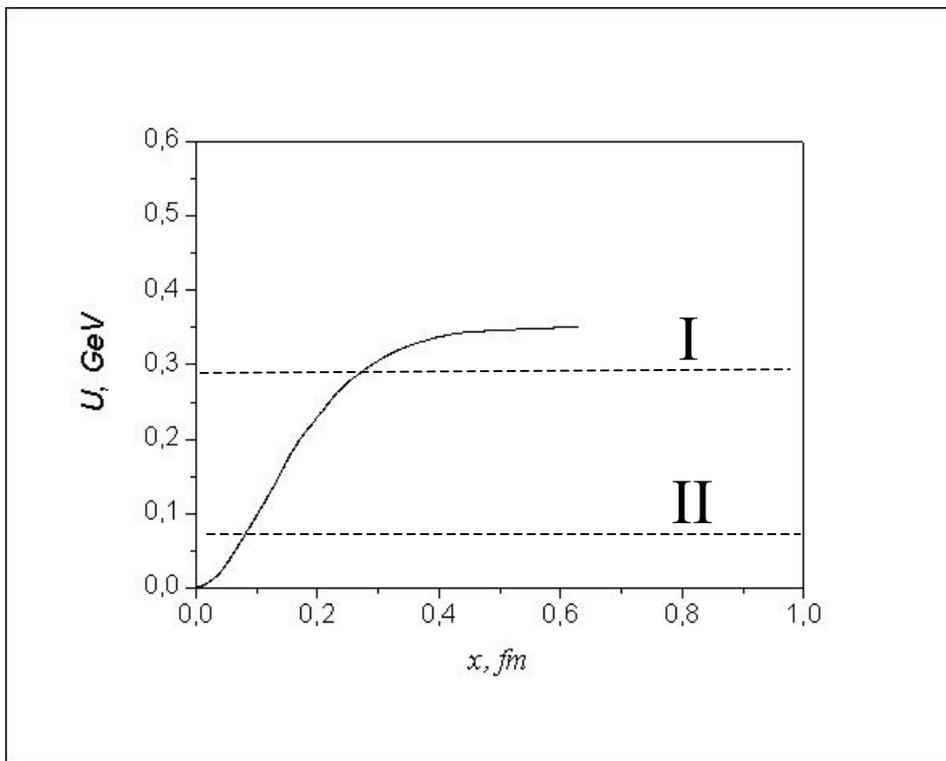
where $M_{Q(\bar{Q})}(x)$ is the dynamical mass of the constituent quark and

$$\rho(x, \vec{r}') = C \left| \varphi_{Q\bar{Q}}(x, \vec{r}') \right| =$$
$$C \left| \varphi_Q(x' + x, y', z') - \varphi_{\bar{Q}}(x' - x, y', z') \right|$$

For simplicity

$$\varphi(\vec{r}) = \frac{(\det \hat{A})^{1/2}}{\pi^{3/2}} \exp \left(- \vec{X}^T \hat{A} \vec{X} \right)$$

