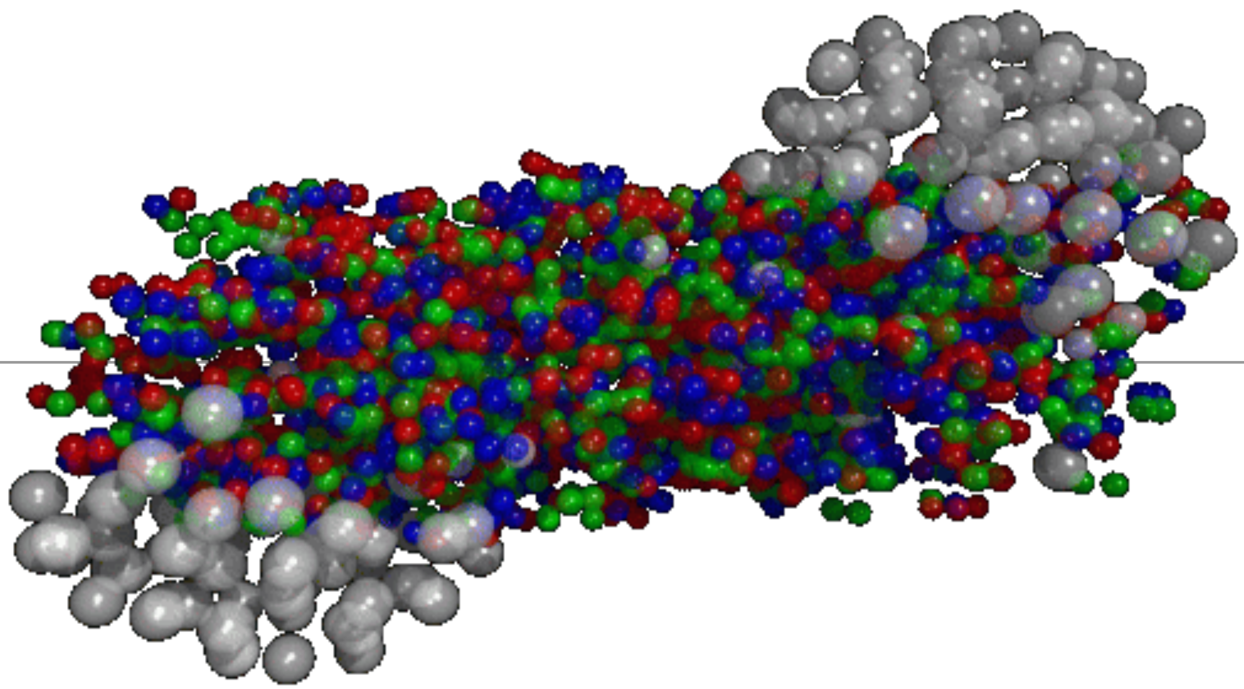
Resonance dynamics and reconstruction in hadronic channels



Marcus Bleicher

Bilateral NICA-JINR-FAIR Workshop
April 2-4th 2012, FIAS, Frankfurt

HIC for FAIR
Helmholtz International Center

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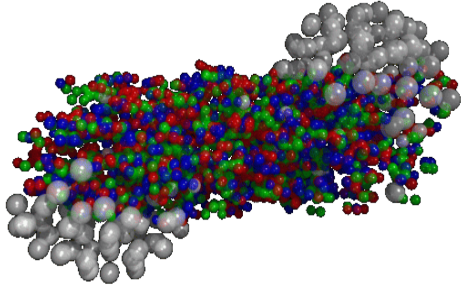
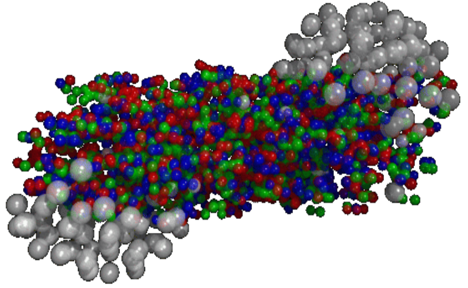


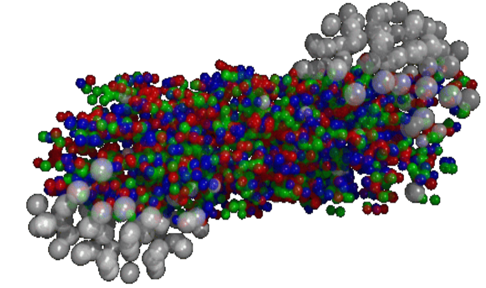
Outline

- **Quick UrQMD reminder**
- **Resonance kinematics**
  - **How deep can we look into heavy ion collisions using resonances?**
    (does high transverse momentum change anything?)
  - **Baryons @ low energies**
  - **a_1**
- **Conclusions**



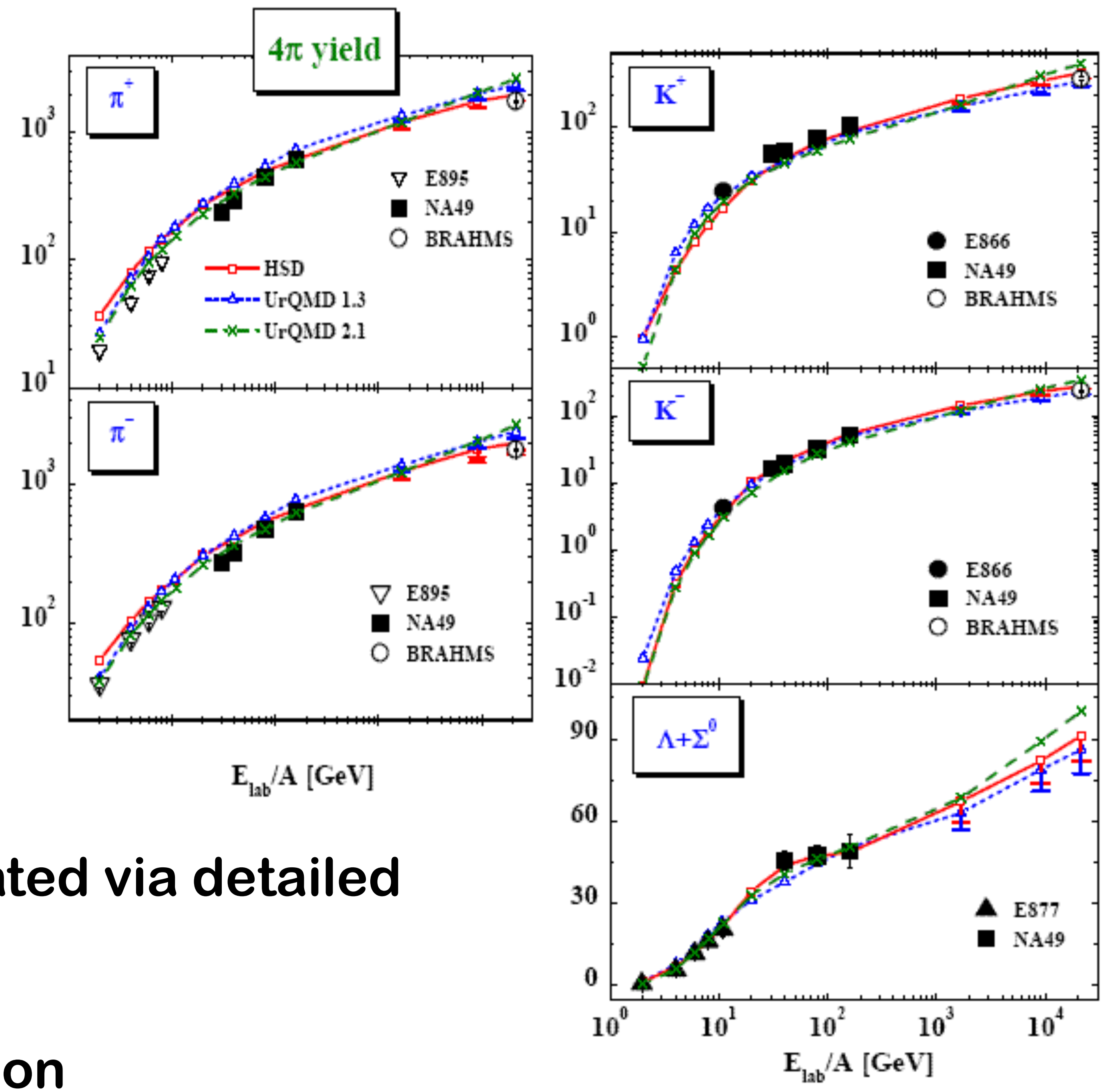
Motivation

- Several physical effects are density driven, e.g.
  - vector meson spectral function broadening
  - chiral phase transition
  - QGP phase transition
  - quarkyonic matter
  - ...



The tool - UrQMD

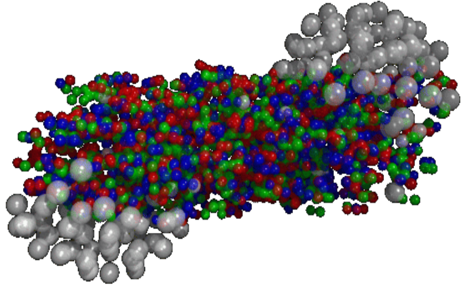
- Ultra Relativistic Quantum Molecular Dynamics
- Non equilibrium transport model
- All hadrons and resonances up to 2.2 GeV included
- Particle production via string excitation and -fragmentation
- Cross sections are fitted to available experimental data or calculated via detailed balance or the additive quark model
- Does account for canonical suppression



No explicit implementation of in-medium modifications!

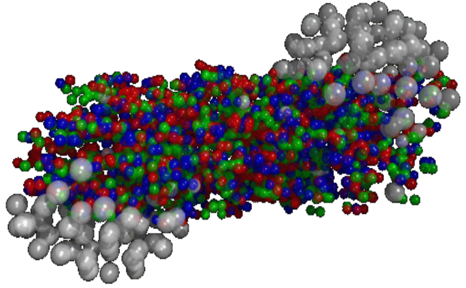
Phys.Rev.C69:054907,2004

Phys.Rev.C74:034902,2006



The tool - UrQMD

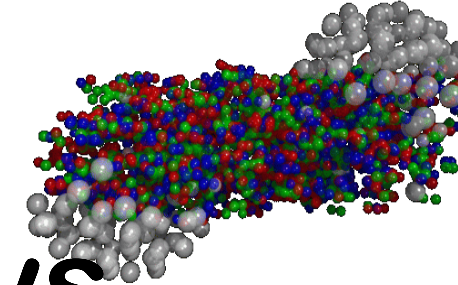
nucleon	Δ	Λ	Σ	Ξ	Ω
N_{938}	Δ_{1232}	Λ_{1116}	Σ_{1192}	Ξ_{1317}	Ω_{1672}
N_{1440}	Δ_{1600}	Λ_{1405}	Σ_{1385}	Ξ_{1530}	
N_{1520}	Δ_{1620}	Λ_{1520}	Σ_{1660}	Ξ_{1690}	
N_{1535}	Δ_{1700}	Λ_{1600}	Σ_{1670}	Ξ_{1820}	
N_{1650}	Δ_{1900}	Λ_{1670}	Σ_{1775}	Ξ_{1950}	
N_{1675}	Δ_{1905}	Λ_{1690}	Σ_{1790}	Ξ_{2025}	
N_{1680}	Δ_{1910}	Λ_{1800}	Σ_{1915}		
N_{1700}	Δ_{1920}	Λ_{1810}	Σ_{1940}		
N_{1710}	Δ_{1930}	Λ_{1820}	Σ_{2030}		
N_{1720}	Δ_{1950}	Λ_{1830}			
N_{1900}		Λ_{1890}			
N_{1990}		Λ_{2100}			
N_{2080}		Λ_{2110}			
N_{2190}					
N_{2200}					
N_{2250}					



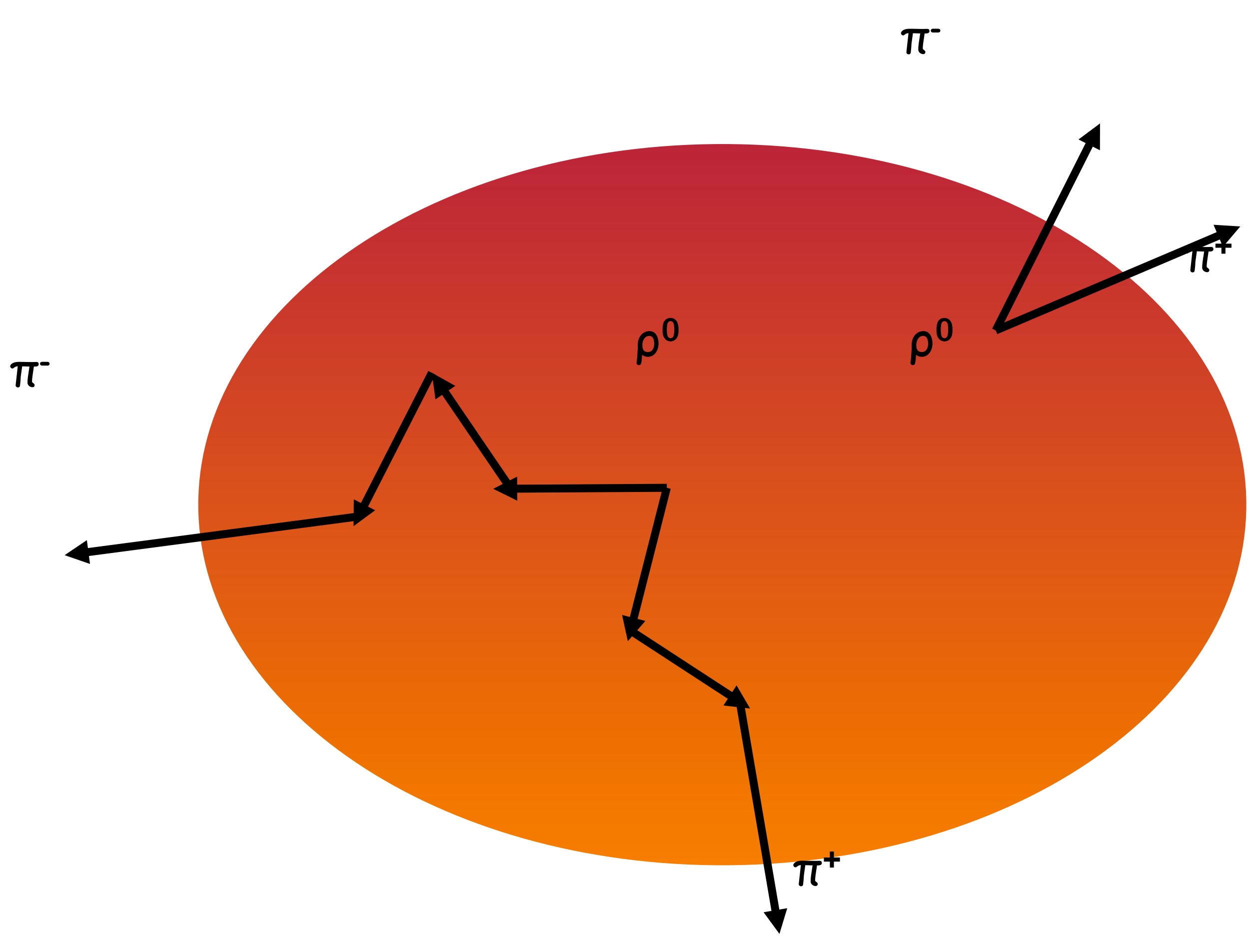
The tool - UrQMD

0^{-+}	1^{--}	0^{++}	1^{++}
π	ρ	a_0	a_1
K	K^*	K_0^*	K_1^*
η	ω	f_0	f_1
η'	ϕ	f_0^*	f_1'
1^{+-}	2^{++}	$(1^{--})^*$	$(1^{--})^{**}$
b_1	a_2	ρ_{1450}	ρ_{1700}
K_1	K_2^*	K_{1410}^*	K_{1680}^*
h_1	f_2	ω_{1420}	ω_{1662}
h_1'	f_2'	ϕ_{1680}	ϕ_{1900}

Dileptonic and hadronic decays

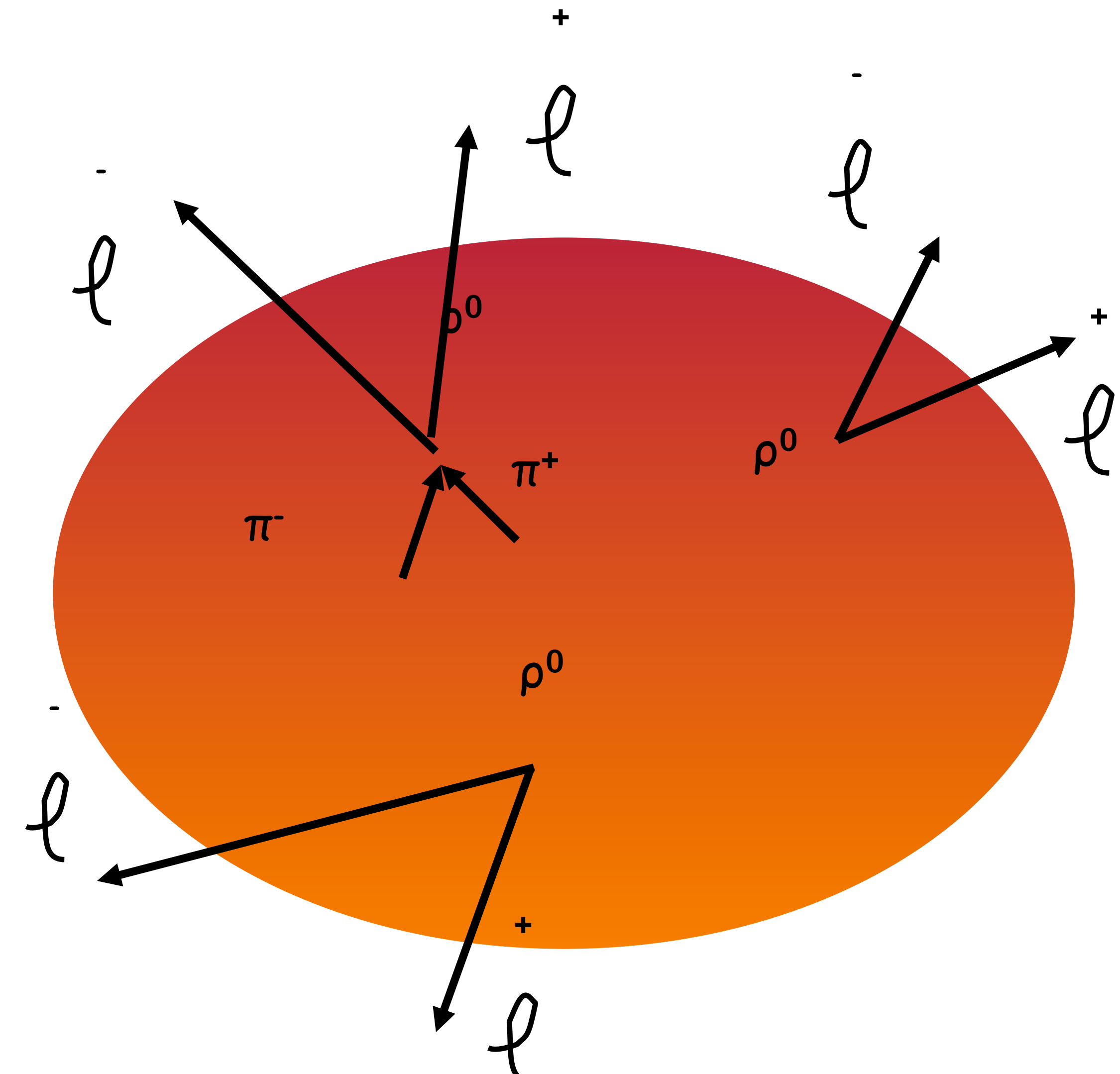


hadronic decay

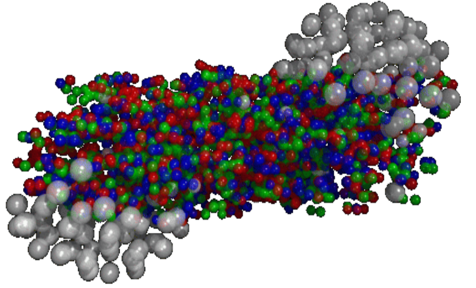


late stage

leptonic decay

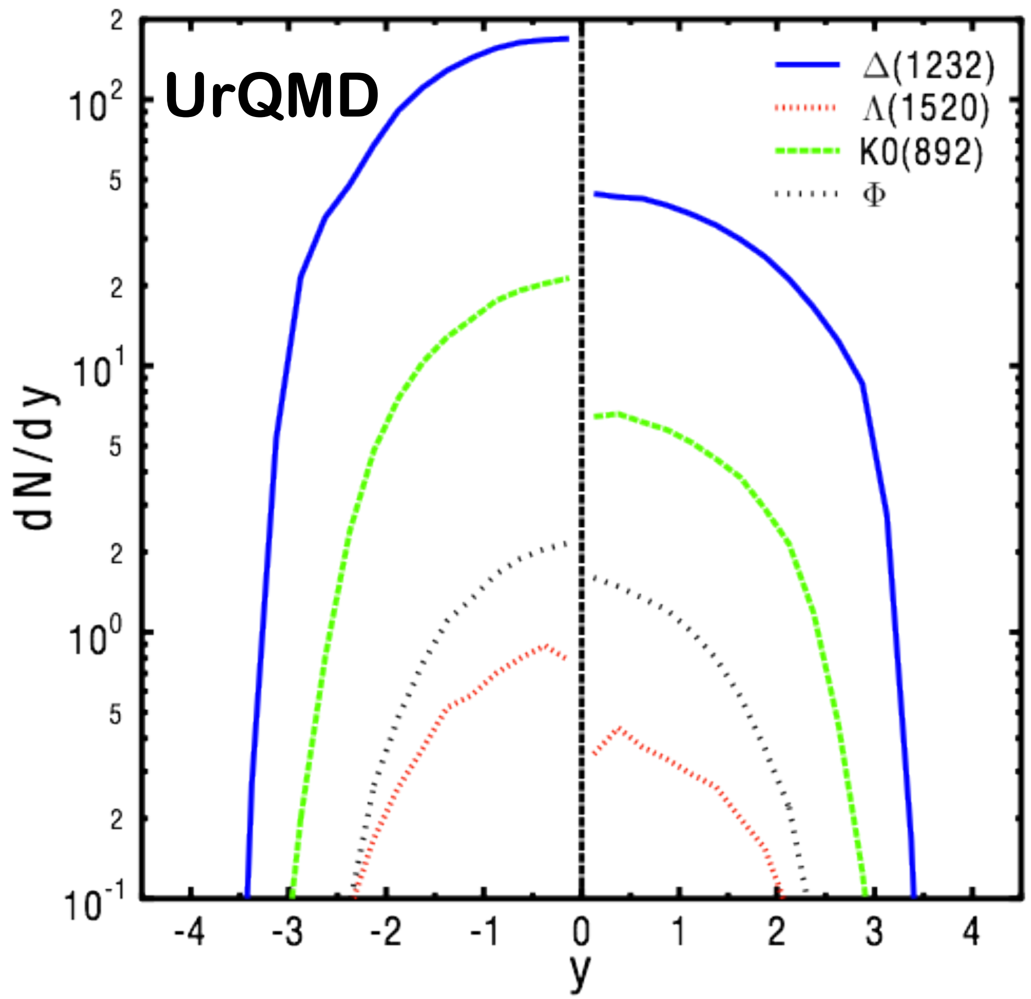


integrated collision

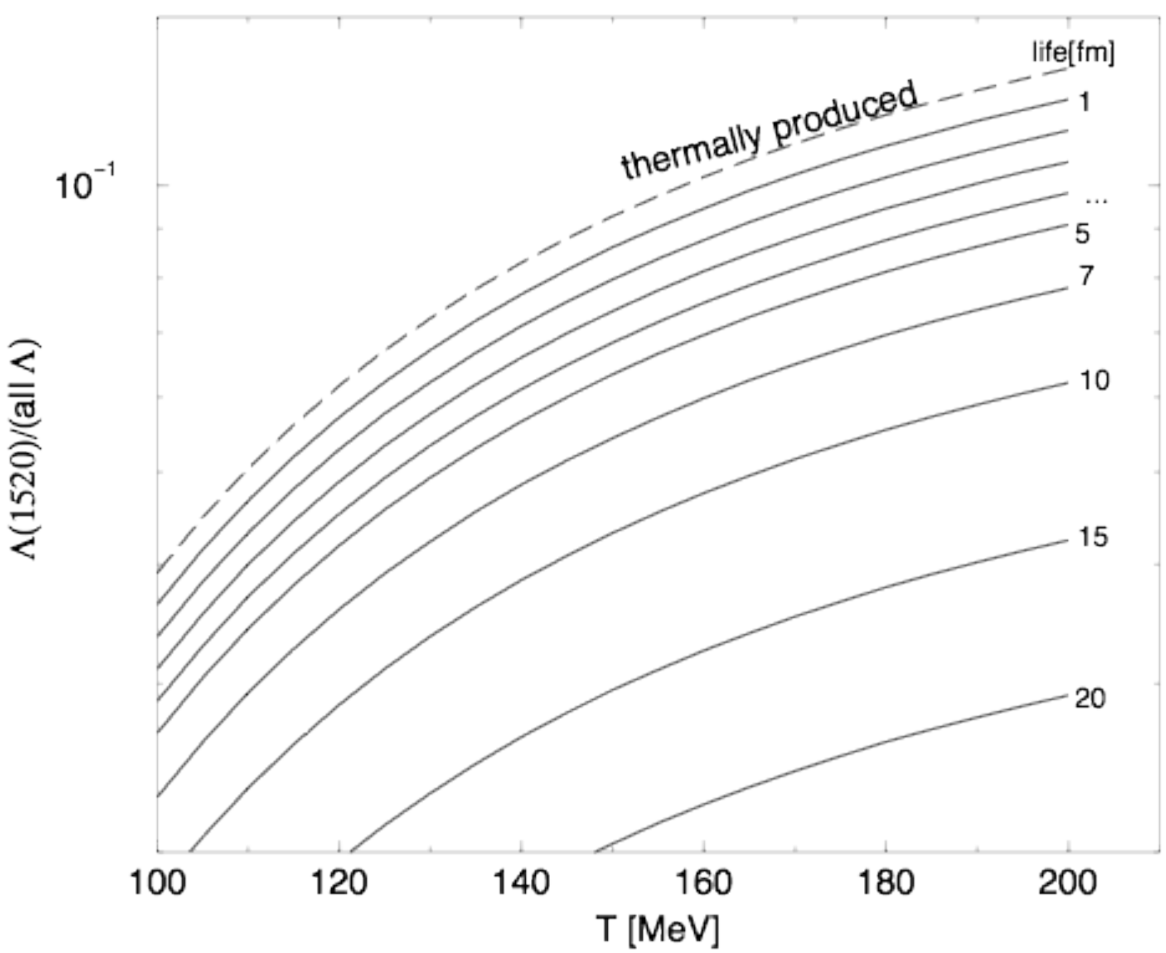


Rescattering

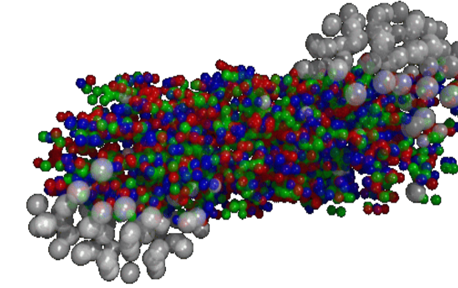
- well known effect, studied in
  - ➡ statistical hadronization models
  - ➡ transport models
  - ➡ hydrodynamical models



Bleicher, Aichelin, Phys.Lett.B530:81-87

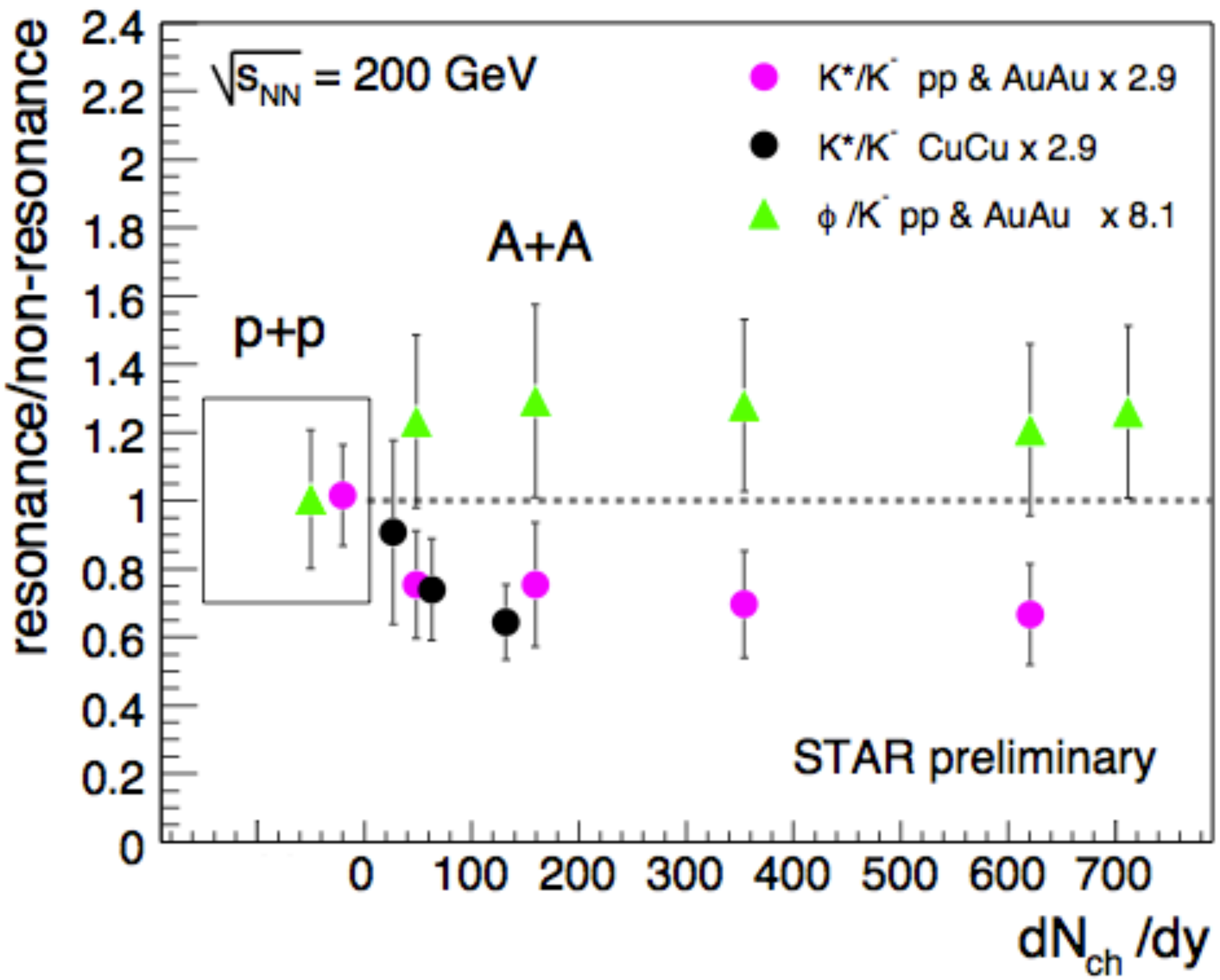


Torrieri, Rafelski, Phys.Lett.B509:239-245

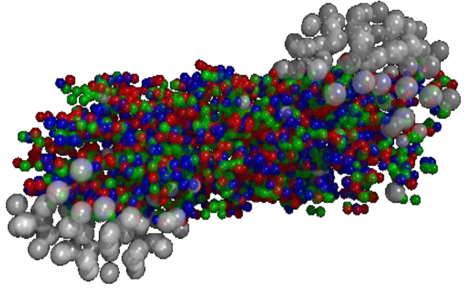


Rescattering

- well known effect, studied in
  - ➡ experiment



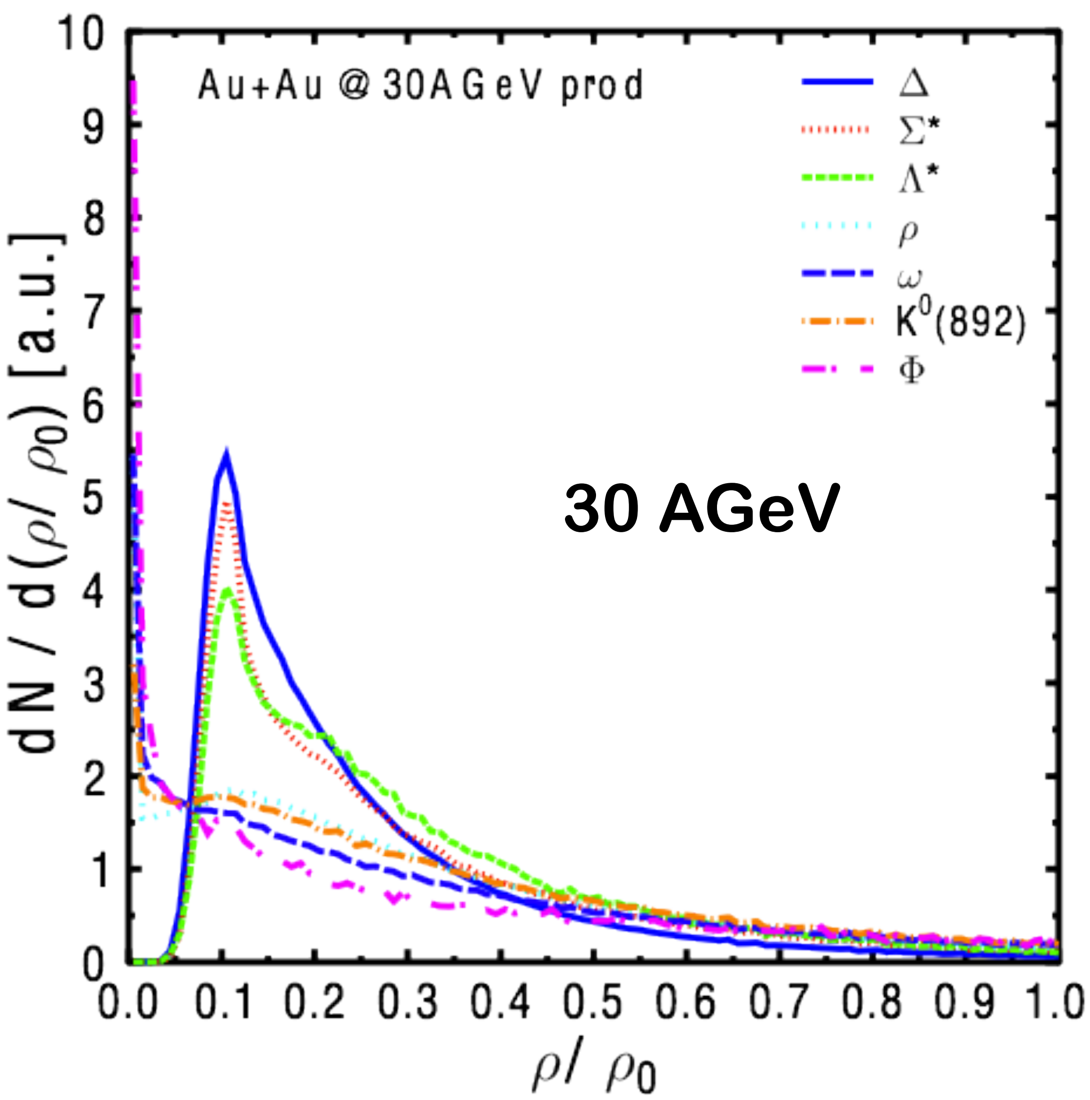
Markert et al. [STAR], J.Phys.G35:044029,2008

