Sivers function: from small to large transverse momenta

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> Oleg Teryaev JINR, Dubna (in collaboration with P.G. Ratcliffe, University of Insubria, Como)

Outline

- Single Spin Asymmetries in QCD Sources of Imaginary Phases
- Unsuppressed by 1/Q twist 3
- Non-universality of Sivers function: Colour modification at large -pT
- Sum rules for effective Sivers function from twist 3 effects in spn-dependent DIS
- Sivers function and GPDs
- Conclusions

Single Spin Asymmetries

Main properties:

- Parity: transverse polarization
- Imaginary phase can be seen Tinvariance or technically - from the imaginary i in the (quark) density matrix
 Various mechanisms – various sources of phases

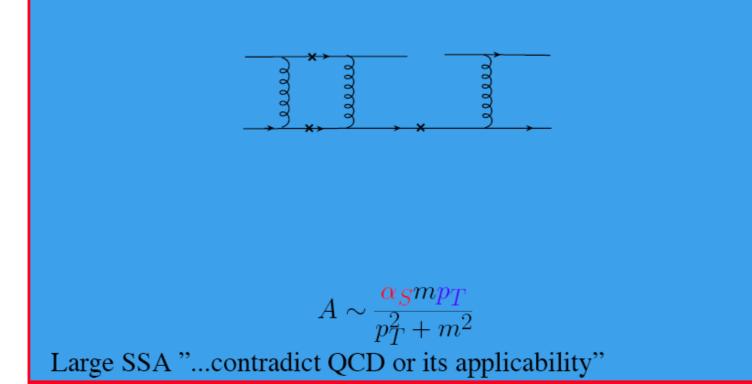
Phases in QCD-I

- QCD factorization soft and hard parts-
- Phases form soft, hard and overlap
- Assume (generalized) optical theorem phase due to on-shell intermediate states – positive kinematic variable (= their invariant mass)
- Hard: Perturbative (a la QED: Barut, Fronsdal (1960), found at JLAB recently):

Kane, Pumplin, Repko (78) Efremov (78)

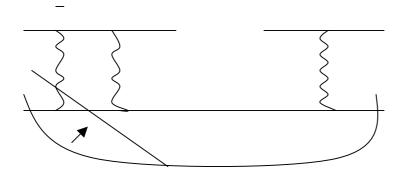
Perturbative PHASES IN QCD

QCD factorization: where to borrow imaginary parts? Simplest way: from short distances - loops in partonic subprocess. Quarks elastic scattering (like q - e scattering in DIS):



Short+ large overlaptwist 3

- Quarks only from hadrons
- Various options for factorization shift of SH separation

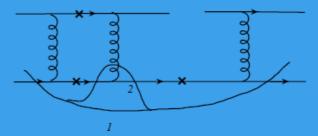


New option for SSA: Instead of 1-loop twist 2

 Born twist 3: Efremov, OT (85, Ferminonc poles); Qiu, Sterman (91, GLUONIC poles)

Twist 3 correlators

Escape: QCD factorization - possibility to shift the borderline between large and short distances



At short distances - Loop \rightarrow Born diagram At Large distances - quark distribution \rightarrow quark-gluon correlator. Physically - process proceeds in the external gluon field of the hadron. Leads to the shift of α_S to non-perturbative domain AND "Renormalization" of quark mass in the external field up to an order of hadron's one

 $rac{lpha_{S}mp_{T}}{p_{T}^{2}+m^{2}}
ightarrow rac{Mb(x_{1},x_{2})p_{T}}{p_{T}^{2}+M^{2}}$

Further shift of phases completely to large distances - T-odd fragmentation functions. Leading twist transversity distribution - no hadron mass suppression.

Phases in QCD –large distances in fragmentation

- Non-perturbative positive variable
- Jet mass-Fragmentation function: Collins(92);Efremov,Mankiewicz, Tornqvist (92),
- Correlated fragmentation: Fracture function: Collins (95), O.T. (98).

Phases in QCD-Large distances in distributions

- Distributions :Sivers, Boer and Mulders no positive kinematic variable producing phase
- QCD: Emerge only due to (initial of final state) interaction between hard and soft parts of the process: "Effective" or "non-universal" SH interactions by physical gluons – Twist-3 (Boer, Mulders, OT, 97)
- Brodsky -Hwang-Schmidt model: the same SH interactions as twist 3 but non-suppressed by Q: Sivers function – leading (twist 2).