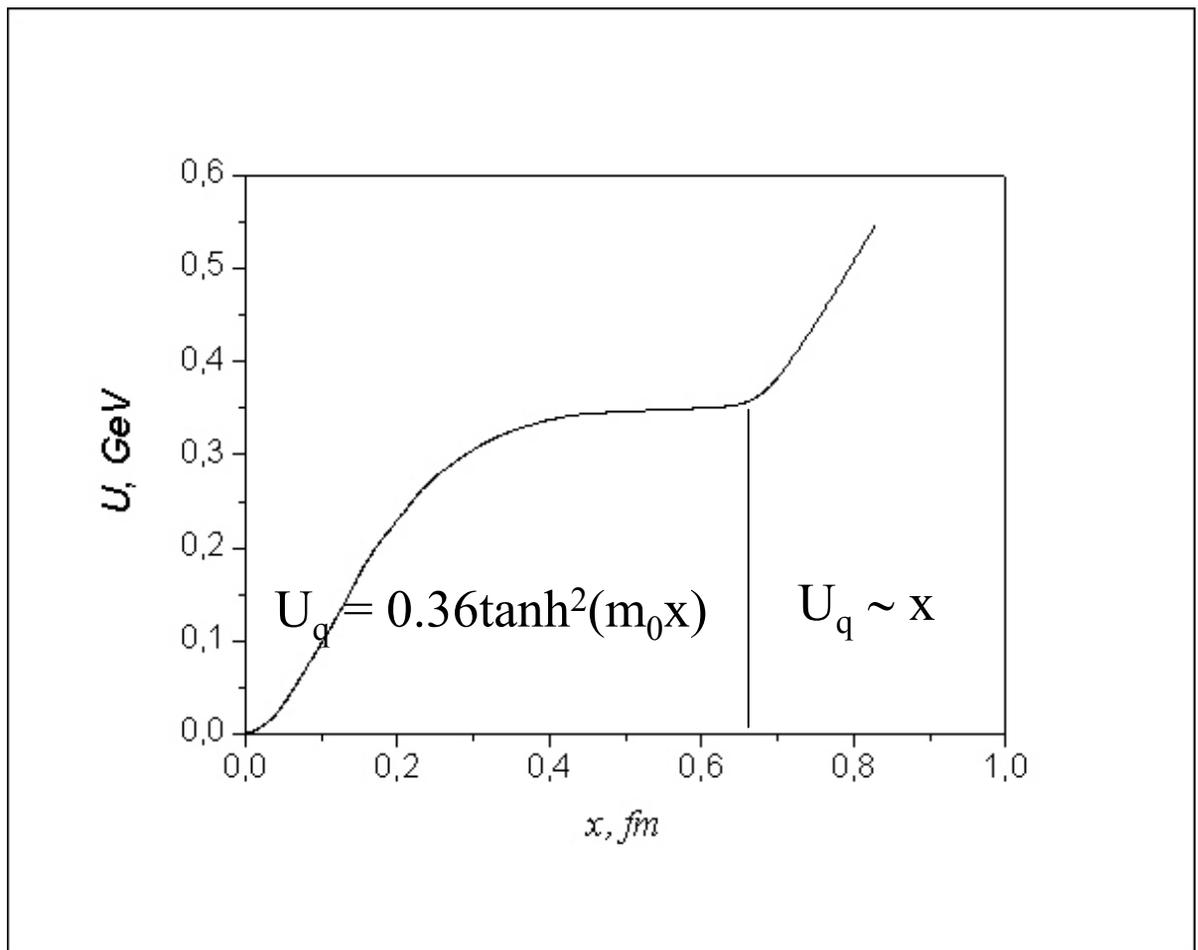
Quark Potential inside Light Hadrons



$U(x) > I$ – constituent quarks

$U(x) < II$ – current(relativistic) quarks

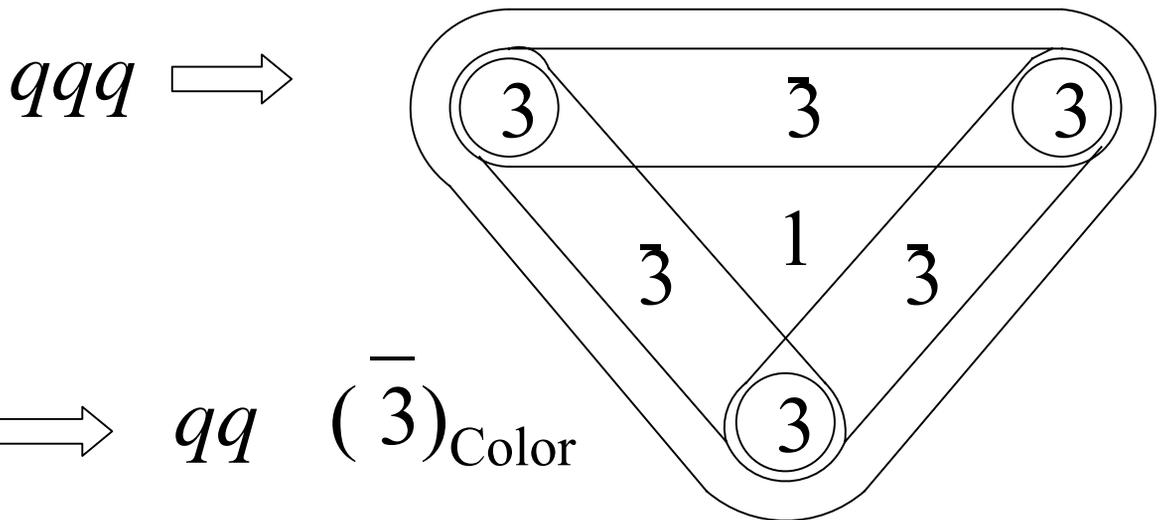
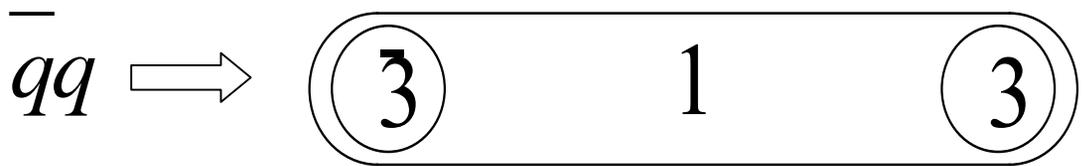
Quark Potential inside Light Hadrons



Generalization to the 3 – quark system (baryons)

