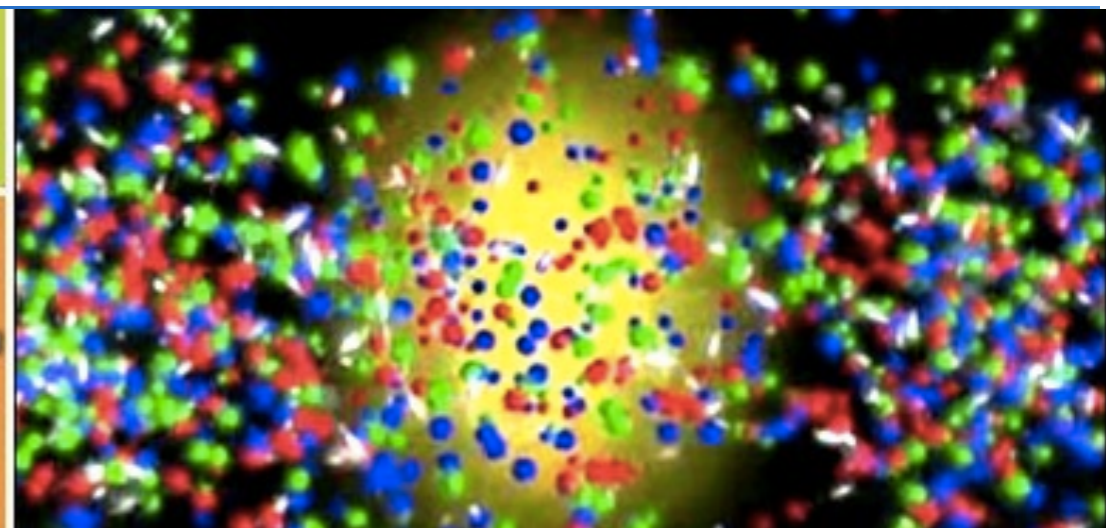


Dubna International Advanced
School of Theoretical Physics

Helmholtz International Summer School

"Physics of Heavy Quarks and Hadrons"

July 15-28, 2013, JINR, Dubna, Russia







The lessons of LHC (Higgs boson and SUSY)

Dmitry Kazakov

BLTP JINR

The Higgs Boson

Questions:

-  Is it the Higgs boson?
-  Is it the SM Higgs boson?
-  Are there deviations from the SM?
-  What are the alternatives?

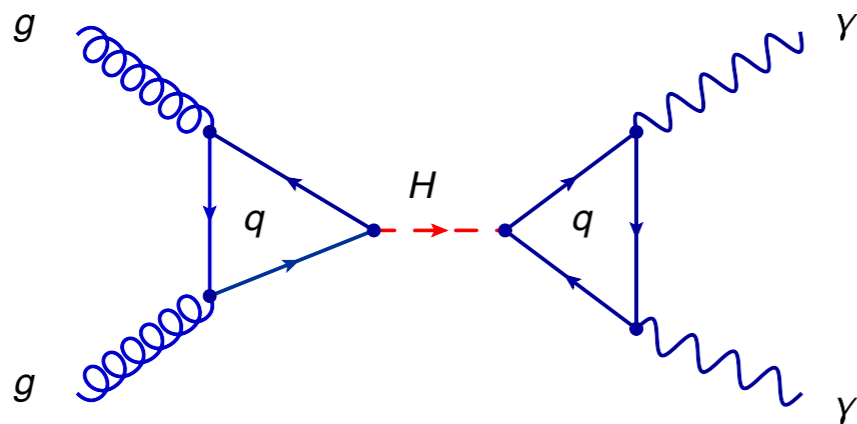
Warning:

Statistics is still low for an ultimate judgement, be patient

Is it the Higgs boson?

Is it 0^+ state?

Production and decay into $\gamma\gamma$



L.Landau'48,C.-N.Yang'50

Spin = ~~0~~, ~~1/2~~, ~~1~~, ~~3/2~~, ~~2~~

Decay into WW

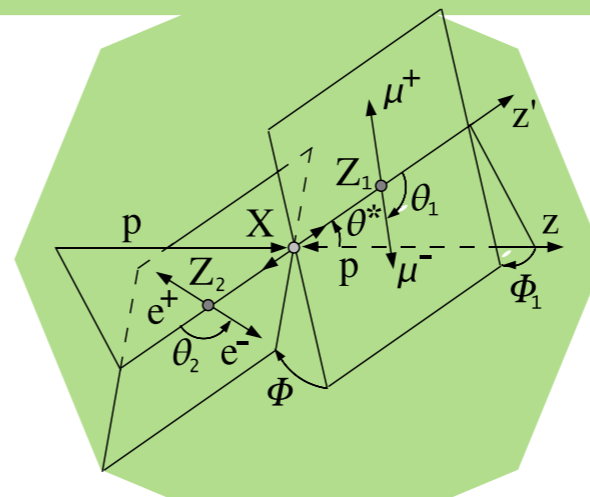
$$\begin{aligned} \mathcal{L} &= \frac{g^2}{4} (v + h(x))^2 W_\mu^+ W^{-\mu} \\ &= m_W^2 W_\mu^+ W^{-\mu} + \frac{2m_W^2}{v} h W_\mu^+ W^{-\mu} + \dots \end{aligned}$$

Decay into ZZ

$$\mathcal{L} = A \frac{\alpha}{4\pi} \frac{1}{M} h F_{\mu\nu} F^{\mu\nu} + B \frac{\alpha}{4\pi} \frac{1}{M} h \epsilon_{\mu\nu\lambda\sigma} F^{\mu\nu} F^{\lambda\sigma}$$

Angular distribution

$$pp \rightarrow ZZ^* \rightarrow l^+ l^- l^+ l^-$$



66% CL scalar

S.Baffoni et al, CMS'12

Y.Gao et al'10, A.De Rujula et al'10

Is it the Higgs boson?