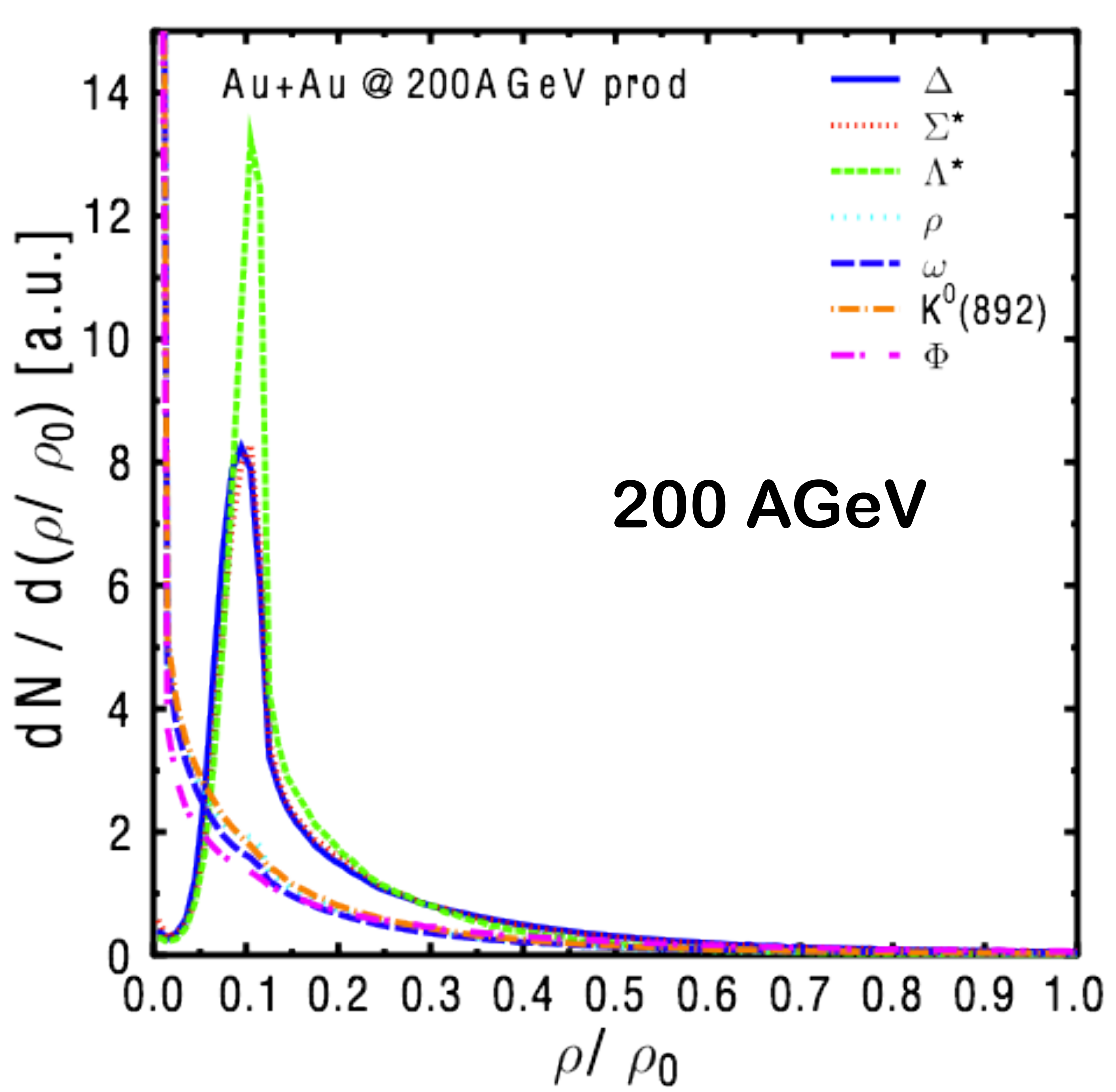


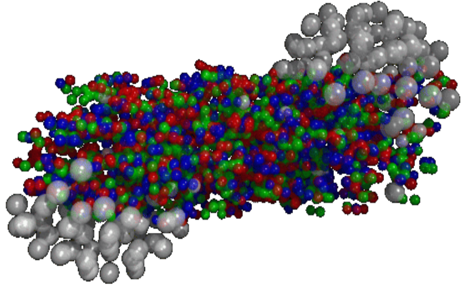
Reach in density

- Normalized density spectrum

- Most resonances originate from very low density



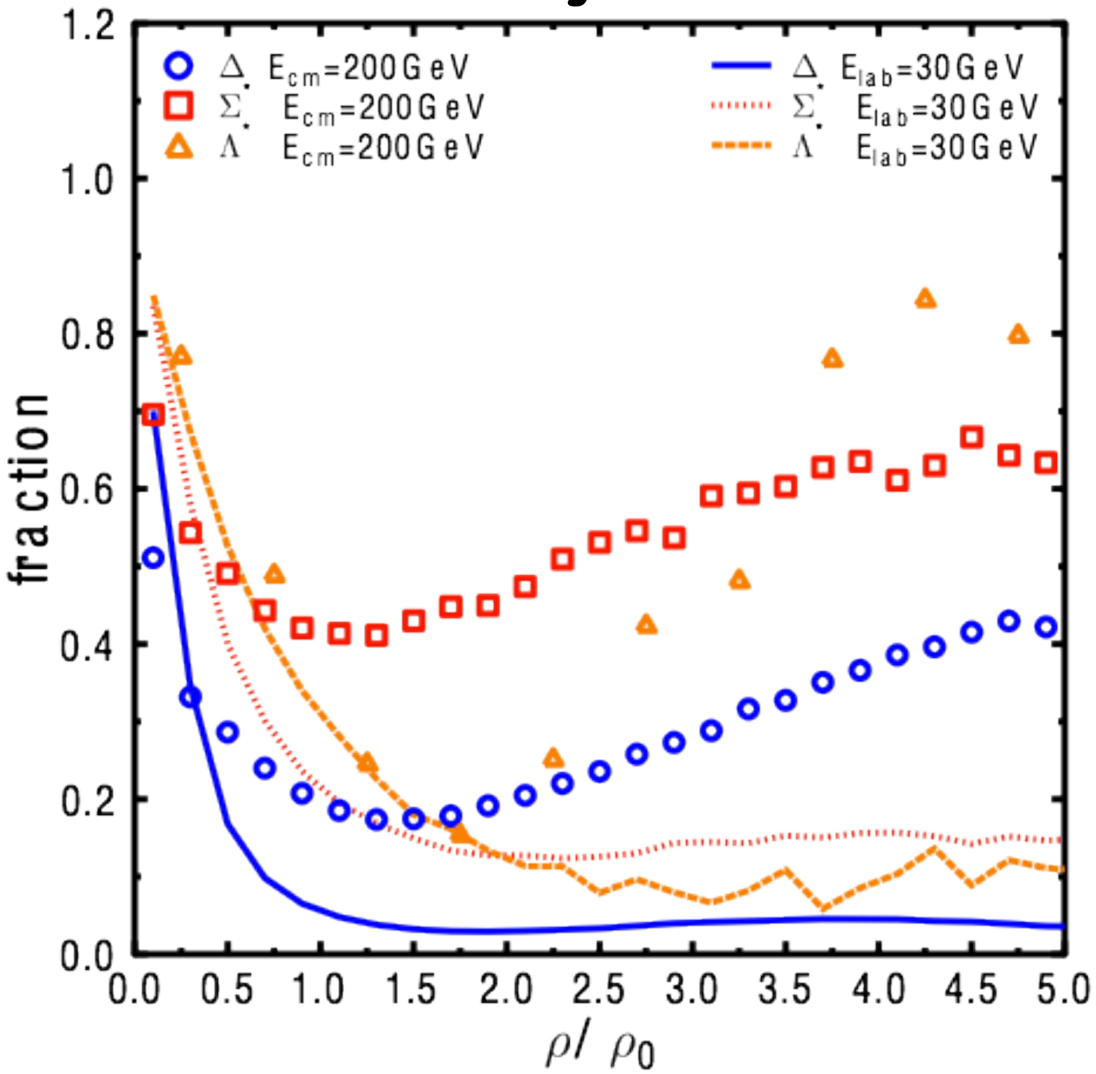




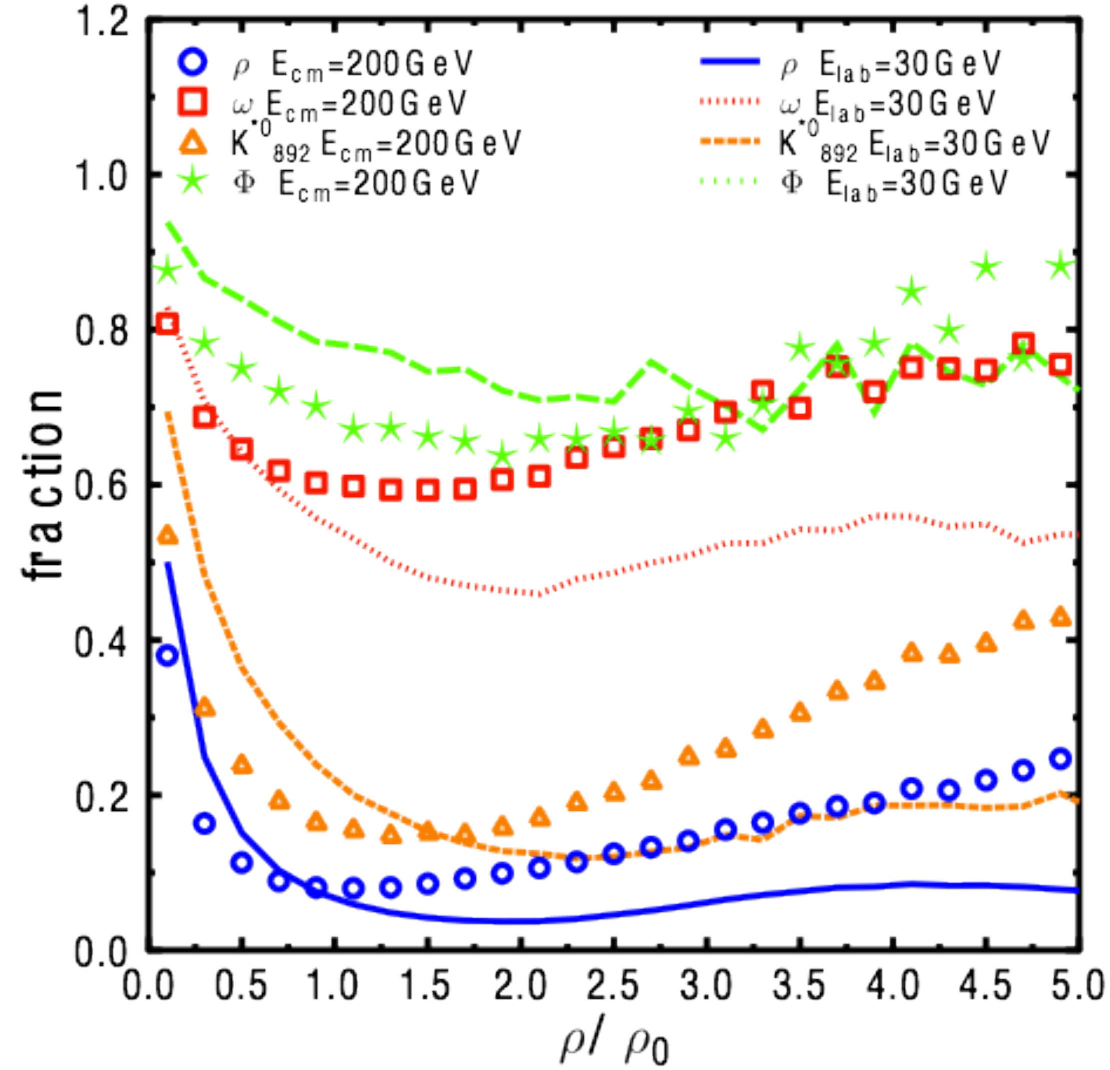
Reconstruction probability

- Probability to reconstruct resonances from a certain density

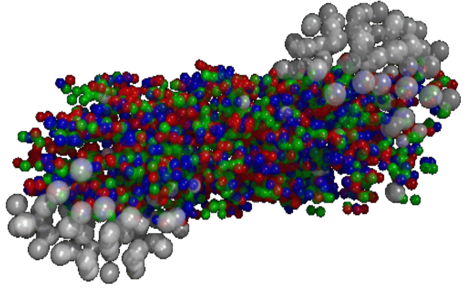
Baryons

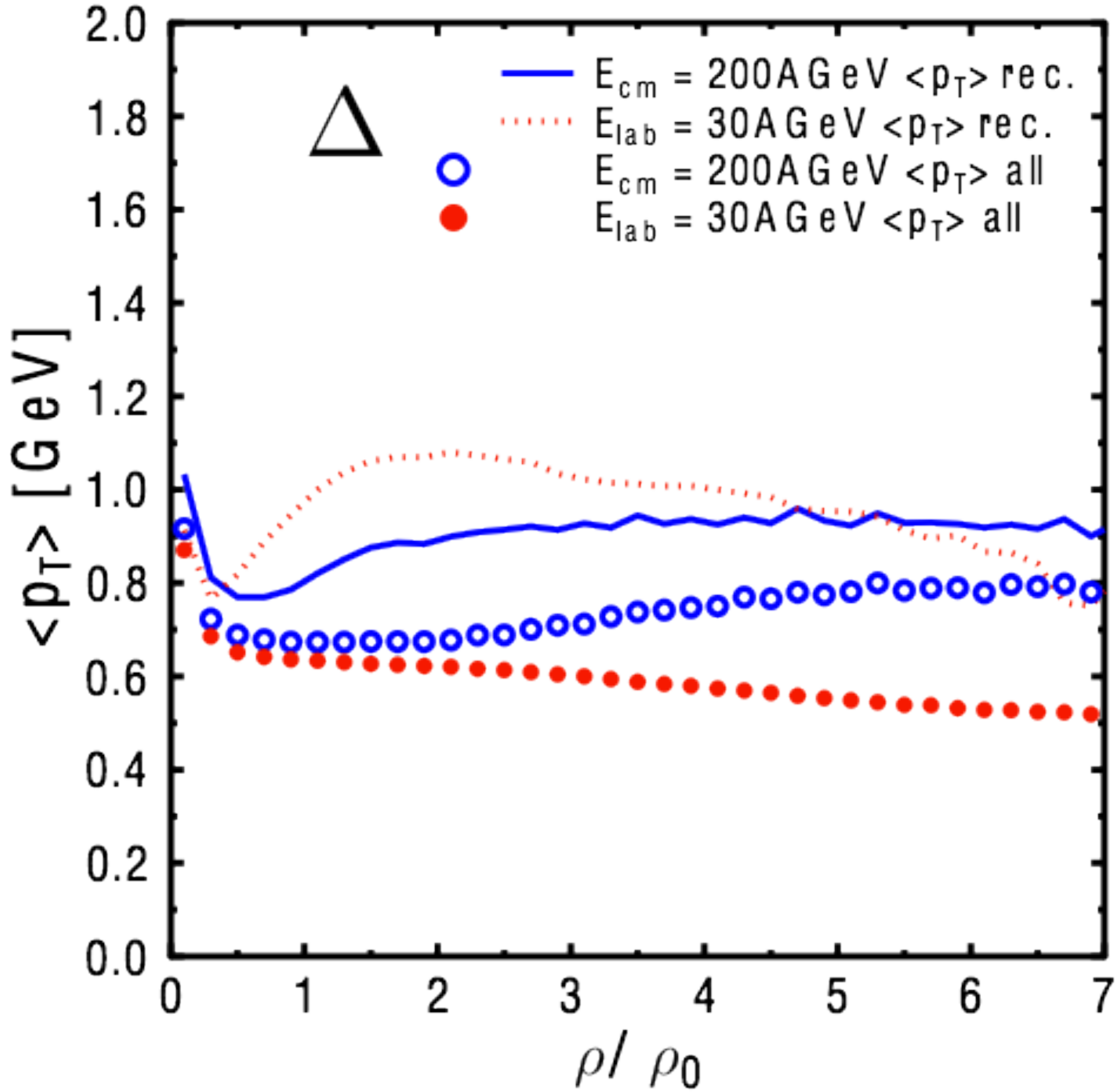


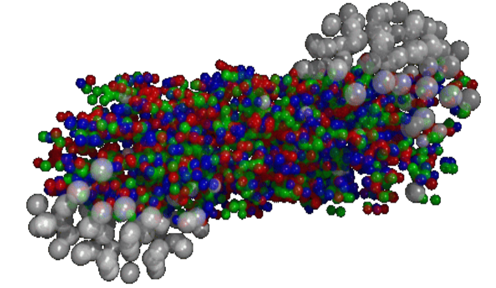
Mesons



p_T dependence



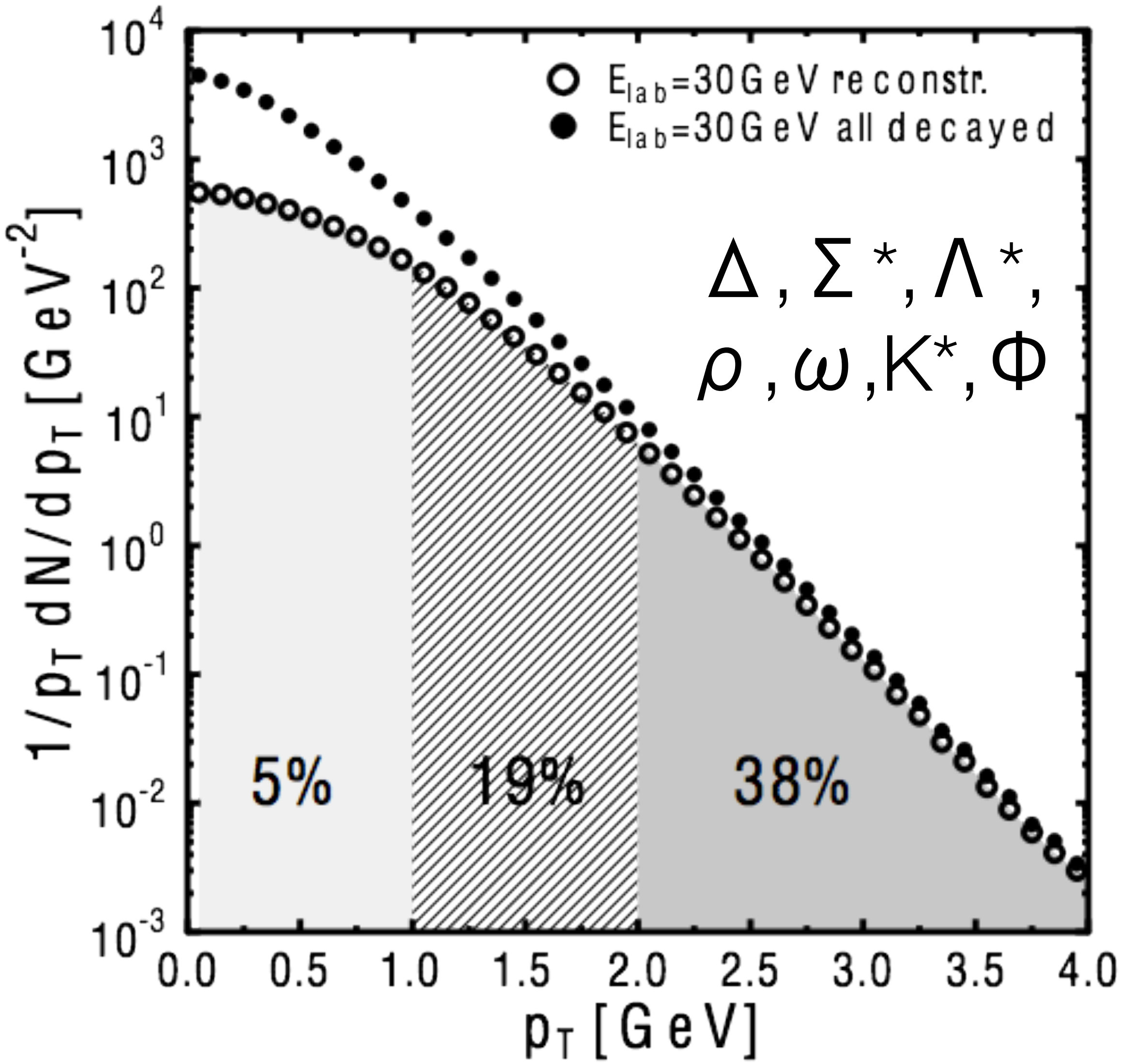


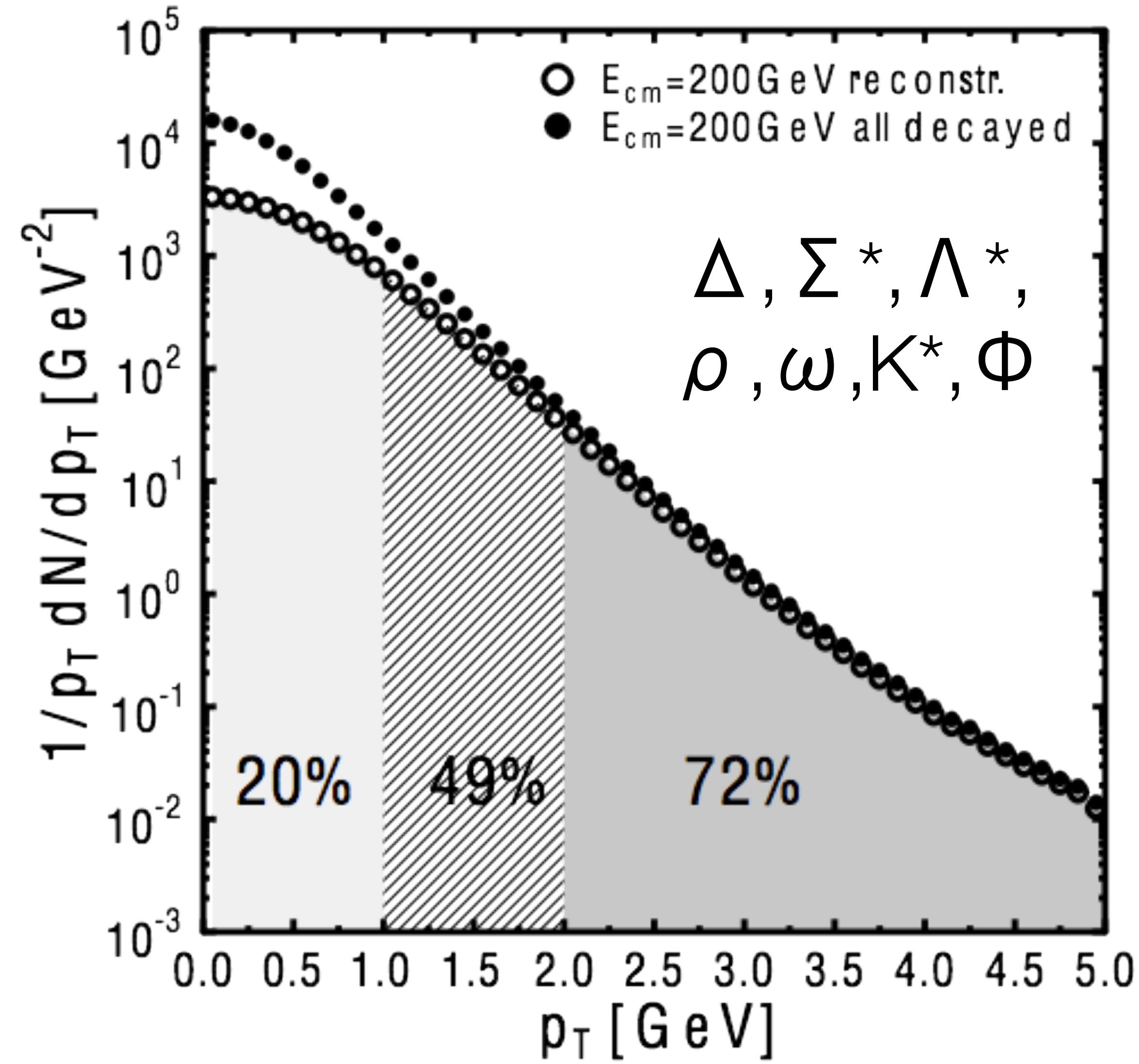


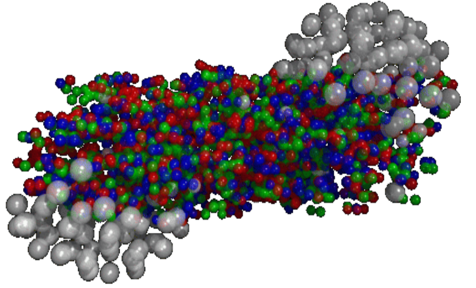
p_T dependence

- **difference in p_T spectrum between observable and all decayed**

- **percentage of reconstructable resonances produced at $\rho > 2\rho_0$ increases with p_T**



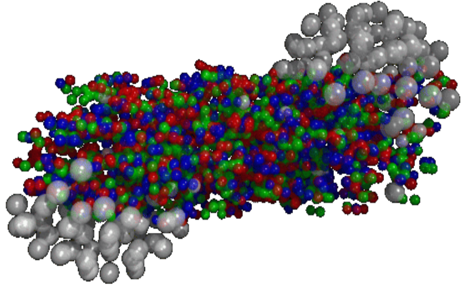




First conclusion

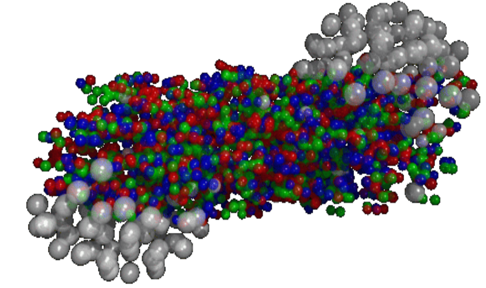
High pT resonances might shed some light on the dense phase of heavy ion collisions!

(but are they really what we want to measure?)



Baryons

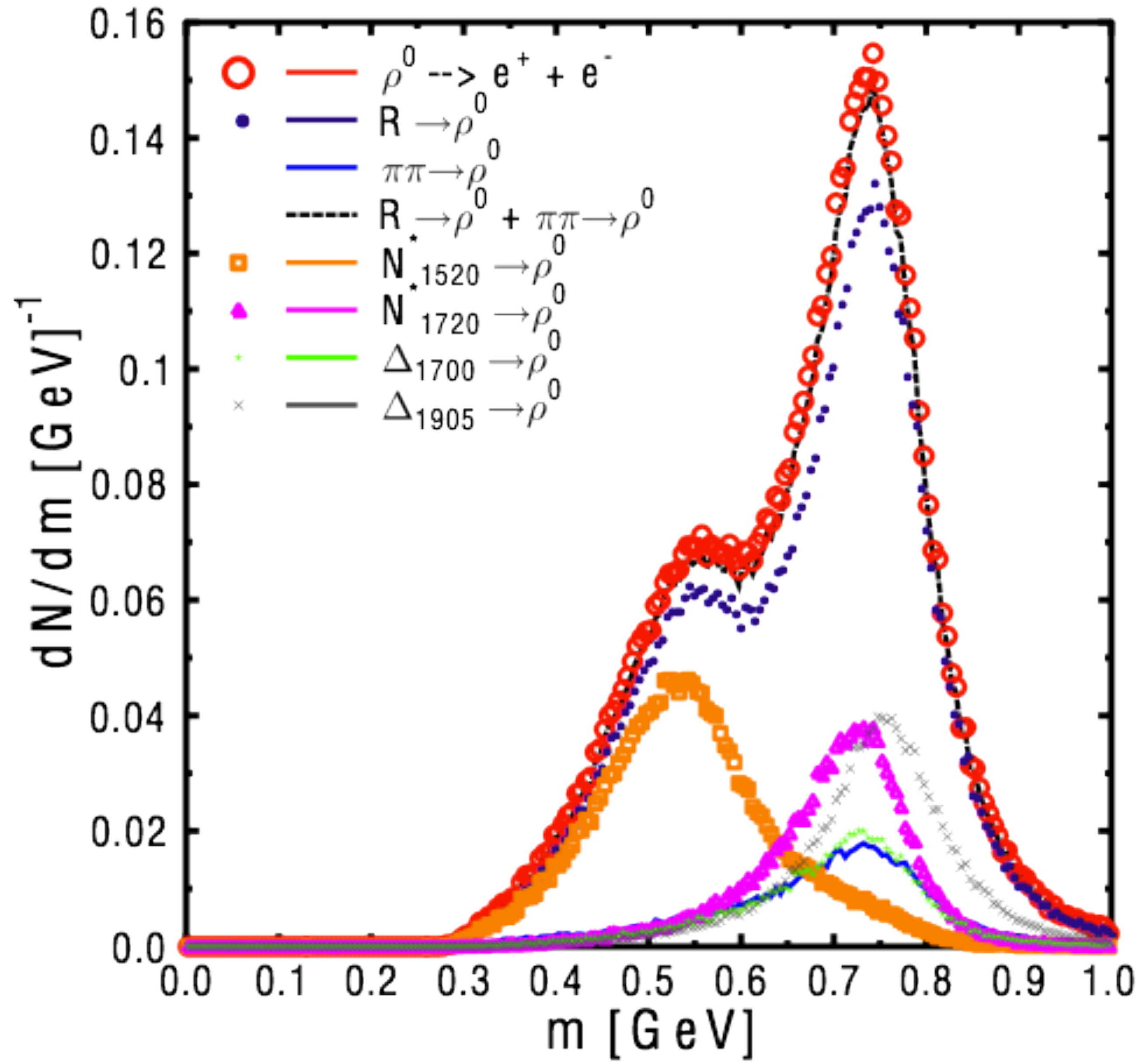
What is the deal about them at low energies?



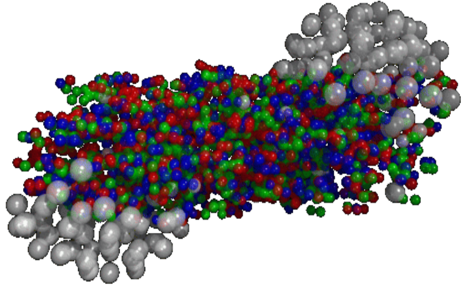
ρ meson in C+C @ 2AGeV

At low energies (~2 AGeV) contributions from baryon resonance decays are dominant.

N^*_{1520} contributes via the decay chain:
$N^*_{1520} \rightarrow N + \rho$
$\rho \rightarrow \pi^+\pi^-$ or $\rho \rightarrow e^+e^-$
to the low mass part of the ρ meson mass spectrum.



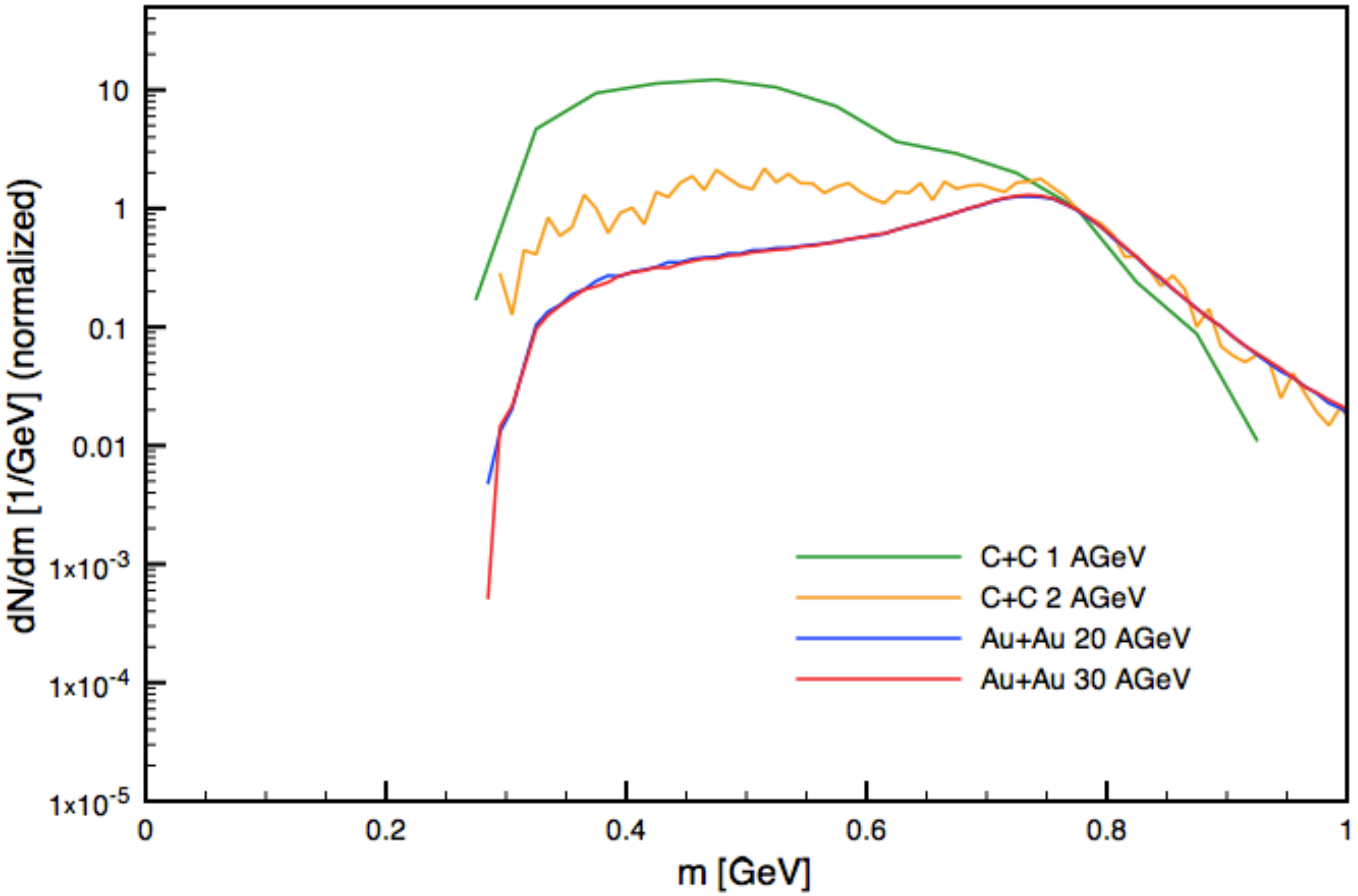
SV, M. Bleicher, Phys.Rev.C74:014902,2006

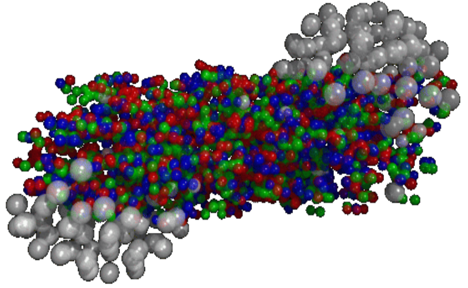


ρ meson mass distribution

At higher energies the contribution from baryonic resonance decays become less important.

Note: All curves normalized to the 770 MeV point.

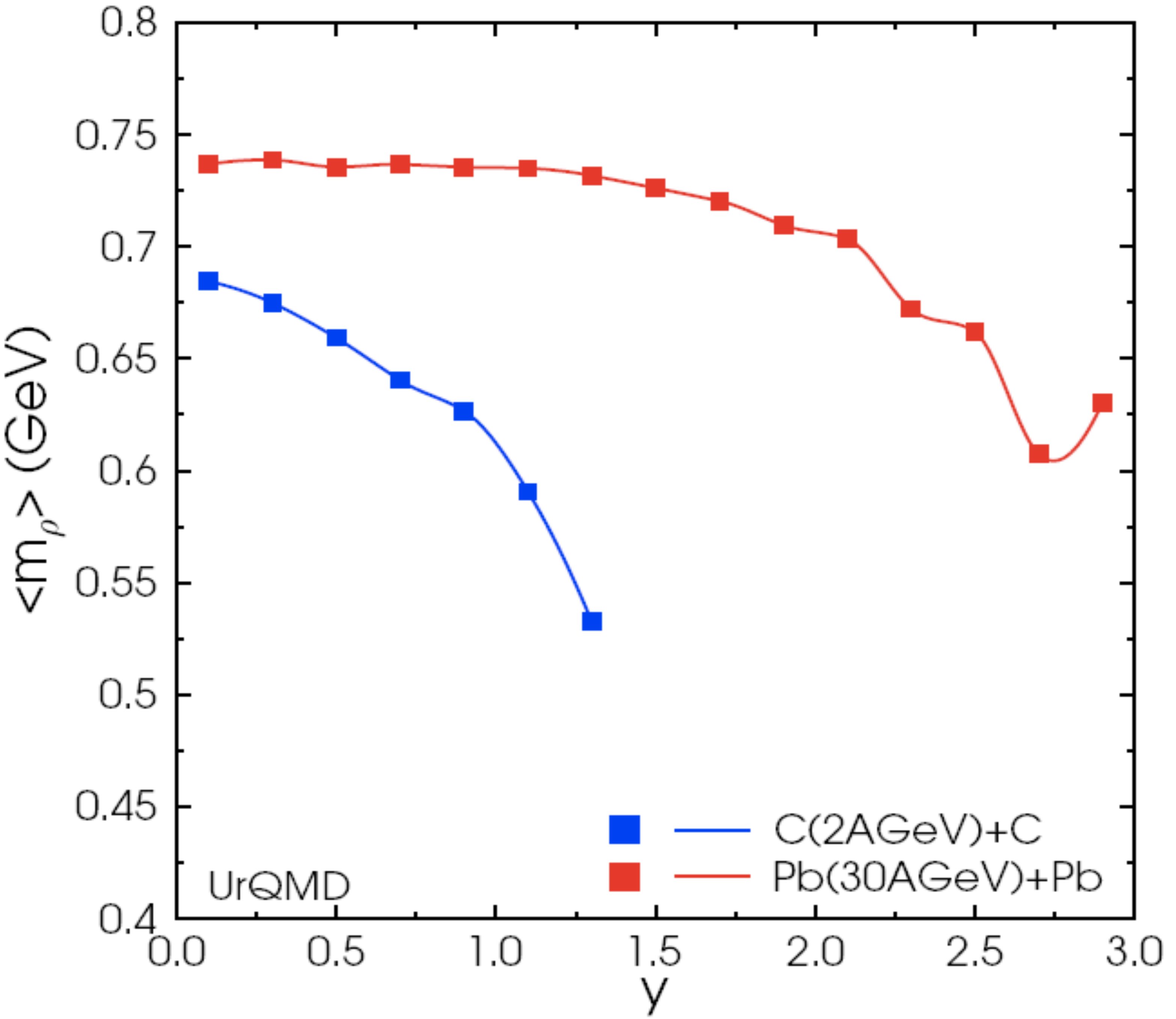


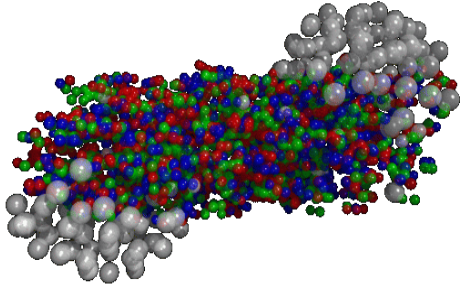


ρ meson at higher energies

Due to the dependence on the baryon density the mass of the ρ meson is rapidity dependent.

The ρ meson mass drops towards higher rapidity.

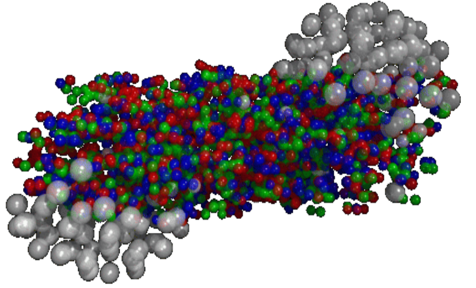




Second conclusion

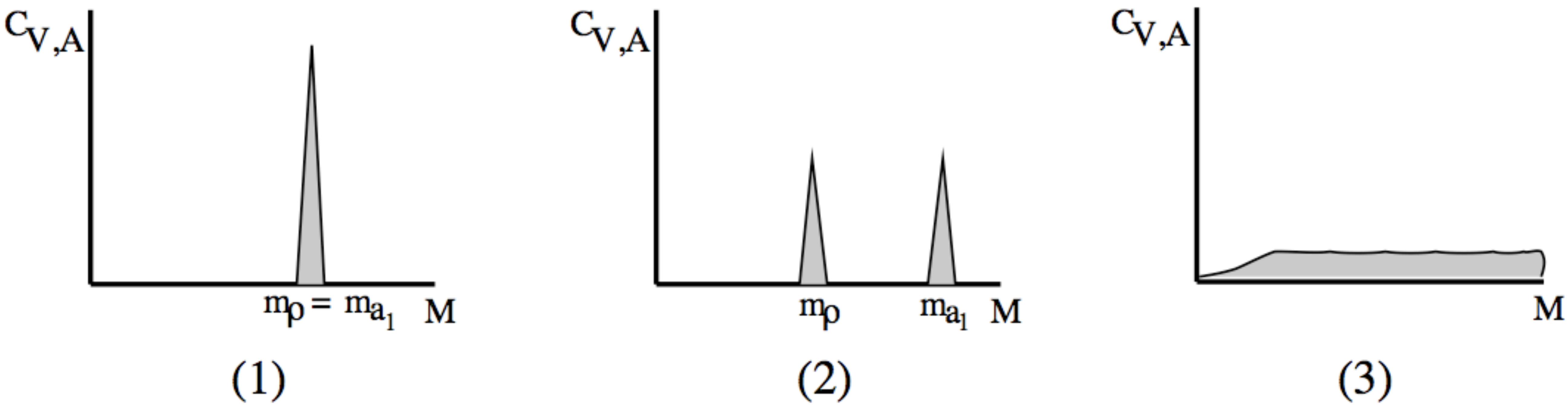
Controlling baryon kinematics is important

(otherwise some spectra seem more interesting than they are)

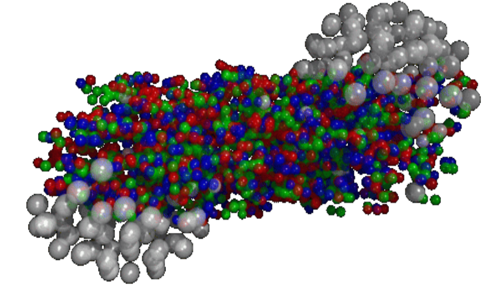


Measuring Chiral Symmetry

- **Can we observe a chirally restored phase? (and how?)**
- **What happens to the ρ meson in the medium? What happens to the a_1 meson?**
- **What can we learn from reasonable hadronic dynamics (without a chirally restored phase)?**



V. Koch, Int.J.Mod.Phys.E6:203-250,1997



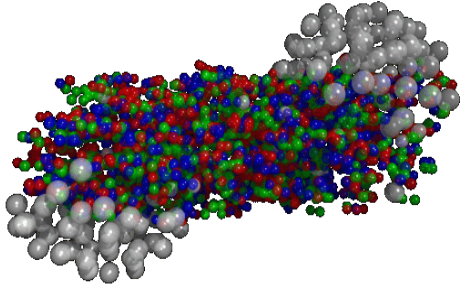
a_1 meson

The a_1 meson mass is expected to be equal to the mass of the ρ meson, in case of chiral symmetry restoration.

Problem:
It is hard to measure.