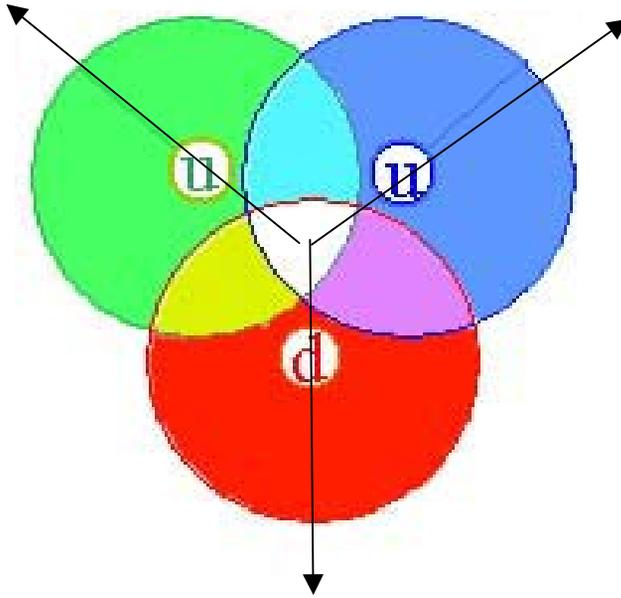
$SU(3)_{Color}$

$3 \iff RGB, \quad \bar{3} \iff CMY$



$\bar{q} \implies qq \quad (\bar{3})_{Color}$

The Proton



One-Quark color wave function

$$\psi(x)_{Color} = \sum_{i=1}^3 a_i(x) |c_i\rangle$$

Where $|c_i\rangle$ are orthonormal states with $i = R, G, B$

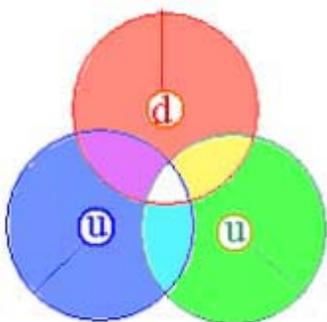
$$\langle c_i | c_j \rangle = \delta_{ij}$$

Nucleon color wave function

$$\psi(x) \rightarrow \frac{1}{\sqrt{6}} \sum_{ijk} e_{ijk} |c_i\rangle |c_j\rangle |c_k\rangle$$

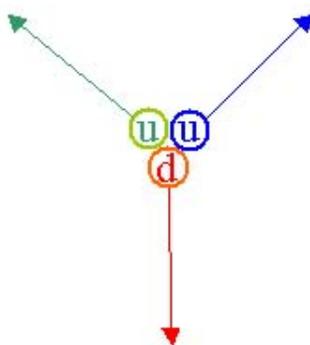
Chiral Symmetry Breaking and its Restoration

$t = 0$
 $x = x_{max}$



Constituent
 Quarks

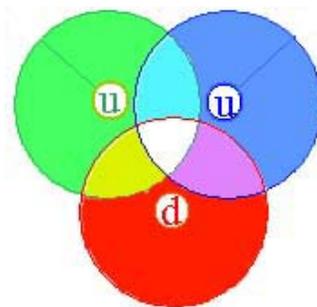
$t = T/4$
 $x = 0$



Current Quarks

Asymptotic Freedom

$t = T/2$
 $x = x_{max}$



Constituent
 Quarks

During the valence quarks oscillations:

$$|B\rangle = c_1 |q_1 q_2 q_3\rangle + c_1 |q_1 q_2 q_3 \bar{q} q\rangle + c_1 |q_1 q_2 q_3 g\rangle + \dots$$

Local gauge invariance in SCQM

Considering each quark separately

$$SU(3)_{\text{Color}} \implies U(1)$$

During valence quarks oscillation destructive interference of their color fields leads to the phase rotation of each quark w.f. in color space:

$$\psi(x)_{\text{Color}} \rightarrow e^{ig\theta(x)} \psi(x)_{\text{Color}}$$

Phase rotation in color space  dressing (undressing) of the quark \equiv the gauge transformation \equiv chiral symmetry breaking (restoration)

$$A_\mu(x) \rightarrow A_\mu(x) + \partial_\mu \theta(x);$$

here

$$A_\mu = (\varphi, 0, 0, 0)$$

Spin in SCQM

1. Now we accept that

$$A_\mu = \{\varphi, \mathbf{A}\}$$

and intersecting \mathbf{E}_{ch} and \mathbf{B}_{ch} create around VQ circulating flow of energy, color analog of the Pointing's vector

$$\mathbf{S}_{\text{ch}} = c^2 \mathbf{E}_{\text{ch}} \times \mathbf{B}_{\text{ch}} .$$

Classical analog of electron spin

– *F. Belinfante 1939; R. Feynman 1964; H. Ohanian 1986; J. Higbie 1988.*

2. Circulating flow of energy carrying along with it hadronic matter is associated with **hadronic matter current**.

3. Total angular momentum created by this Pointing's vector

$$\mathbf{s}_Q = \mathbf{L}_g = (\dots) \int_a^\infty d^3r [\mathbf{r} \times (\mathbf{E}_{\text{ch}} \times \mathbf{B}_{\text{ch}})]$$

is associated with the total spin angular momentum of the constituent quark.