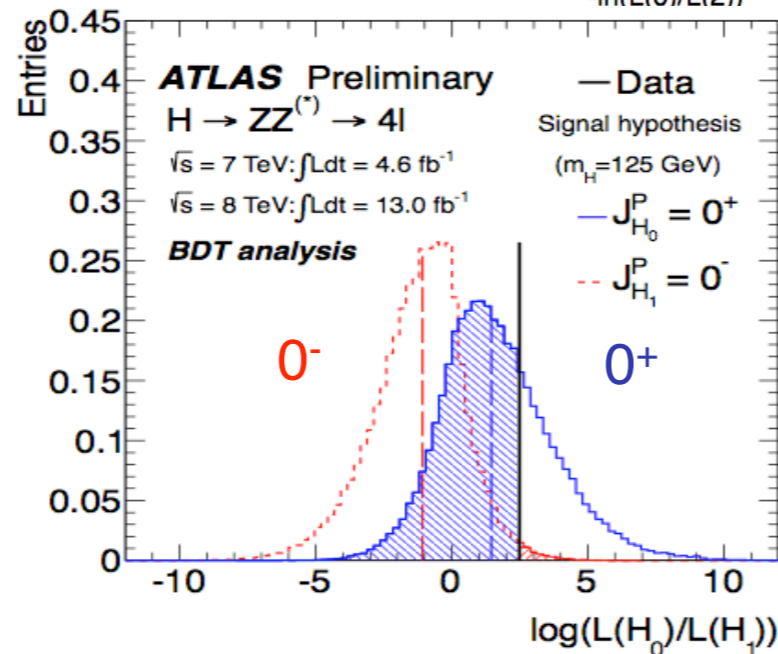
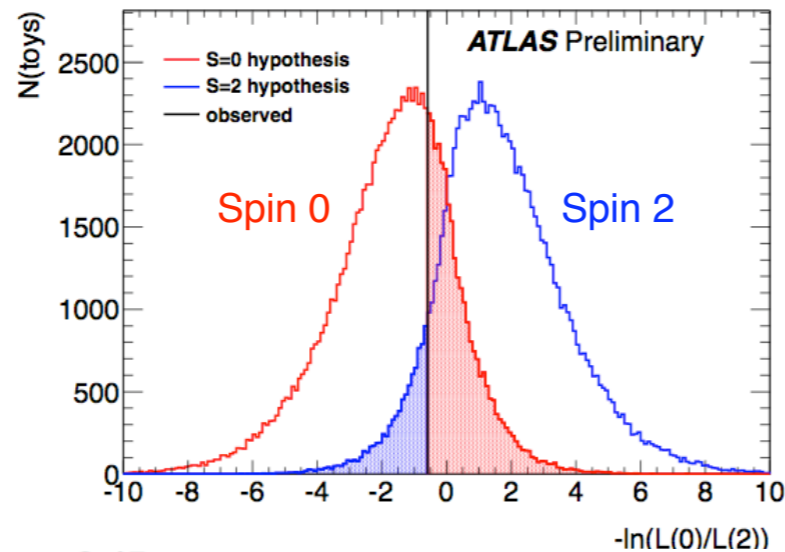
Is it 0^+ state?

Preliminary results ATLAS/CMS

$$pp \rightarrow ZZ^* \rightarrow l^+l^-l^+l^-$$

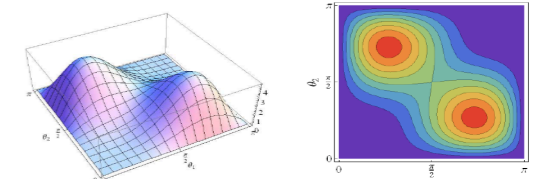
1) $gg \rightarrow X \rightarrow \gamma\gamma$ $qq \rightarrow X \rightarrow \gamma\gamma$
 spin-0: flat in $\cos\theta^*$
 spin-2: quartic in $\cos\theta^*$

$$\frac{d\sigma}{d\Omega} \propto \frac{1}{4} + \frac{3}{2}\cos^2\theta + \frac{1}{4}\cos^4\theta$$



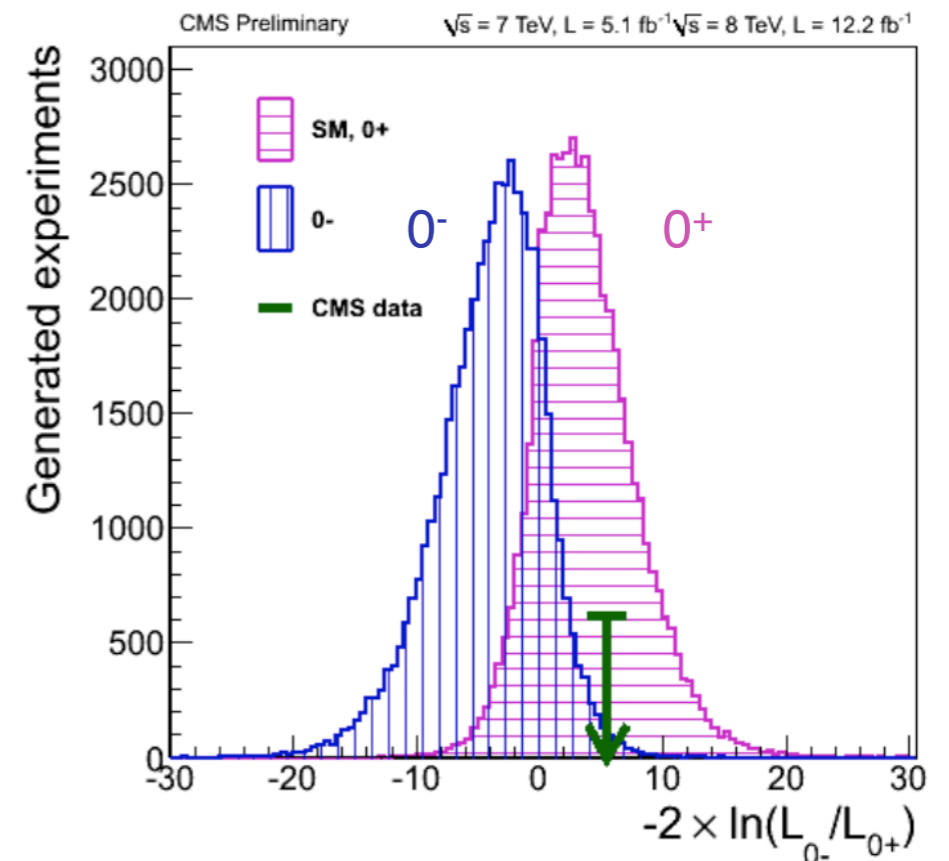
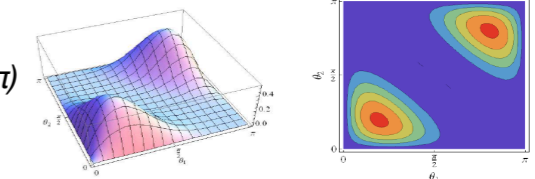
G. Tonelli, CERN/INFN/UNIFI