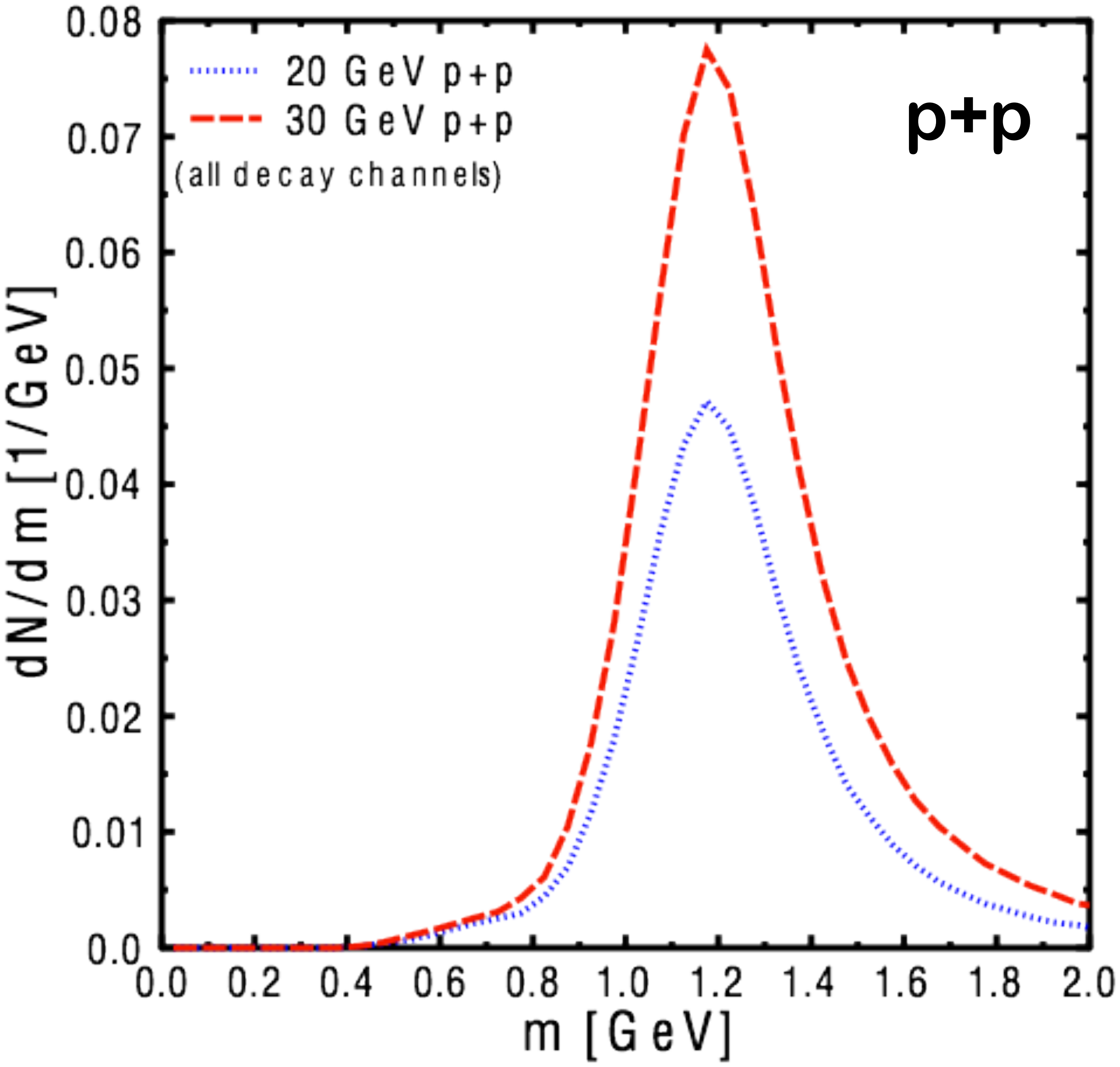
$a_1(1260)$ DECAY MODES

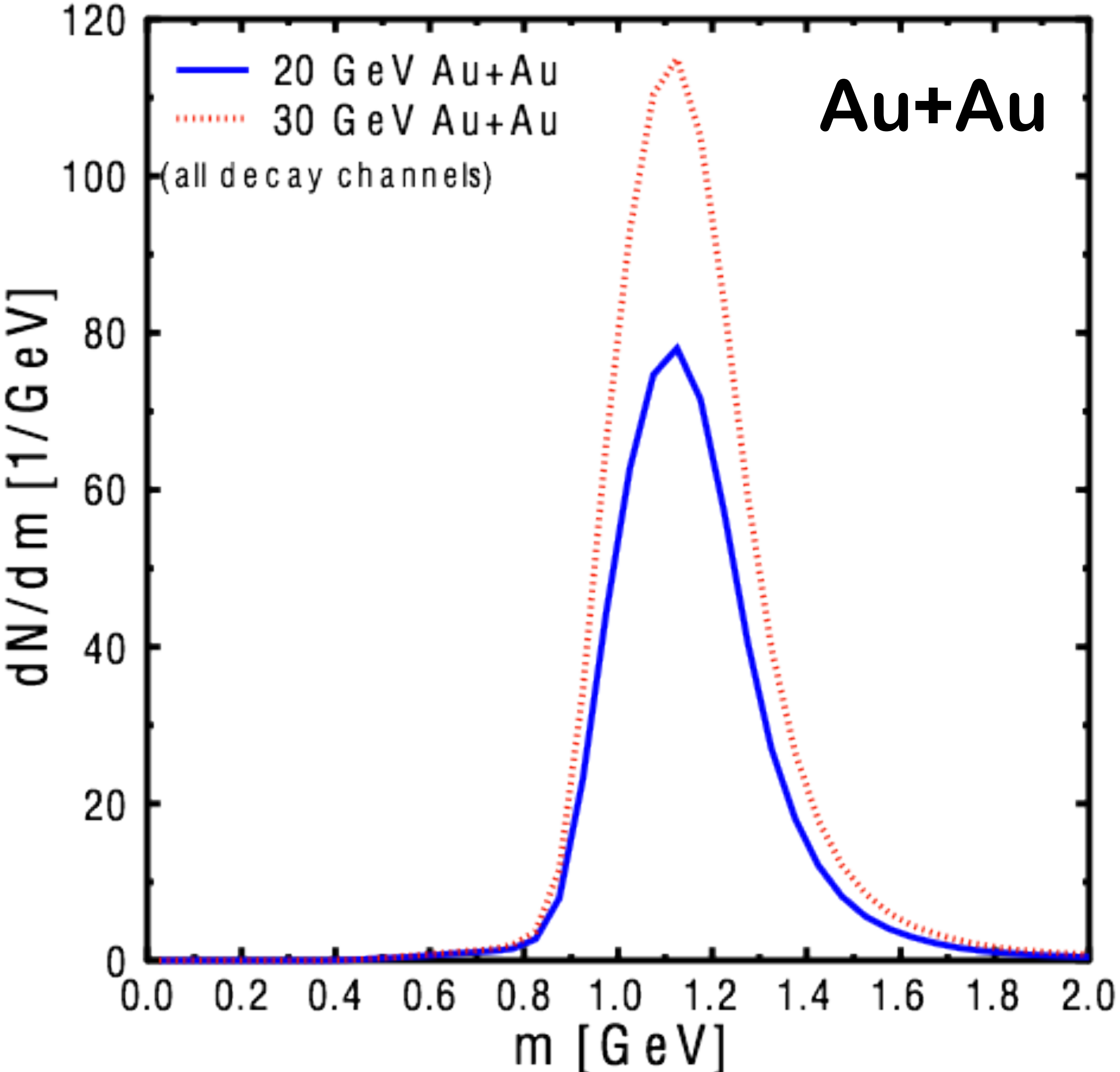
	Mode	Fraction (Γ_i/Γ)
Γ_1	$\pi^+\pi^-\pi^0$	
Γ_2	$\pi^0\pi^0\pi^0$	
Γ_3	$(\rho\pi)_{S-\text{wave}}$	seen
Γ_4	$(\rho\pi)_{D-\text{wave}}$	seen
Γ_5	$(\rho(1450)\pi)_{S-\text{wave}}$	seen
Γ_6	$(\rho(1450)\pi)_{D-\text{wave}}$	seen
Γ_7	$\sigma\pi$	seen
Γ_8	$f_0(980)\pi$	not seen
Γ_9	$f_0(1370)\pi$	seen
Γ_{10}	$f_2(1270)\pi$	seen
Γ_{11}	$K\overline{K}^*(892)+$ c.c.	seen
Γ_{12}	$\pi\gamma$	seen



a_1 meson

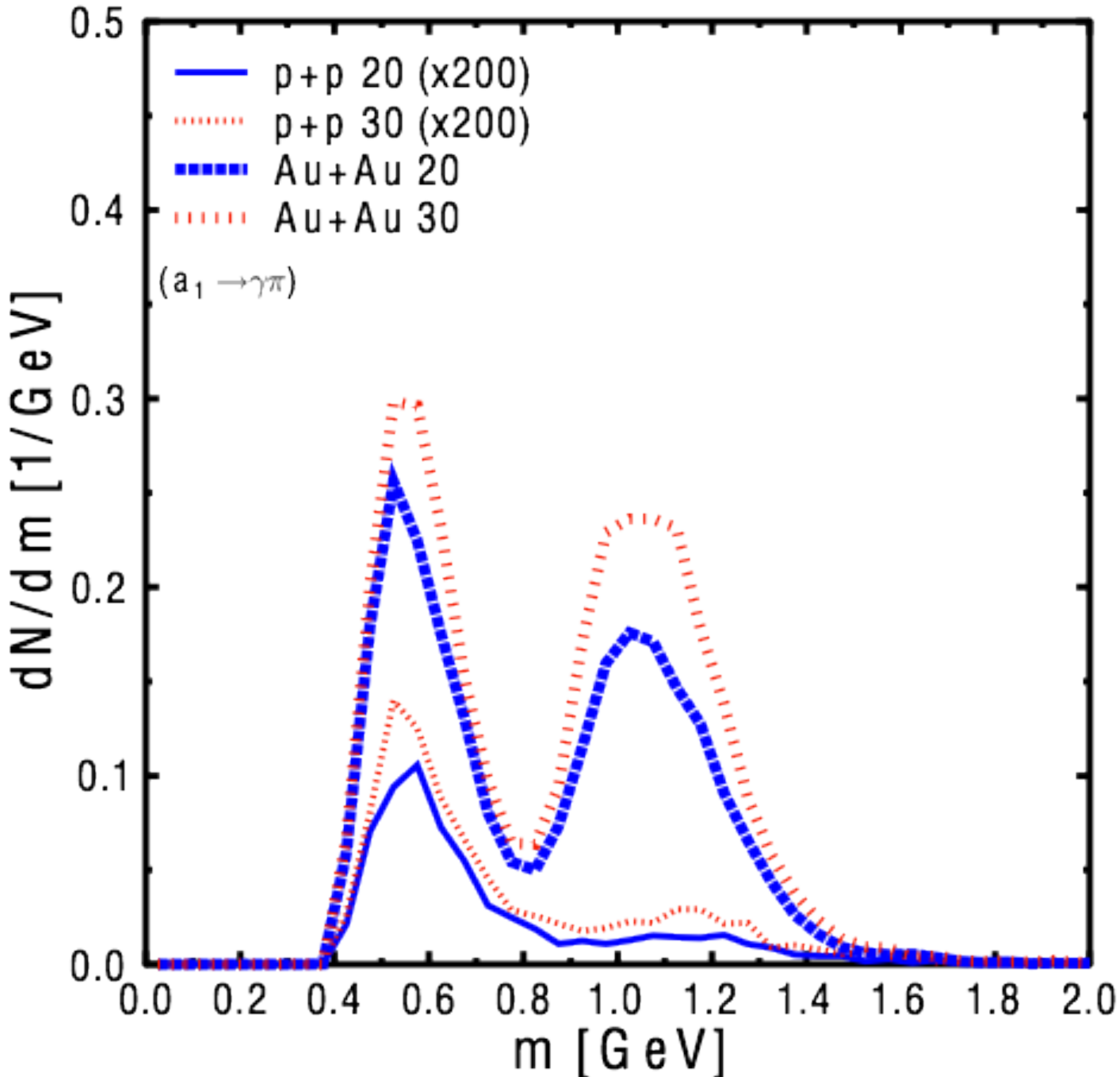
Idea: Check the mass distribution from the transport code.

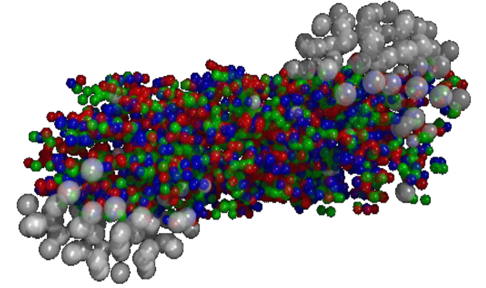




a_1 meson

Next: trigger on the decay channel $a_1 \rightarrow \gamma\pi$ (assumed width = 640keV)





a_1 meson

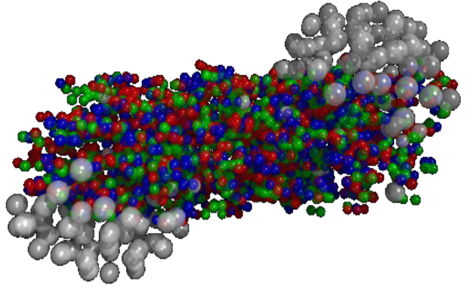
→ **Mass dependent branching ratios**

$$\Gamma_{i,j}(M) = \Gamma_R^{i,j} \frac{M_R}{M} \left(\frac{\langle p_{i,j}(M) \rangle}{\langle p_{i,j}(M_R) \rangle} \right)^{2l+1} \frac{1.2}{1 + 0.2 \left(\frac{\langle p_{i,j}(M) \rangle}{\langle p_{i,j}(M_R) \rangle} \right)^{2l}}$$

Low mass a_1 favors $\gamma\pi$ decay, not $\rho\pi$

Trigger on $a_1 \rightarrow \gamma\pi$ = trigger on low mass a_1 mesons

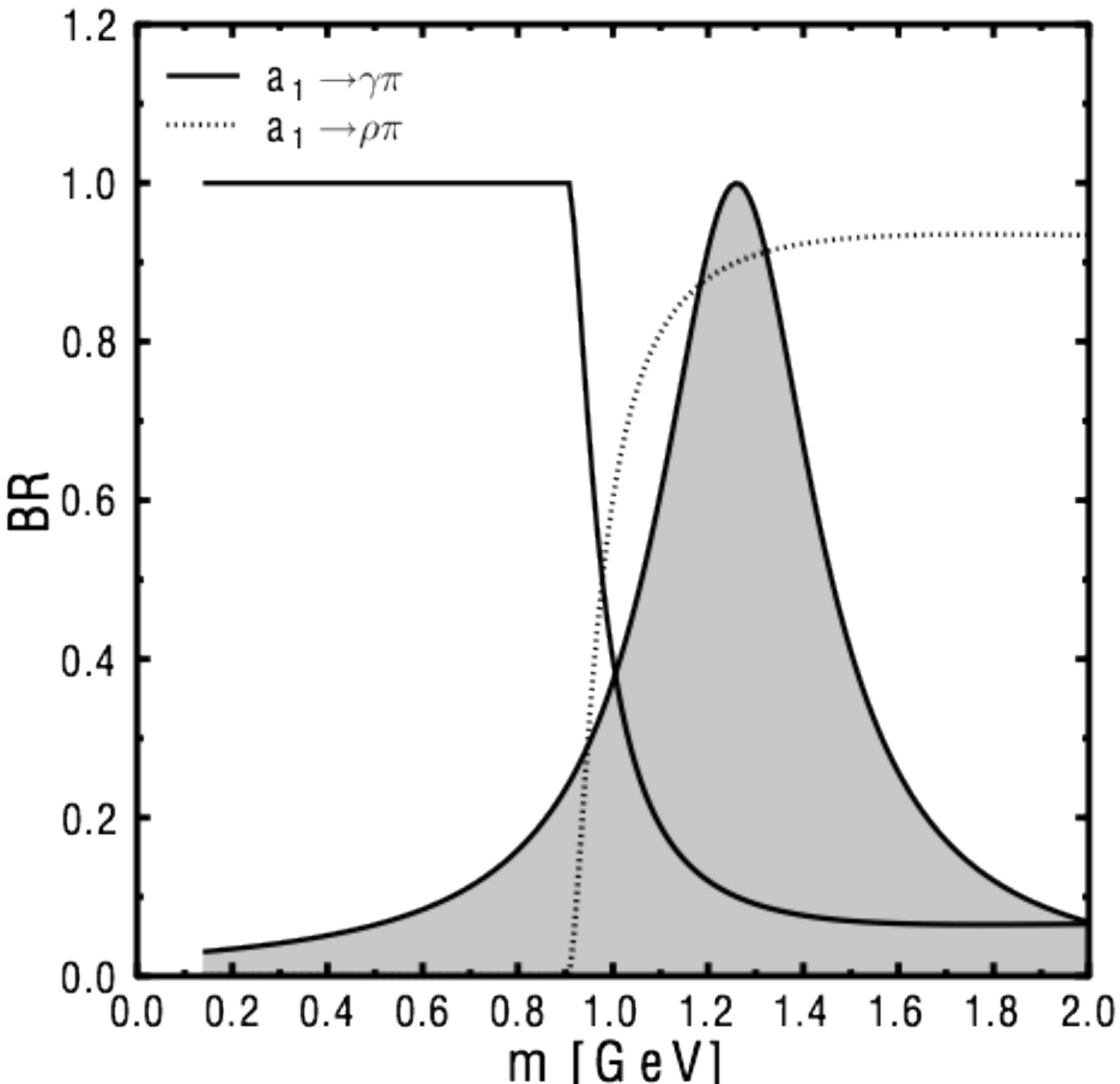
H. Sorge, Phys.Rev.C52:3291,1995

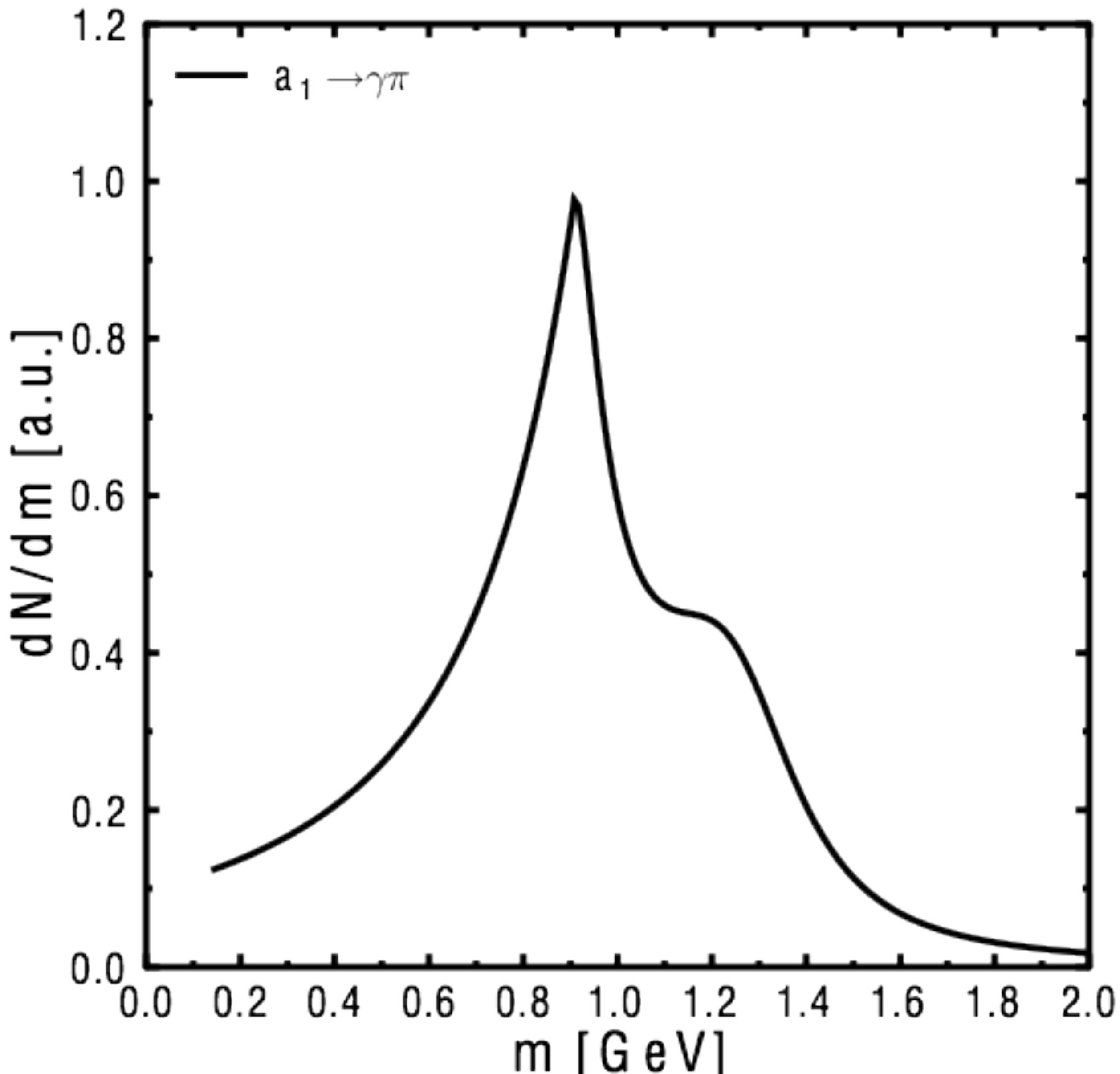


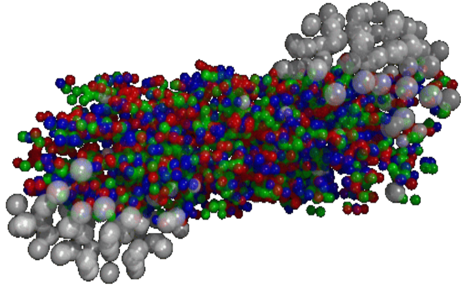
a_1 meson – a sketch!

Below 900 MeV $\gamma\pi$ decay is dominant, $\rho\pi$ is kinematically suppressed.

Branching ratio folded with BW distribution

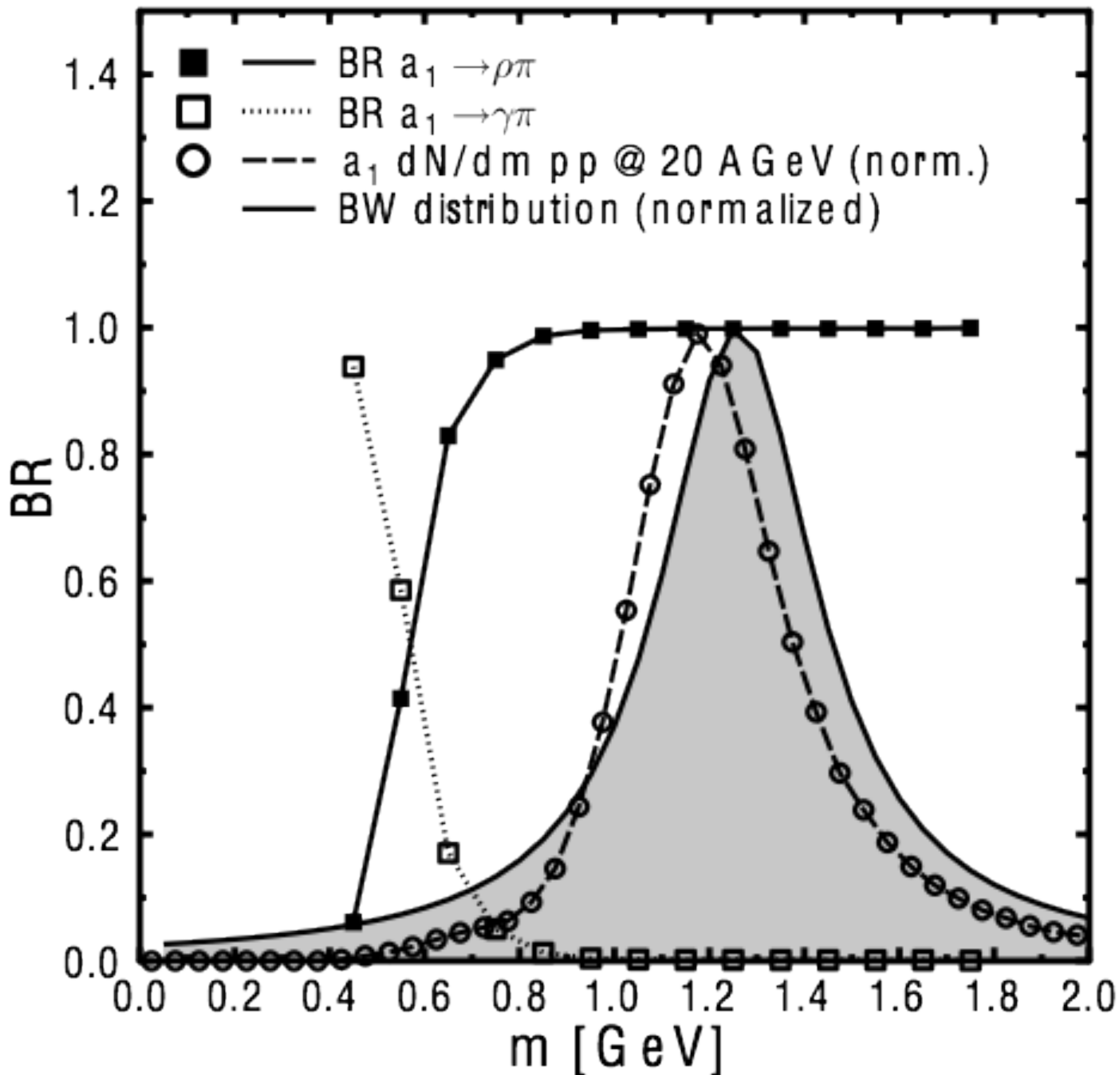




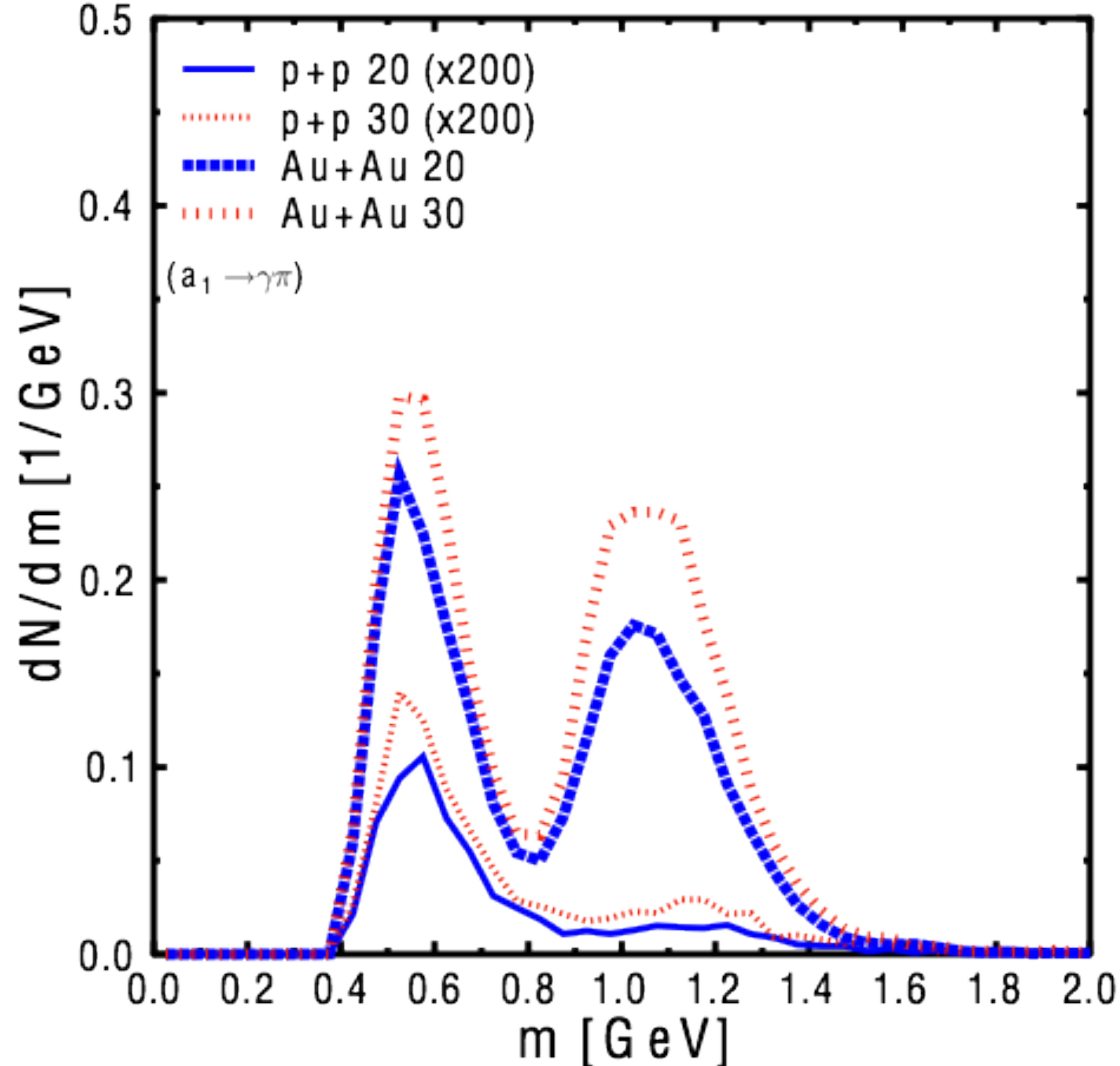


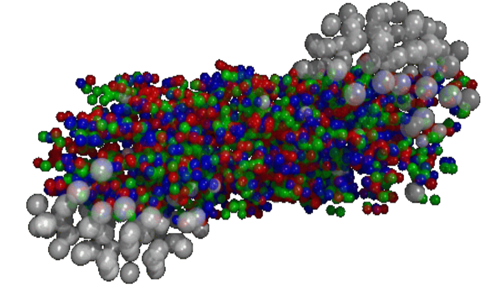
a_1 meson

Full model calculation



a_1 meson


p+p 20 (x200)
p+p 30 (x200)
Au+Au 20
Au+Au 30
$(a_1 \rightarrow \gamma\pi)$
dN/dm [1/GeV]
m [GeV]
0.0
0.1
0.2
0.3
0.4
0.5
0.0
0.2
0.4
0.6
0.8
1.0
1.2
1.4
1.6
1.8
2.0



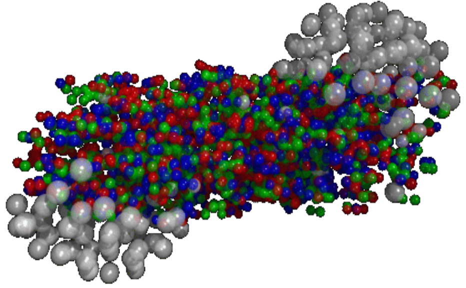
Take home messages

- Experimentally reconstructable resonances are sensitive to the high density region if measured at high p_T
  (however, medium modification gets weaker)

- Beware of baryon kinematics
  (i.e. even calculations w/o explicit medium modification have a medium)

- $a_1 \rightarrow \gamma\pi$ might not be the golden channel
  (mass dependent branching ratios may surprise you)

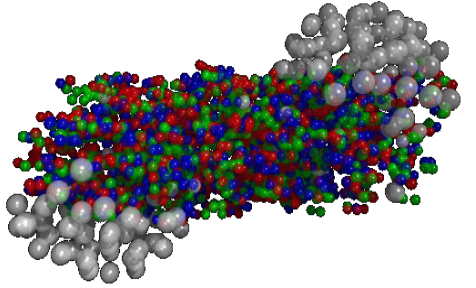
the end



Take home messages

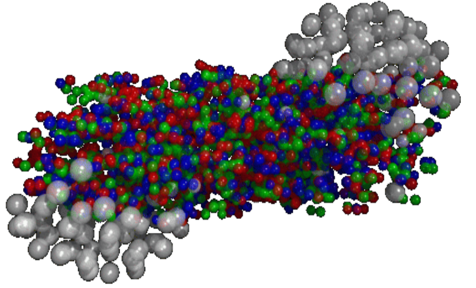
- Experimentally reconstructable resonances are not sensitive to the high density region unless measured at high p_T
- Beware of baryons kinematics
- $a_1 \rightarrow \gamma\,\pi$ might not be the golden channel

Thanks!



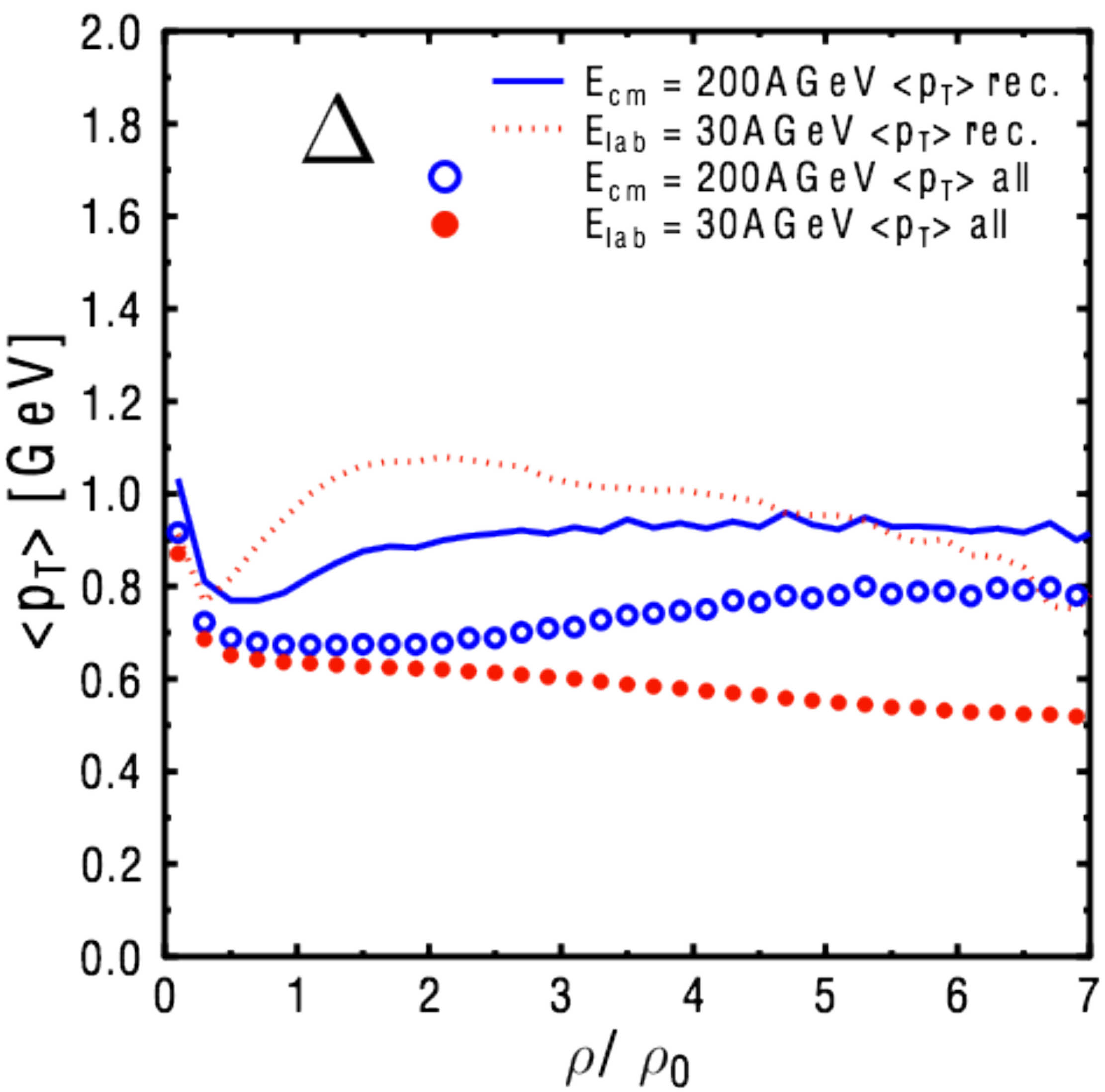
Backup

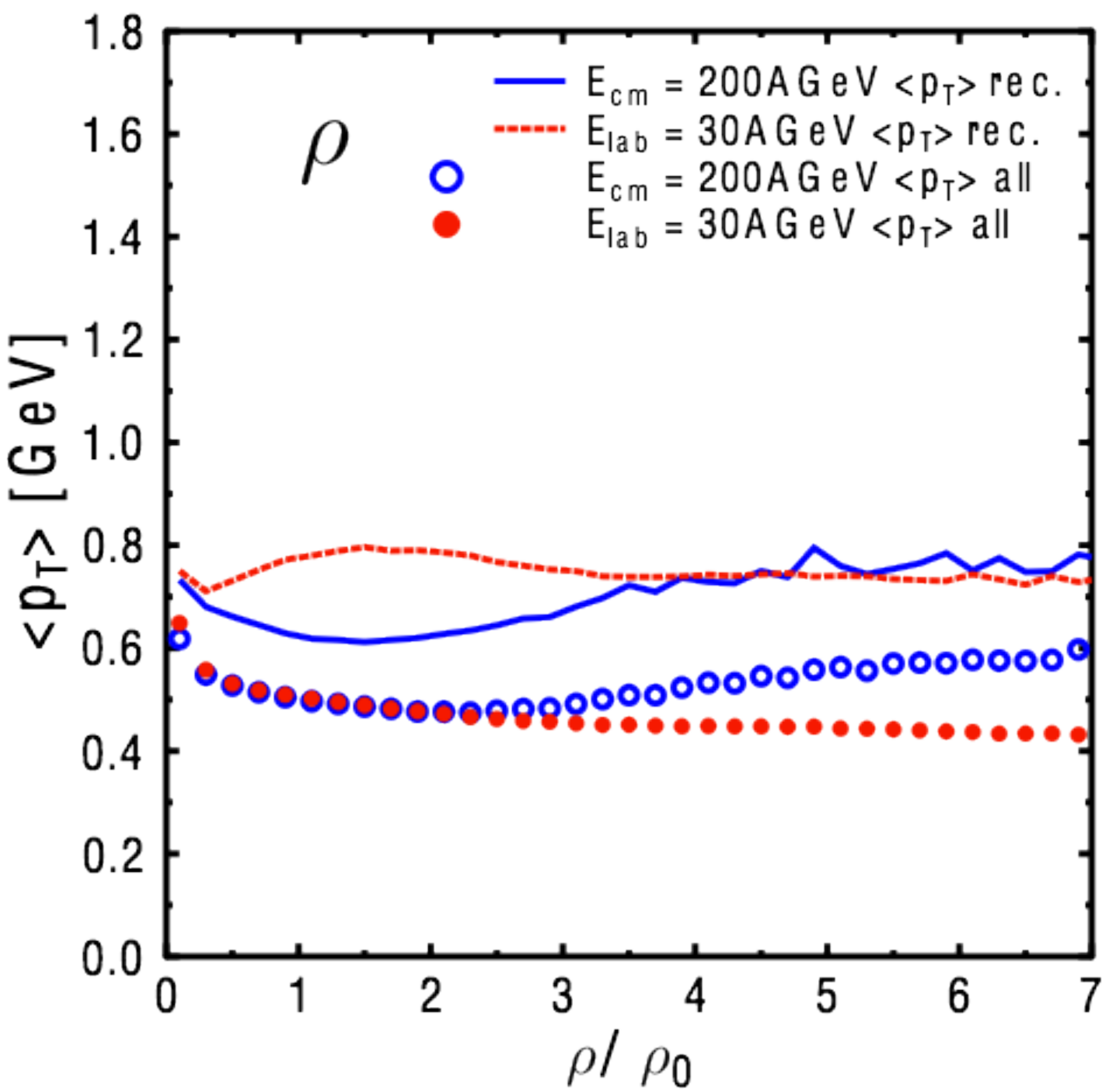
BACKUP SLIDES



p_T dependence

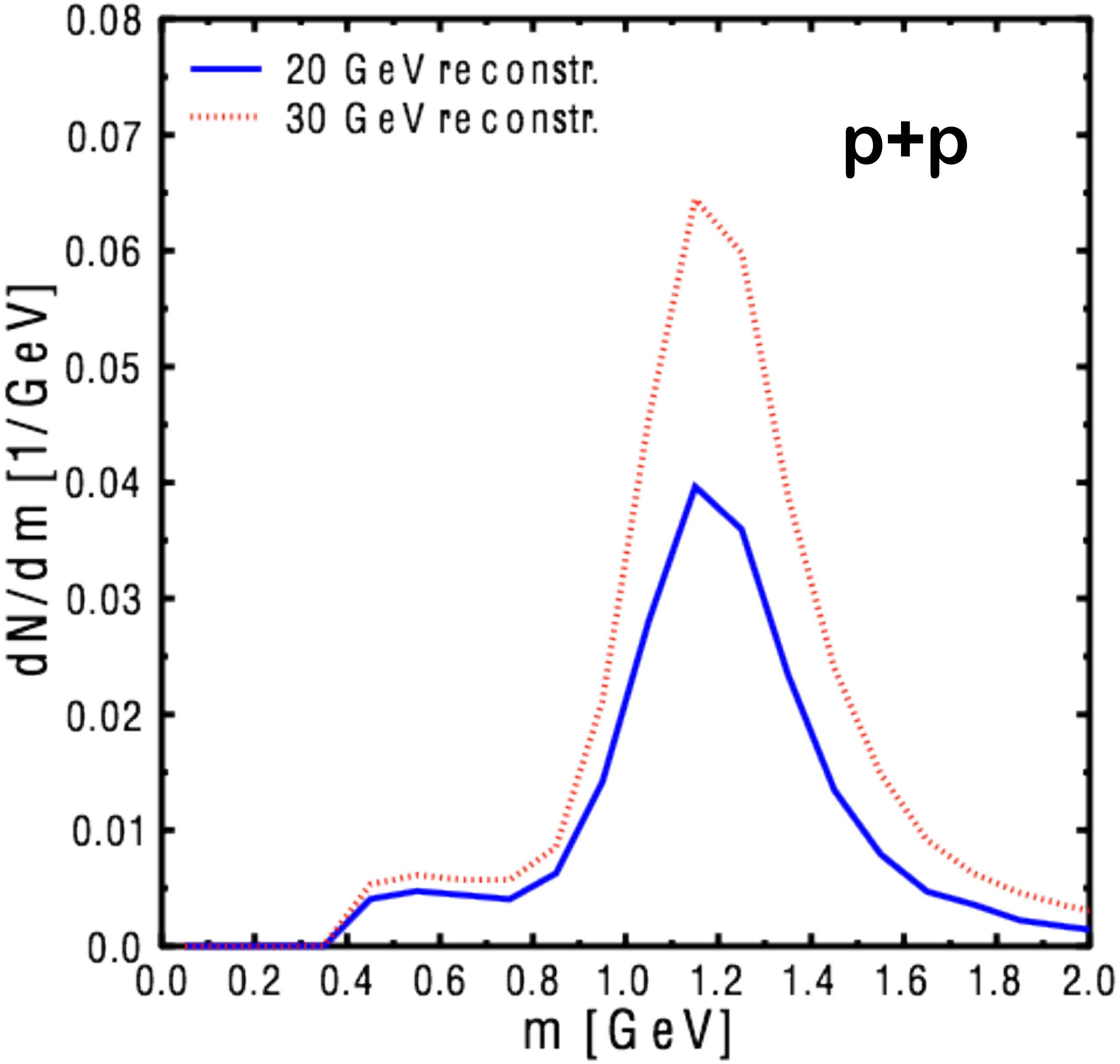
- average transverse momentum depends on density
- reconstructable resonances have higher p_T

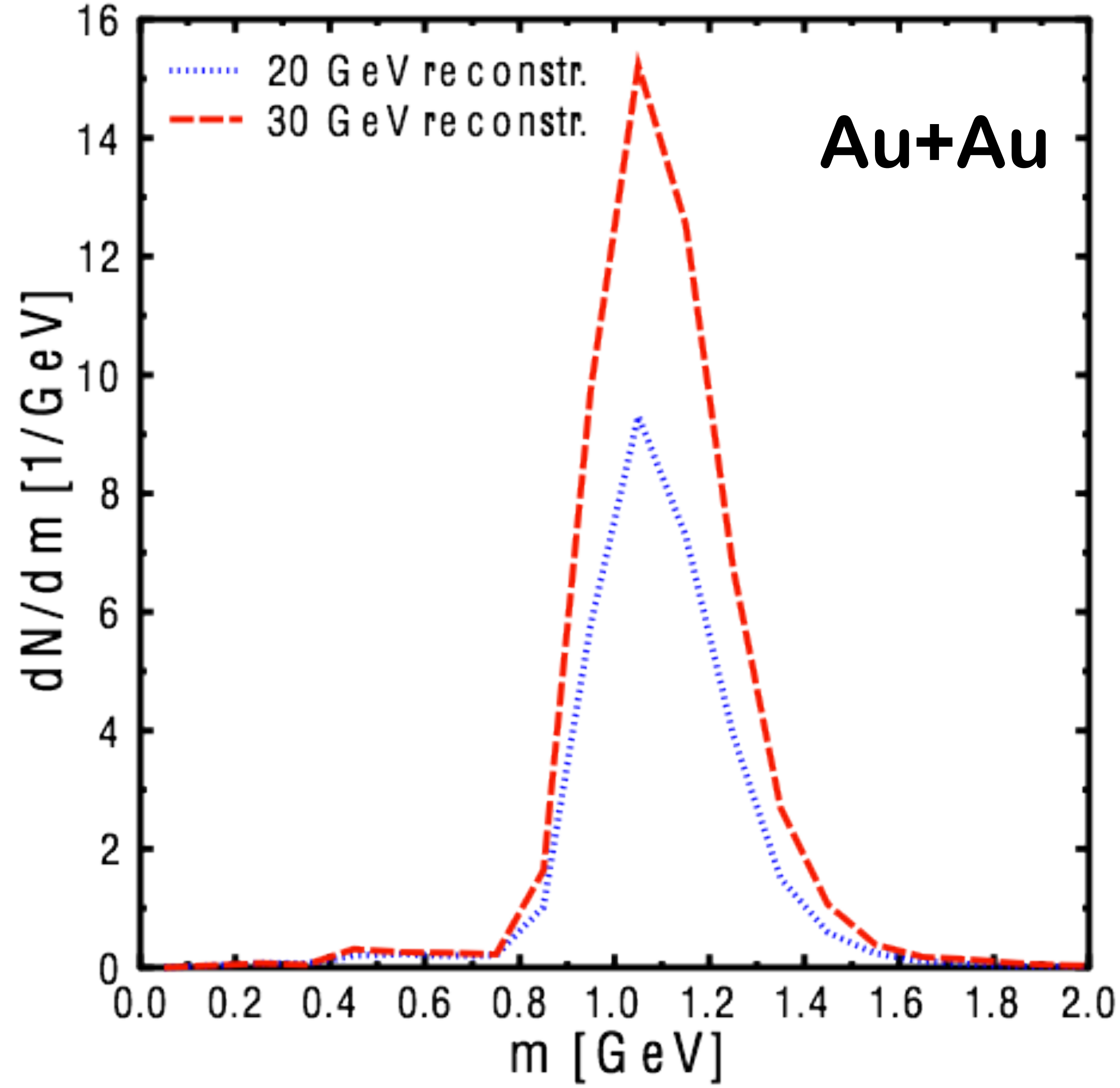




a_1 meson

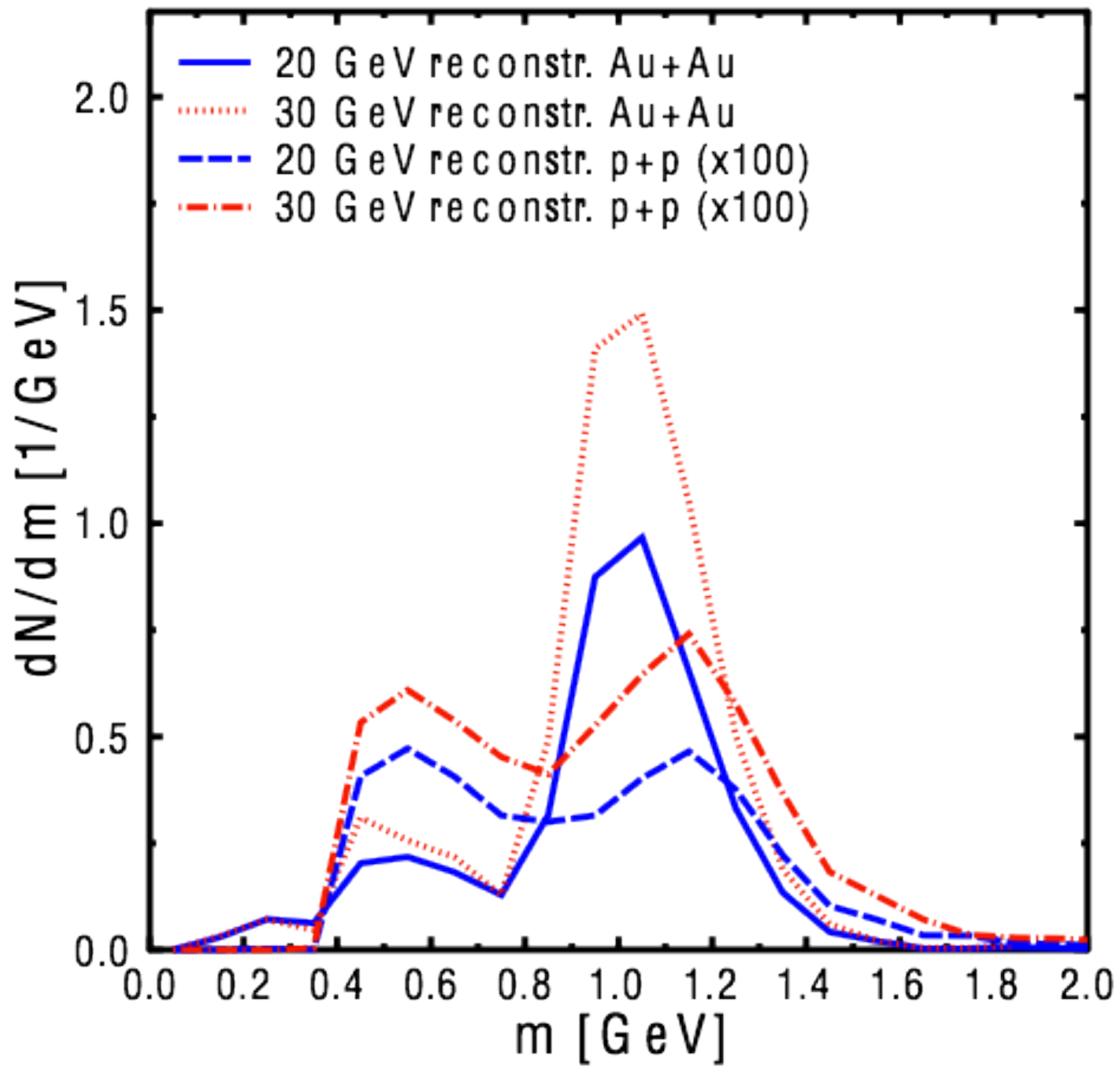
Idea: Check the mass distribution from the transport code.