4. Quark oscillations lead to changing of the values of \mathbf{E}_{ch} and \mathbf{B}_{ch} : at the origin of oscillations they are concentrated in a small space region around VQ. As a result hadronic current is concentrated on a narrow shell with small radius.
5. Quark spins are perpendicular to the plane of oscillation.
6. Analogue from hydrodynamics:

Helmholtz laws for velocity field

$$((\partial \xi) / (\partial t)) + \nabla \times (\xi \times \mathbf{v}) = 0,$$

$$\xi = \nabla \times \mathbf{v},$$

$$\nabla \cdot \mathbf{v} = 0,$$

lead to

$$\oint_{\sigma} \mathbf{v} \cdot d\mathbf{r} = s = \text{const.}$$

$$v \propto 1/r$$

The velocity field of the hadronic matter around the valence quarks is irrotational.

Parameters of SCQM

1. Mass of Constituent Quark

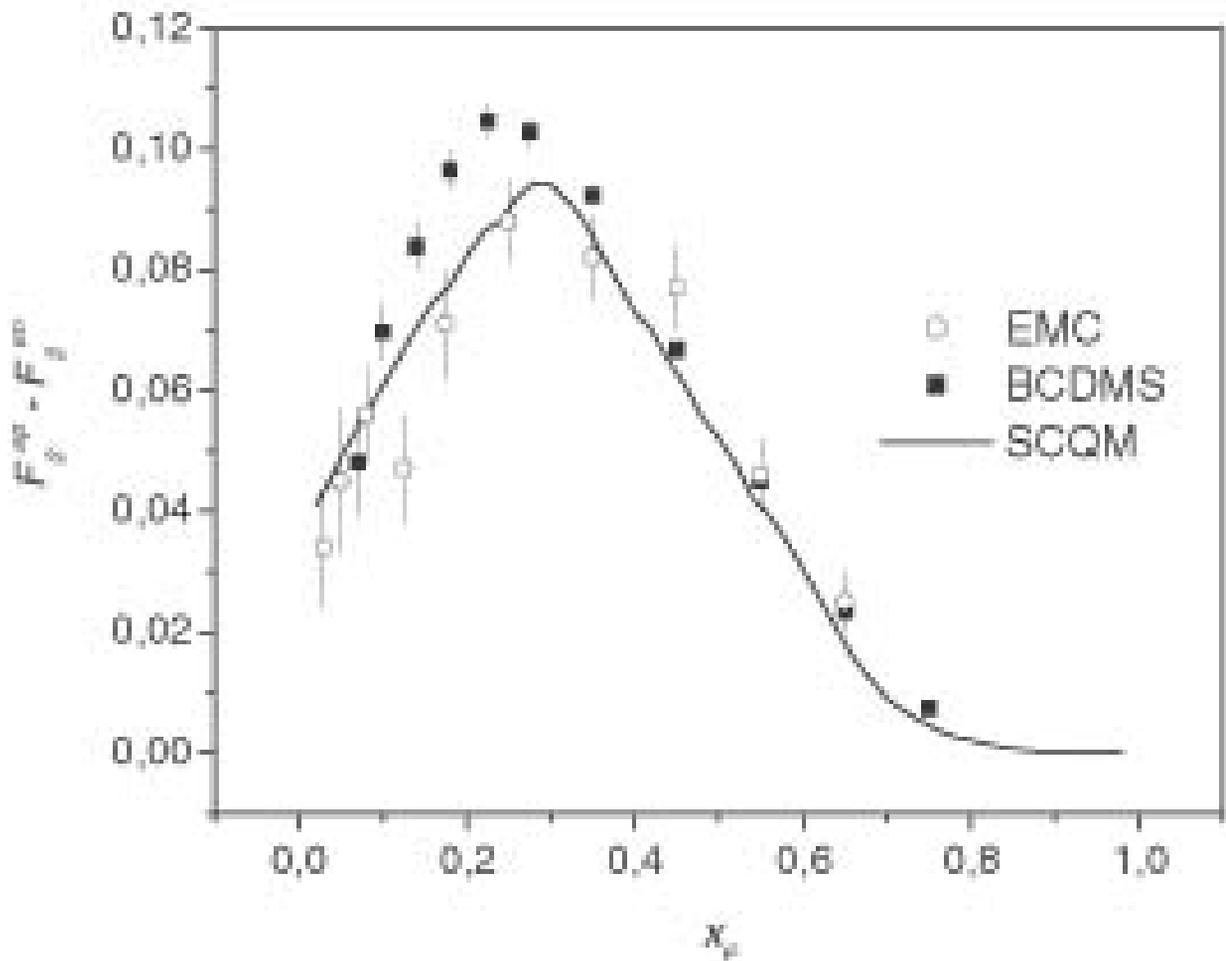
$$M_{Q(\bar{Q})}(x_{\max}) = \frac{1}{3} \left(\frac{m_{\Delta} + m_N}{2} \right) \approx 360 \text{ MeV},$$

2. Maximal Displacement of Quarks: $x_{\max} = 0.64 \text{ fm}$,

3. Constituent quark sizes (parameters of gaussian distribution): $\sigma_{x,y} = 0.24 \text{ fm}$, $\sigma_z = 0.12 \text{ fm}$

Parameters 2 and 3 are derived from the calculations of Inelastic Overlap Function (IOF) and σ_{in} in $\bar{p} p$ and pp – collisions.