Polar angle distribution for $X_2 \rightarrow W^+W^-$



Polar angle distribution for $X_0 \rightarrow W^+W^-$

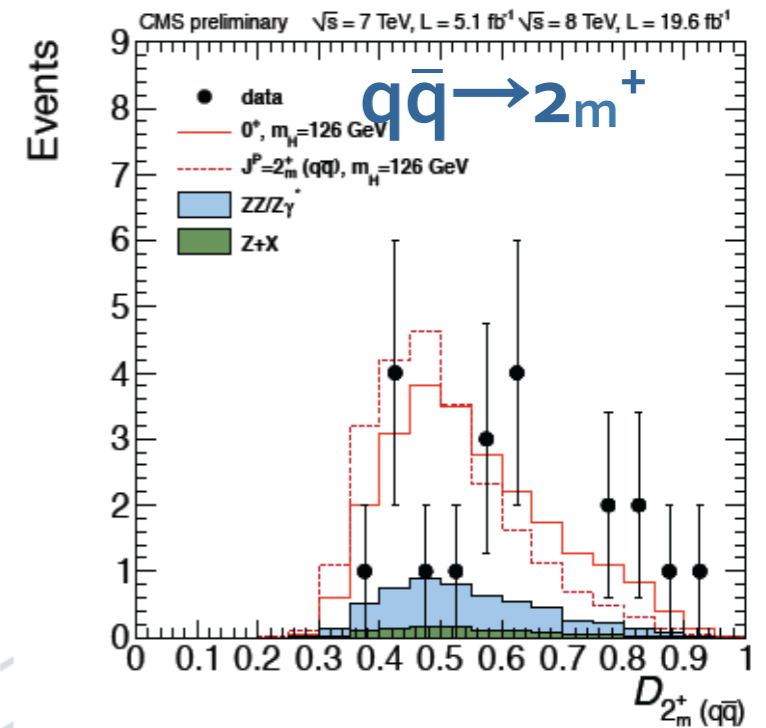
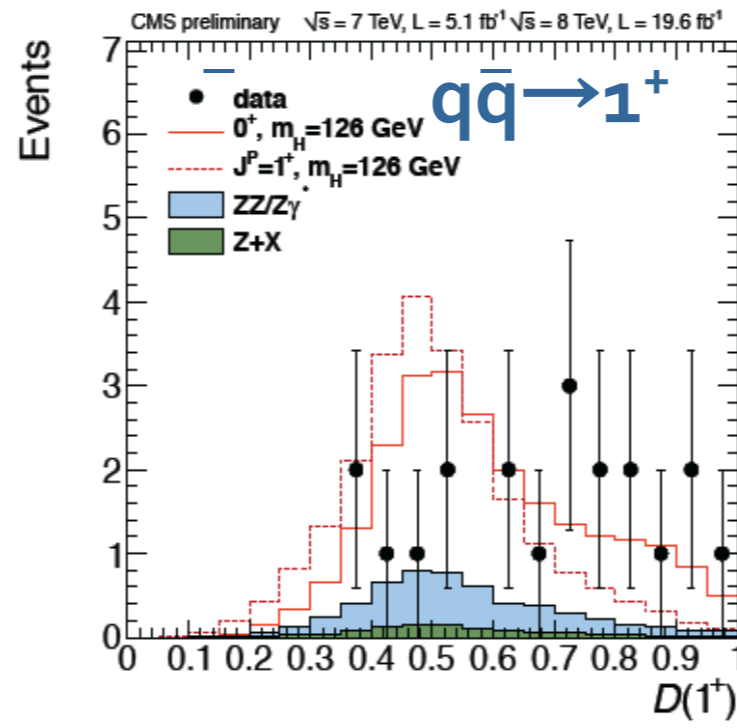
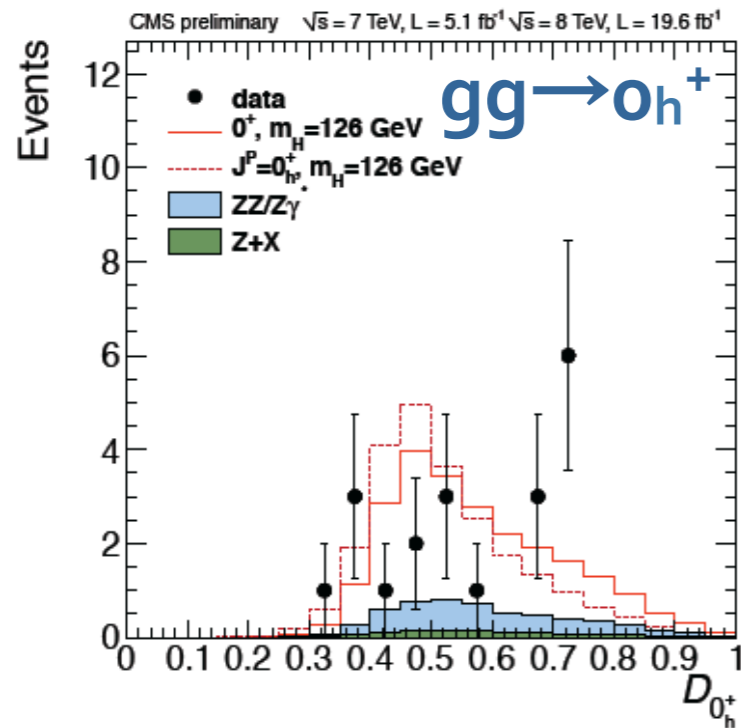
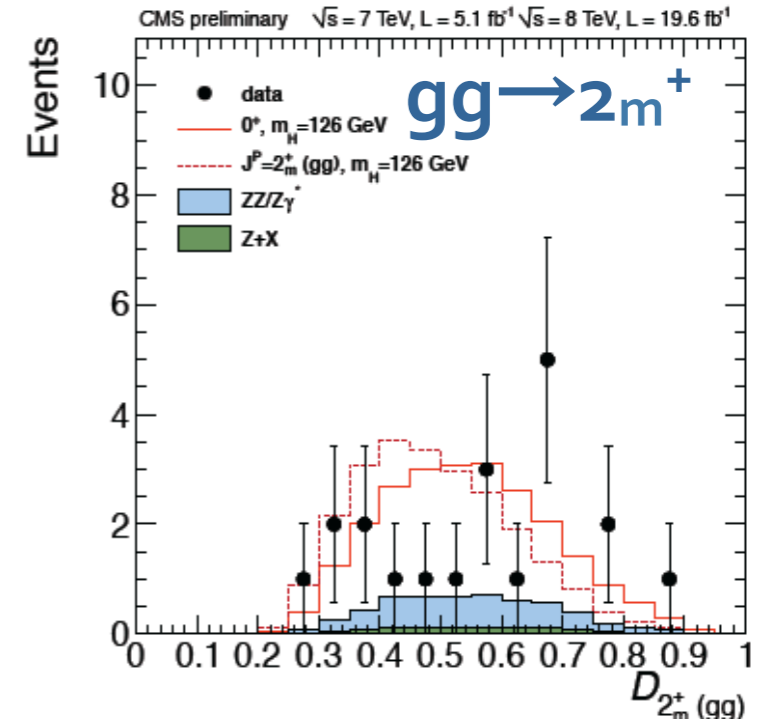
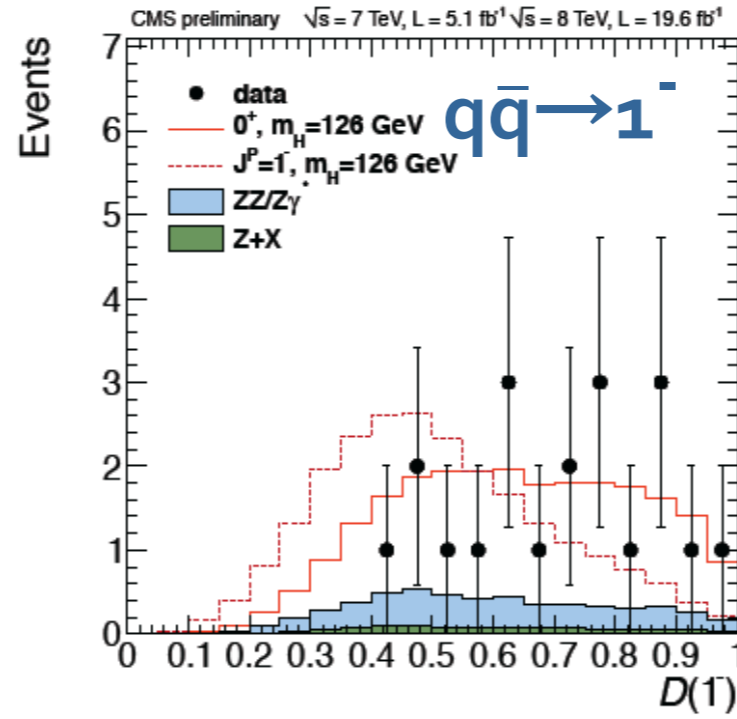
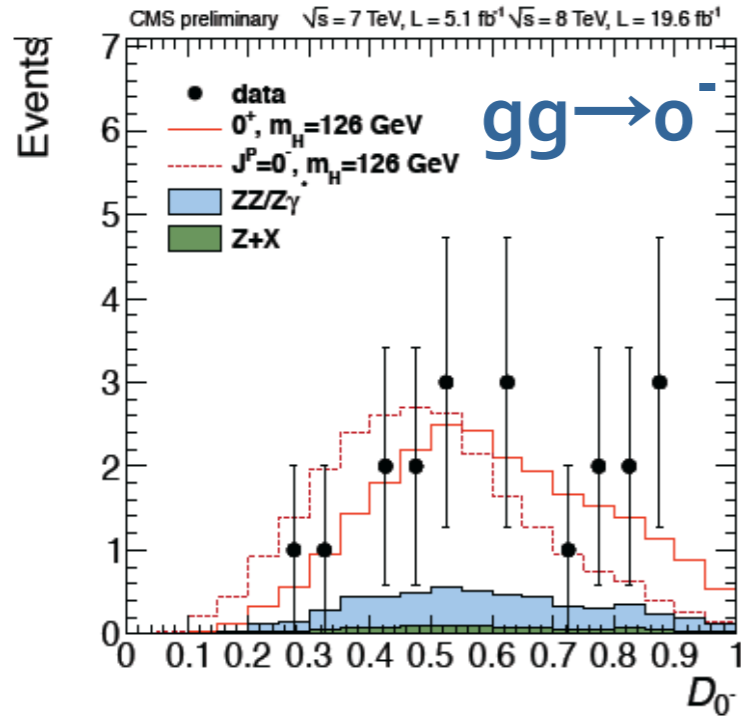
(for $\phi = \pi$)