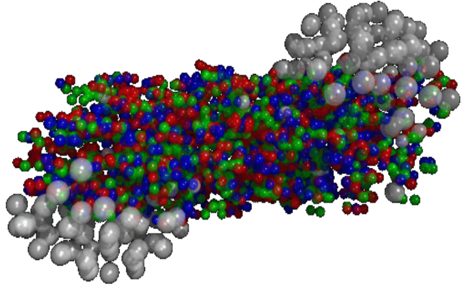




a_1 meson

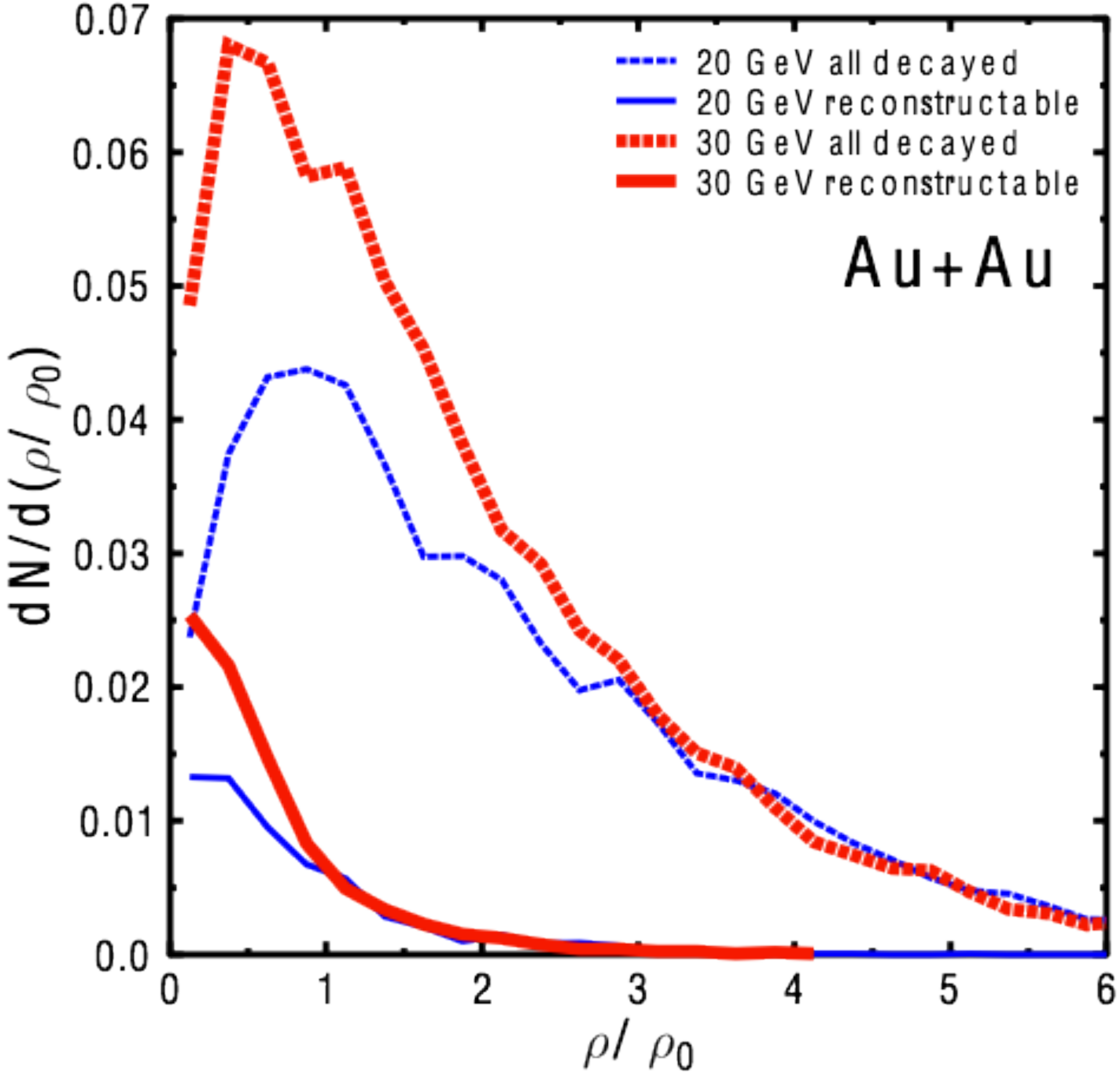
Next: trigger on the decay channel $a_1 \rightarrow \gamma\,\pi$. (assumed BR = 0.1)

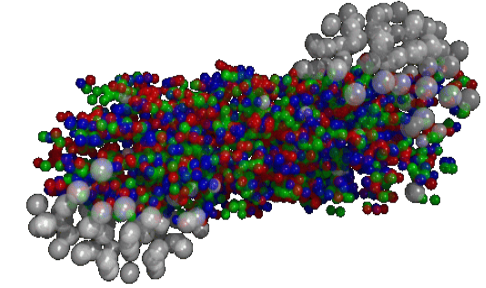




a_1 meson - density

Density at the point of decay of the a_1 meson



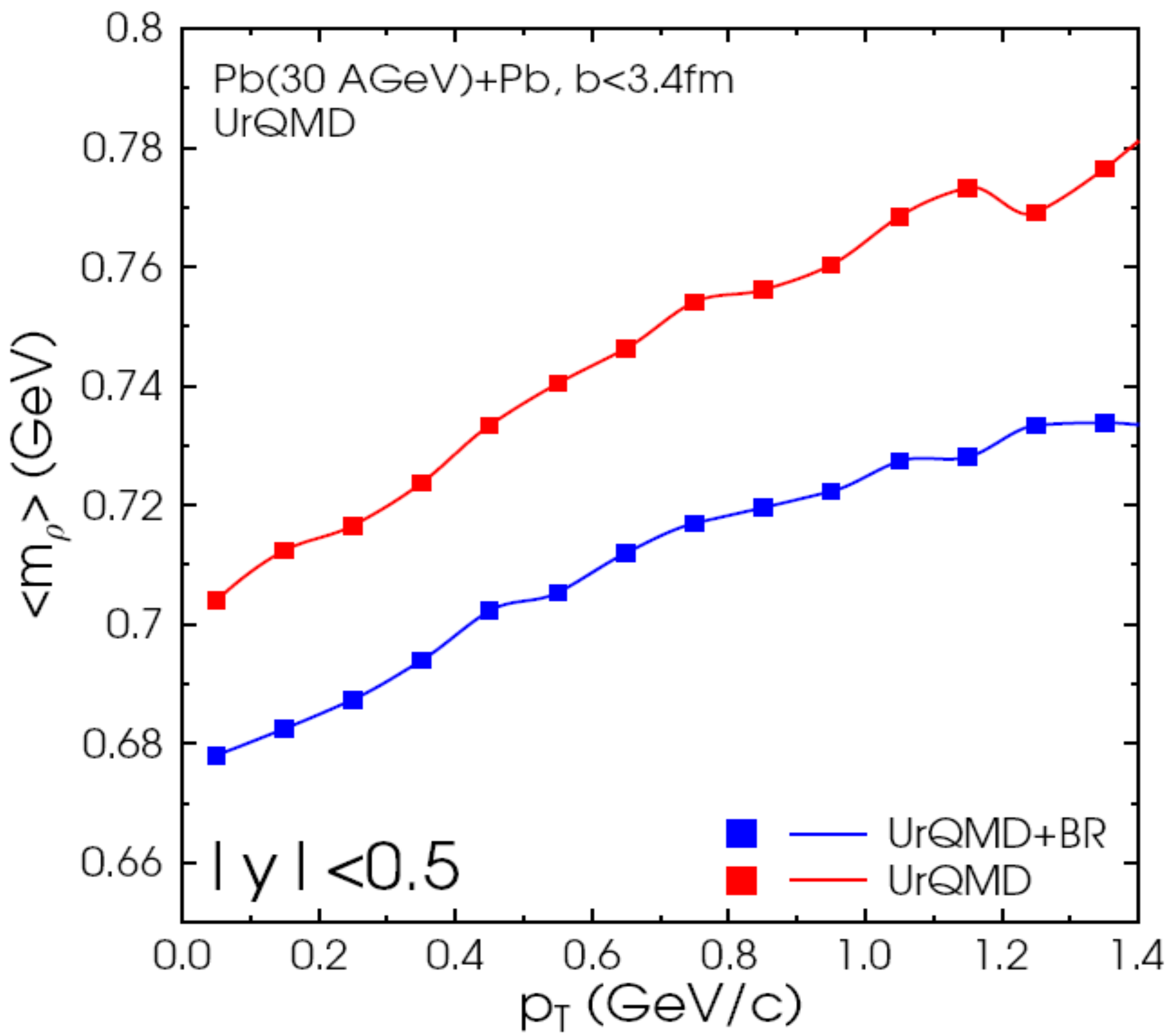


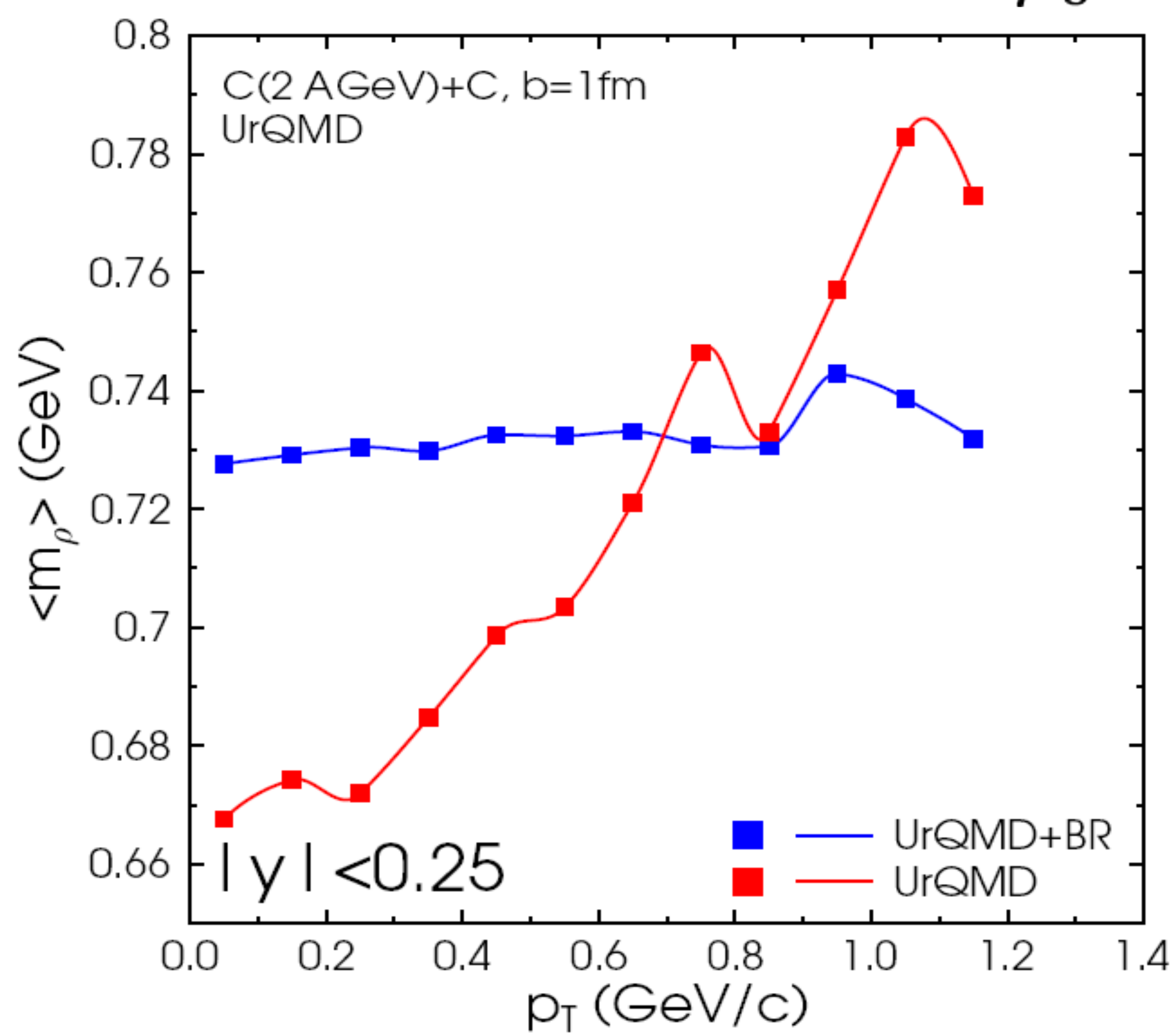
ρ meson at higher energies

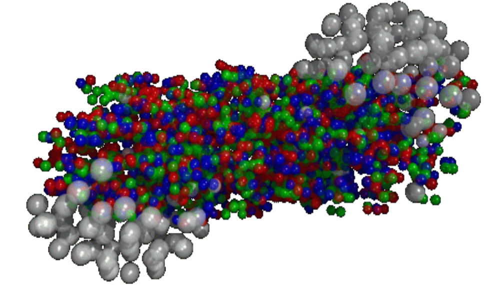
The ρ meson mass drops towards lower transverse momentum.

BR:

$$m^* = m_\rho(1 - 0.15\frac{\rho_B}{\rho_0})$$

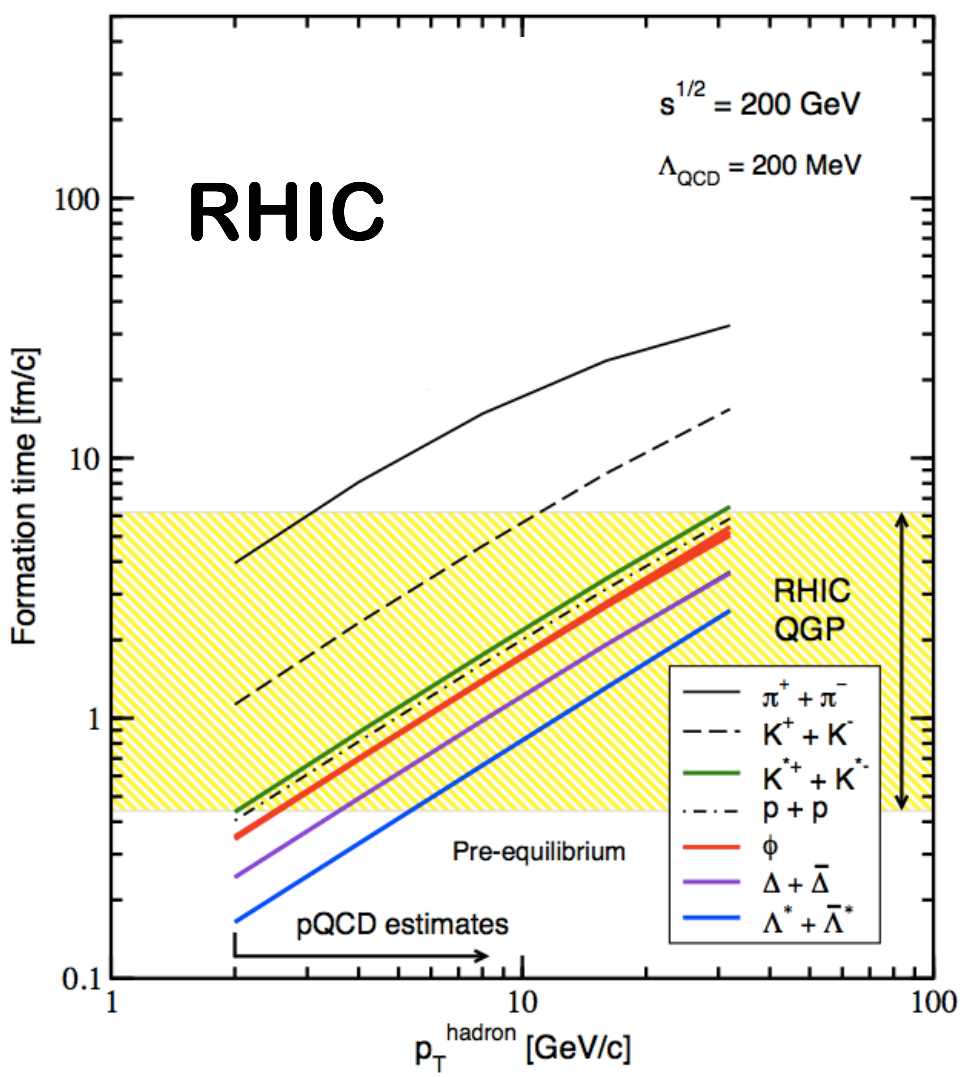


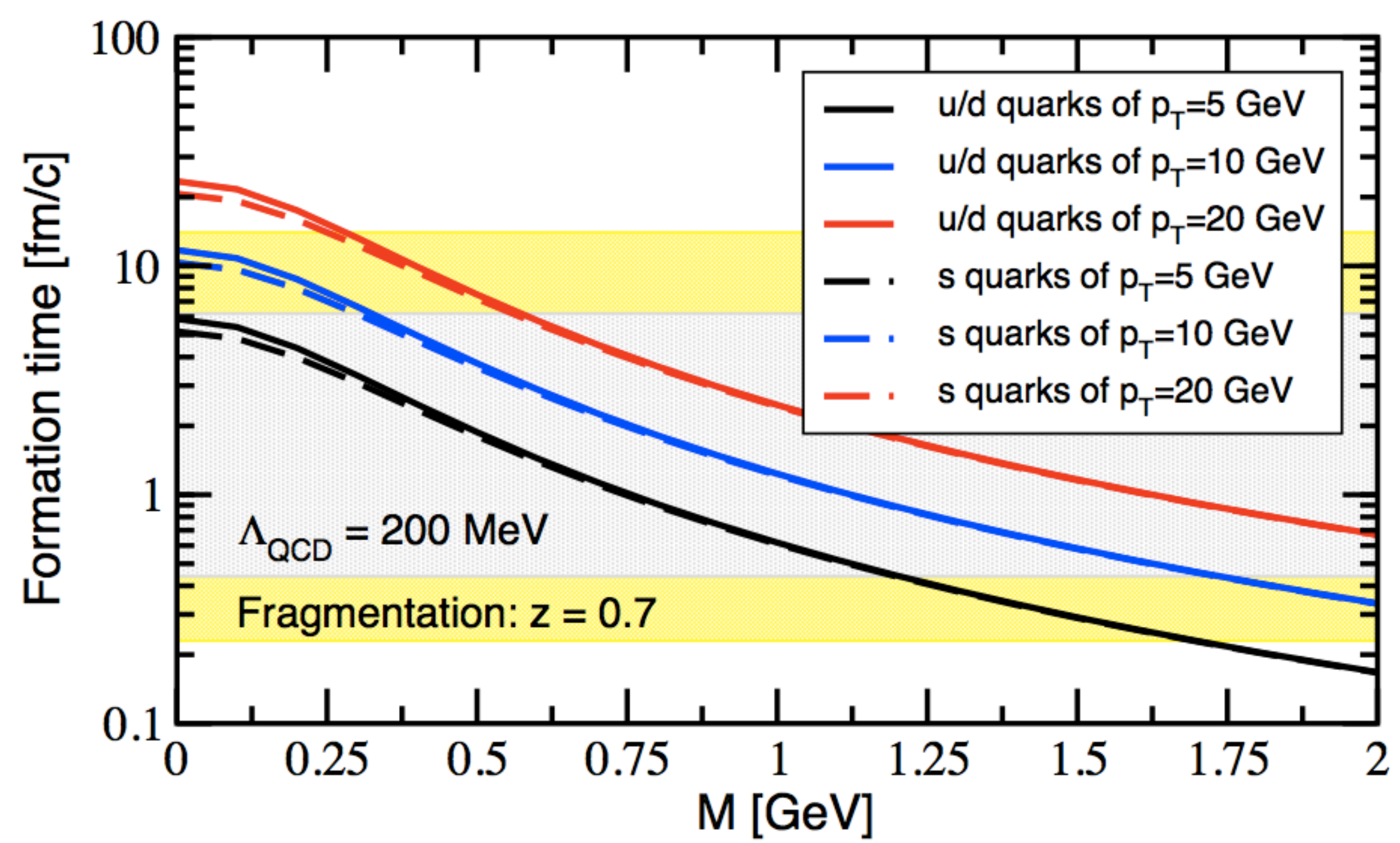




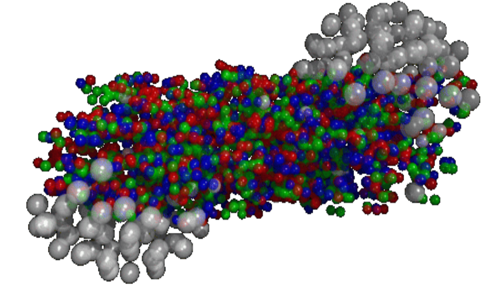
Formation time

- **formation time is mass and p_T dependent**
- **shaded areas indicate the estimated lifetime of the partonic phase**





Markert, Bellwied, Vitev,
Phys.Lett.B669:92-97,2008



Cross sections

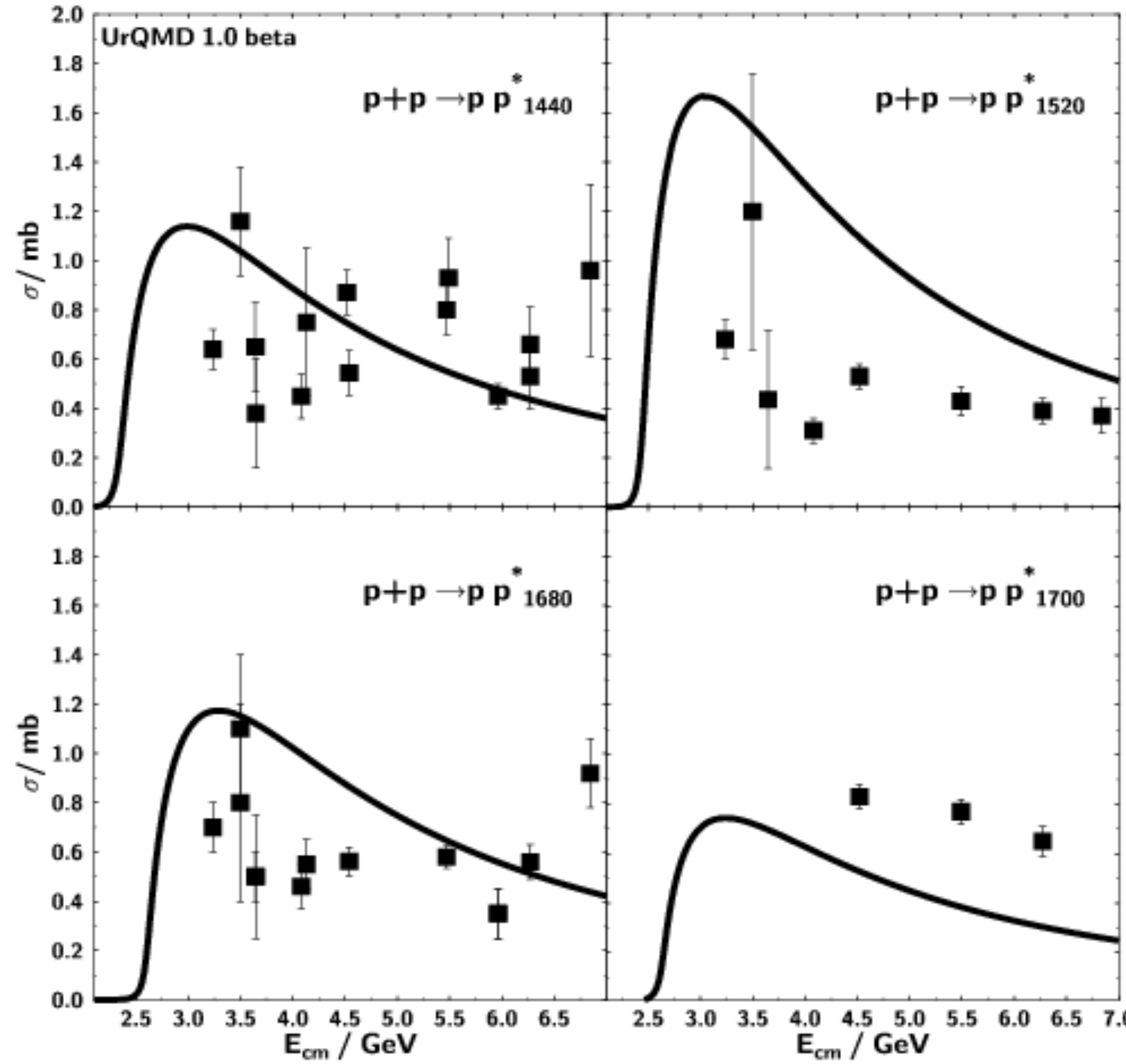
$$\sigma_{1,2\to 3,4}(\sqrt{s}) \sim (2s_3+1)(2s_4+1)\frac{<p_{3,4}>}{<p_{1,2}>}\frac{1}{\sqrt{s}}|M(m_3,m_4)|^2$$

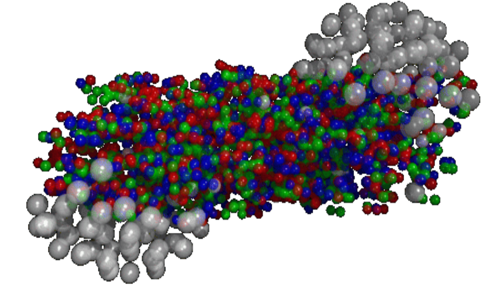
Global fit with the same kind of matrix element for 5 channels

$$NN \to NN^*, N\Delta^*, \Delta\Delta, \Delta N^*, \Delta\Delta^*$$

$$|M(m_3,m_4)|^2 = A\frac{1}{(m_4-m_3)^2(m_4+m_3)^2}$$

Data from elementary reactions are needed as an input into theory! (HADES?)





Cross sections

$$NN \rightarrow N\Delta_{1232}, NN^*, N\Delta^*, \Delta_{1232}\Delta_{1232}, \Delta_{1232}N^*, \text{ and } \Delta_{1232}\Delta^*$$

parametrized via

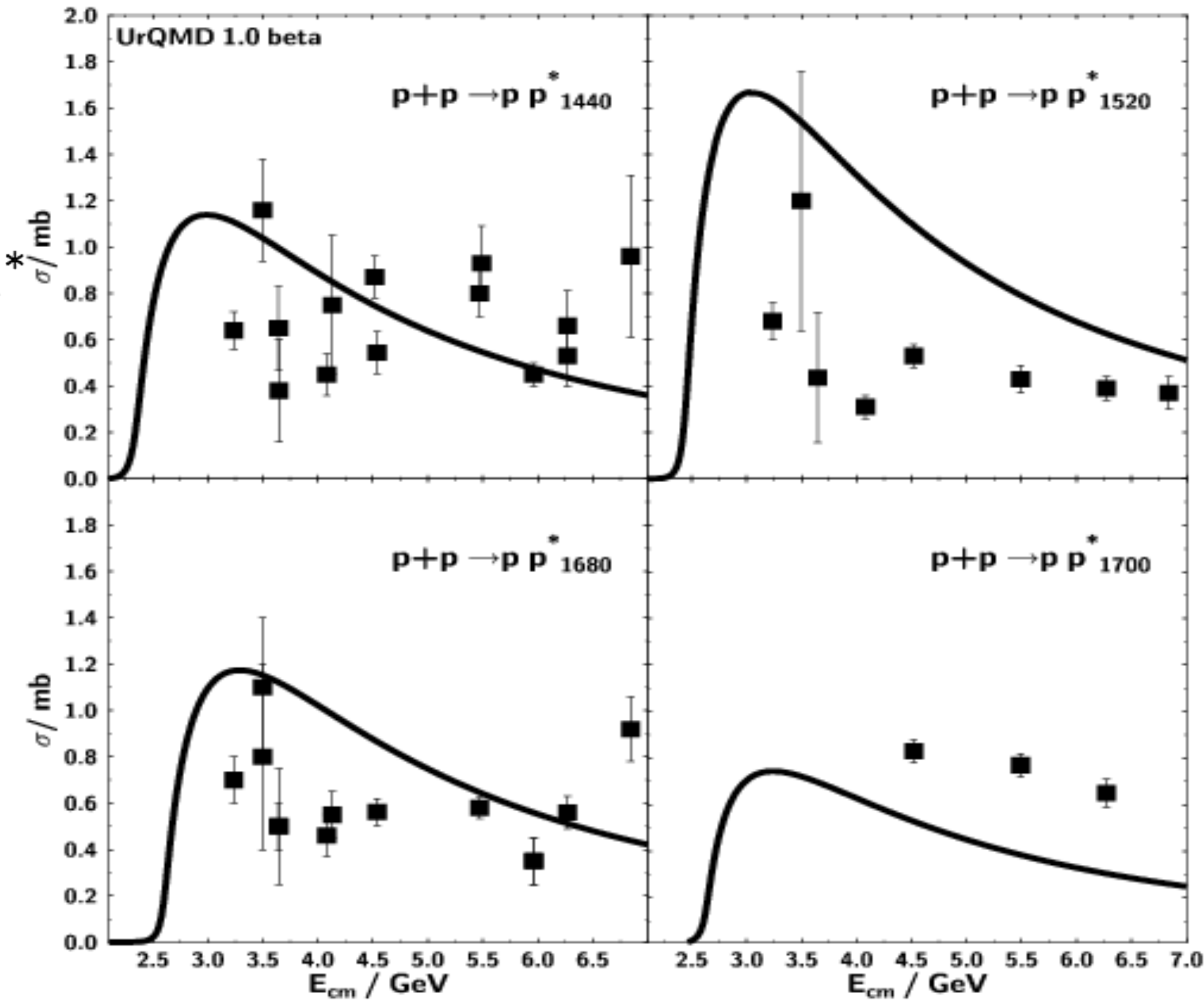
$$\sigma_{1,2\rightarrow 3,4}(\sqrt{s}) \sim (2s_3+1)(2s_4+1)\frac{<p_{3,4}>}{<p_{1,2}>}\frac{1}{\sqrt{s}}|M(m_3,m_4)|^2$$

Global fit with the same kind of matrix element of

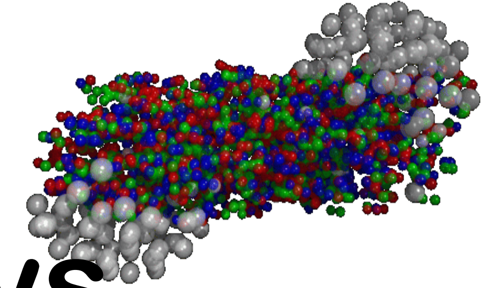
$$NN \rightarrow NN^*, N\Delta^*, \Delta_{1232}\Delta_{1232}, \Delta_{1232}N^*, \text{ and } \Delta_{1232}\Delta^*$$

$$|M(m_3,m_4)|^2 = A\frac{1}{(m_4-m_3)^2(m_4+m_3)^2}$$

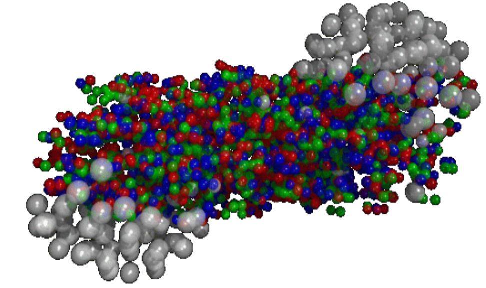
Data from elementary reactions are needed as an input into theory! (HADES?)



Dileptonic and hadronic decays

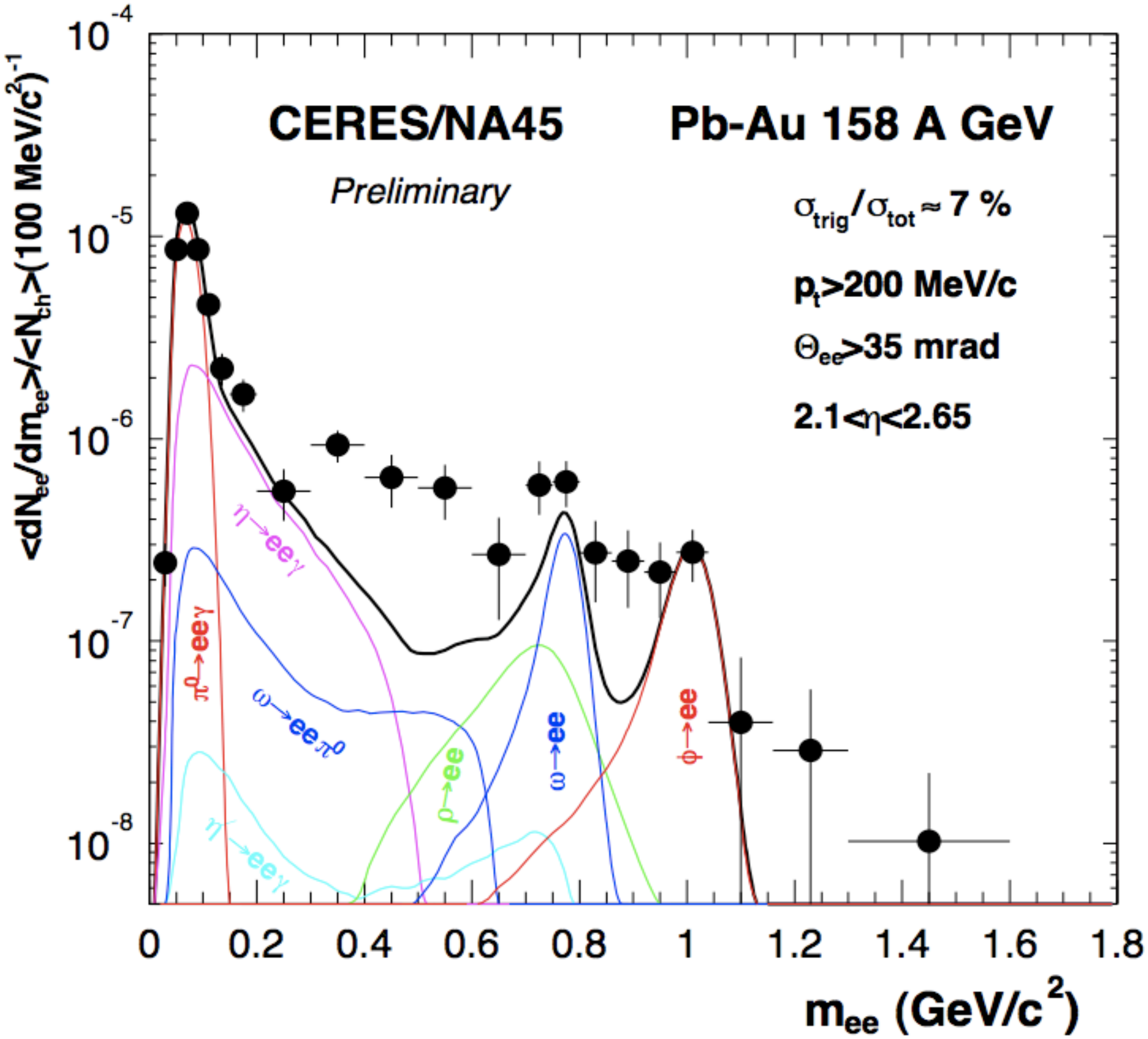


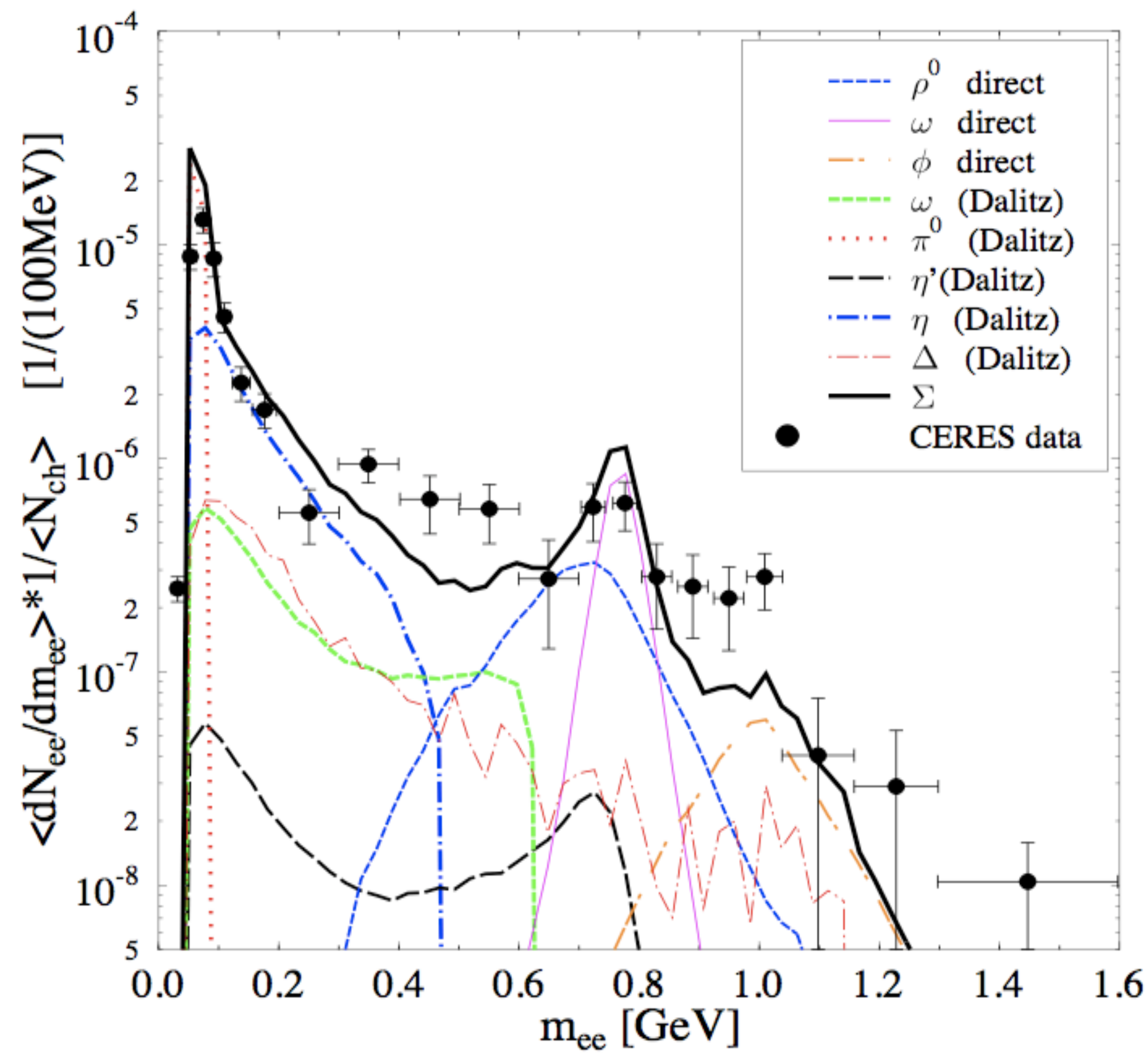
Dileptons	Hadrons
do not interact strongly with the surrounding medium	suffer from final state interactions
originate from various sources in various mass regions (note: Dalitz decays)	originate from various sources in various mass regions
Typical branching ratios on the order of 10^{-4} - 10^{-5}	Typical branching ratios on the order of 0.1 - 1
when measured reflect the integrated collision history	when measured reflect the late stage (after freezeout) of the collision



Dilepton spectra

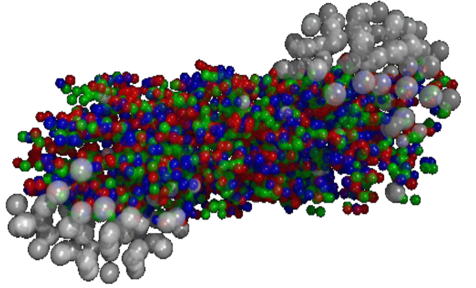
Certainly not a perfect fit, yet doing better than the hadronic cocktail.