Structure Function of Valence Quarks in Proton



Summary on SCQM

- Quarks and gluons inside hadrons are strongly correlated;
- Valence quarks do not orbit inside nucleons
- Constituent quarks are identical to vortical solitons.
- Hadronic matter distribution inside hadrons is fluctuating quantity resulting in interplay between constituent and current quarks.
- **Hadronic matter distribution inside the nucleon is deformed; it is oblate in relation to the spin direction.**

What is “Spin crisis”

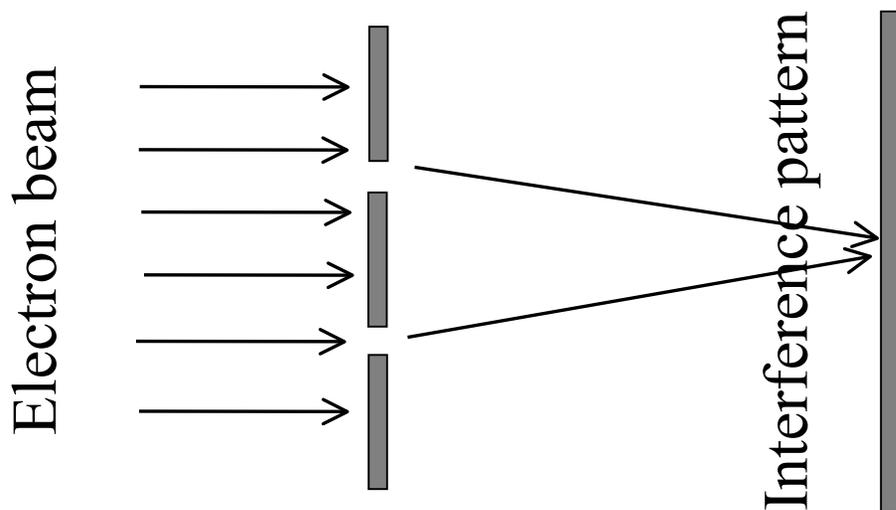
- Quark-vortex is singularity (hole) in vacuum that means nonsimple-connectedness of vacuum structure.
- In DIS the transverse size of virtual photon emitted by incident lepton is given by $d^2 \propto 1/Q^2$.

- $$s_q = \oint_{\sigma} \mathbf{v} \cdot d\mathbf{r}, \quad \mathbf{v} \propto \frac{1}{r}$$

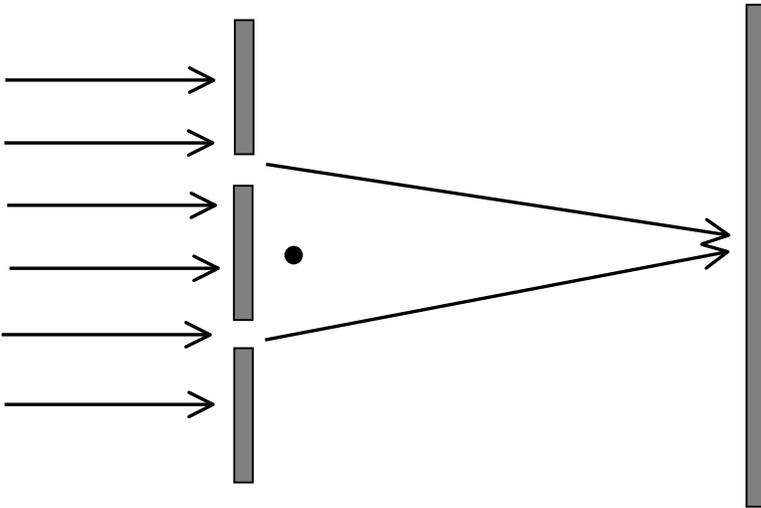
- If the wavelength of the virtual photon in DIS is large enough to cover the center of the vortex then the circulation integral is nonzero and equal to $s_q = 1/2$; otherwise $s_q = 0$.