Preliminary results on scalar/pseudo-scalar; observation consistent with 0^+ (0^-) within **0.5 (2.45) σ** .

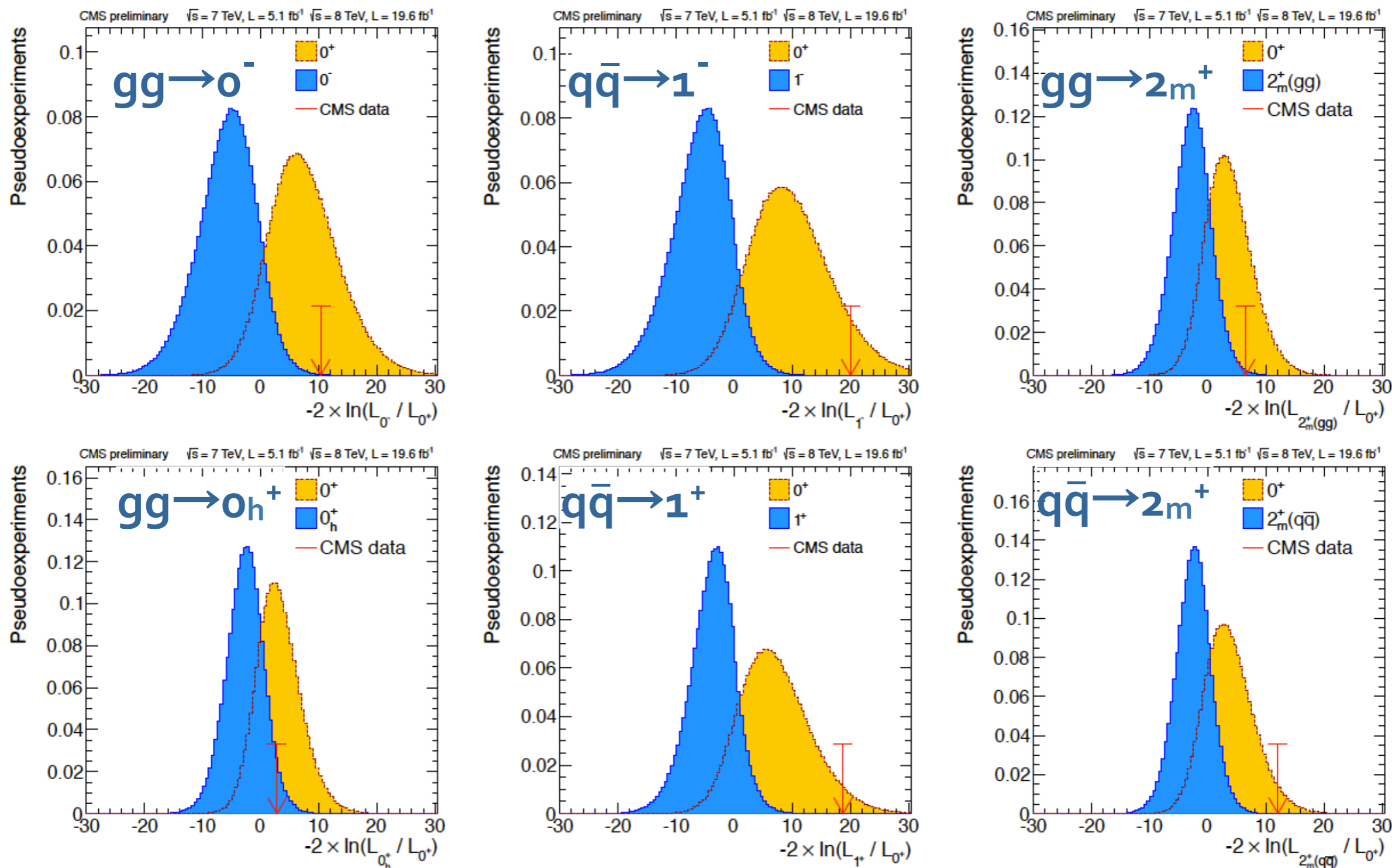


D_{JP} distributions (with $D_{bkg} > 0.5$)