CERES Collaboration, Nucl. Phys A 774, 43, 2006

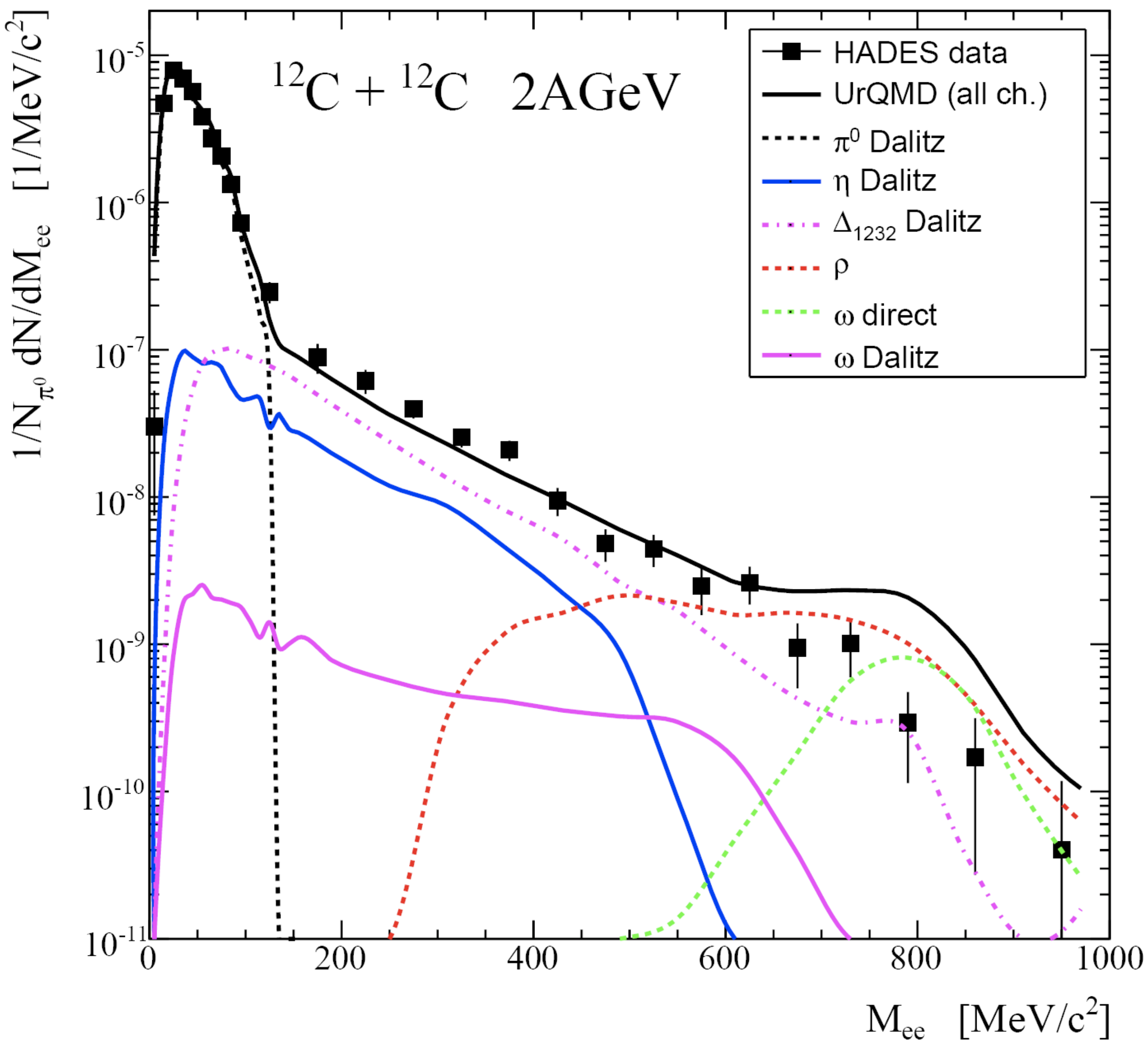
D. Schumacher, diploma thesis, 2006



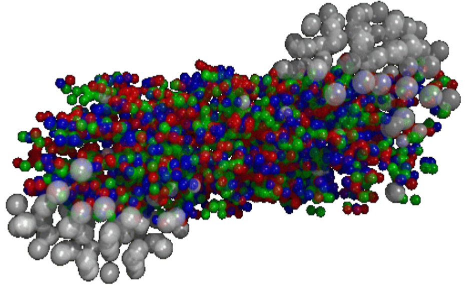
Dilepton spectra

The comparison to the data gives a reasonable explaination of the 400-600 MeV region.

Note: No explicit in-medium modifications have been implemented.



HADES collaboration, Phys. Rev. Lett.98, 052302, 2007
D. Schumacher, S.Vogel, M. Bleicher, in preparation



Dilepton approaches

1) Shining

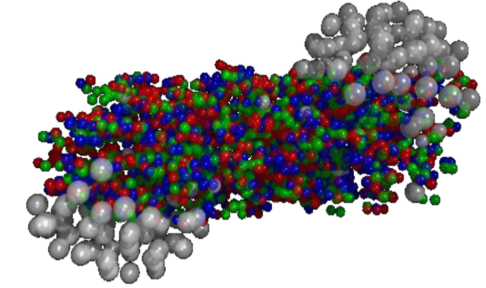
- Evaluate lifetime of the resonance, weight accordingly

2) Full weight only when resonance decays - ignore absorbed resonances

- Weight decayed resonance with vacuum width / BR

3) Full weight when absorbed/decayed

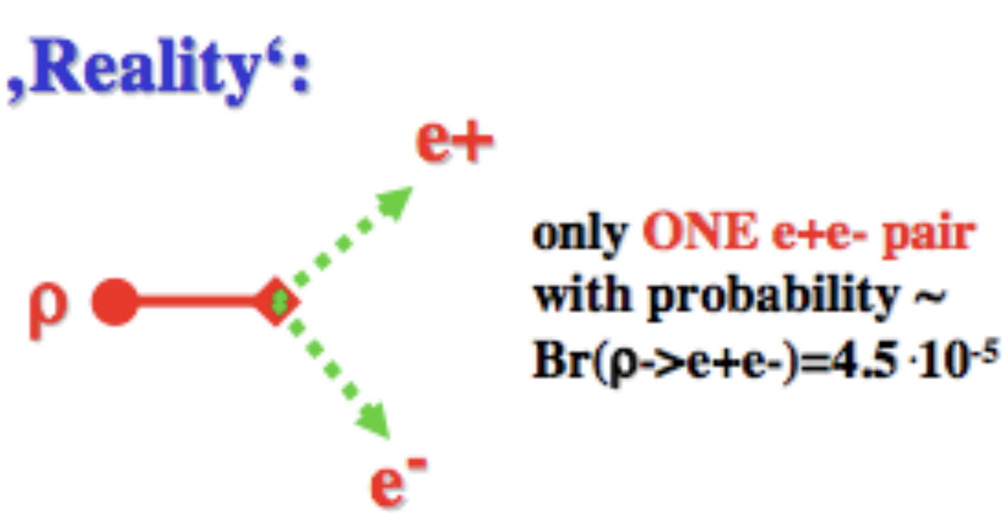
- Weight all decayed/absorbed resonances with vacuum width / BR (most optimistic approach)

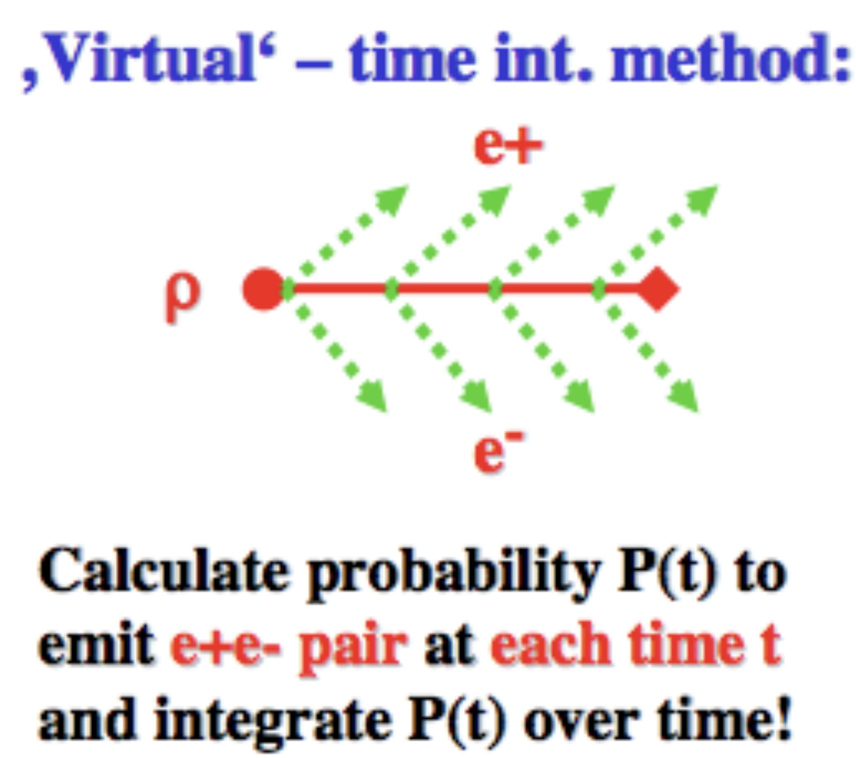


Dilepton via shining

- Dileptons are emitted throughout the whole resonance evolution ("time-integration method" or "shining")

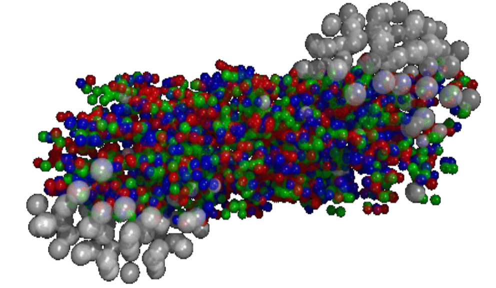
- $\Rightarrow$Resonances do not have to decay to emit dileptons





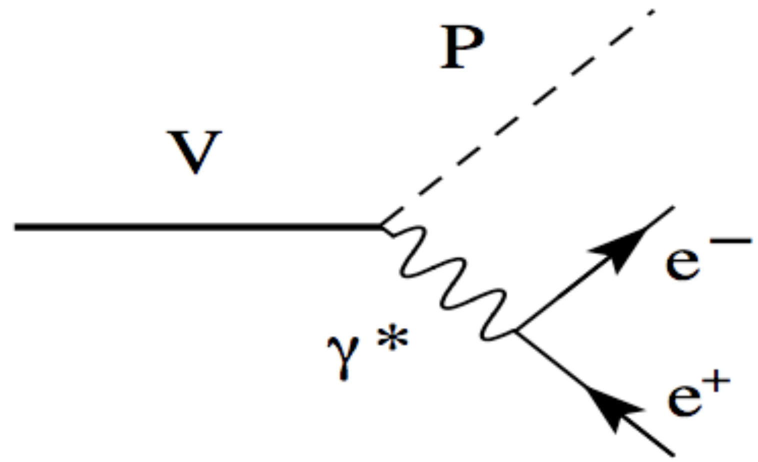
Pictures from:
http://www.th.physik.uni-frankfurt.de/~brat/hsd.html

Heinz and Lee, Nucl.Phys.A544:503-508,1992
Ko and Li, Nucl.Phys.A582:731-748,1995

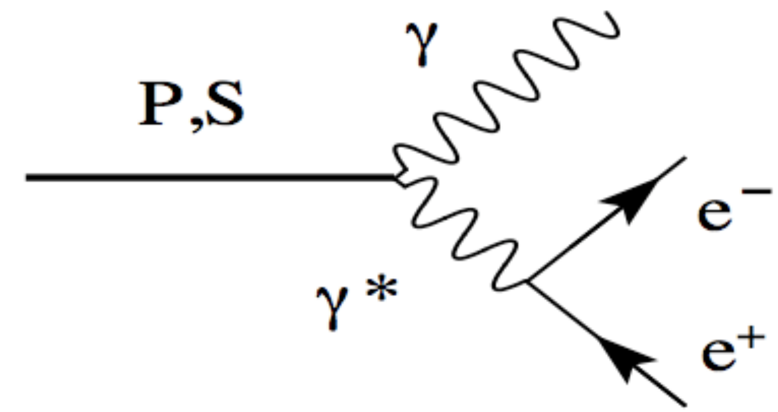


Dilepton sources

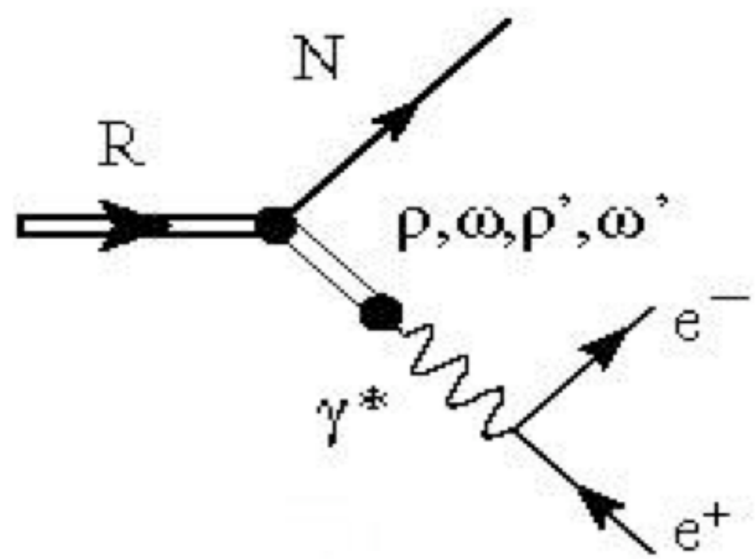
- **Bremsstrahlung**
- **Dalitz decays**
- π^0, η , η ' , ω, Δ ...
- **Direct decays**
- ρ, ω, ϕ ...
- **Thermal dileptons**



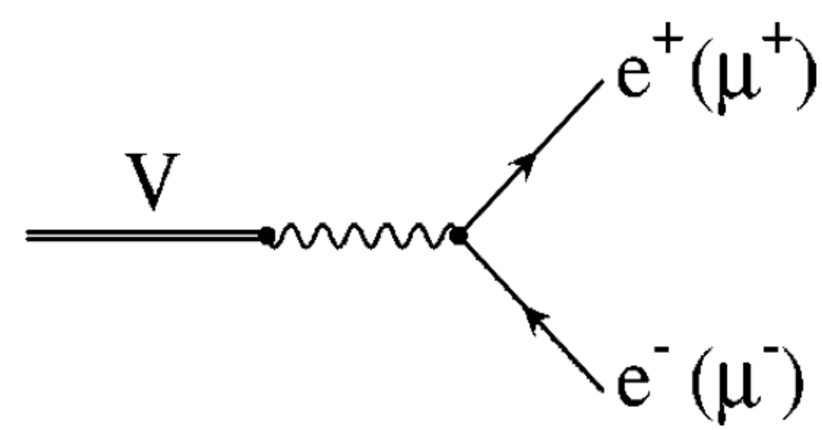
Dalitz decay (vector meson)



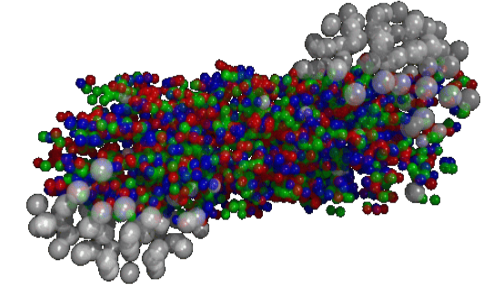
(pseudoscalar meson)



Dalitz decay (Δ)



Direct decay



Dilepton sources

Pseudoscalar mesons:

$$\frac{dN_{A\to\gamma e^+e^-}}{dM} = \frac{4\alpha}{3\pi M}\sqrt{1-\frac{4m_e^2}{M^2}}\left(1+\frac{2m_e^2}{M^2}\right)\left(1-\frac{M^2}{m_A^2}\right)^3 \times |F_{AB}(M^2)|^2\frac{\Gamma_{A\to 2\gamma}}{\Gamma_{tot}}\langle N_A\rangle$$

Vector mesons:

$$\frac{dN_{A\to Be^+e^-}}{dM} = \frac{2\alpha}{3\pi M}\sqrt{1-\frac{4m_e^2}{M^2}}\left(1+\frac{2m_e^2}{M^2}\right)|F_{AB}(M^2)|^2\frac{\Gamma_{A\to 2\gamma}}{\Gamma_{tot}}\langle N_A\rangle \times \left(\left(1+\frac{M^2}{m_A^2-m_B^2}\right)^2-\left(\frac{2m_AM}{m_A^2-m_B^2}\right)^2\right)^{3/2}$$

Direct decays:

$$BR(V\to e^+e^-) = \frac{\Gamma_{V\to e^+e^-}(M)}{\Gamma_{tot}}$$

$$\Gamma_{V\to e^+e^-}(M) = \frac{\Gamma_{V\to e^+e^-}(m_V)}{m_V}\frac{m_V^4}{M^3}\sqrt{1-\frac{4m_e^2}{M^2}}\left(1+2\frac{m_e^2}{M^2}\right)$$

Dilepton sources

Δ baryon:

$$\frac{dN_{e^+e^-}}{dM} = \int \frac{dN_{\Delta\to Ne^+e^-}}{dM}(M_\Delta)\frac{dN_\Delta}{dM_\Delta}dM_\Delta = \int \frac{2\alpha}{3\pi M}\frac{\Gamma(M_\Delta, M)}{\Gamma^{tot}_{\Delta 0}}\frac{dN_\Delta}{dM_\Delta}dM_\Delta$$

$$\Gamma(M_\Delta, M) = \frac{\lambda^{1/2}(M^2, m_N^2, M_\Delta^2)}{16\pi M_\Delta^2} m_N [2\mathcal{M}_t(M, M_\Delta) + \mathcal{M}_l(M, M_\Delta)]$$

$$\lambda(m_A^2, m_1^2, m_2^2) = (m_A^2 - (m_1 + m_2)^2)(m_A^2 - (m_1 - m_2)^2)$$

$$\mathcal{M}_l = (efg)^2 \frac{m_\Delta^2}{9m_N} M^2 4(m_\Delta - m_N - q_0)$$

$$\mathcal{M}_t = (efg)^2 \frac{m_\Delta^2}{9m_N}[q_0^2(5m_\Delta - 3(q_0 + m_N)) - M^2(m_\Delta + m_N + q_0)]$$

L.G. Landsberg, Phys.Rept.128:301-376,1985
P. Koch, Z. Phys. C57:283-304, 1993

Dilepton via shining

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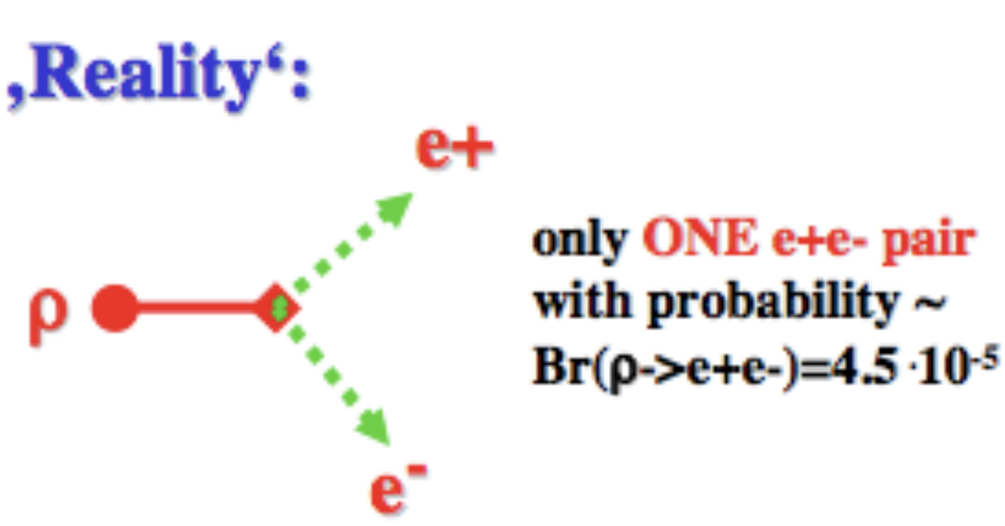
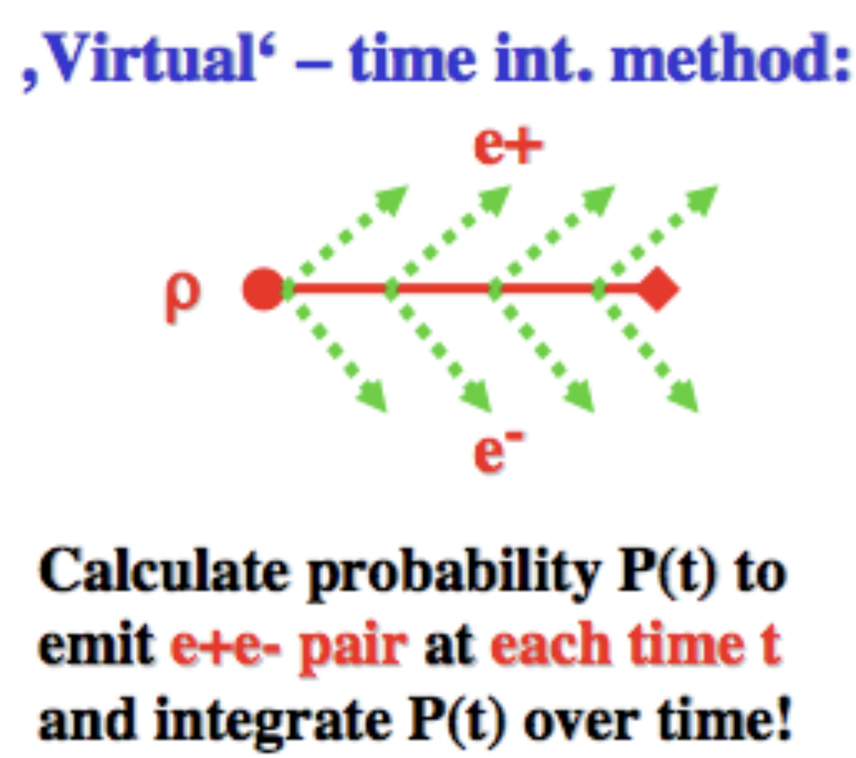
- ⇒Resonances do not have to decay to emit dileptons





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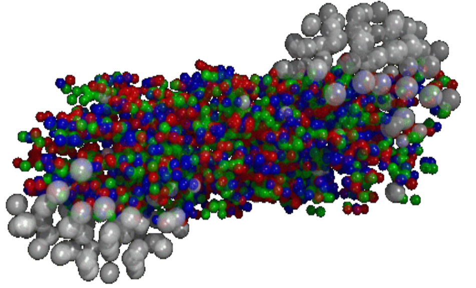


Dilepton approaches

1) Shining

- Evaluate lifetime of the resonance, weight accordingly

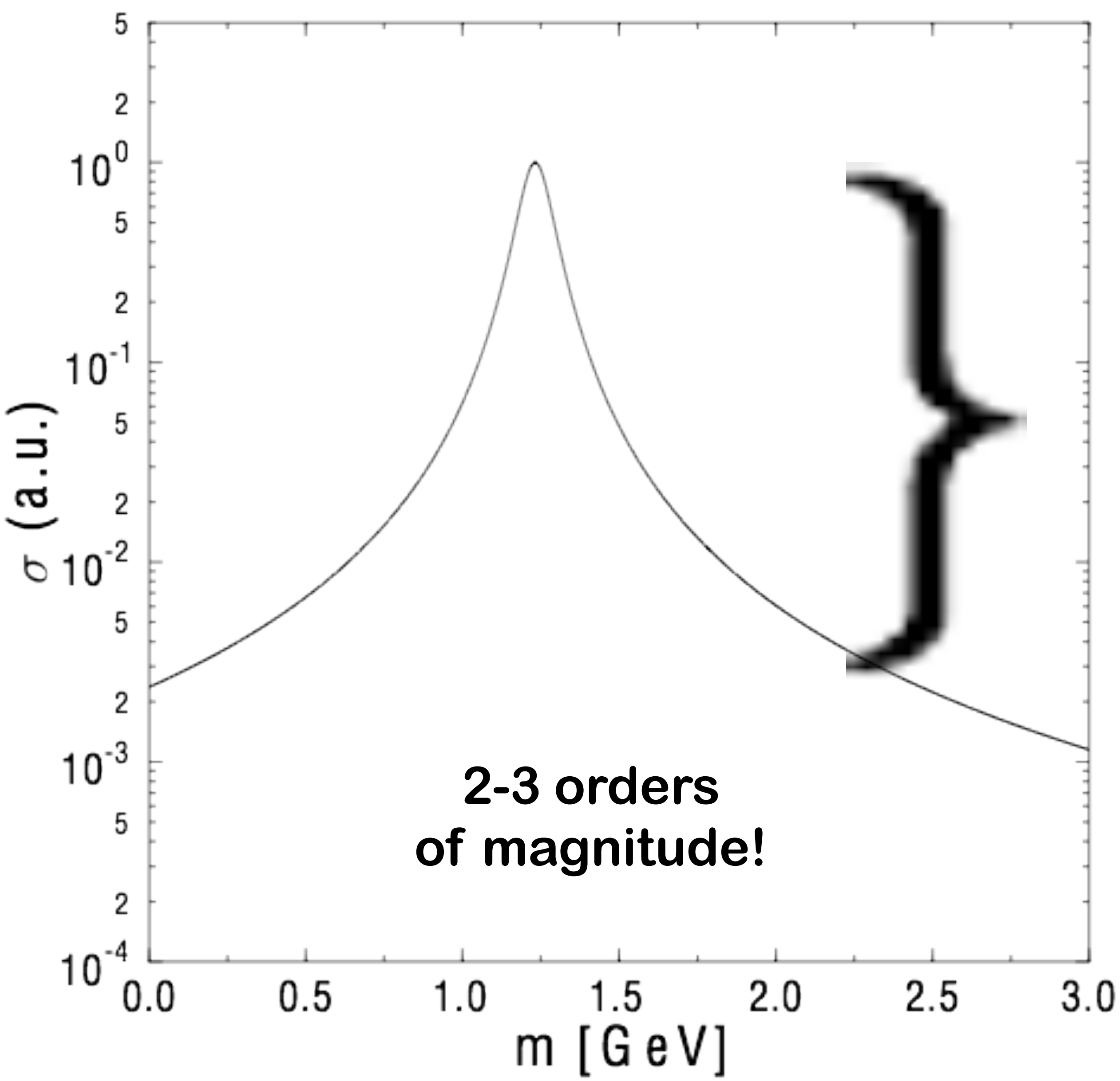
2) Full weight only when resonance decays - ignore absorbed resonances

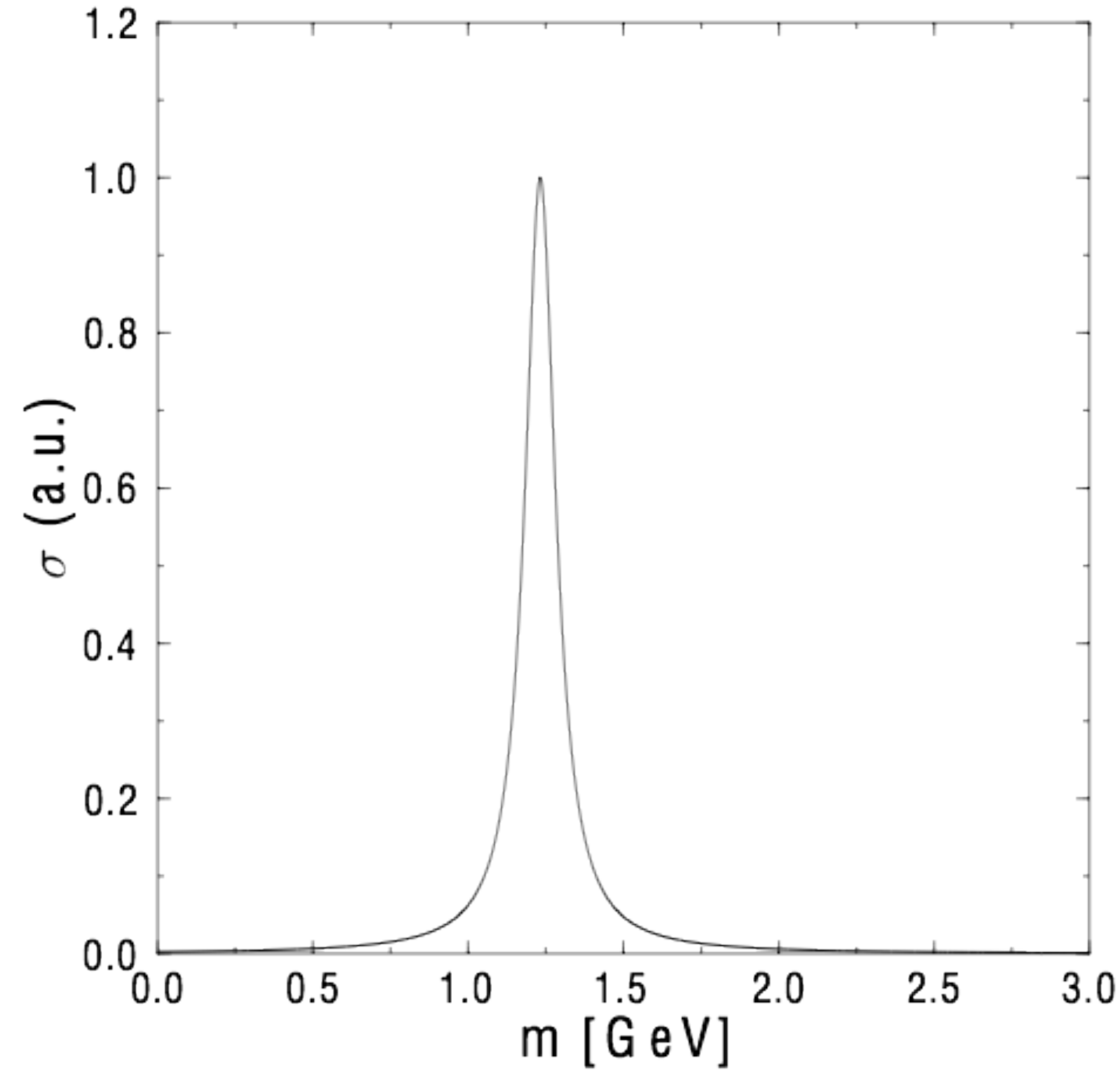
- Weight decayed resonance with vacuum width / BR

3) Full weight when absorbed/decayed

- Weight all decayed/absorbed resonances with vacuum width / BR (most optimistic approach)

Breit-Wigner

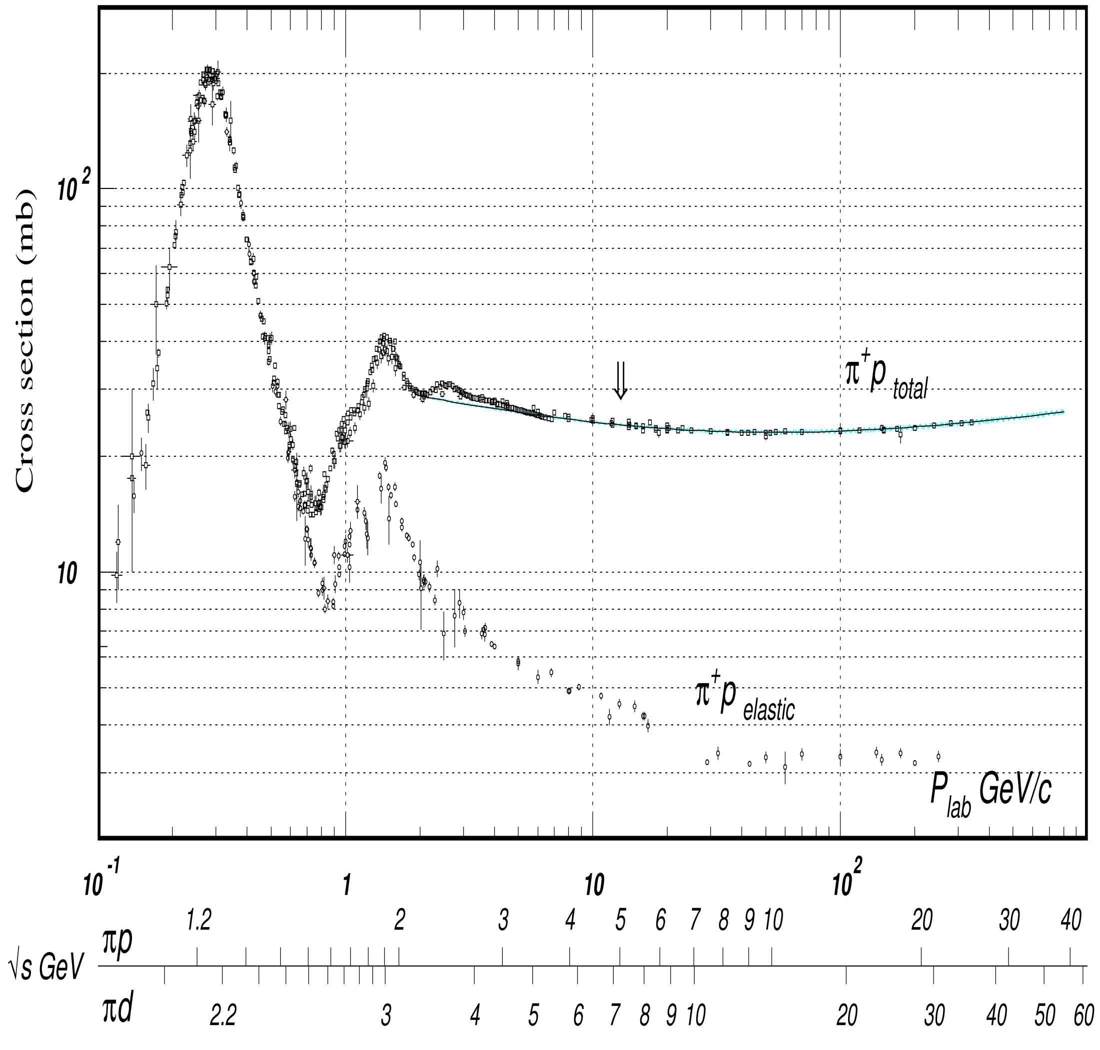




From pole to 2.2 GeV the Breit-Wigner provides a difference of 2-3 order of magnitude!

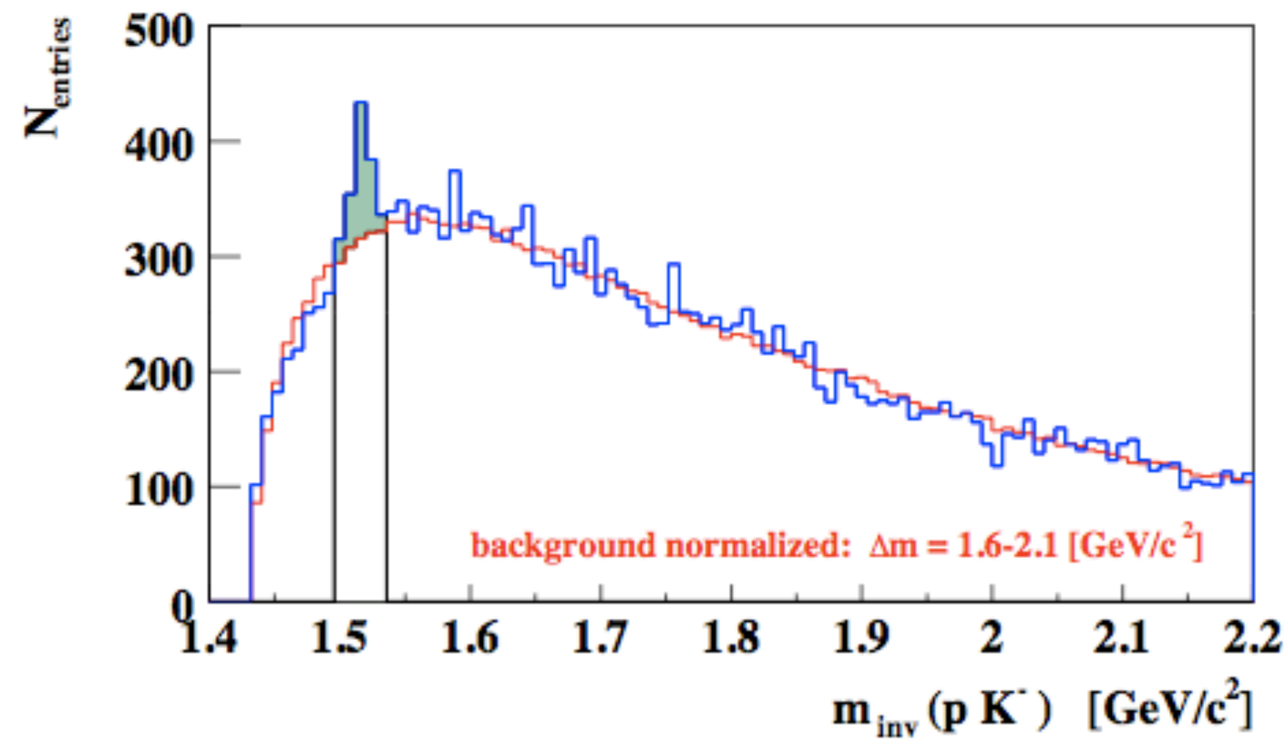
Measuring resonances

- Since resonances decay on timescales of several fm they cannot be measured directly
- Resonances are measured via their decay products, cross section follows a Breit-Wigner law

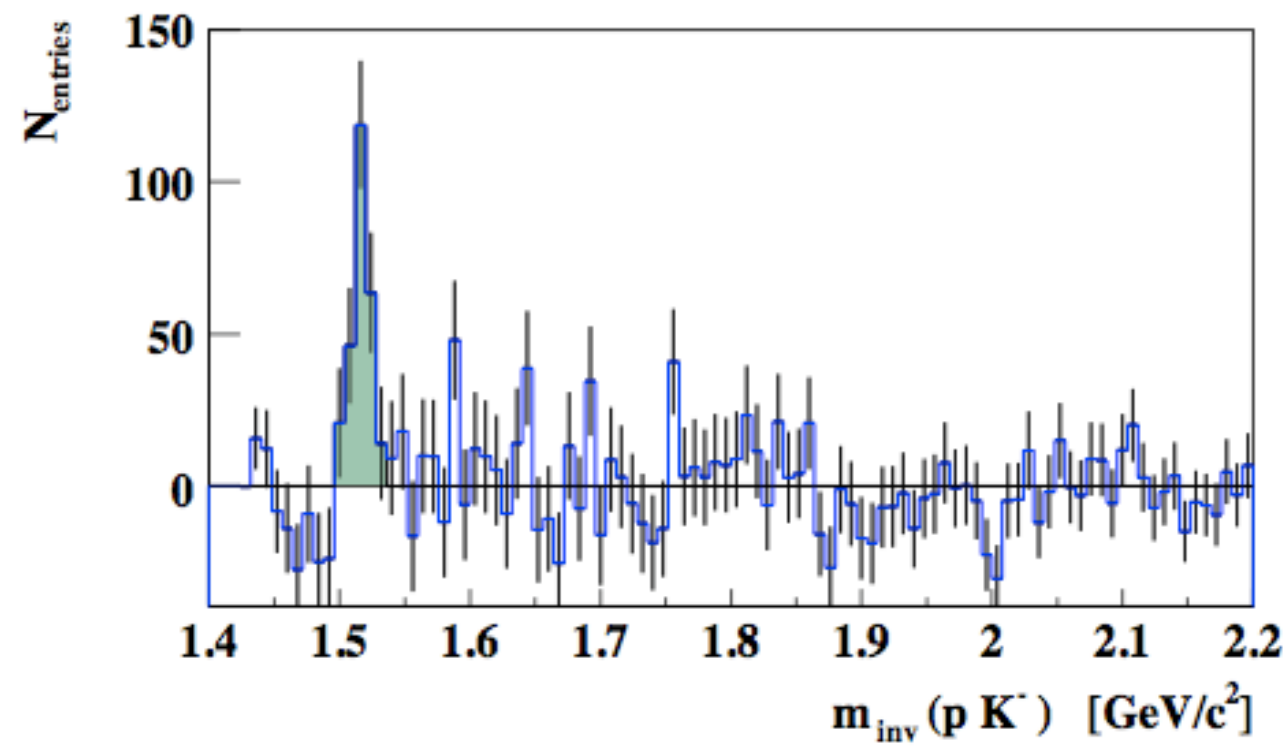


Measuring resonances in p+p

Correlate all protons and kaons in the event, plot invariant mass.