Aaronov-Bohm effect

Double-slit experiment



Double-slit experiment with thin solenoid



Dirac phase $\varphi_{\text{D}} = \frac{e}{\hbar c} \int \mathcal{A}^\alpha dr_\alpha.$

is dependent on the transport path. Phase transfer along the closed contour is determined by the field flux Φ through this contour:

$$\varphi_{\text{D}} = \frac{e}{\hbar c} \oint \mathcal{A} d\mathbf{r} = \frac{e}{\hbar c} \int \mathcal{B} ds = \frac{e}{\hbar c} \Phi.$$

According to SCQM

▪ In DIS

$$Q^2 \longrightarrow 0 \quad s_q \longrightarrow 1/2$$

$$Q^2 \longrightarrow \infty \quad s_q \longrightarrow 0$$

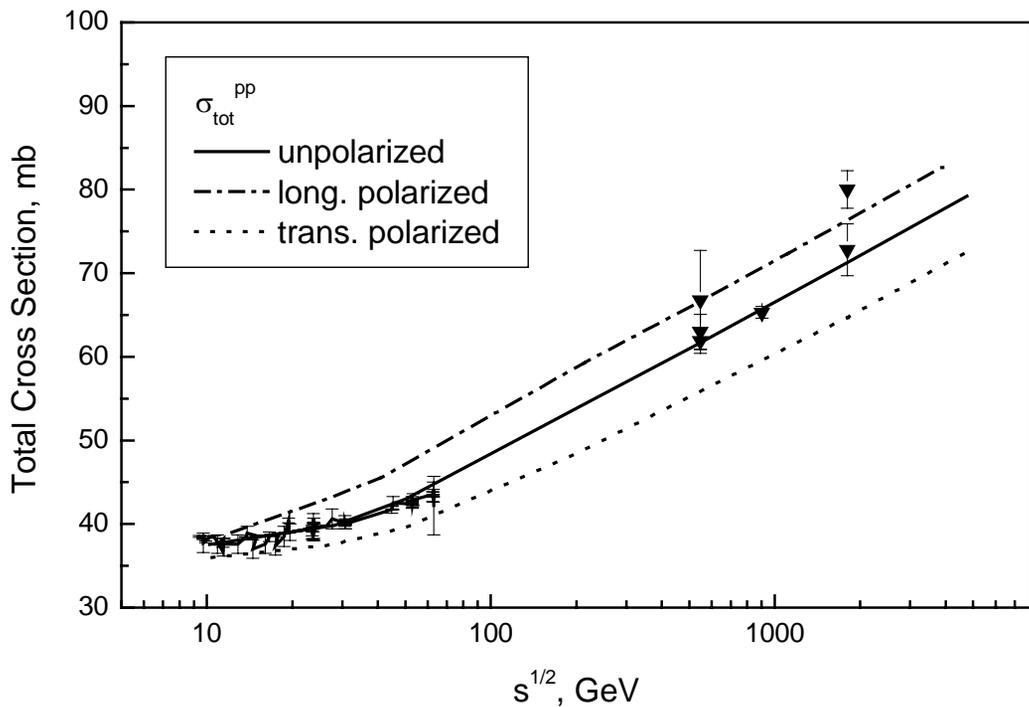
Our calculations

$$\text{At } Q^2 = 3(\text{GeV}/c)^2 \quad \Delta\Sigma = 0.18$$

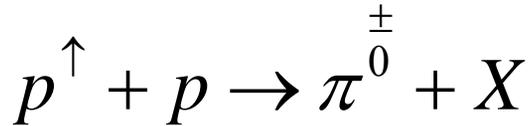
Spin Effect in Soft process

Proposal for forward elastic pp-scattering on RHIC

$$\sigma_T^{tot} < \sigma_L^{tot}$$



Single Spin Asymmetry in proton – proton collisions



- In the factorized parton model

$$\sigma_{\pi/p} \approx f_q \otimes \sigma_{q \rightarrow q'} \otimes D_{\pi/q'}$$

$$\frac{d\sigma_\uparrow}{d^3\bar{p}} = \int dx \cdot d^2k_\perp f_{q/p}(x, \bar{k}_\perp) \int dy \cdot f_{r/p}(y) \times$$