What is "Leading" twist?

- Practical Definition Not suppressed as M/Q
- However More general definition: Twist 3 may be suppresses

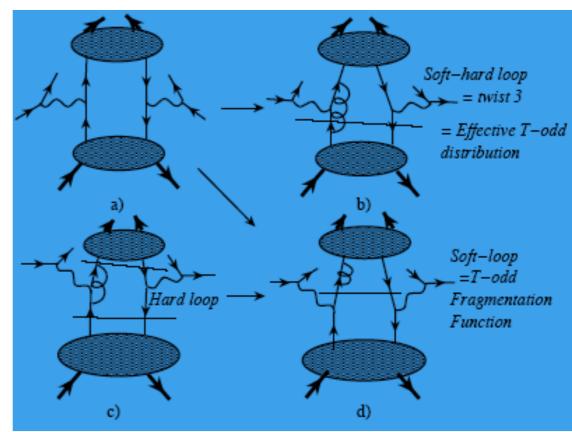
as M/P_T

Twist 3 may contribute at leading order in 1/Q !

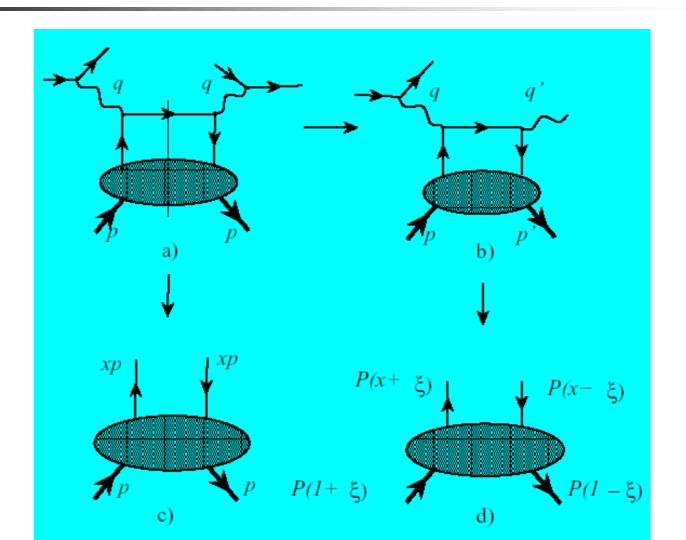
Does this happen indeed?? – Explicit calculation for the case when $Q >> P_T$ May be interesting for experimental studies

Sources of Phases in SIDIS and Drell-Yan processes

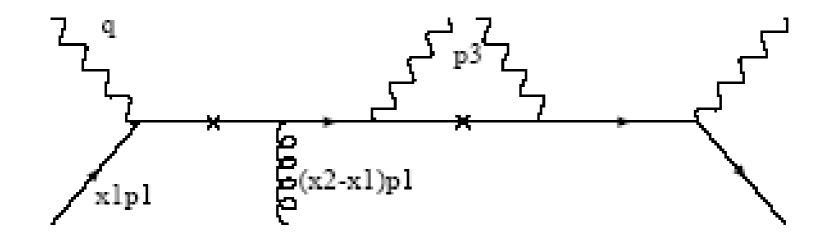
- a) Born no SSA
- b) -Sivers (can be only effective)
- -for both SIDIS and DY
- c) Perturbative
- d) Collins (SIDIS) or (effective) Boer (DY)



Final Pion -> Photon: SIDIS ->
SIDVCS (clean, easier than exclusive)
- analog of DVCS



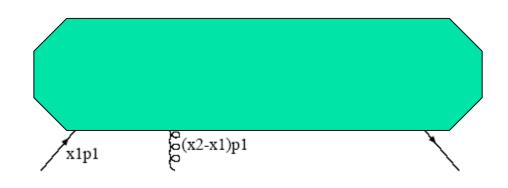
Twist 3 partonic subprocesses for SIDVCS



Real and virtual photons most clean tests of QCD

- Both initial and final real :Efremov, O.T. (85)
- Initial quark/gluon, final real : Efremov, OT (86, fermionic poles); Qui, Sterman (91, GLUONIC poles)
- Initial real, final-virtual (or quark/gluon) Korotkiian, O.T. (94)
- Initial –virtual, final-real: O.T., Srednyak (05; smooth transition from fermionic via hard to GLUONIC poles).