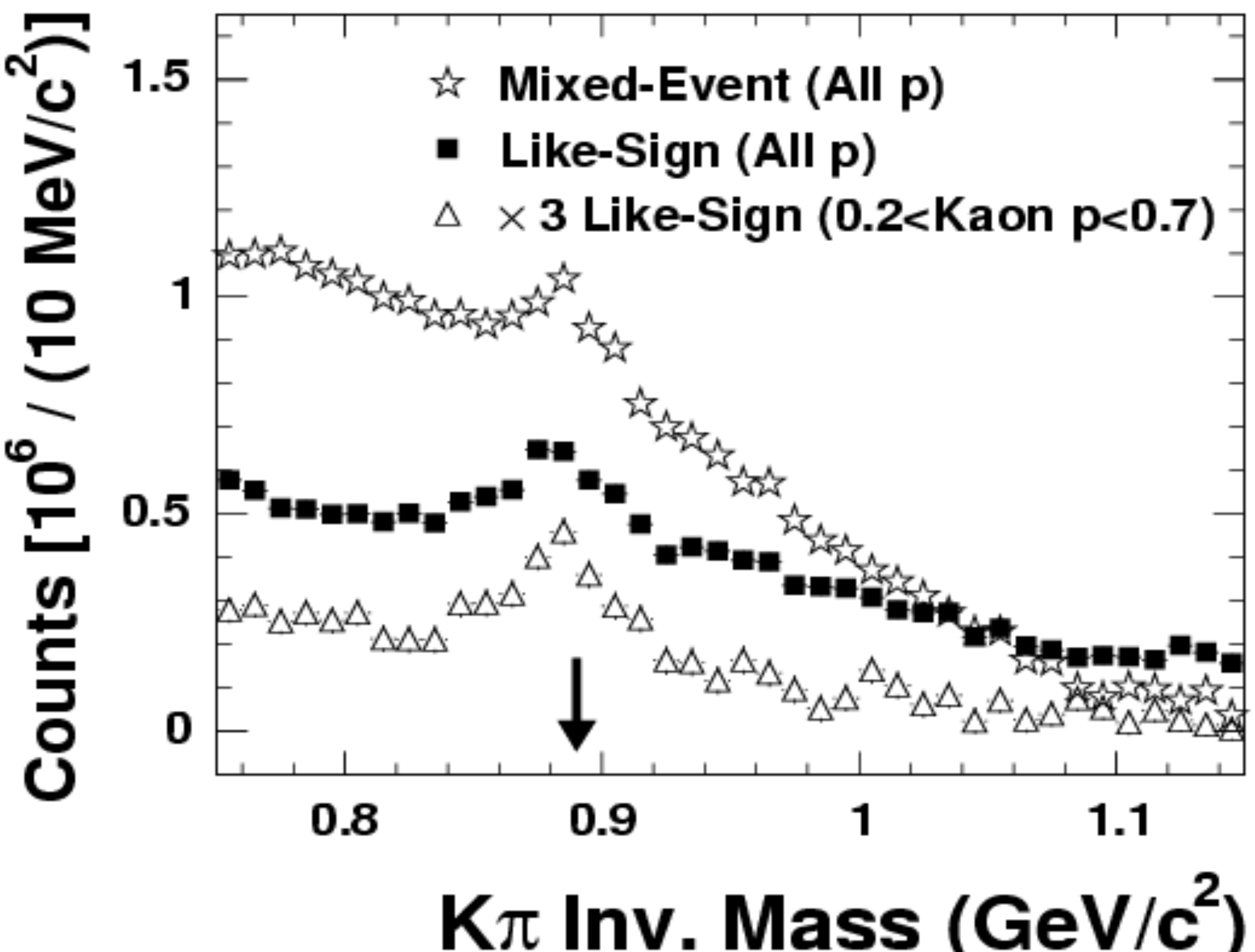
Lots of uncorrelated pairs → background subtraction needed

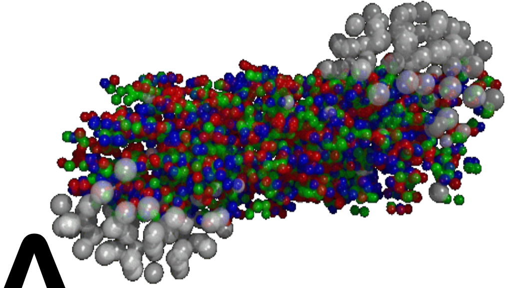


Still a visible peak, but not as clear as before.

Measuring resonances in A+A

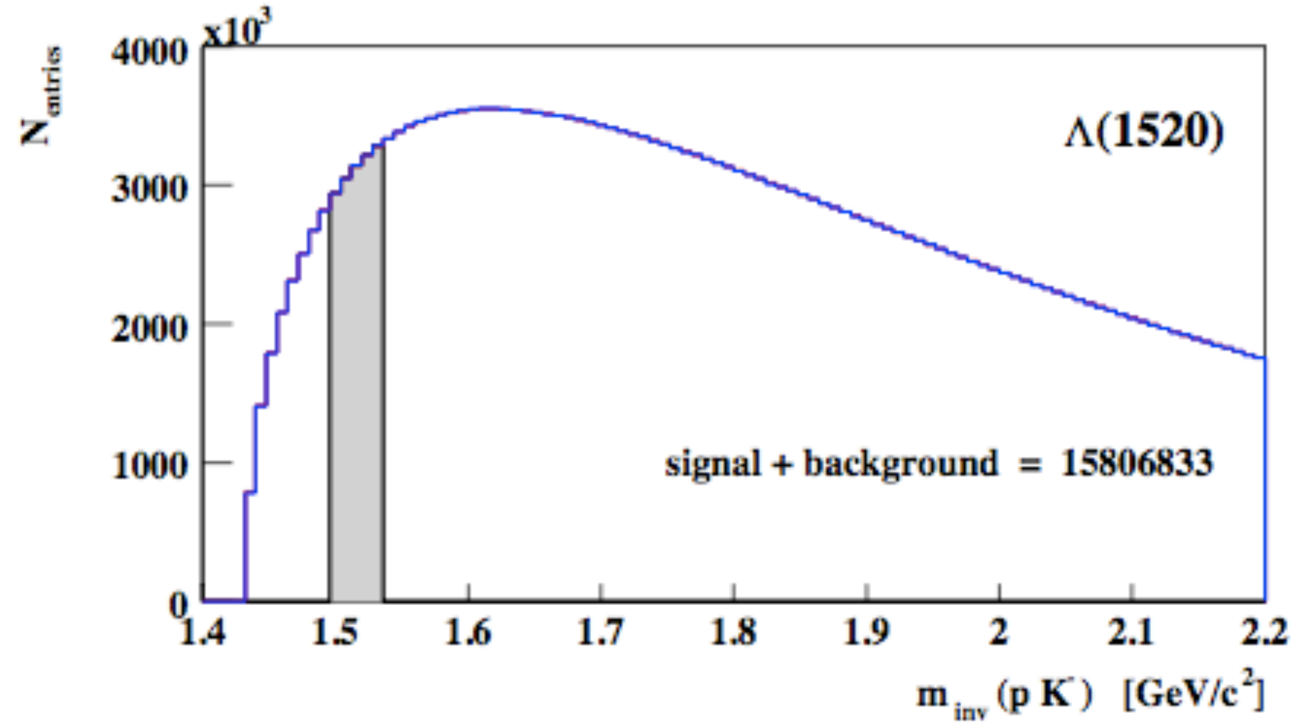
Different methods to subtract the background lead to slightly different results.



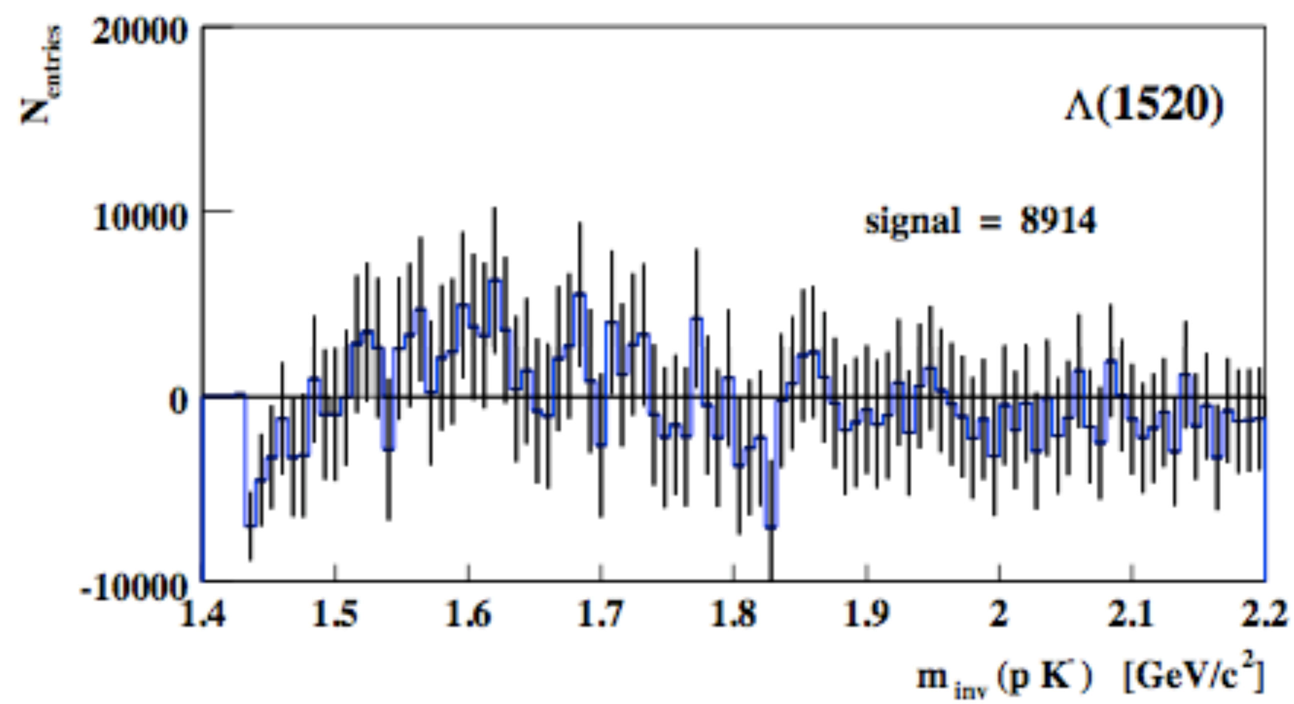


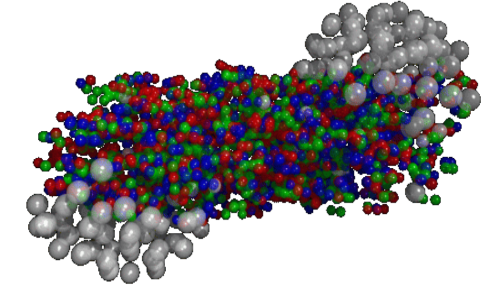
Measuring resonances in A+A

Correlate all protons and kaons in the event, plot invariant mass.



Peak?





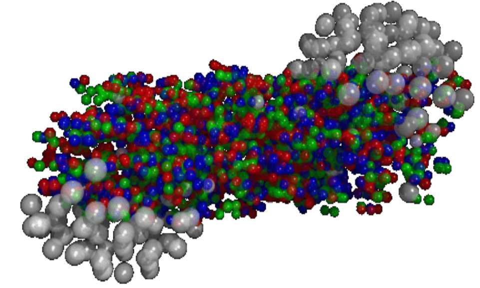
Density calculation

- **Lorentz-transform the CF density to the frame where the three-current vanishes (Eckart frame)**

$$\vec{\beta}_{CF} = \frac{\sum_{j=1}^{N} \left(\frac{\vec{p_j}}{E_j}\right) \cdot P_j}{\sum_{j=1}^{N} P_j}$$

$$\begin{pmatrix} \gamma & -\beta_x\gamma & -\beta_y\gamma & -\beta_z\gamma \\ -\beta_x\gamma & 1+(\gamma-1)\frac{\beta_x^2}{\beta^2} & (\gamma-1)\frac{\beta_x\beta_y}{\beta^2} & (\gamma-1)\frac{\beta_x\beta_z}{\beta^2} \\ -\beta_y\gamma & (\gamma-1)\frac{\beta_y\beta_x}{\beta^2} & 1+(\gamma-1)\frac{\beta_y^2}{\beta^2} & (\gamma-1)\frac{\beta_y\beta_z}{\beta^2} \\ -\beta_z\gamma & (\gamma-1)\frac{\beta_z\beta_x}{\beta^2} & (\gamma-1)\frac{\beta_z\beta_y}{\beta^2} & 1+(\gamma-1)\frac{\beta_z^2}{\beta^2} \end{pmatrix}$$

- **The zero-component of the transformed four-current is the relevant density**

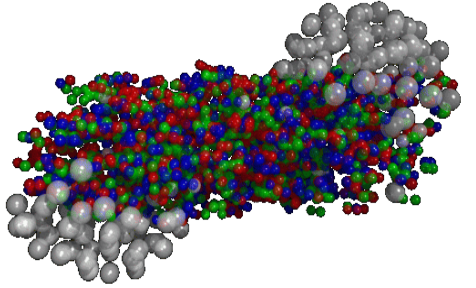


Density calculation

- Local baryon density is the zeroth component of the baryon four-current $j^{\mu} = (\rho_B, \vec{j})$ when the baryon is at rest
- UrQMD calculates in the Computational Frame (CF), which is usually the CMS (due to symmetry)
- $j^{\mu}_{CF} = (\rho_{B_{CF}}, \vec{j}_{CF})$ can be calculated as a sum over Gaussians

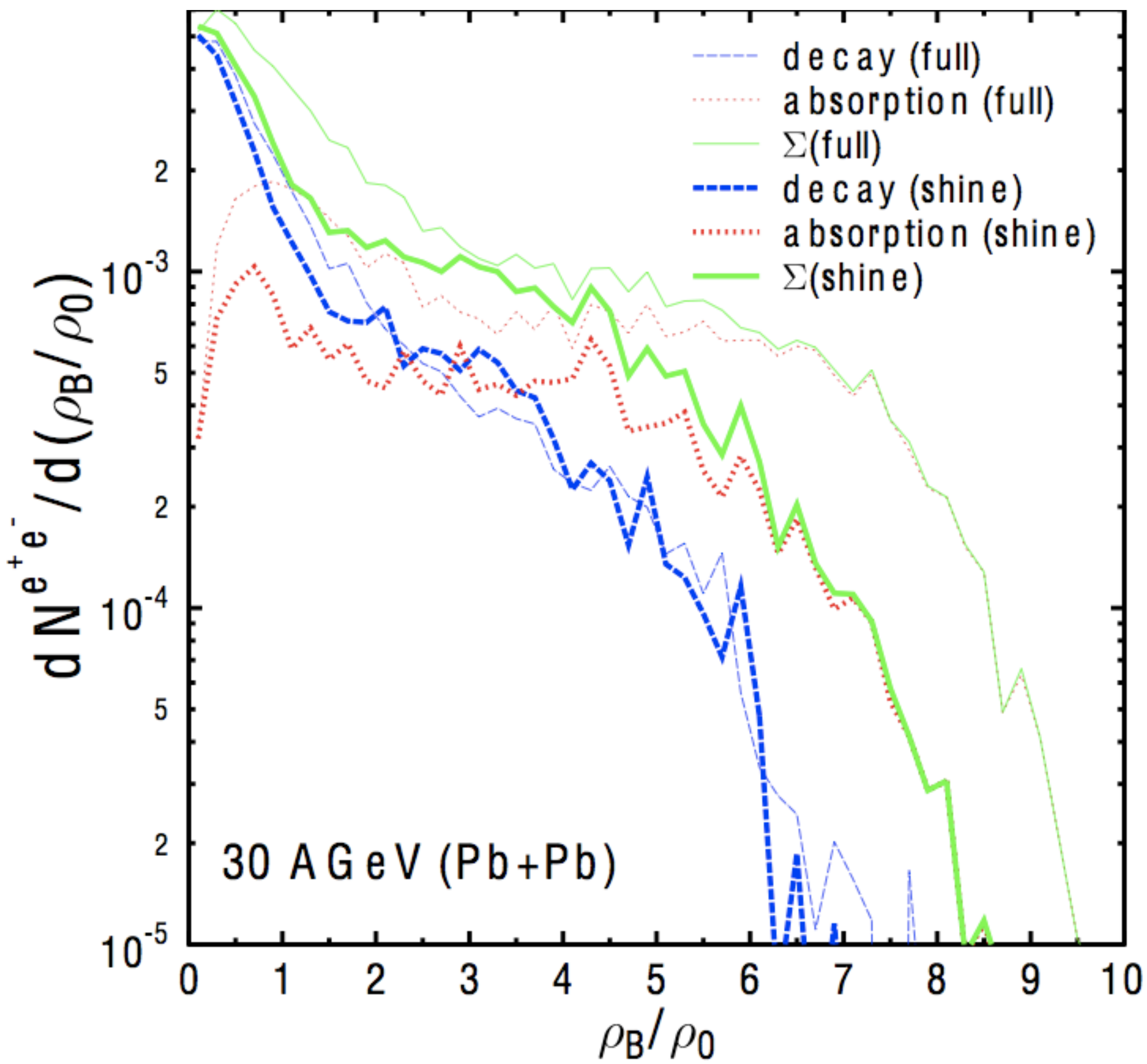
$$\rho_{CF}(\vec{r_i}) \quad = \sum_{j=1}^{N} \left(\frac{1}{\sqrt{2\pi}\sigma}\right)^3 \gamma_z \, \mathrm{e}^{\left(-\frac{(x-x_0)^2+(y-y_0)^2+(z-z_0)^2\gamma_z^2}{2\sigma^2}\right)}$$

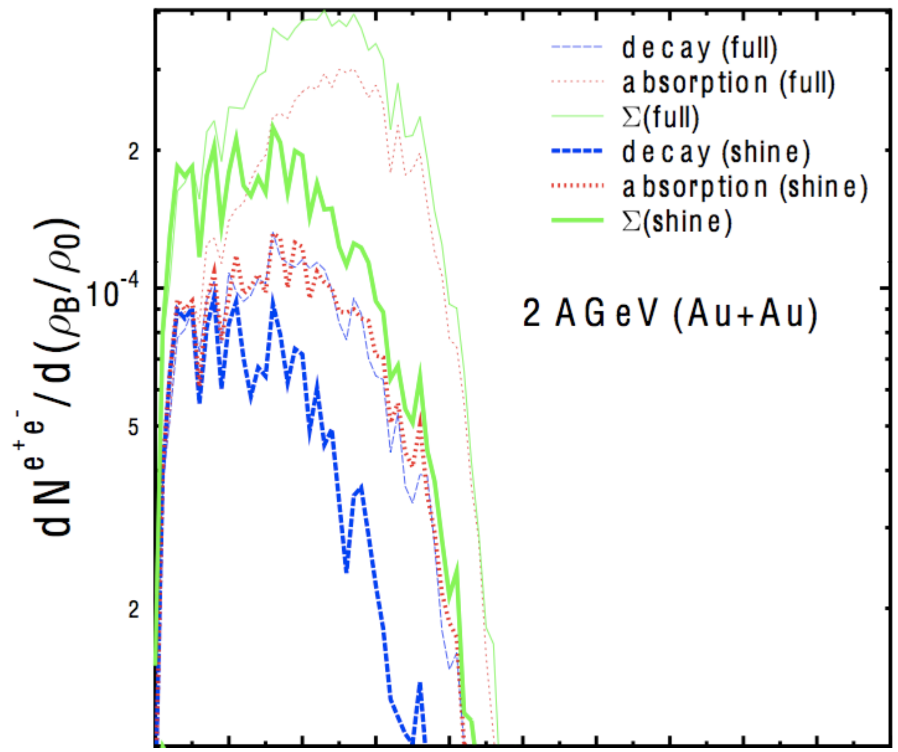
$$= \sum_{j=1}^{N} P_j$$

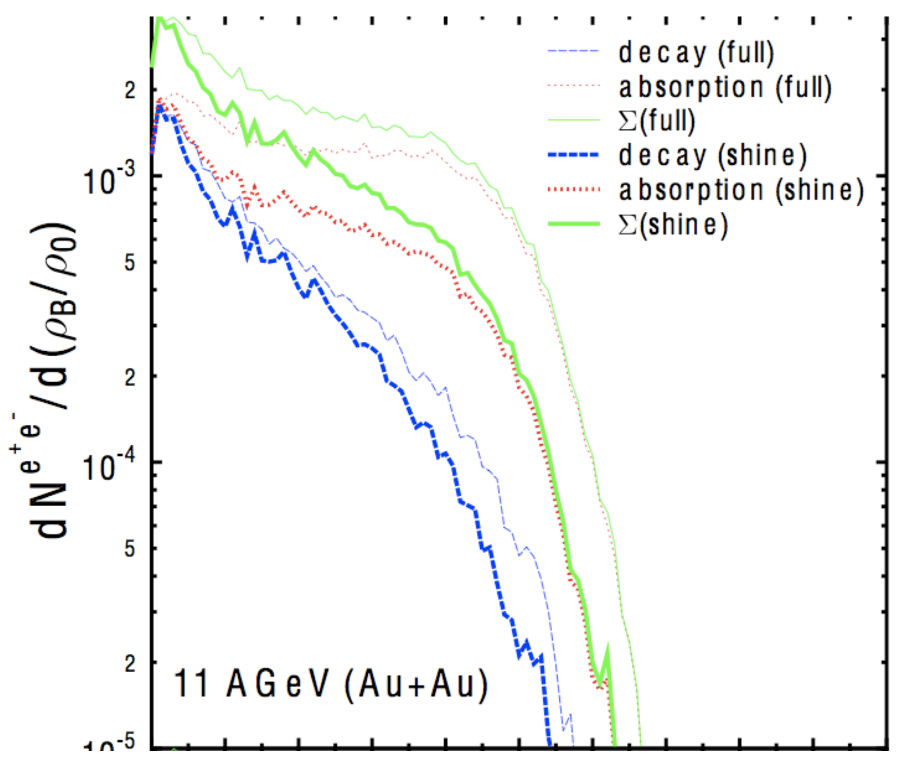


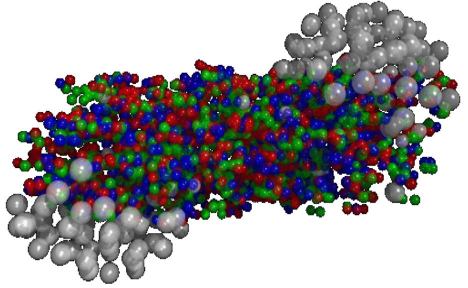
Dileptons

Even in the most optimistic approach dileptons only reach out to 2-3 ρ_0
Shining approach only reaches out to 1-2 ρ_0



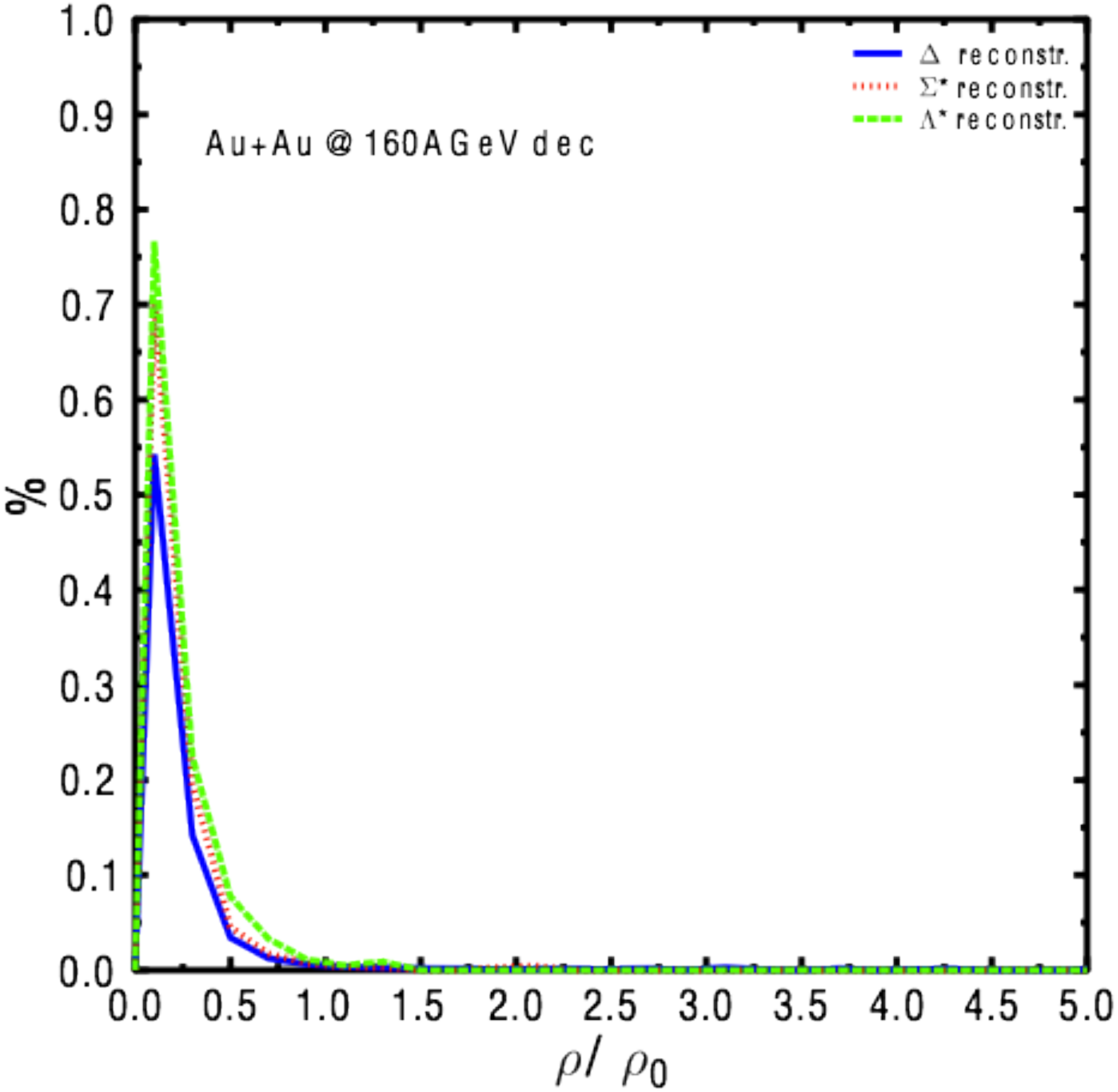


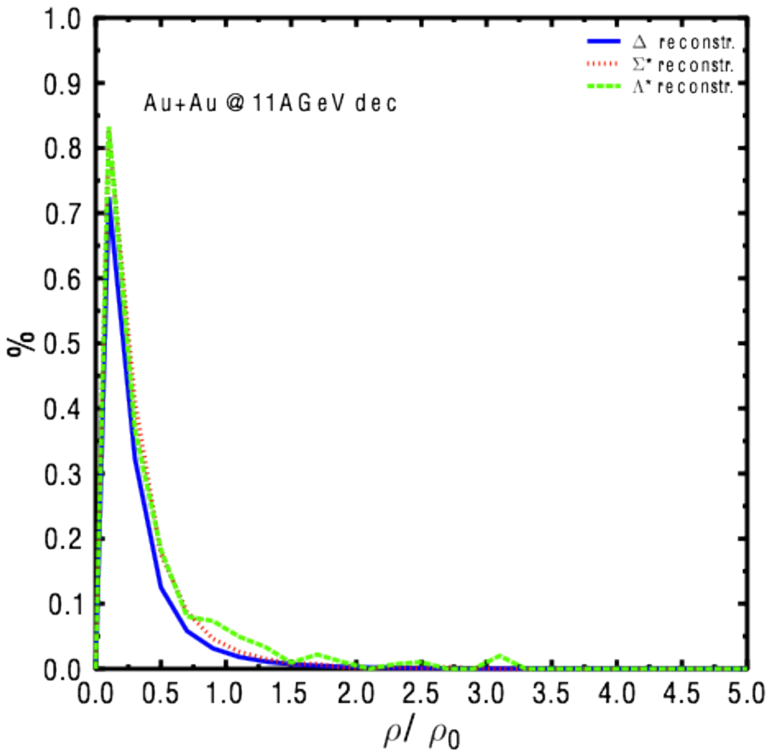


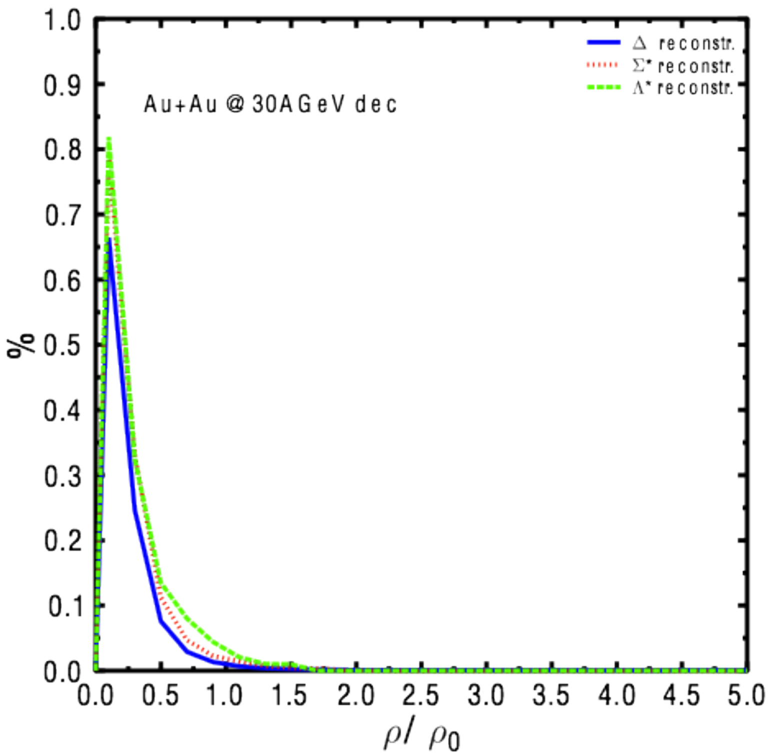


Baryon resonances @ decay

Density evaluated at point of decay
Only resonances that decay at low density are reconstructable

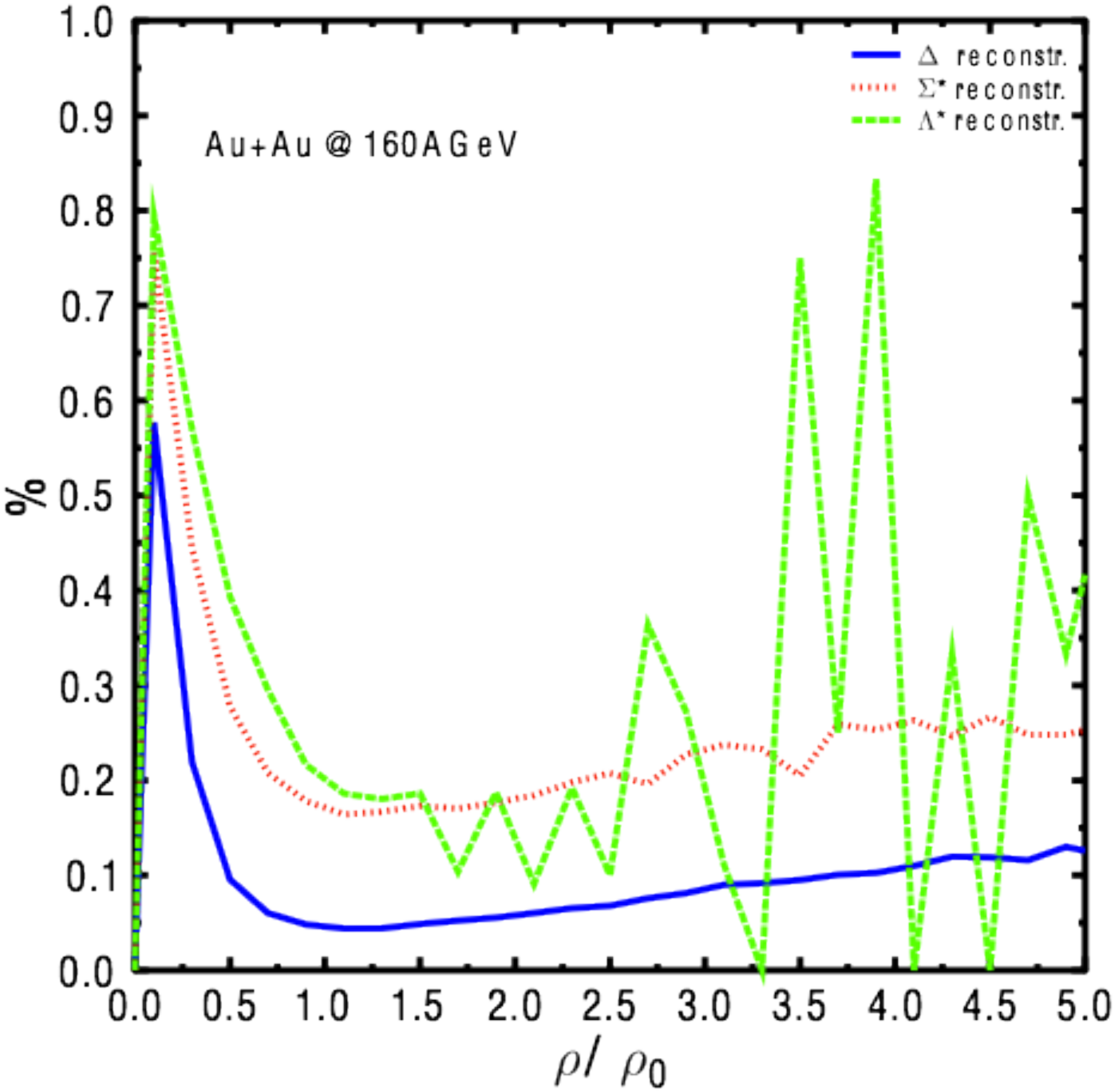


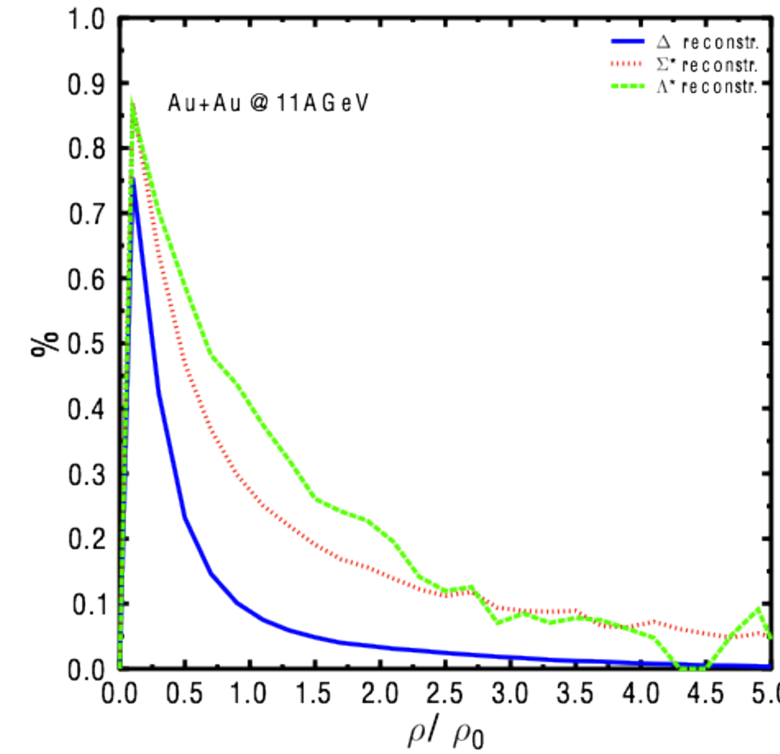


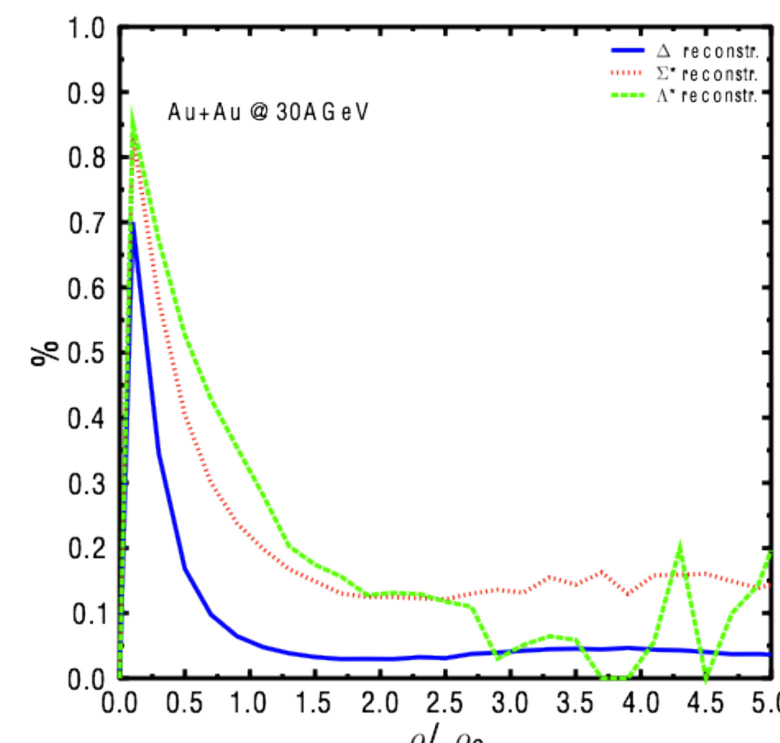


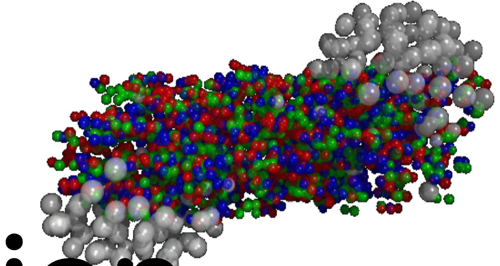
Baryon resonances @ production

Density evaluated at point of production
At higher density resonances become more reconstructable again (p_T effect?)
Roughly scales with lifetime of the resonance