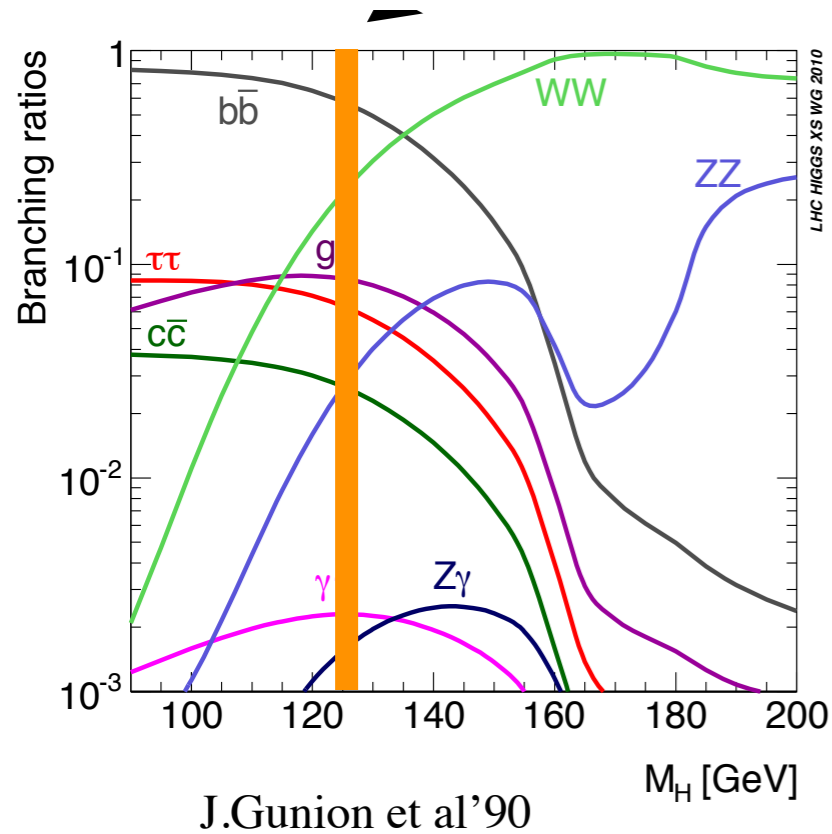


Spin-parity: test statistics



Decay modes

for 126 GeV Higgs



$b\bar{b}$	56%	$\tau^+\tau^-$	6.2%	$\gamma\gamma$	0.23%
WW^*	23%	ZZ^*	2.9%	γZ	0.16%
gg	8.5%	$c\bar{c}$	2.8%	$\mu^+\mu^-$	0.02%

1. $\gamma\gamma$ decay mode: Observed (4.5 σ in ATLAS, 4.1 σ in CMS) .
2. ZZ^* decay mode: Observed (3.6 σ in ATLAS, 3.2 σ in CMS) .
3. WW^* decay mode: Observed (2.8 σ in ATLAS, 1.6 σ in CMS) .
4. $b\bar{b}$ decay mode: Observed (2.8 σ in CDF/D0 combination)
5. $\tau\bar{\tau}$ decay mode: Not observed yet

All these decays are consistent with the Higgs boson

Is it quantitatively consistent with the SM?

Is the Higgs Standard Model-like or not?

$$\mu = \sigma \cdot BR / (\sigma \cdot BR)|_{SM}$$

Conclusion (so far)

1. too much $\gamma\gamma$, $\mu_{ATLAS} = 1.8 \pm 0.5$
 $\mu_{CMS} = 1.56 \pm 0.43$
2. too few $b\bar{b}$, $\mu \sim 0.7$
3. no $\tau\bar{\tau}$
4. WW, ZZ about all right

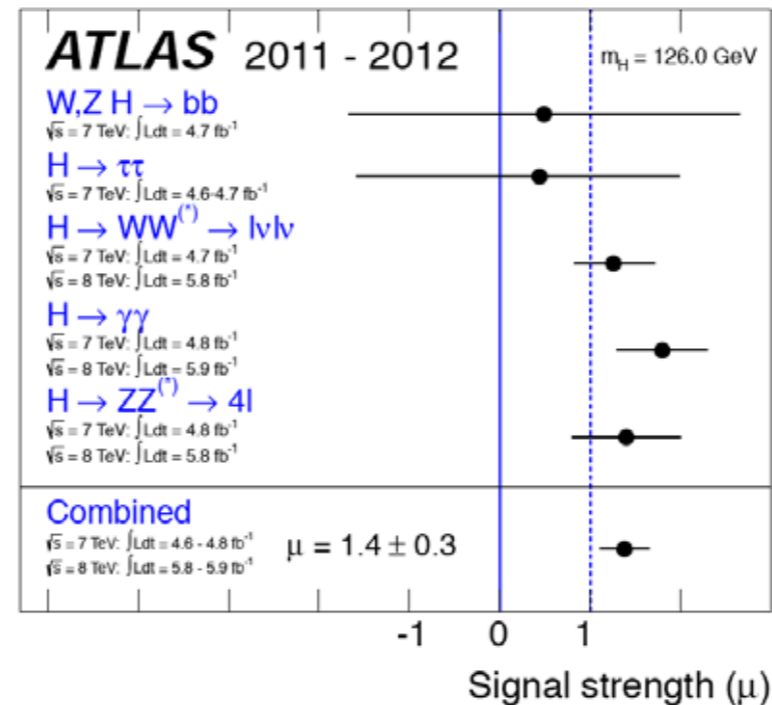
$$\begin{aligned} \mu(ZZ^{(*)} \rightarrow 4l)_{ATLAS} &= 1.4 \pm 0.6, \\ \mu(ZZ^{(*)} \rightarrow 4l)_{CMS} &= 0.7^{+0.4}_{-0.3}, \\ \mu(WW^{(*)} \rightarrow l\nu l\nu)_{ATLAS} &= 1.3 \pm 0.5, \\ \mu(WW^{(*)} \rightarrow l\nu l\nu)_{CMS} &= 0.6 \pm 0.4. \end{aligned}$$