Quark-gluon correlators



- Non-perturbative NUCLEON structure physically mean the quark scattering in external gluon field of the HADRON.
- Depend on TWO parton momentum fractions
- For small transverse momenta quark momentum fractions are close to each other- gluonic pole; probed if :
 Q >> P_T>> M

$$x_2 - x_1 = \delta = \frac{p_T^2 x_B}{Q^2 z}$$

Cross-sections at low transverse momenta:

$$d\sigma_{total} = f(x_{Bj}) 8Q^2 \frac{x_{Bj}^2 (1 + (1 - y)^2)(1 + (1 - z)^2)}{y^2 z \delta}$$
(12)

$$d\sigma_{ax1x2} = b_A(x_{Bj}, x_2) 8M p_T \frac{x_{Bj}(1 + (1 - y)^2)(2 - z)}{y^2(1 - z)\delta} s_T \sin(\phi_s^h)$$
(13)

$$d\sigma_{vx1x2} = b_V(x_{Bj}, x_2) 8M p_T \frac{x_{Bj}(1 + (1 - y)^2)(1 + (1 - z)^2)}{y^2 z(1 - z)\delta} s_T sin(\phi_s^h)$$
(14)

$$d\sigma_{a0x2} = -b_A(0, x_2) 8M p_T \frac{x_{Bj}^2 (2(1-y)(1-2z) + y^2(1-z))}{y^2 z^2 \delta} s_T sin(\phi_s^h)$$

(14) - non-suppressed for large Q if Gluonic pole exists=effective Sivers function; spin-dependent looks like unpolarized (soft gluon)

$$A \propto \frac{2M p_T \varphi_V(\chi_B)}{m_T^2 \chi_B q(\chi_B)} S_T \sin \phi_h^s$$

Effective Sivers function

- Needs (soft) talk of large and short distances
- Complementary to gluonic exponential, when longitudinal (unsuppressed by Q, unphysical) gluons get the physical part due to transverse link (Belitsky, Ji, Yuan)
- We started instead with physical (suppressed as 1/Q) gluons, and eliminated the suppression for gluonic pole.
- Another support Ji, Qiu, Vogelsang, Yuan in DY and SIDIS – PERTURBATIVE Sivers function from twist 3

Other way - NP Sivers and gluonic poles at large PT (P.G. Ratcliffe, OT, hep-ph/0703293)

Sivers factorized (general!) expression

$$d\Delta\sigma = \int d^2k_T dx \, f_S(x,k_T) \, \operatorname{Tr}[\gamma^{\rho} H(x,k_T)] \frac{\epsilon^{\rho s P k_T}}{M}$$

Expand in kT = twist 3 part of Sivers

$$d\Delta\sigma = \int dx\, f_S(x,k_T)\,\, {\rm Tr} \! \left[\gamma^\rho \frac{\partial H(x,k_T=0)}{\partial k_T^\alpha} k_T^\alpha \right] \epsilon^{\rho s P k_T} \label{eq:sigma_static_static}$$

From Sivers to twist 3 - II

• Angular average : $\langle k_T^{\mu} k_T^{\nu} \rangle = -\frac{g_T^{\mu\nu}}{2} \langle k_T^2 \rangle$

$$g_T^{\mu\nu} = g^{\mu\nu} - P^{\mu}n^{\nu} - n^{\mu}P^{\nu}$$

- As a result $d\Delta \sigma = -M \int dx f_S^{(1)}(x) \operatorname{Tr} \left[\gamma^{\rho} \frac{\partial H(x, k_T = 0)}{\partial k_T^{\alpha}} \right]$ $f_S^{(1)}(x) = \int d^2 k_T f_S(x, k_T) \frac{k_T^2}{2M^2} \cdot \frac{\left(\epsilon^{\rho s P \alpha} - P^{\alpha} \epsilon^{\rho s P n}\right)}{2M^2}$
- M in numerator sign of twist 3. Higher moments – higher twists. KT dependent function – resummation of higher twists

From Sivers to gluonic poles -

Final step – kinematical identity

$$\epsilon^{\rho s P \alpha} = P^{\alpha} \epsilon^{\rho s P n} - P^{\rho} \epsilon^{\alpha s P n}$$

Two terms are combined to one

$$d\Delta\sigma = M \int dx \, f_S^{(1)}(x) \, \operatorname{Tr}\left[\gamma \cdot P \frac{\partial H(x, k_T = 0)}{\partial k_T^{\alpha}}\right] \epsilon^{\alpha s P n}$$

- Key observation exactly the form of Master Formula for gluonic poles (Koike et al, 2007)
- Non Suppression as 1/Q seen!