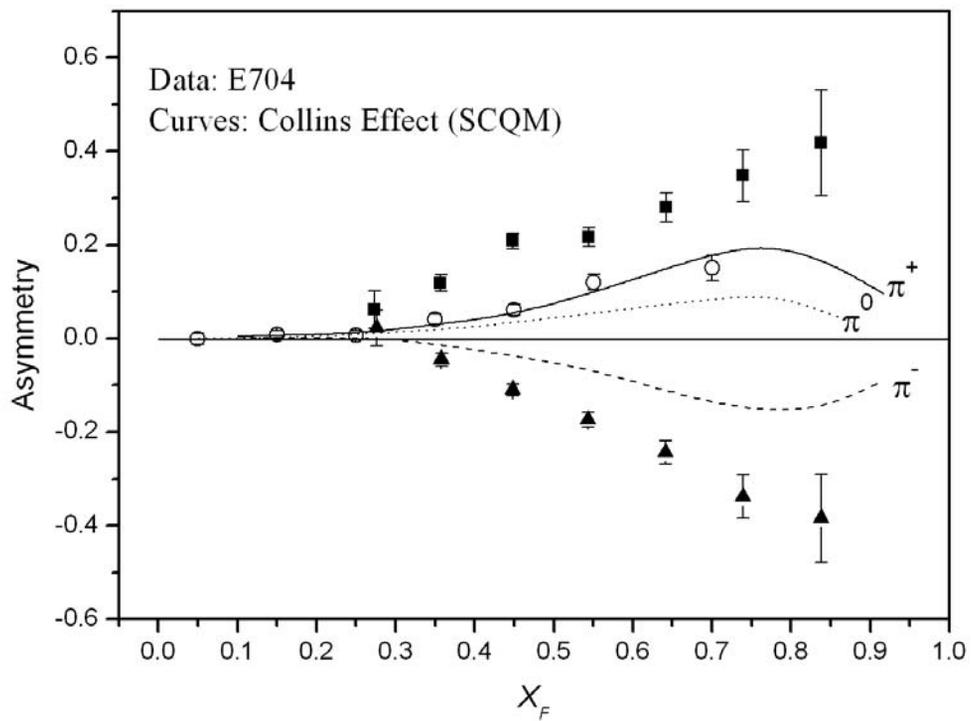
$$\int d\cos\theta \cdot d\varphi \frac{d\sigma(\bar{P}_q, \theta, \varphi)}{d\Omega} \times$$

$$\int dz \cdot d^2\bar{h}_\perp D_{h/q'}(\bar{P}_{q'}, z, \bar{h}_\perp) \cdot \delta^3(\bar{p} - zk' - \bar{h}_\perp)$$

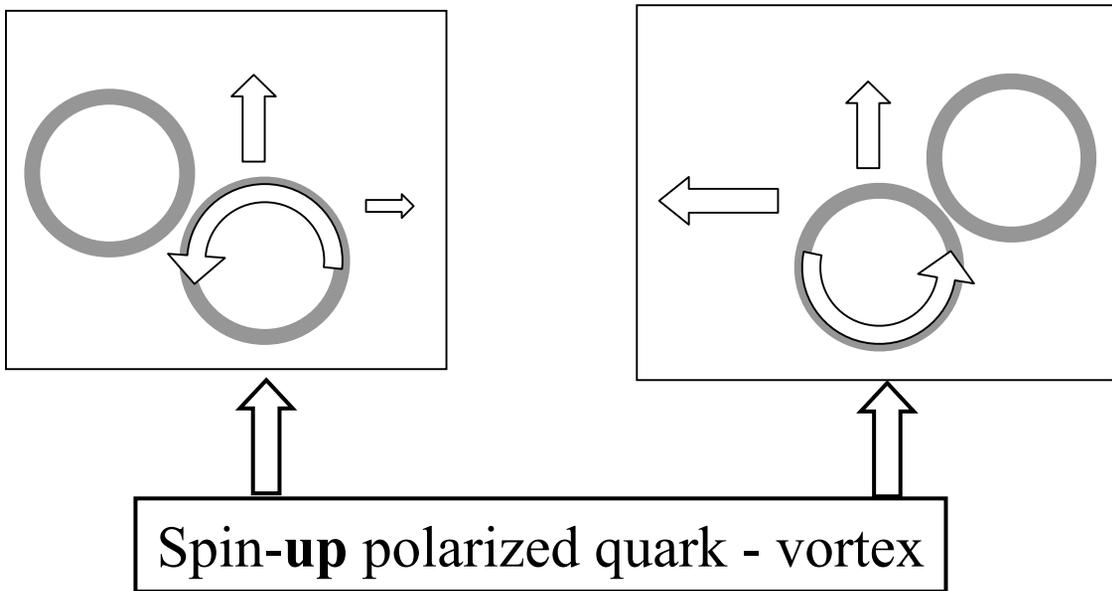
where

$$\vec{P}_q = \frac{\Delta_\perp f_{q\uparrow/p\uparrow}}{f_{q/p}} = \frac{f_{q\uparrow/p\uparrow} - f_{q\downarrow/p\uparrow}}{f_{q/p}}$$