Theoretical suggestions

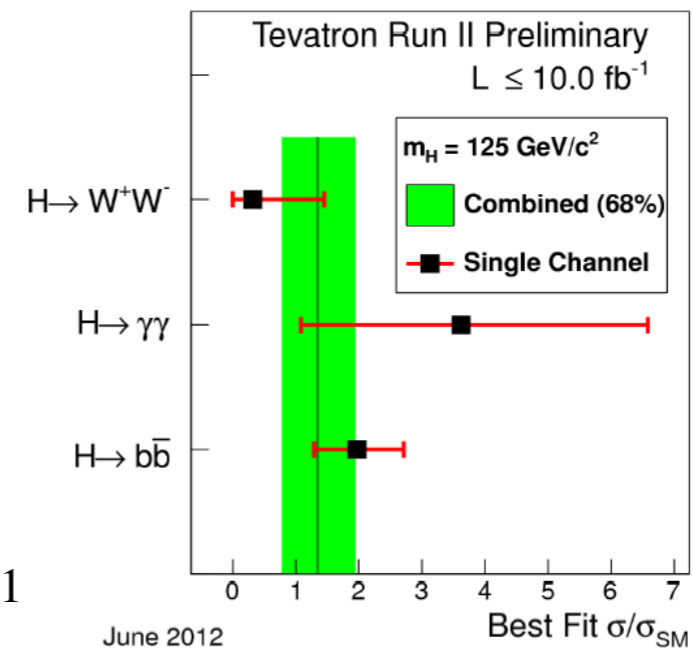
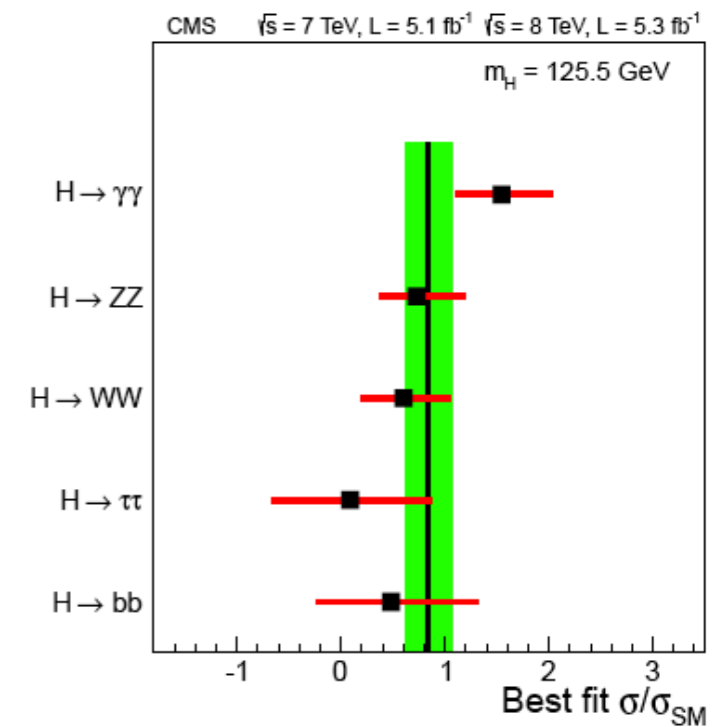
1. Extended Higgs sector
2. Colour singlet matter particles (stau)
3. New coloured matter particles (stop)
4. Strongly interacting Higgs sector

U.Ellwanger'11,
J.Gunion et al'12
M.Carena et al'11
Z. Kang et al'11
J.Espinosa et al'12
F.Goertz et al'11