Effective Sivers function

 Follows the expression similar to BMT, BMP, JVY

$$x f_{s}^{(1)}(x) = \sum C_{i} \frac{1}{2M} T_{j}(x, x),$$

- Up to Colour Factors !
- Defined by colour charge (natural for low energy theorems!): Collins sign rule: ISI -> -FSI holds because of quark -> antiquark (cf. Abelian charge – Collins&Qiu (arXiv:0705.2141)

What are these factors?

- SIDIS at large pT : -1/6 for mesons from quark, 3/2 from gluon fragmentation (kaons?)
- DY at large pT (PAX): 1/6 in quark antiquark annihilation, - 3/2 in gluon Compton subprocess – Collins sign rule more elaborate – involve crossing of distributions and fragmentations - special role of PION DY (COMPASS).
- Hadronic pion production more complicated studied for P-exponentials by Amsterdam group + VW
- FSI for pions from quark fragmentation
- -1/6 x (non-Abelian Compton) +1/8 x (Abelian Compton)

How to pass from high to low PT

- Hard poles in correlators (become soft ay small PT – c.f. SIDVCS)
- Low pT cannot distinguish fragmentation from quarks and gluons: 3/2-1/6 = 4/3 (Abelian)
- Strong transverse momentum dependence, very different for mesons from quark and gluon fragmentation

Colour flow

Quark at large PT:-1/6

- Gluon at large PT : 3/2
- Low PT combination of quark and gluon: 4/3 (absorbed to definition of Sivers function)
- Similarity to colour transparency phenomenon

Twist 3 factorization (Efremov, OT '84, Ratcliffe, Qiu, Sterman)

 Convolution of soft (S) and hard (T) parts

$$d\sigma_s = \int dx_1 dx_2 \frac{1}{4} Sp[S_\mu(x_1, x_2)T_\mu(x_1, x_2)]$$

 Vector and axial correlators: define hard process for both double (g₂) and single asymmetries

$$T_{\mu}(x_1, x_2) = \frac{M}{2\pi} (\hat{p}_1 \gamma^5 s_{\mu} b_A(x_1, x_2) - i \gamma_{\rho} \epsilon^{\rho \mu s p_1} b_V(x_1, x_2))$$

Twist 3 factorization -II

Non-local operators for quark-gluon correlators

 $b_A(x_1, x_2) = \frac{1}{M} \int \frac{d\lambda_1 d\lambda_2}{2\pi} e^{i\lambda_1 (x_1 - x_2) + i\lambda_2 x_2} \langle p_1, s | \bar{\psi}(0) \hat{n} \gamma^5 (D(\lambda_1) s) \psi(\lambda_2) | p_1, s \rangle,$

 $b_{V}(x_{1},x_{2}) = \frac{i}{M} \int \frac{d\lambda_{1} d\lambda_{2}}{2\pi} e^{i\lambda_{1}(x_{1}-x_{2})+i\lambda_{2}x_{2}} \epsilon^{\mu s p_{1}n} \langle p_{1},s | \bar{\psi}(0) \hat{n} D_{\mu}(\lambda_{1}) \psi(\lambda_{2}) | p_{1},s \rangle$

Symmetry properties (from Tinvariance)

$$b_A(x_1, x_2) = b_A(x_2, x_1), \ b_V(x_1, x_2) = -b_V(x_2, x_1)$$

Twist-3 factorization -III

Singularities

$$b_A(x_1, x_2) = \varphi_A(x_1)\delta(x_1 - x_2) + b_A^r(x_2, x_1),$$

$$b_V(x_1, x_2) = \frac{\varphi_V(x_1)}{x_1 - x_2} + b_V^r(x_1, x_2)$$

- Very different: for axial Wandzura-Wilczek term due to intrinsic transverse momentum
- For vector-GLUONIC POLE (Qiu, Sterman '91)
 large distance background