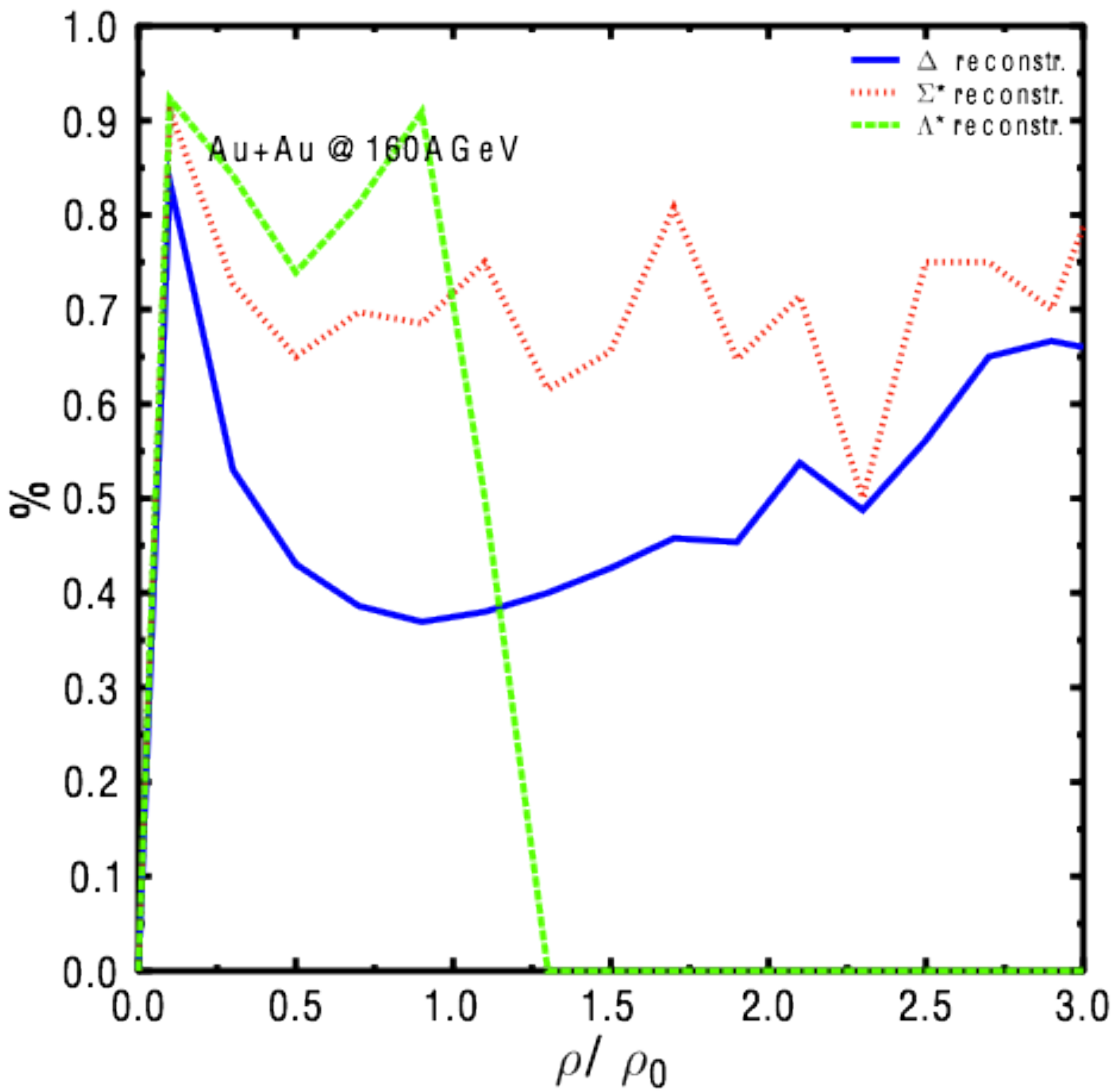


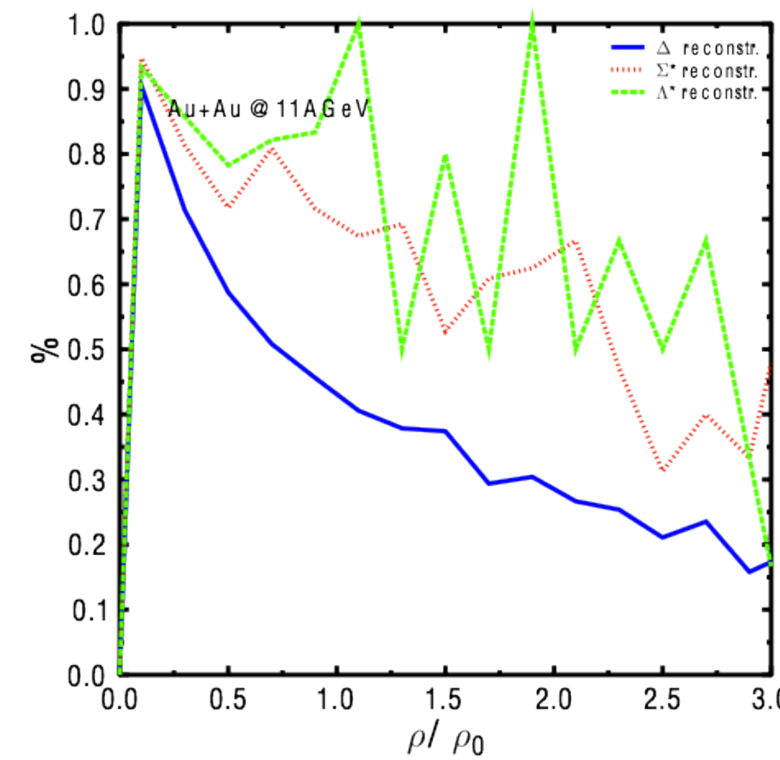


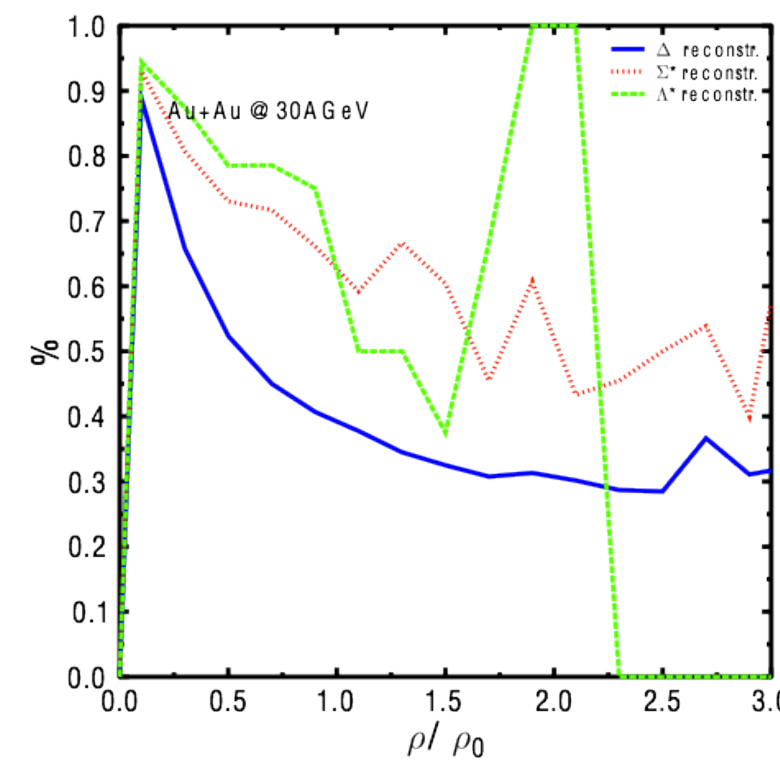


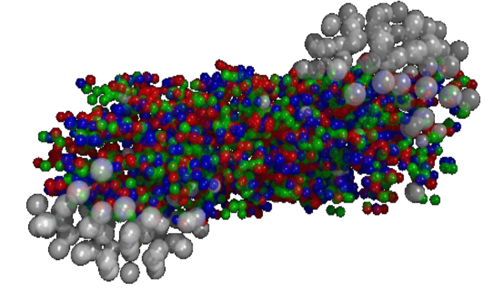
Baryon resonances @ production

p_T cut applied (p_T > 2GeV)
"Reconstructability" increases at higher density
Problem: High p_T resonances might not be the most interesting to study

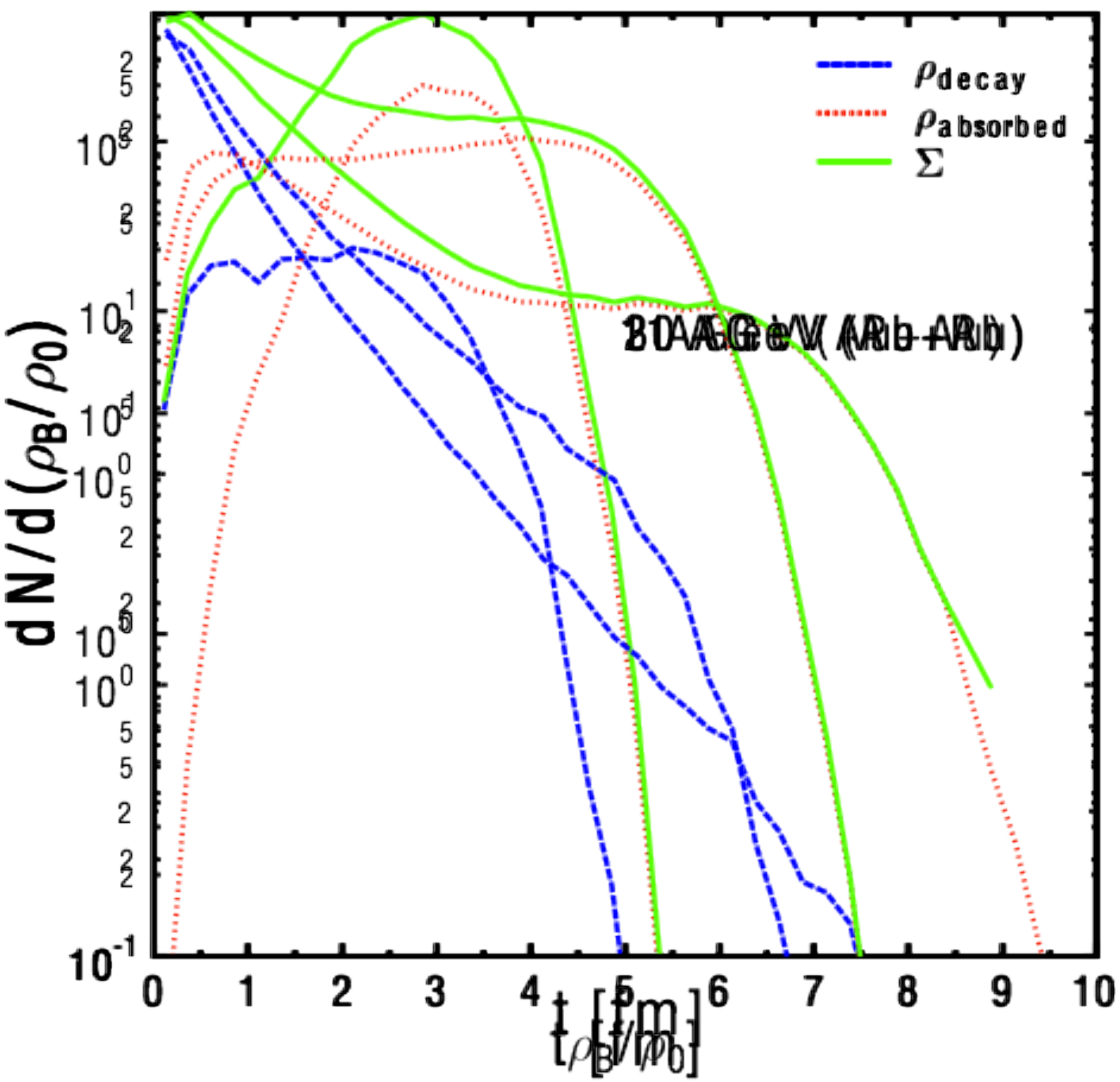


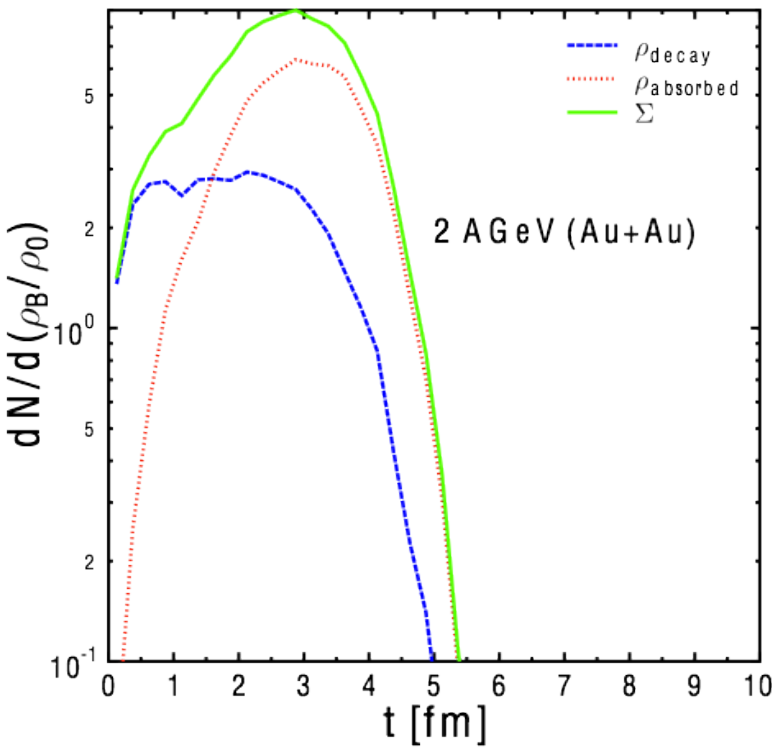


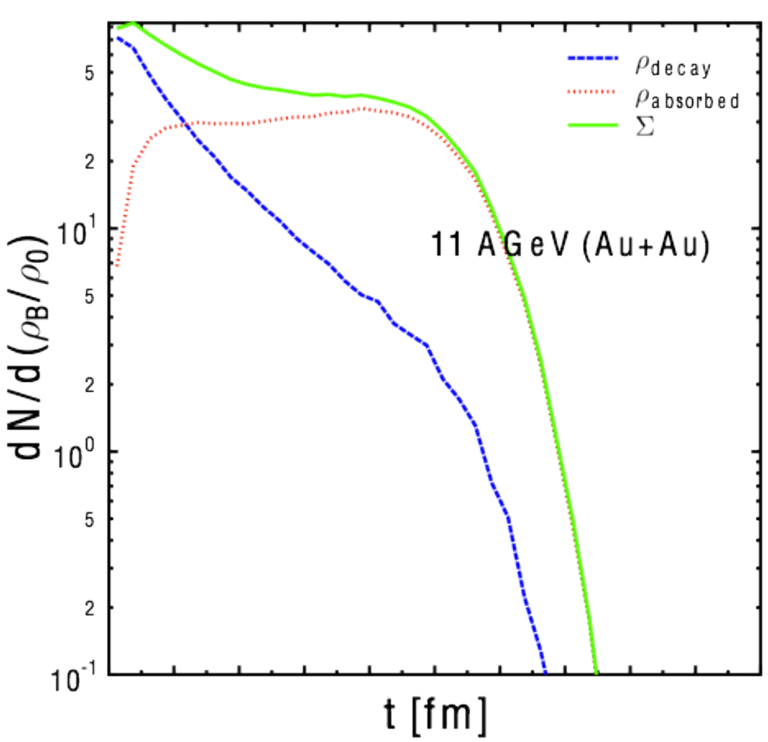


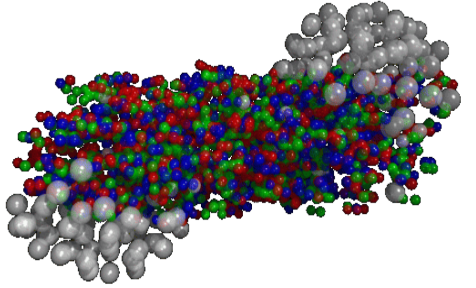


Density distribution



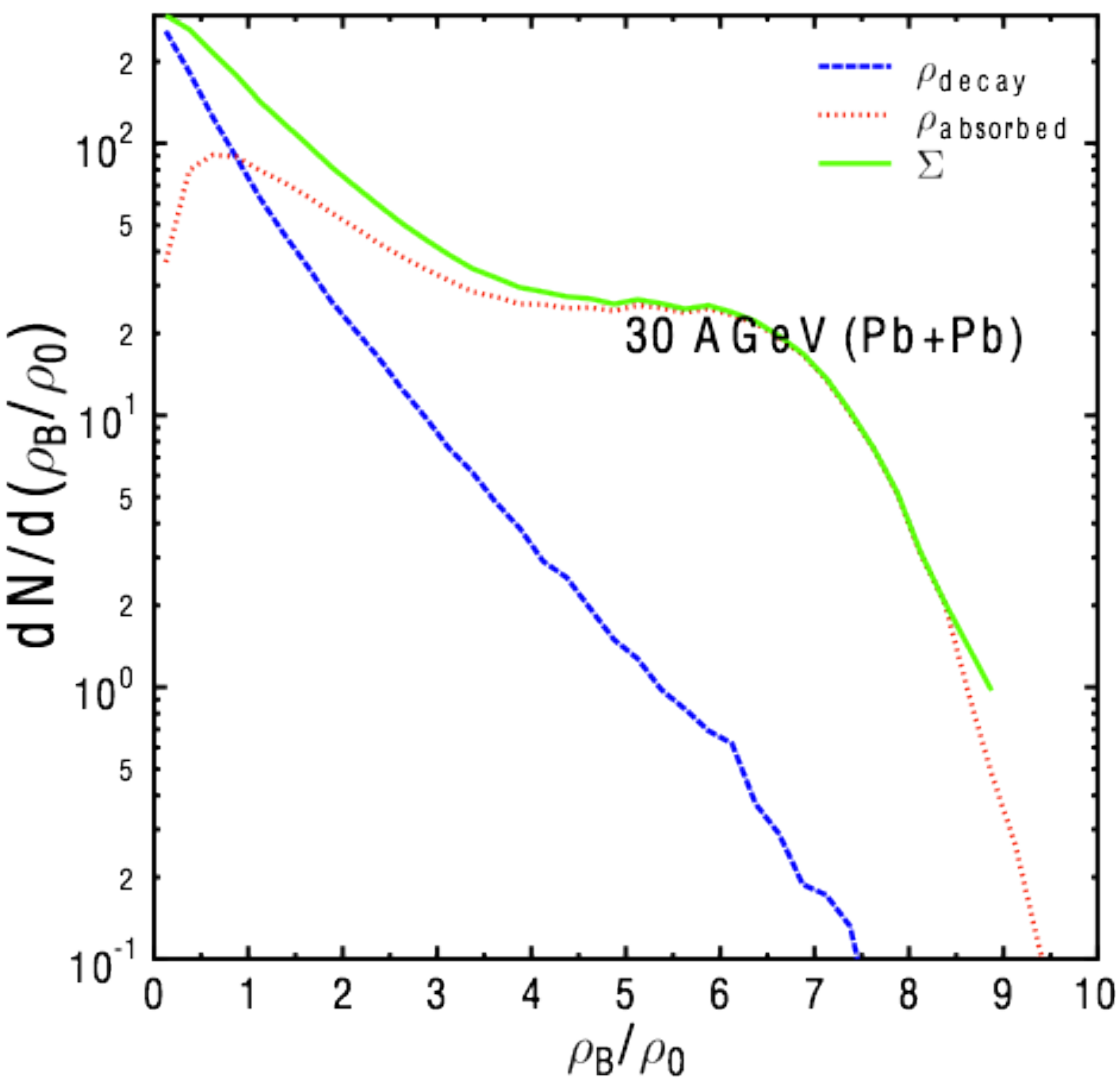


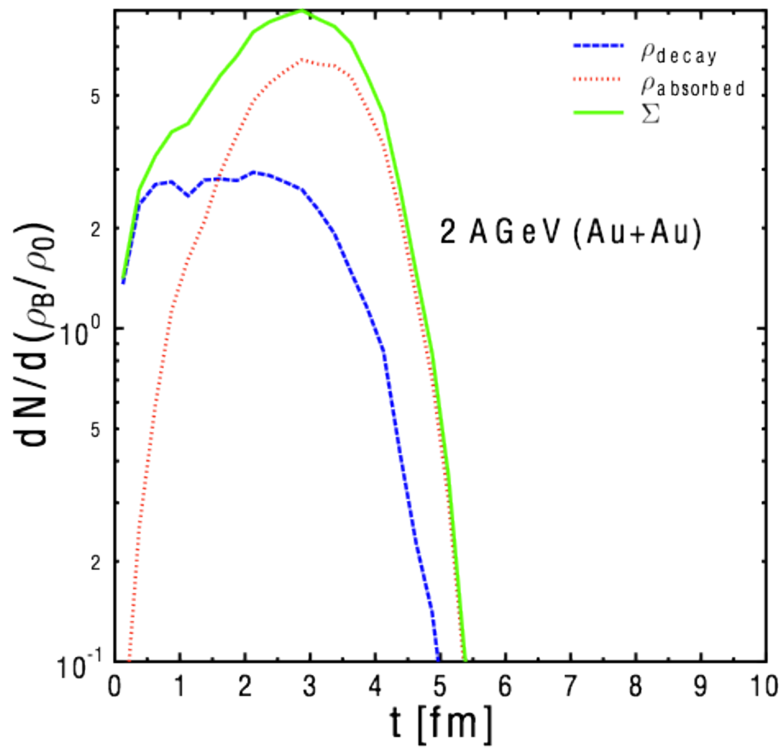


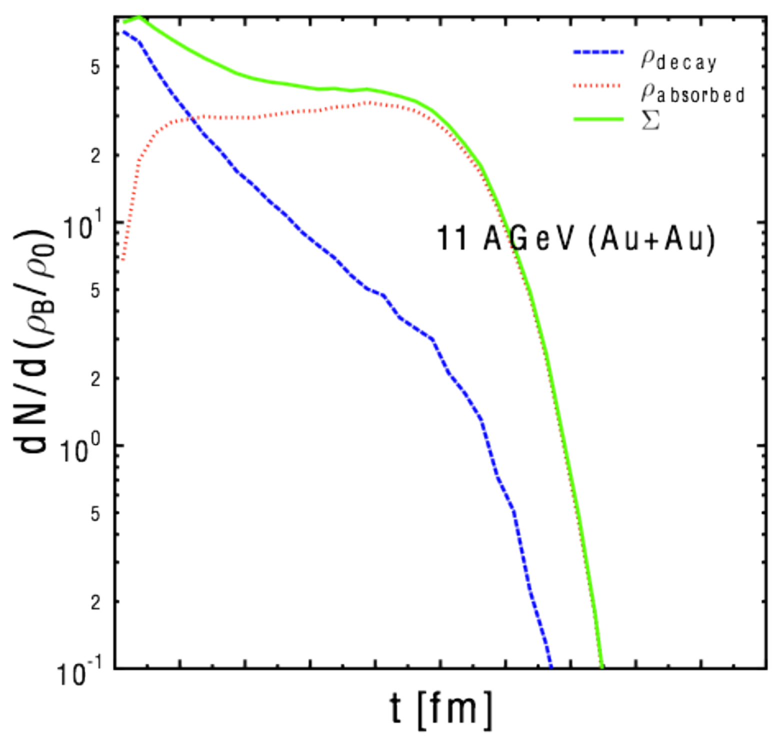


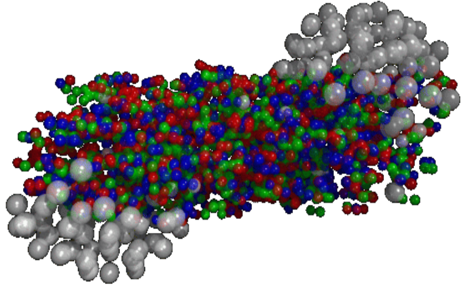
Density distribution

- ρ decays do not reach out to high density
- most resonances at such densities are re-absorbed

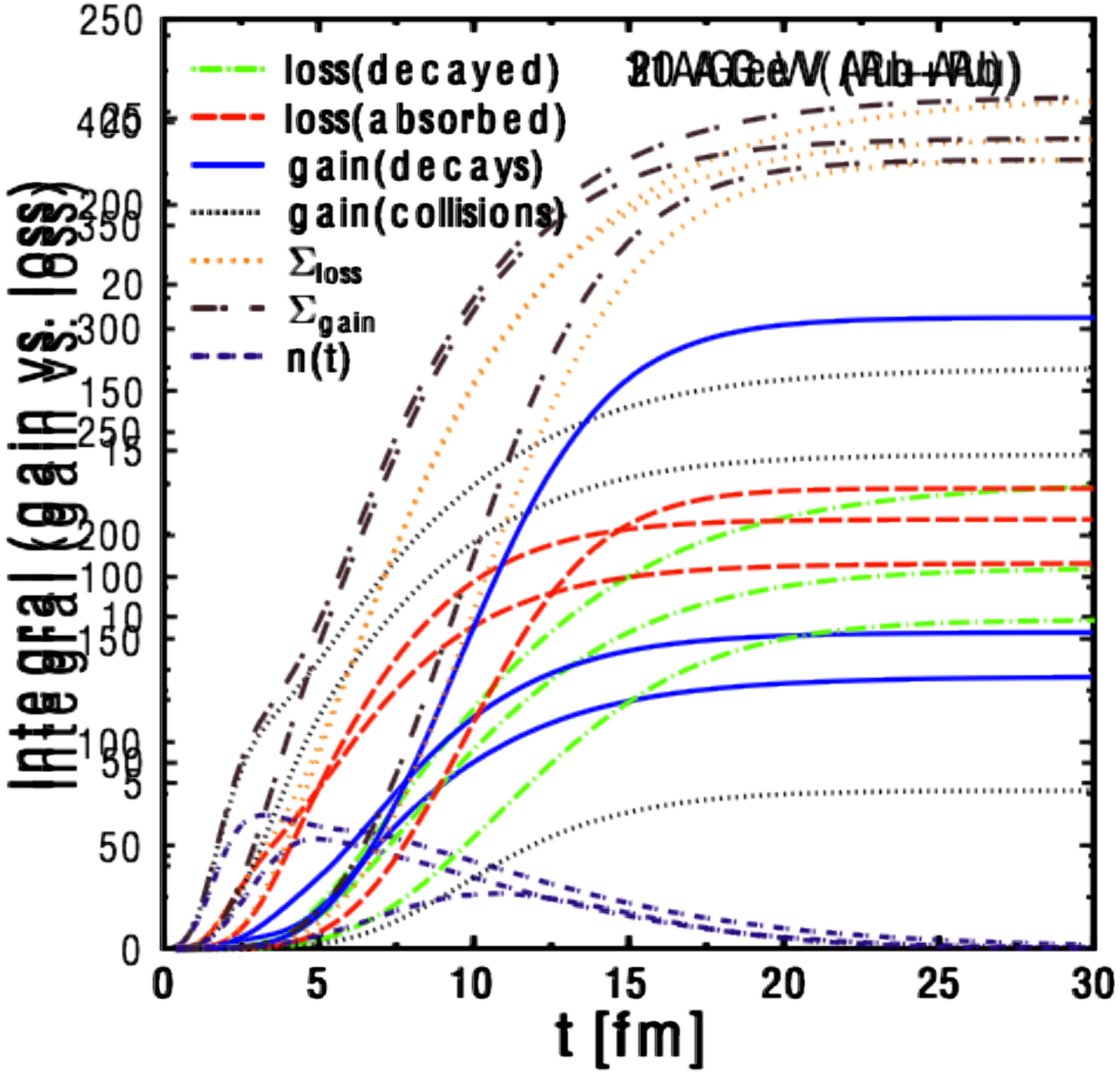


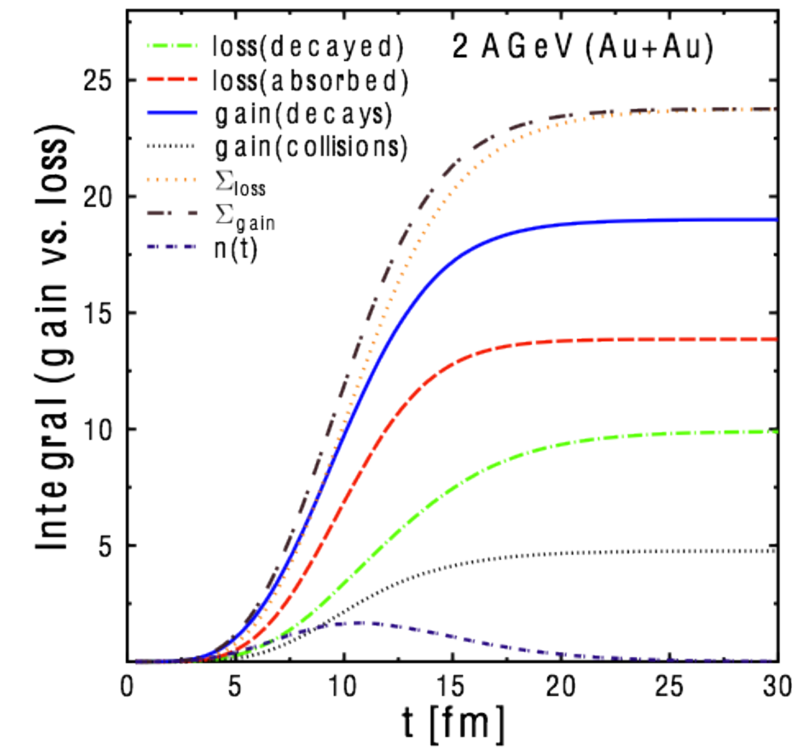


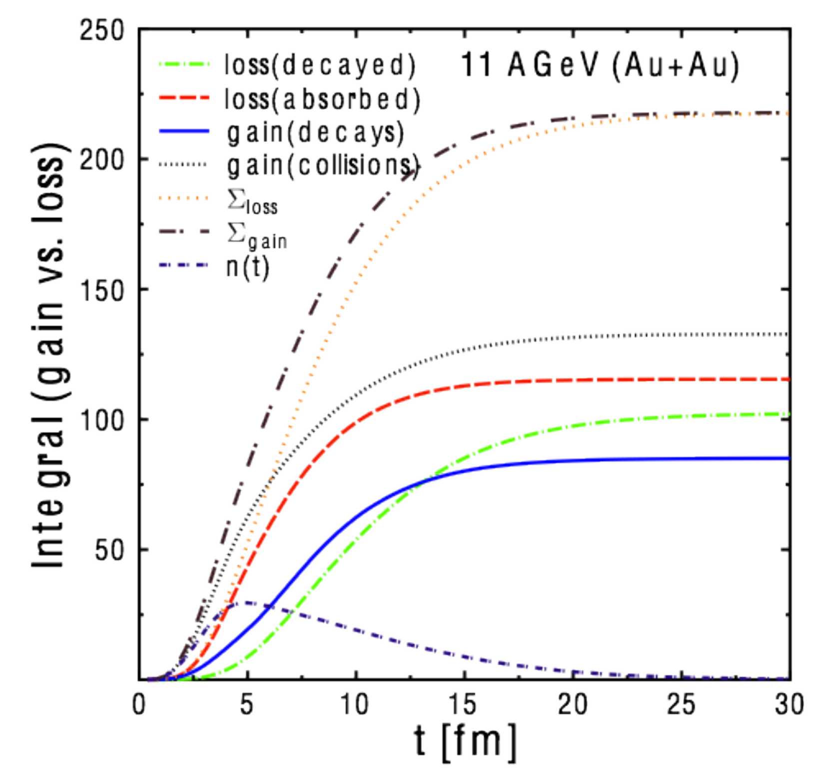


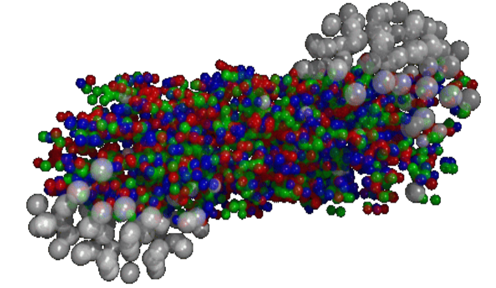


Integral values







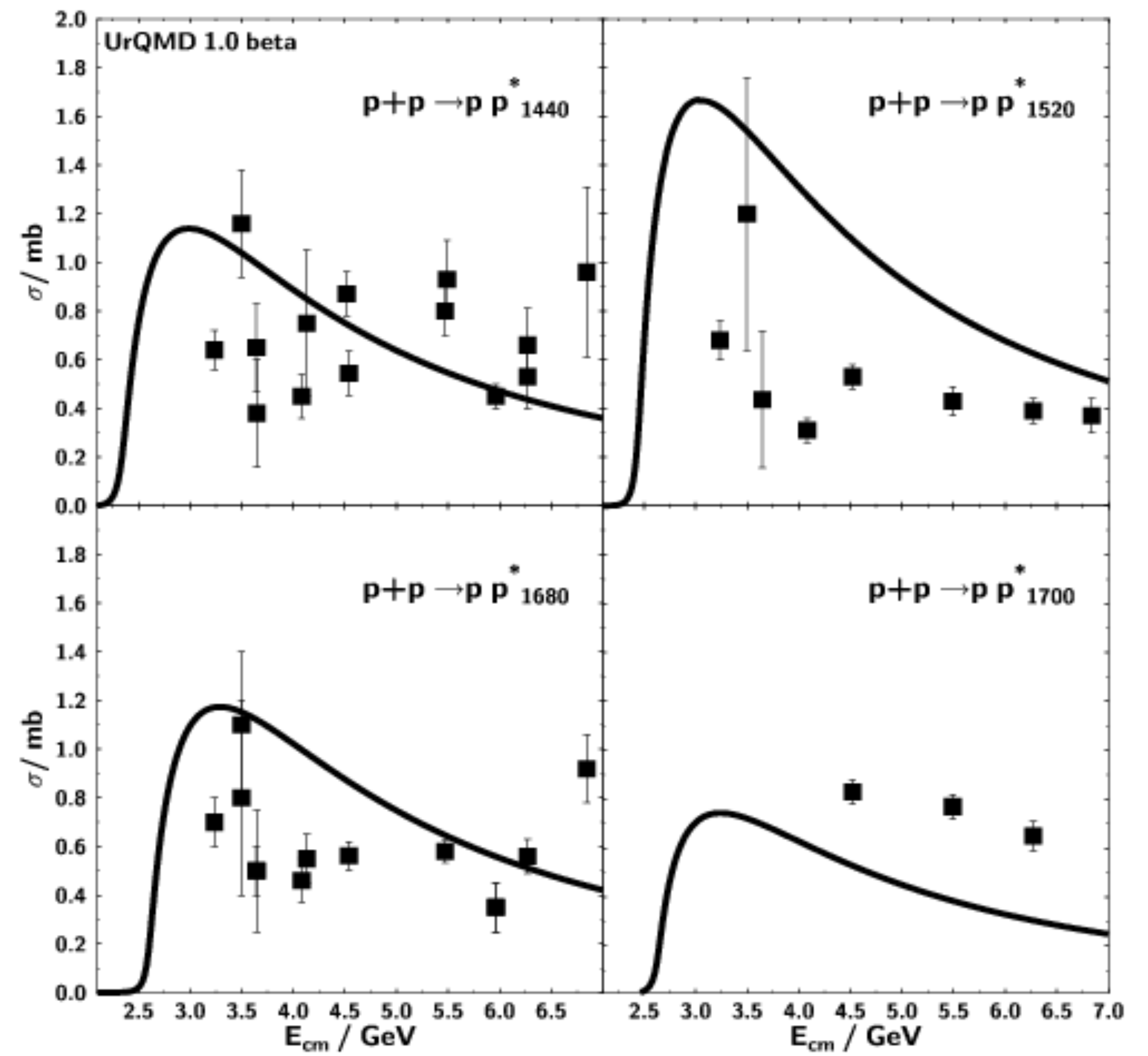


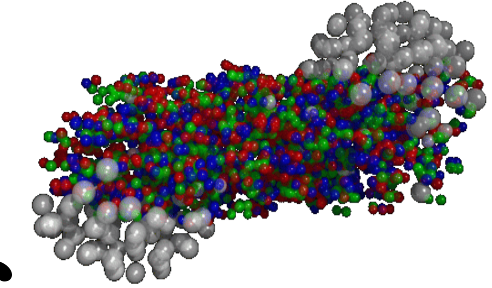
Cross sections

$$\sigma_{1,2\to3,4}(\sqrt{s}) \sim (2s_3+1)(2s_4+1)\frac{< p_{3,4} >}{< p_{1,2} >}\frac{1}{\sqrt{s}}|M(m_3,m_4)|^2$$

Global fit with the same kind of matrix element

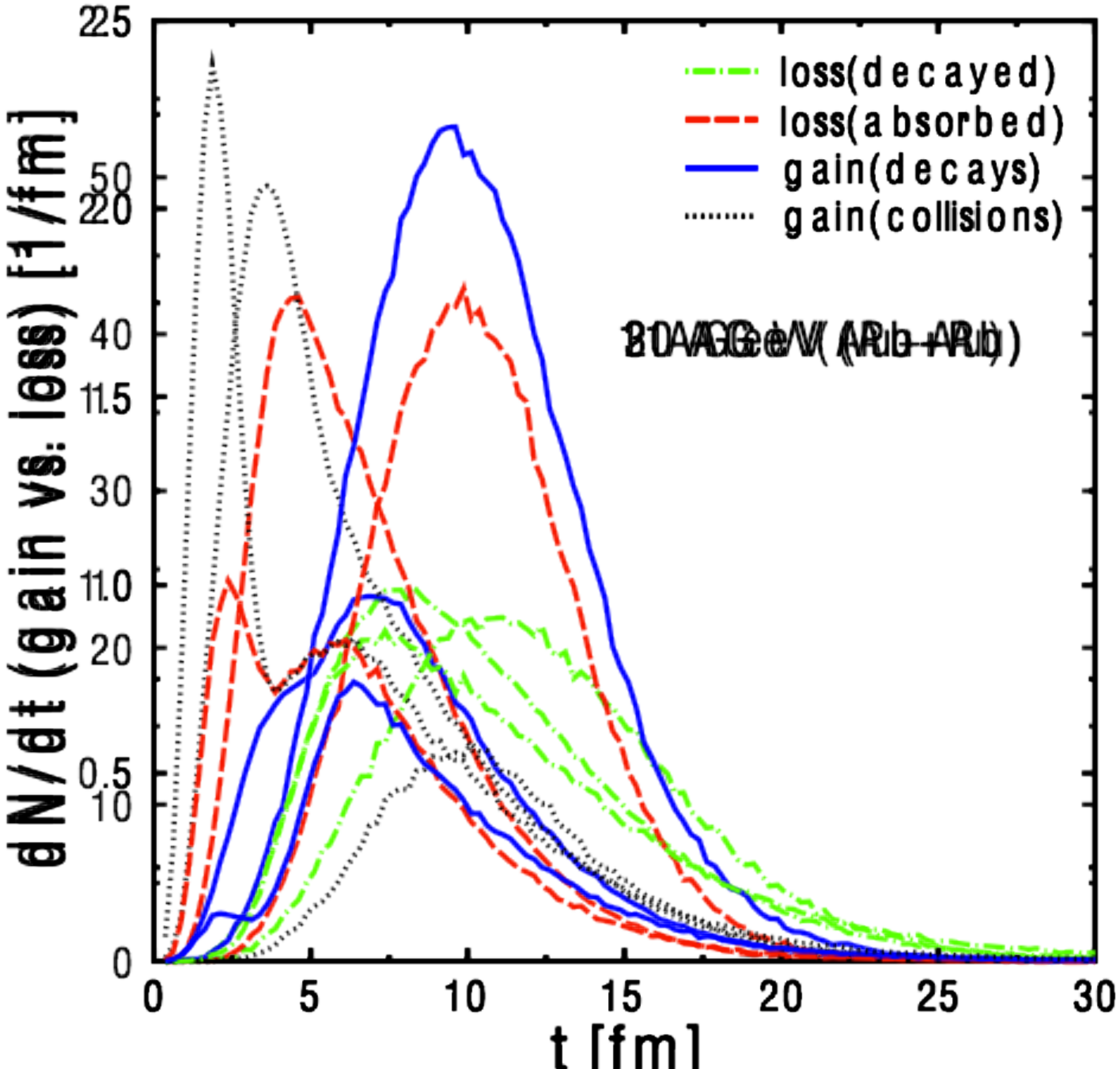
$$|M(m_3,m_4)|^2 = A\frac{1}{(m_4-m_3)^2(m_4+m_3)^2}$$

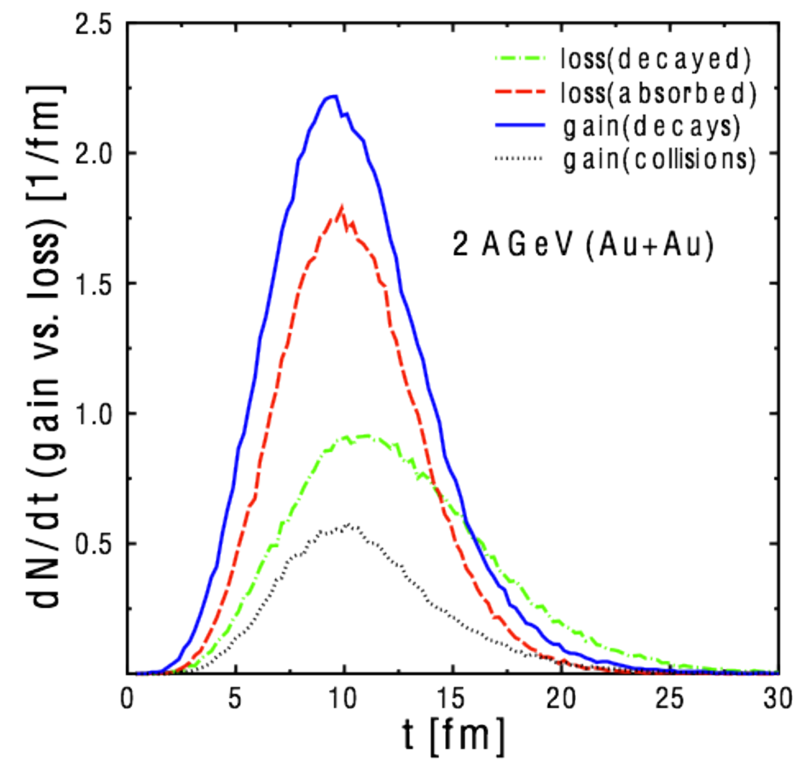


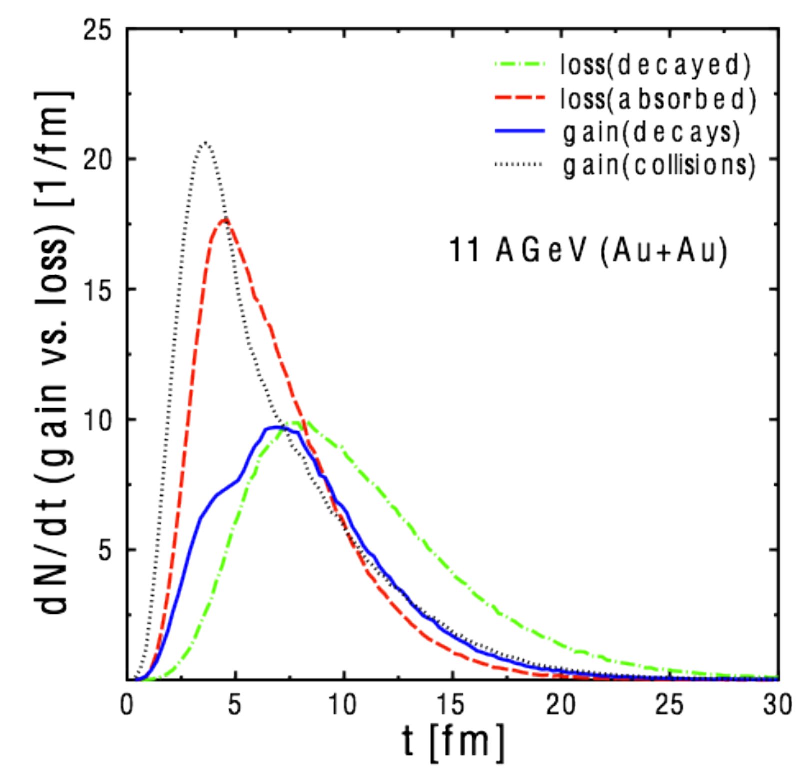


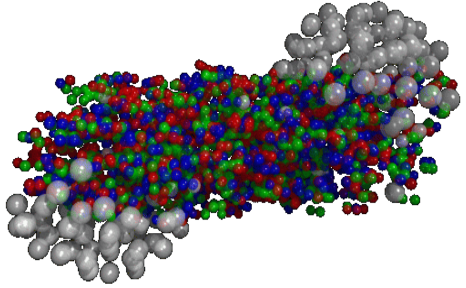
Gain/Loss rates of ρ mesons

SiS energies: Most gain from decay
AGS and FAIR energies: More gain from collisions
In the early stage - loss by absorption dominant
Decays set in in the late stage