Collins Effect in SSA



Orbital Angular Momentum in SSA



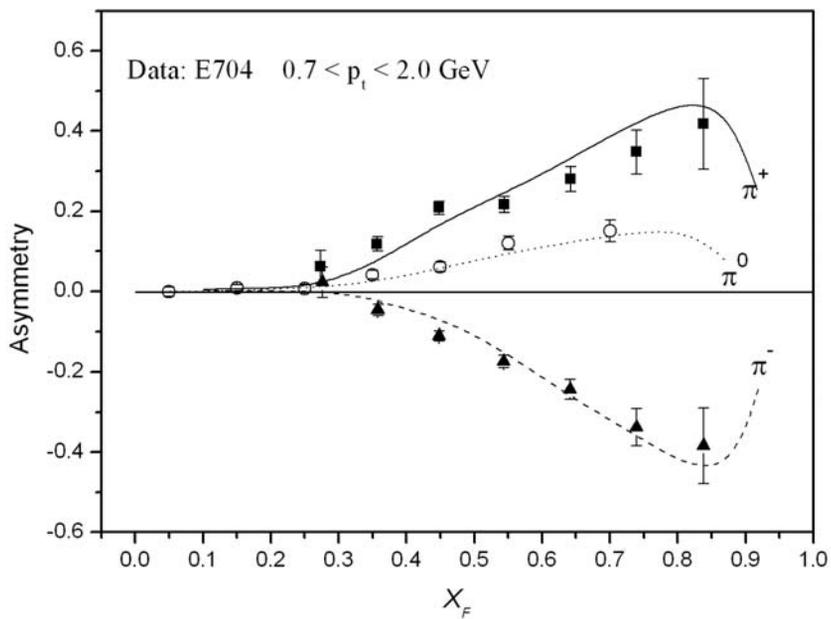
Chou & Yang 1976

*Hadronic matter current distributions inside
polarized hadrons and nuclei*

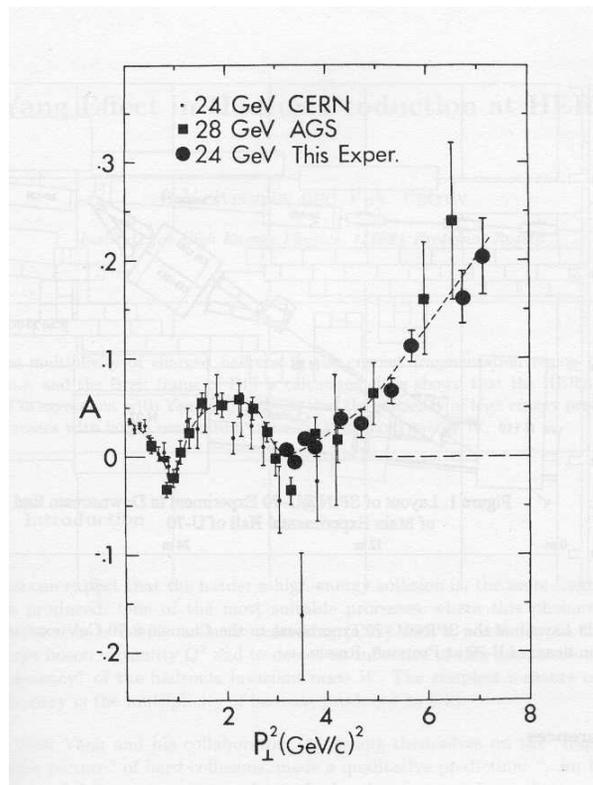
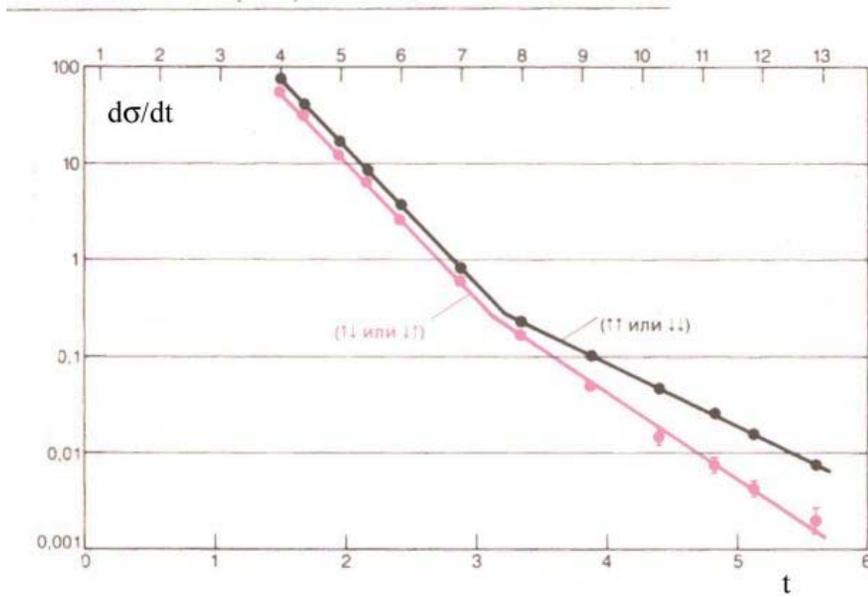
$$\text{Opacity} \propto P_{eff}^{\alpha} \rho(b_x, b_y, b_z)$$

$$P_{eff}^{\alpha} \cong P_{in}^{\alpha} (1 - \alpha v_z)$$

Collins & Orbital Angular Momentum in SSA



Experiments with Polarized Protons



Visualization of Hadron Structure and Scattering Processes

- Pion
- rho – meson
- Nucleon
- Soft scattering of polarized protons
(proposal for RHIC)
- Role of the orbital angular momentum
of chiral (quark-antiquark and gluon)
condensate in SSA and Double Spin
Asymmetry.