Sum rules

 EOM + n-independence (GI+rotational invariance) –relation to (genuine twist
 3) DIS structure functions

$$\begin{split} &\int_{0}^{1} x^{n} \bar{g}_{2}(x) dx = \int_{0}^{1} x^{n} (\frac{n}{n+1} g_{1}(x) + g_{2}(x)) dx = \\ &- \frac{1}{\pi(n+1)} \int_{|x_{1}, x_{2}, x_{1} - x_{2}| \leq 1} dx_{1} dx_{2} \sum_{f} e_{f}^{2} [\frac{n}{2} b_{V}(x_{1}, x_{2}) (x_{1}^{n-1} - x_{2}^{n-1}) + \\ &b_{A}^{r}(x_{1}, x_{2}) \phi_{n}(x_{1}, x_{2})], \quad \phi_{n}(x, y) = \frac{x^{n} - y^{n}}{x - y} - \frac{n}{2} (x^{n-1} - y^{n-1}), \quad n = 0, 2... \end{split}$$



To simplify – low moments

$$\int_0^1 x^2 \hat{g}_2(x) dx = -\frac{1}{3\pi} \int_{|x_1, x_2, x_1 - x_2| \le 1} dx_1 dx_2 \sum_f e_f^2 b_V(x_1, x_2) (x_1 - x_2)$$

Especially simple – if only gluonic pole kept:

$$\int_0^1 x^2 \bar{g}_2(x) dx = -\frac{1}{3\pi} \int_{|x_1, x_2, x_1 - x_2| \le 1} dx_1 dx_2 \sum_f e_f^2 \varphi_V(x_1)$$
$$= -\frac{1}{3\pi} \int_{-1}^1 dx_1 \sum_f e_f^2 \varphi_V(x_1) (2 - |x_1|)$$

Gluonic poles and Sivers function

- Gluonic poles effective Sivers functions-Hard and Soft parts talk, but SOFTLY
- Implies the sum rule for effective Sivers function (soft=gluonic pole dominance assumed in the whole allowed x's region of quark-gluon correlator)

$$x f_{T}(x) = \frac{1}{2M}T(x,x) = \frac{1}{4}\phi_{v}(x)$$

$$\int_{0}^{1} dx x^{2} \bar{g}_{2}(x) = \frac{4}{3\pi} \int_{0}^{1} dx x f_{T}(x)(2-x)$$

Compatibility of SSA and DIS

- Extractions of Sivers function: "mirror" u and d
- First moment of EGMMS = 0.0072 (0.0042 0.014)
- Twist -3 similar for neutron and proton (0.005) and of the same sign – nothing like mirror picture seen –but supported by colour ordering!
- Current status: Scale of Sivers function seems to be reasonable, but flavor dependence differs qualitatively.
- Inclusion of pp data, global analysis including gluonic (=Sivers) and fermionic poles

Relation of Sivers function to GPDs

- Qualitatively similar to Anomalous Magnetic Moment (Brodsky et al)
- Quantification : weighted TM moment of Sivers PROPORTIONAL to GPD E (hep-ph/0612205): $x f_{T}(x) \sim xE(x)$
- Burkardt SR for Sivers functions is now related to Ji SR for E and, in turn, to Equivalence Principle

$$\sum_{q,G} \int dxx f_T(x) = \sum_{q,G} \int dxx E(x) = 0$$

How gravity is coupled to nucleons?

- Current or constituent quark masses ?-neither!
- Energy momentum tensor like electromagnertic current describes the coupling to photons

Equivalence principle

- Newtonian "Falling elevator" well known and checked
- Post-Newtonian gravity action on SPIN known since 1962 (Kobzarev and Okun') – not yet checked
- Anomalous gravitomagnetic moment iz ZERO or
- Classical and QUANTUM rotators behave in the SAME way

Gravitational formfactors

 $\langle p'|T^{\mu\nu}_{q,g}|p\rangle = \bar{u}(p') \Big[A_{q,g}(\Delta^2) \gamma^{(\mu} p^{\nu)} + B_{q,g}(\Delta^2) P^{(\mu} i \sigma^{\nu)\alpha} \Delta_{\alpha}/2M] u(p) + B_{q,g}(\Delta^2) P^{(\mu} i \sigma^{\nu)\alpha} \Delta_{\alpha}/2M] u(p)$

• Conservation laws - zero Anomalous Gravitomagnetic Moment : $\mu_G = J$ (g=2) $P_{e,q} = A_{e,q}(0)$

$$\begin{split} P_{q,g} &= A_{q,g}(0) & A_{q}(0) + A_{q}(0) = 1 \\ J_{q,g} &= \frac{1}{2} \left[A_{q,g}(0) + B_{q,g}(0) \right] & A_{q}(0) + B_{q}(0) + A_{g}(0) + B_{g}(0) = 1 \end{split}$$