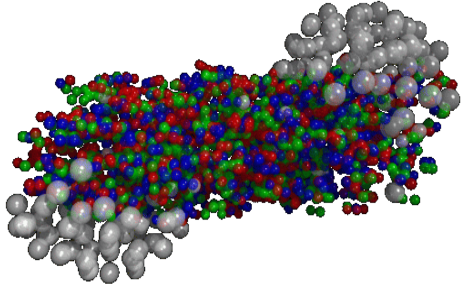






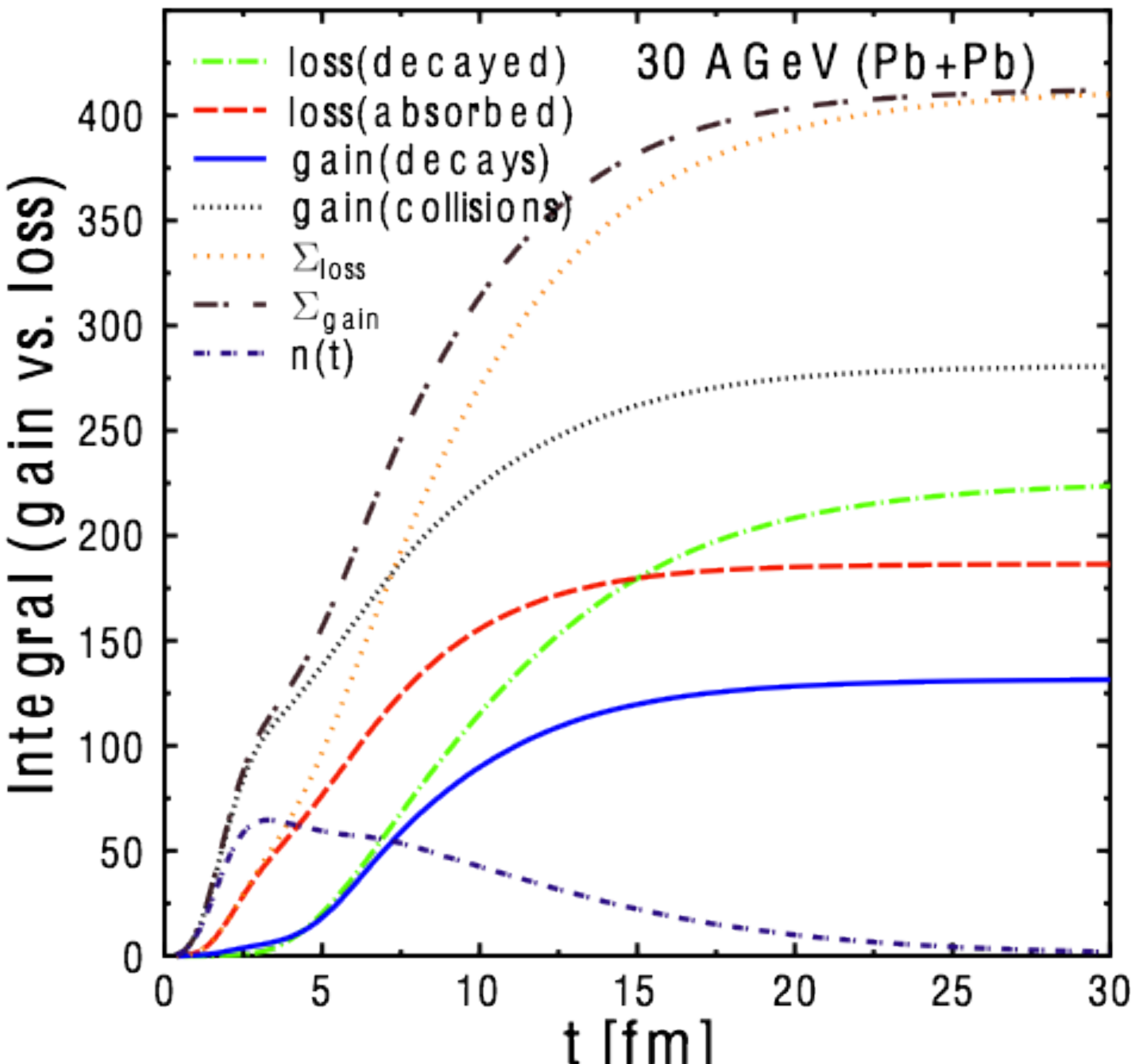
Gain/Loss terms

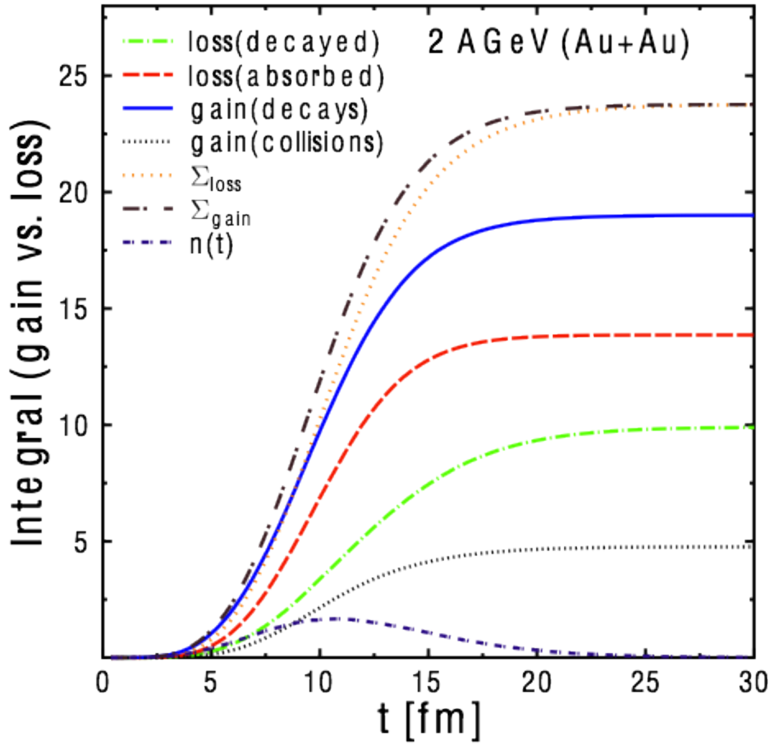
- Resonances can stem from two processes
  - Collisions (e.g. $\pi\pi \rightarrow \rho$)
  - Decays of heavier resonances (e.g. $N^*_{1520} \rightarrow N + \rho$)
- Resonances can be destroyed by two processes
  - Decays (e.g. $\rho \rightarrow e^+e^-$)
  - Absorption (e.g. $N + \rho \rightarrow N^*_{1520}$)

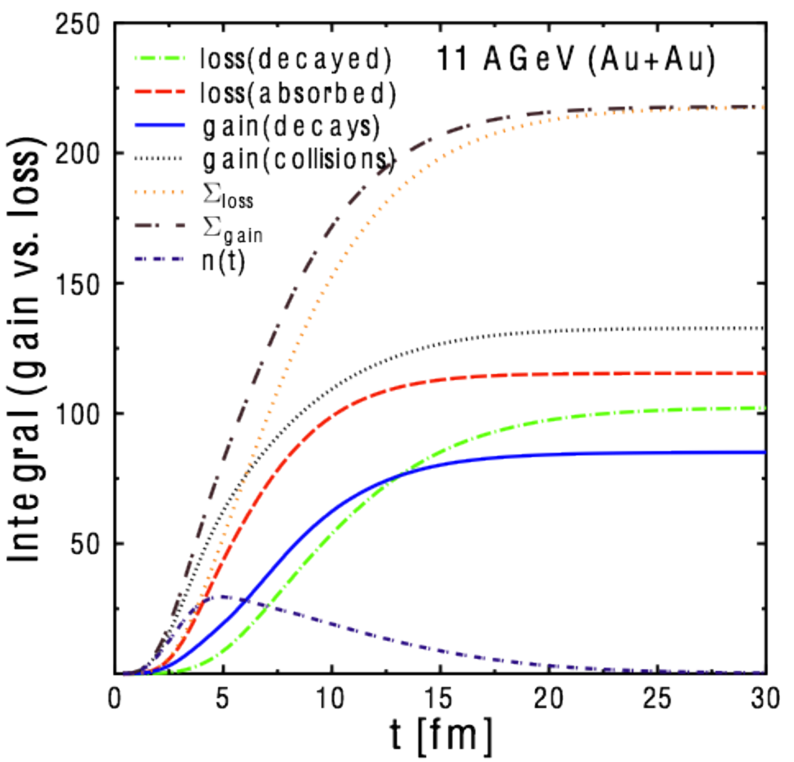


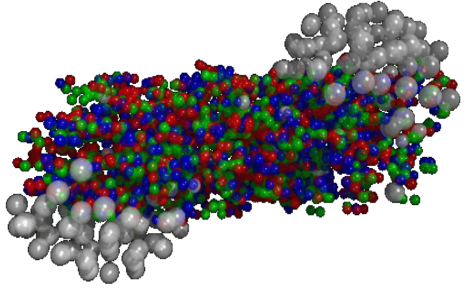
Integral values

- Consistency check: Sum of gain and collision agree
- Difference gives the number of resonances in the system









a_1 meson

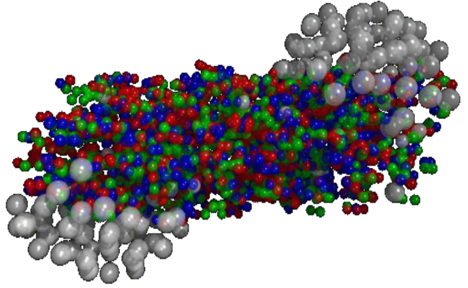
What about the other channels?

Experimentally not feasible:

Higher mass resonances are either not known or the decay channel analyses contradict each other (further exp. studies certainly useful!).

$a_1(1260)$ DECAY MODES

	Mode	Fraction (Γ_i/Γ)
Γ_1	$\pi^+\pi^-\pi^0$	
Γ_2	$\pi^0\pi^0\pi^0$	
Γ_3	$(\rho\pi)_{S-\text{wave}}$	seen
Γ_4	$(\rho\pi)_{D-\text{wave}}$	seen
Γ_5	$(\rho(1450)\pi)_{S-\text{wave}}$	seen
Γ_6	$(\rho(1450)\pi)_{D-\text{wave}}$	seen
Γ_7	$\sigma\pi$	seen
Γ_8	$f_0(980)\pi$	not seen
Γ_9	$f_0(1370)\pi$	seen
Γ_{10}	$f_2(1270)\pi$	seen
Γ_{11}	$K\overline{K}^*(892)+$ c.c.	seen
Γ_{12}	$\pi\gamma$	seen



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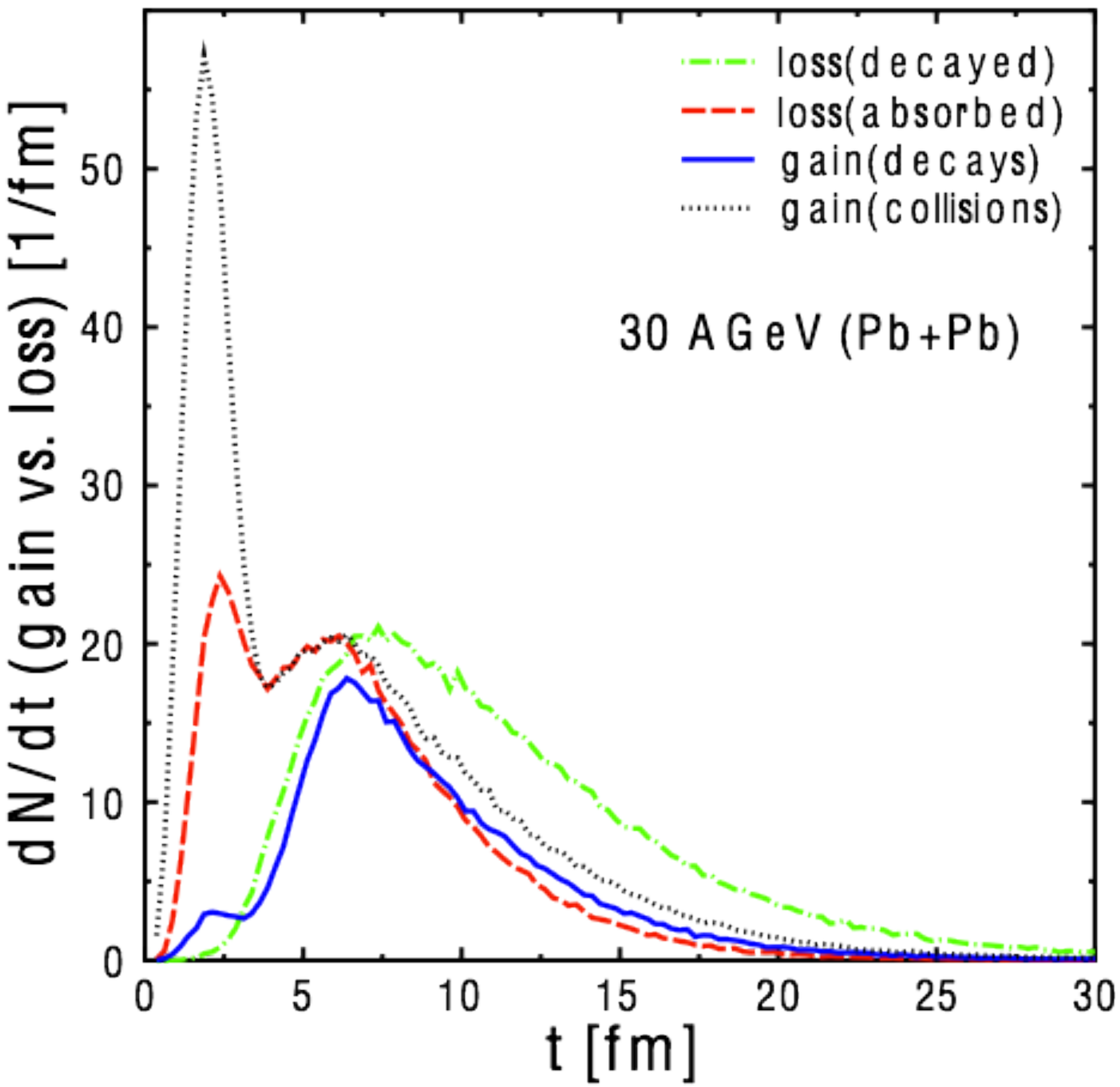
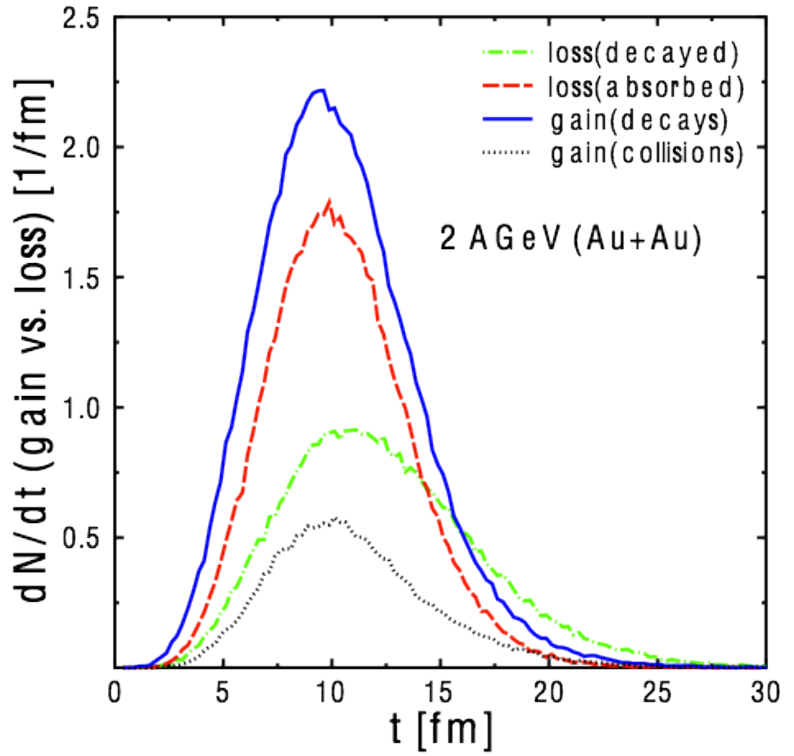
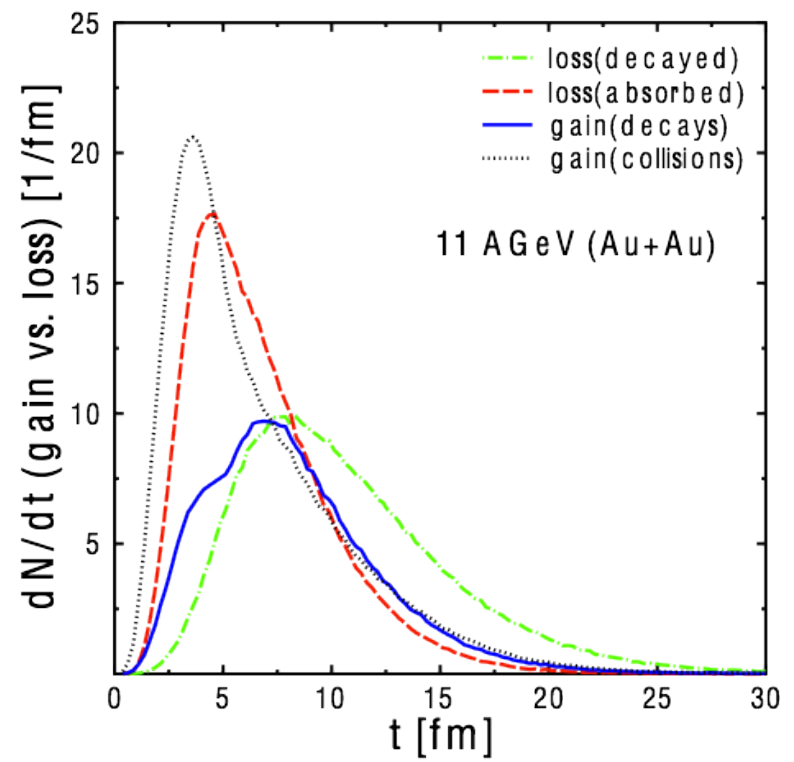
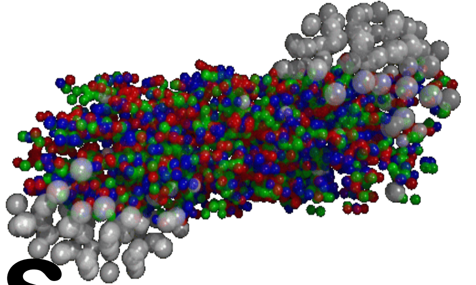


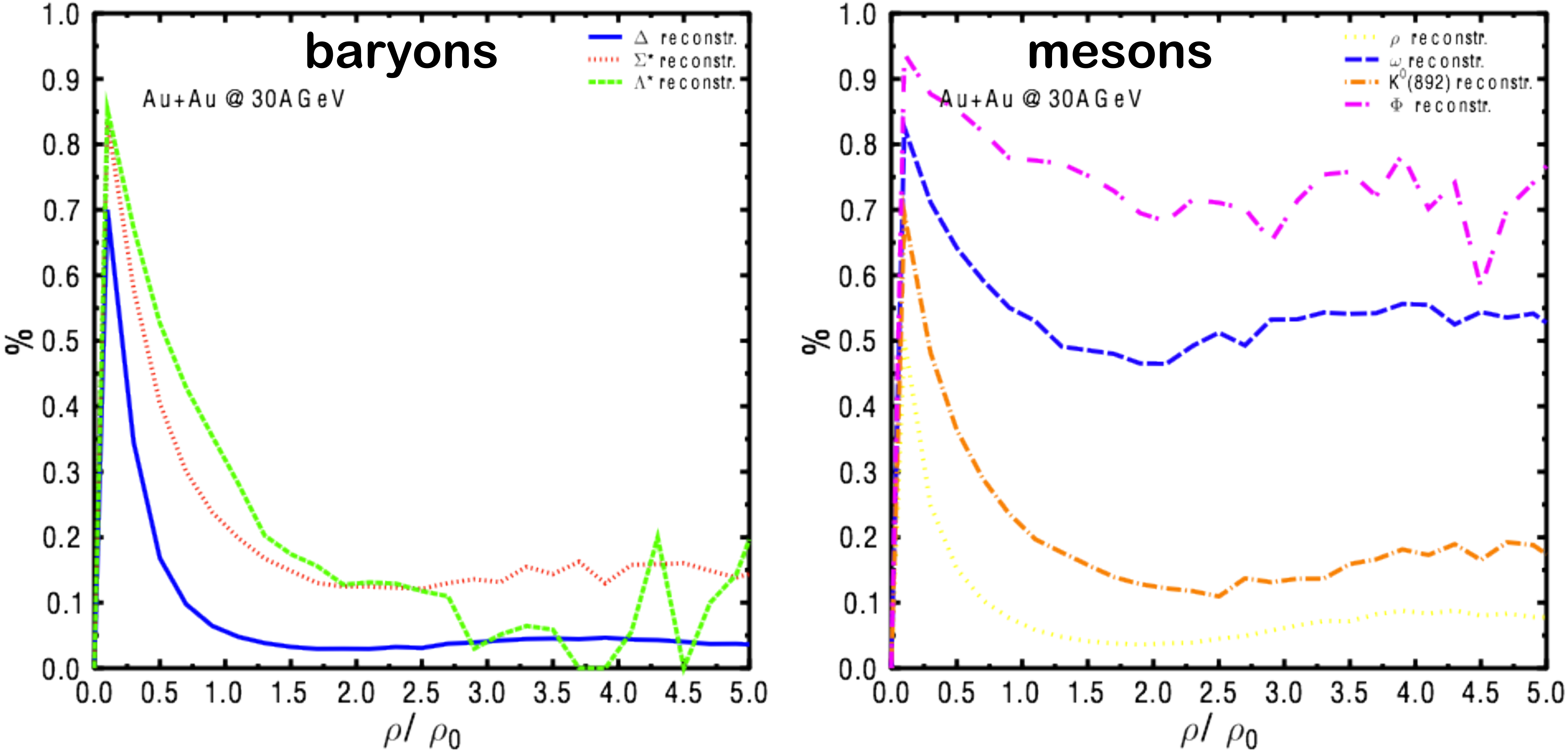




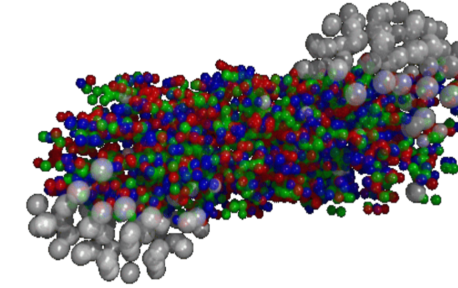
Percentage of rec. resonances



Which percentage can you reconstruct at a given density?

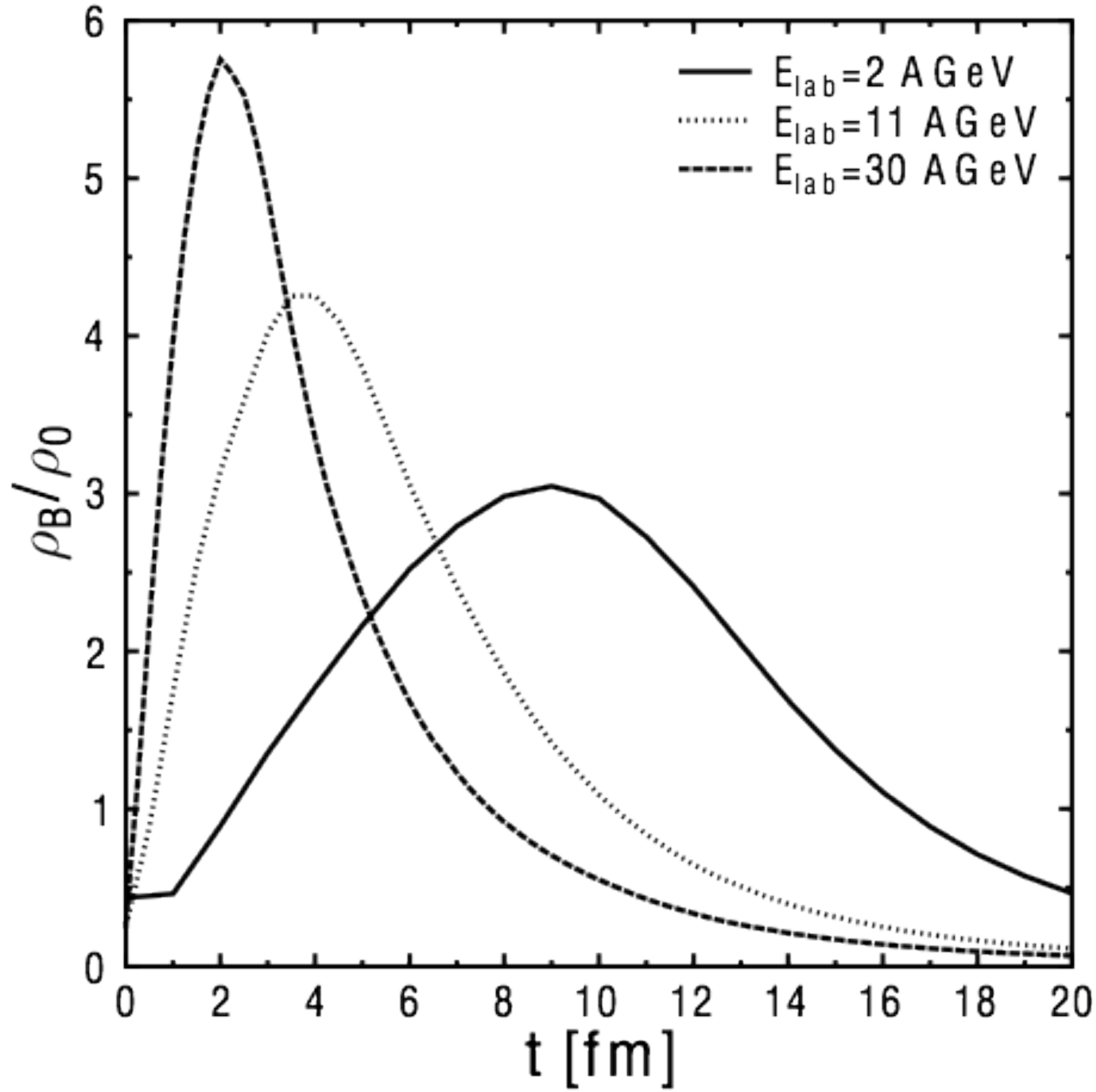


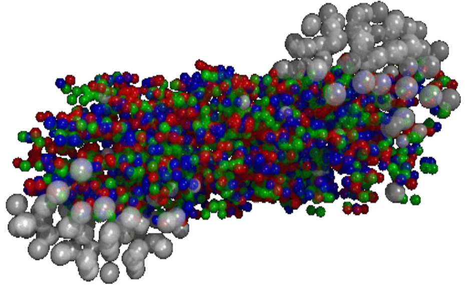
Clear ordering with lifetime, saturates at higher densities.



Time evolution of ρ_B

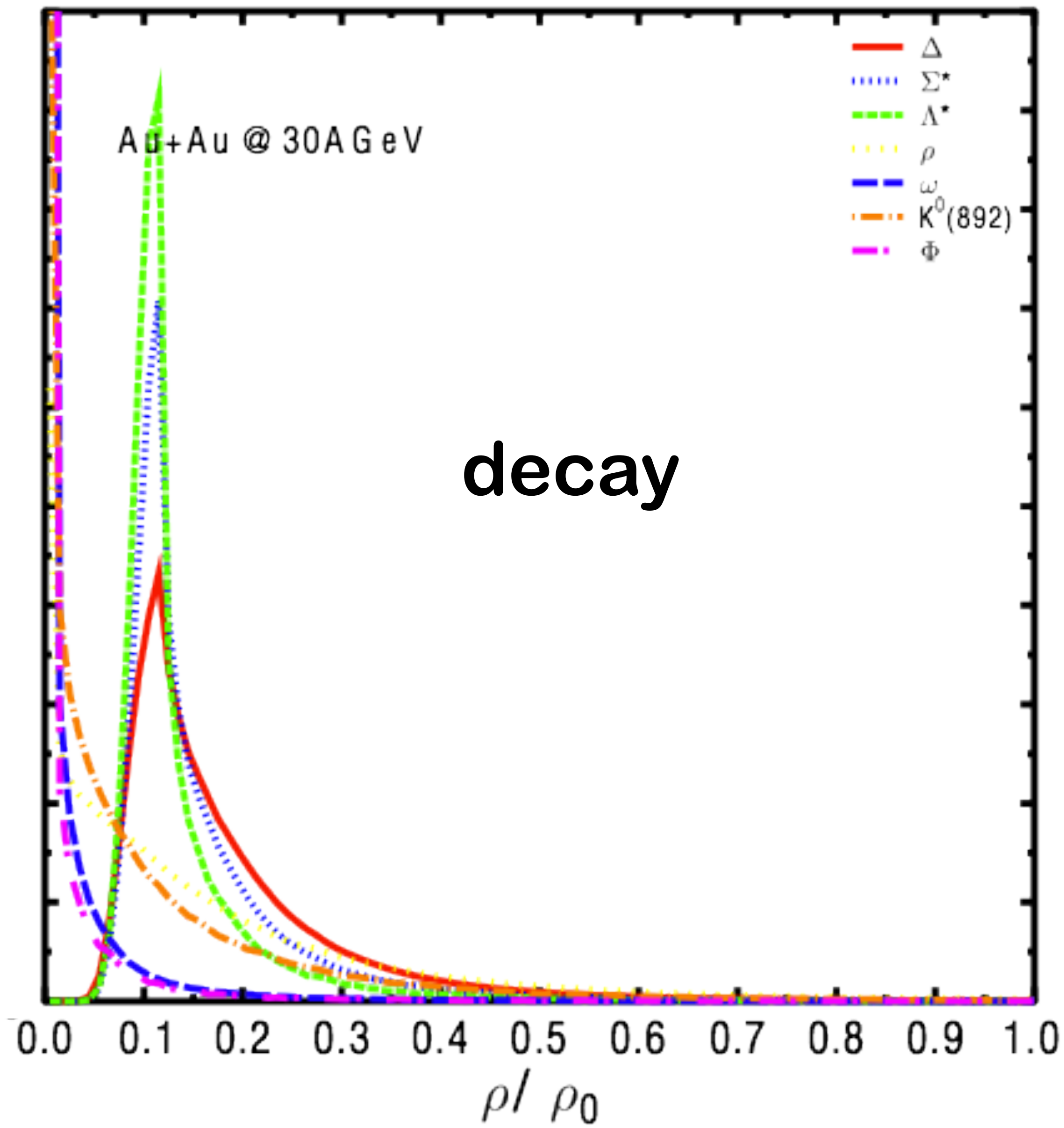
- Central Pb+Pb (Au+Au) collisions
- Averaged over all hadron positions

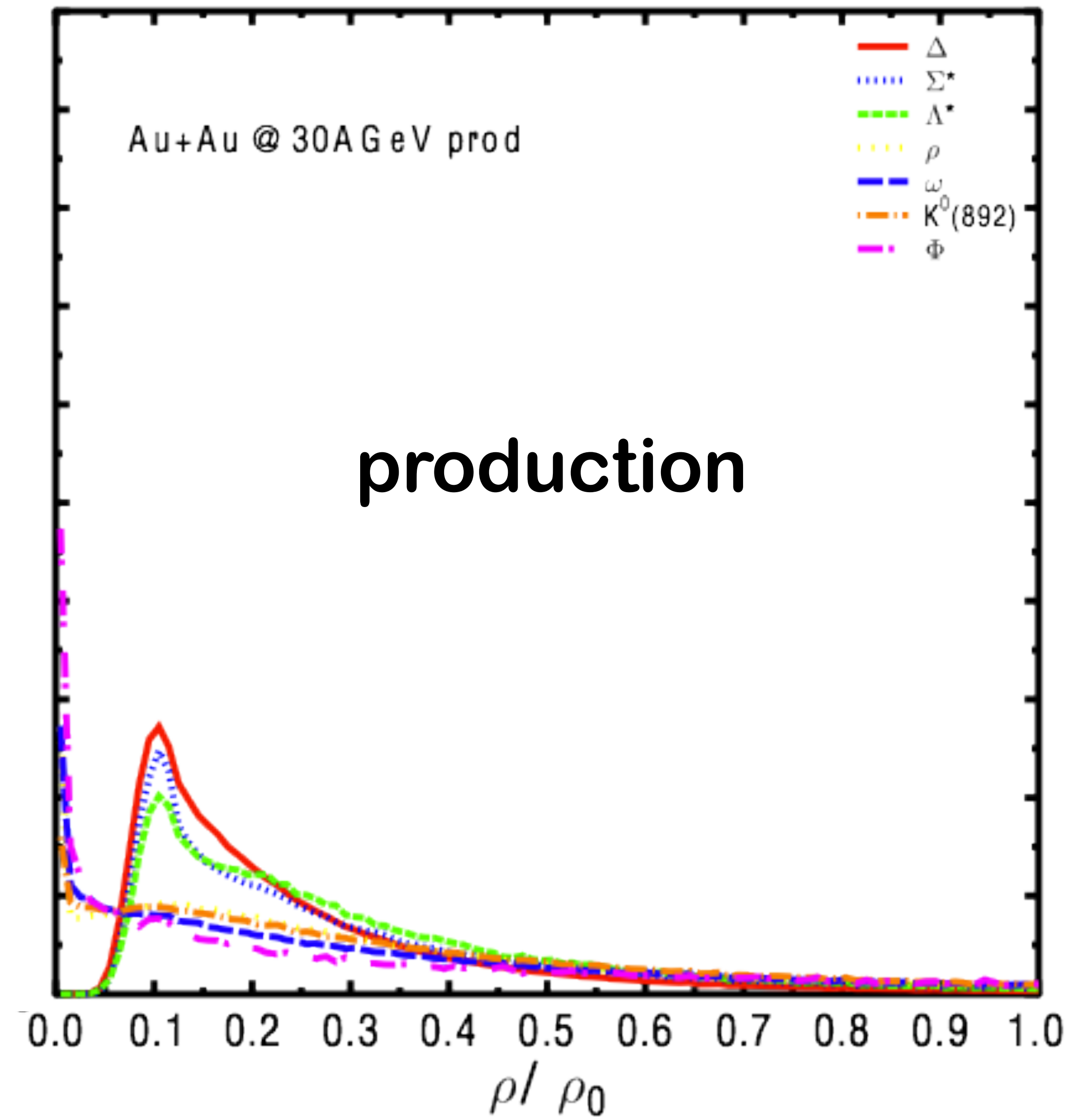




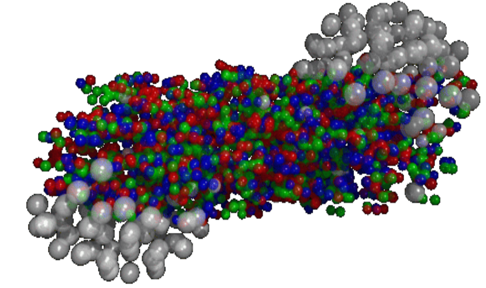
Density?

What percentage of resonances originates from what density?





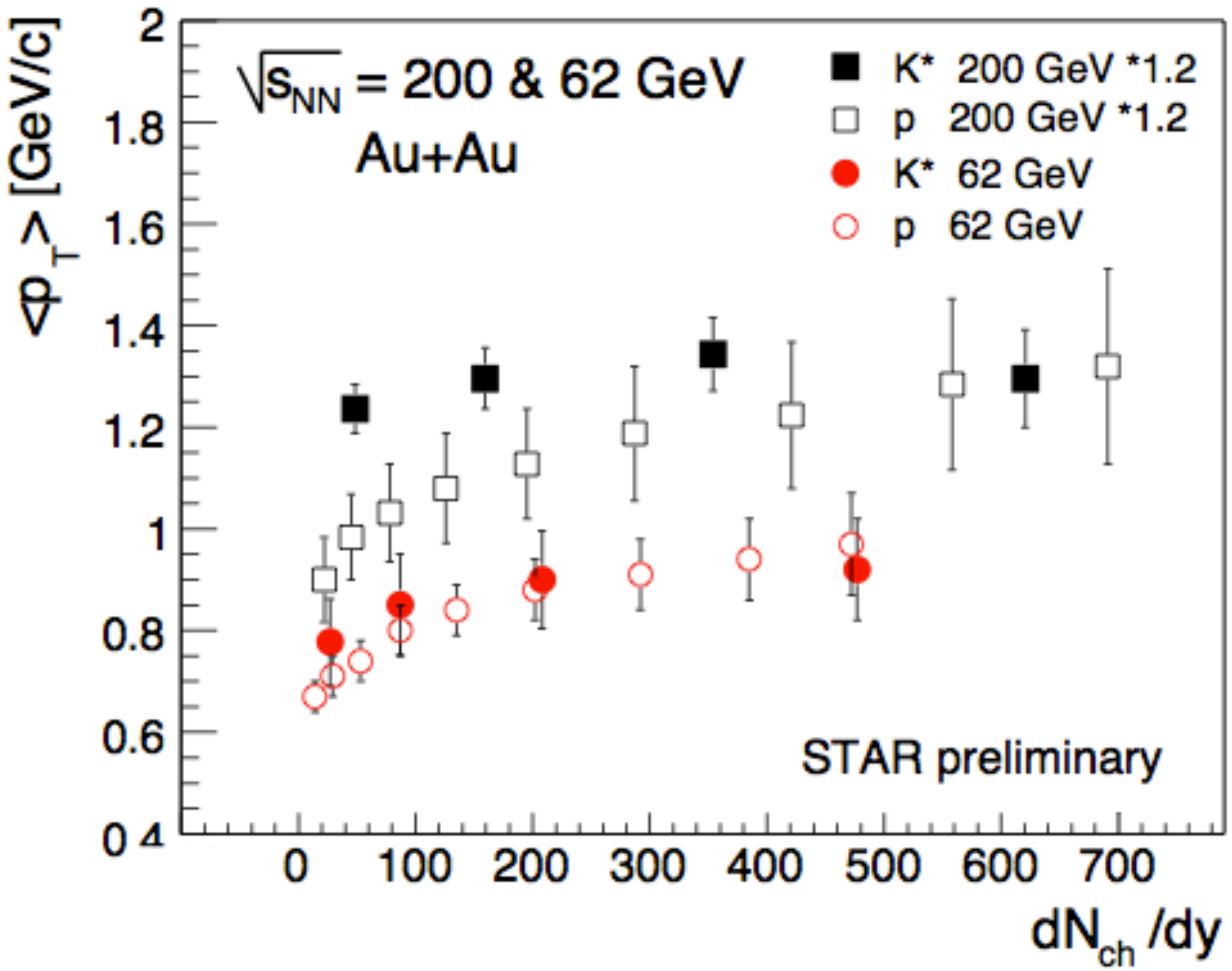
Most resonances are from the low density region of the collision.



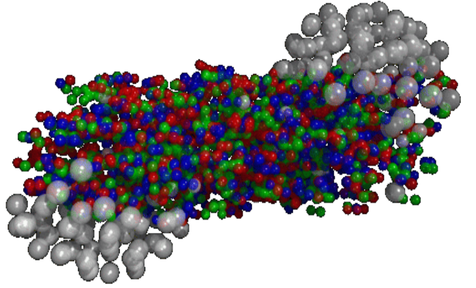
p_T dependence (exp)

- average transverse momentum depends on centrality as well

- more central ⇒ more transverse momentum



Markert et al. [STAR], J.Phys.G35:044029,2008



Motivation

- Before including those density-driven effects into theoretical models one should check:
  - the maximum density which is reached in heavy ion collisions
  - the behaviour of the system without any medium effects
  - from what stage is the information one can gather experimentally from? and how?