G. Aad et al. [ATLAS]'12



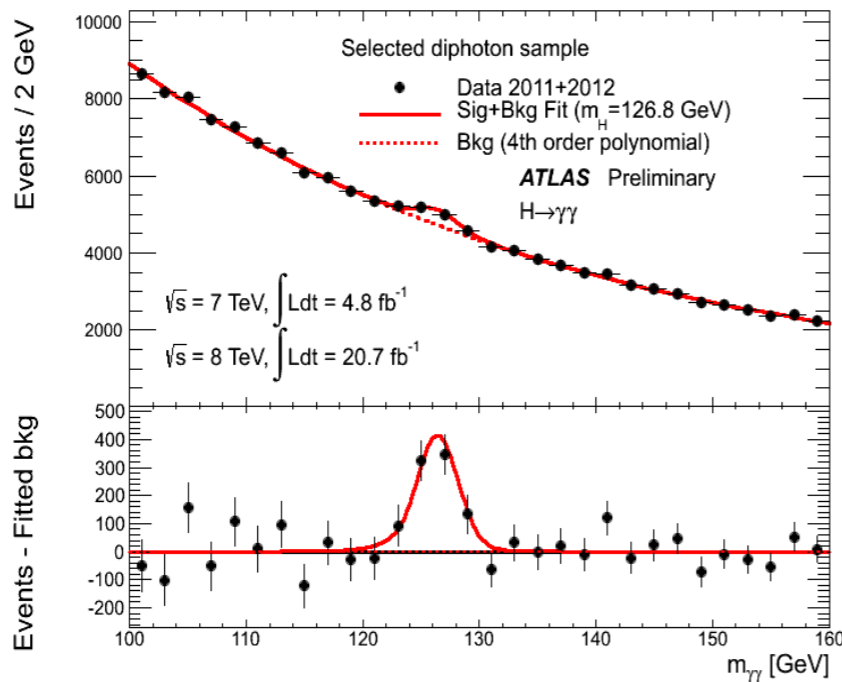
S. Chatrchyan et al. [CMS]'12



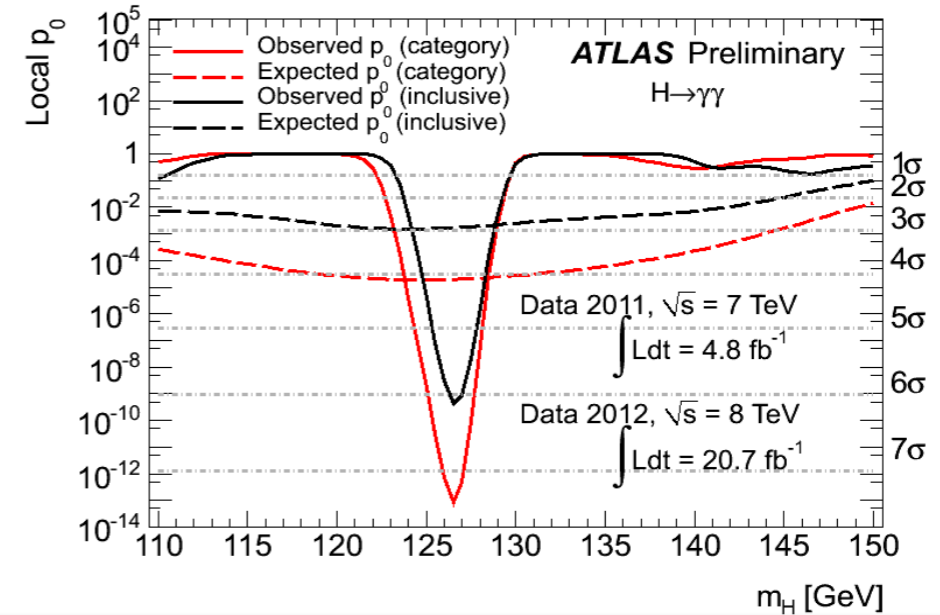
J. Haley et al, CDF/D0'12

Higgs $\rightarrow \gamma\gamma$

Update from ATLAS

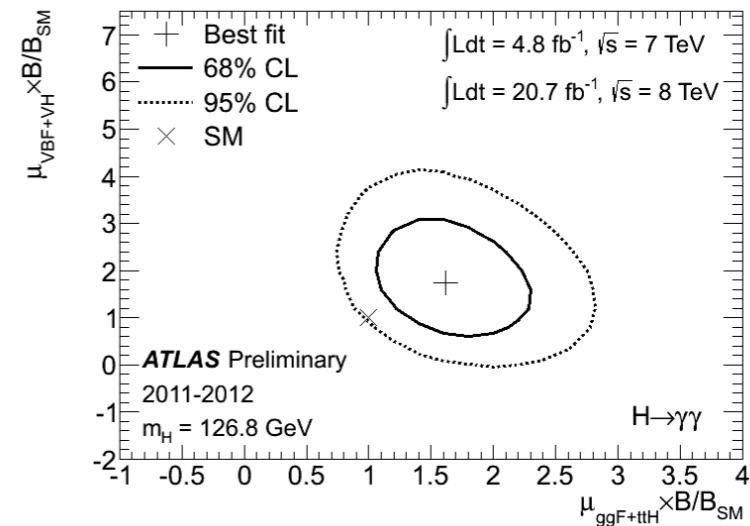


**Mass window ~125 GeV
with 90% signal: S/B~3%**



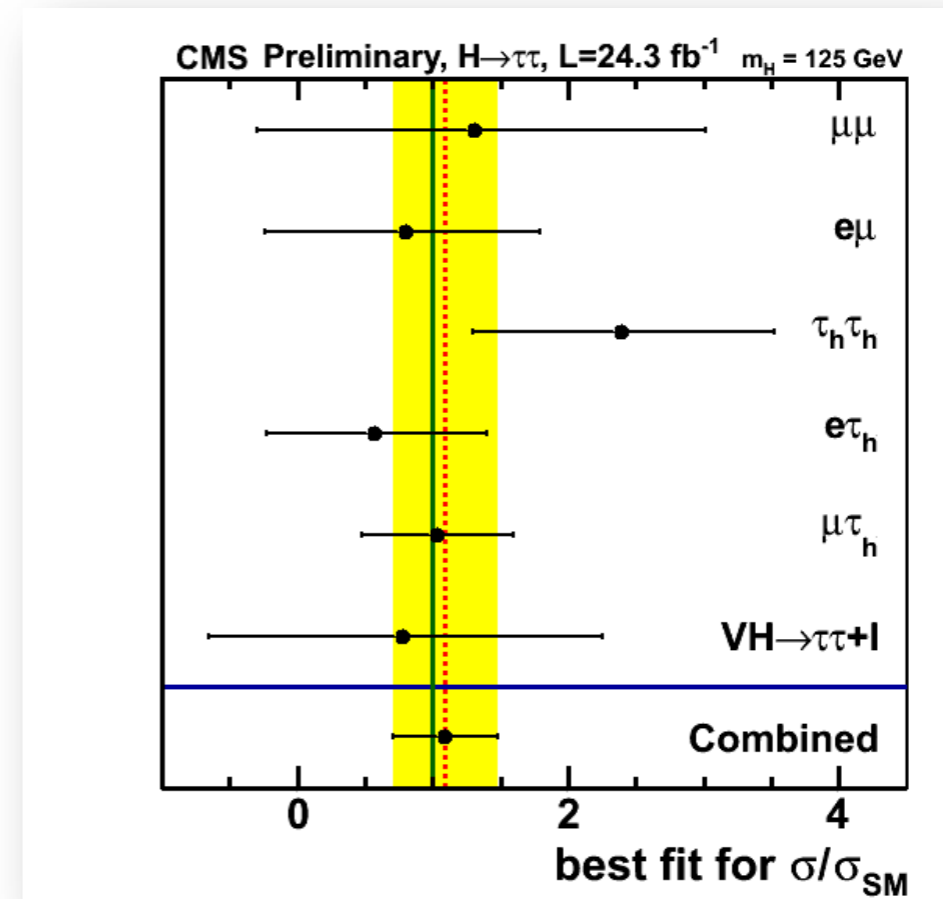
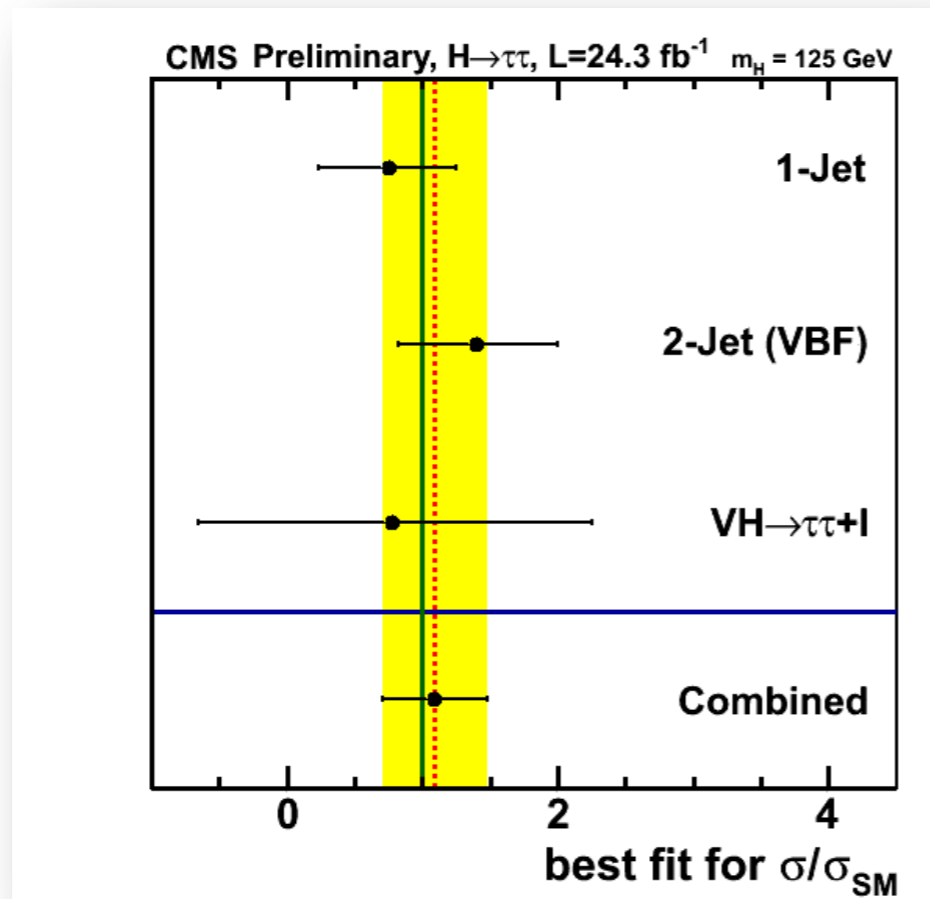
Significance; obs: 7.4σ; exp: 4.1σ