- May be extracted from high-energy experiments/NPQCD calculations
- Describe the partition of angular momentum between quarks and gluons
- Describe interaction with both classical and TeV gravity – similar t-dependence to EM FF

Electromagnetism vs Gravity

- Interaction field vs metric deviation
 - $M = \langle P' | J^{\mu}_{q} | P \rangle A_{\mu}(q) \qquad \qquad M = \frac{1}{2} \sum_{q,G} \langle P' | T^{\mu\nu}_{q,G} | P \rangle h_{\mu\nu}(q)$
- Static limit

$$\langle P|J^{\mu}_{q}|P\rangle = 2e_{q}P^{\mu} \qquad \qquad \sum_{q,G} \langle P|T^{\mu\nu}_{i}|P\rangle = 2P^{\mu}P^{\nu} \\ h_{00} = 2\phi(x)$$

$$M_0 = \langle P | J^{\mu}_q | P \rangle A_{\mu} = 2e_q M \phi(q) \qquad M_0 = \frac{1}{2} \sum_{q,G} \langle P | T^{\mu\nu}_i | P \rangle h_{\mu\nu} = 2M \cdot M \phi(q)$$

Mass as charge – equivalence principle

Gravitomagnetism

Gravitomagnetic field – action on spin – $\frac{1}{2}$ from $M = \frac{1}{2} \sum_{q,G} \langle P' | T^{\mu\nu}_{q,G} | P \rangle h_{\mu\nu}(q)$

$$\vec{H}_J = \frac{1}{2} rot \vec{g}; \ \vec{g}_i \equiv g_{0i}$$
 spin dragging twice
smaller than EM

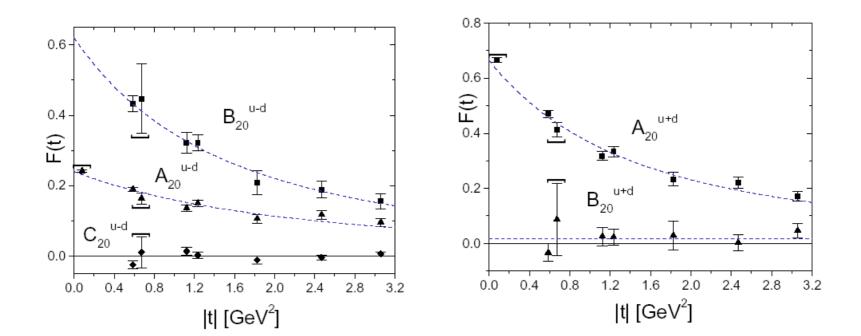
- Lorentz force similar to EM case: factor $\frac{1}{2}$ cancelled with 2 from $h_{00} = 2\phi(x)$ Larmor frequency same as EM $\vec{H}_L = rot\vec{g}$
- Orbital and Spin momenta dragging the same Equivalence principle $\omega_J = \frac{\mu_G}{I}H_J = \frac{H_L}{2} = \omega_L$

Sivers function and Extended Equivalence principle

- Second moment of E zero SEPARATELY for quarks and gluons –only in QCD beyond PT (OT, 2001) supported by lattice simulations etc.. ->
- Gluon Sivers function is small! (COMPASS, STAR, Brodsky&Gardner)
- BUT: gluon orbital momentum is NOT small: total about 1/2, if small spin – large (longitudinal) orbital momentum
- Gluon Sivers function should result from twist 3 correlator of 3 gluons: remains to be proved!

Generalization of Equivalence principle

 Various arguments: AGM 0 separately for quarks and gluons – most clear from the lattice (LHPC/SESAM, confirmed recently)



CONCLUSIONS

- Sivers and other TMD functions contain infinite tower of twists starting from 3 – special role of moments
- Colour charge of initial/final partons crucial NO factorization in naive sense (cf Abelian model of Collips&Qiu)
- Transverse momentum dependence of Sivers SSA in SIDIS and DY (PAX) is a new sensitive test of QCD
- Relation of Sivers function to twist 3 in DIS: Reasonable magnitude, but problems with flavor dependence.Bochum results with suppressed singlet twist 3 supported!
- Relation of Sivers to GPD's link to Nucleon Spin and Equivalence Principle
- Problems: evolution (no WW for Sivers) and SR from twist 3.