Mass: 126.8 ± 0.2(stat) ± 0.7(sys) GeV



Higgs- \rightarrow $\tau\tau$

- Yields by “type” and by decay channel



- Consistent picture across channels and categories
- Combined best-fit $\hat{\mu}$ of 1.1 ± 0.4

Signal Strength

- Parameter of interest : μ (global)

$\Rightarrow \mu = 1.43 \pm 0.16$ (stat) ± 0.14 (sys)

Council Dec 2012 $\mu = 1.35 \pm 0.19$ (stat) ± 0.15 (sys)

- Consistency tests

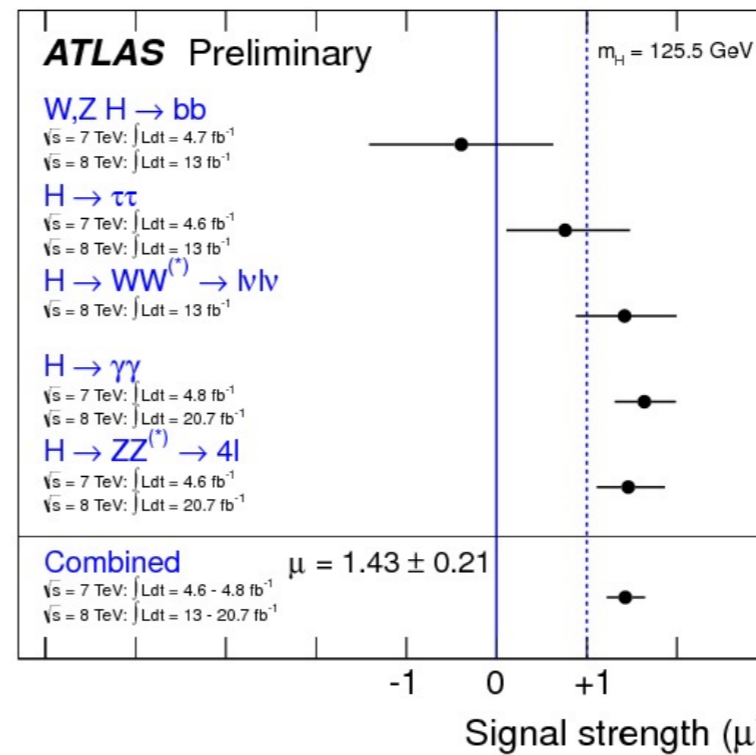
- global μ with SM: 3%
 - 11% with rectangular QCD scale and parton dist functions

- 5 μ_i with SM: 8%
- 5 μ_i with 1.43: 32%

- μ, m_H contours

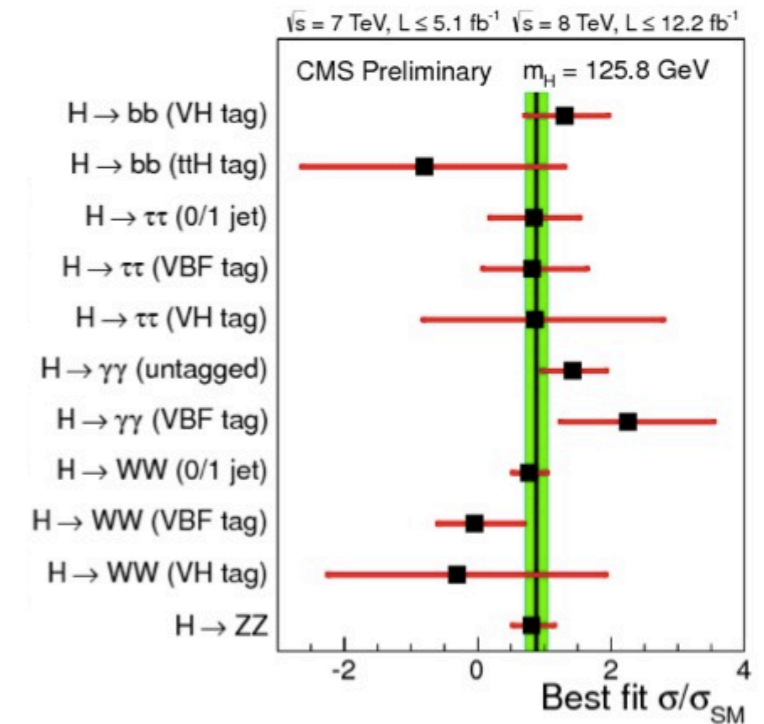
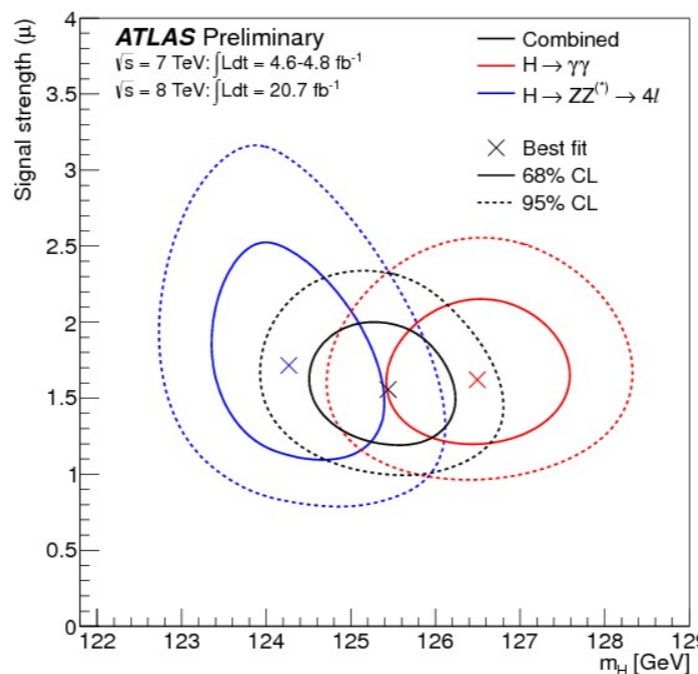
- $\gamma\gamma$
- 4l
- combined

Bruno Mansoulié (IRFU-CEA), Moriond-EW, March 2013



at $m_H = 125.5$

[124.5-126.5]:
 $\mu \pm 4\%$



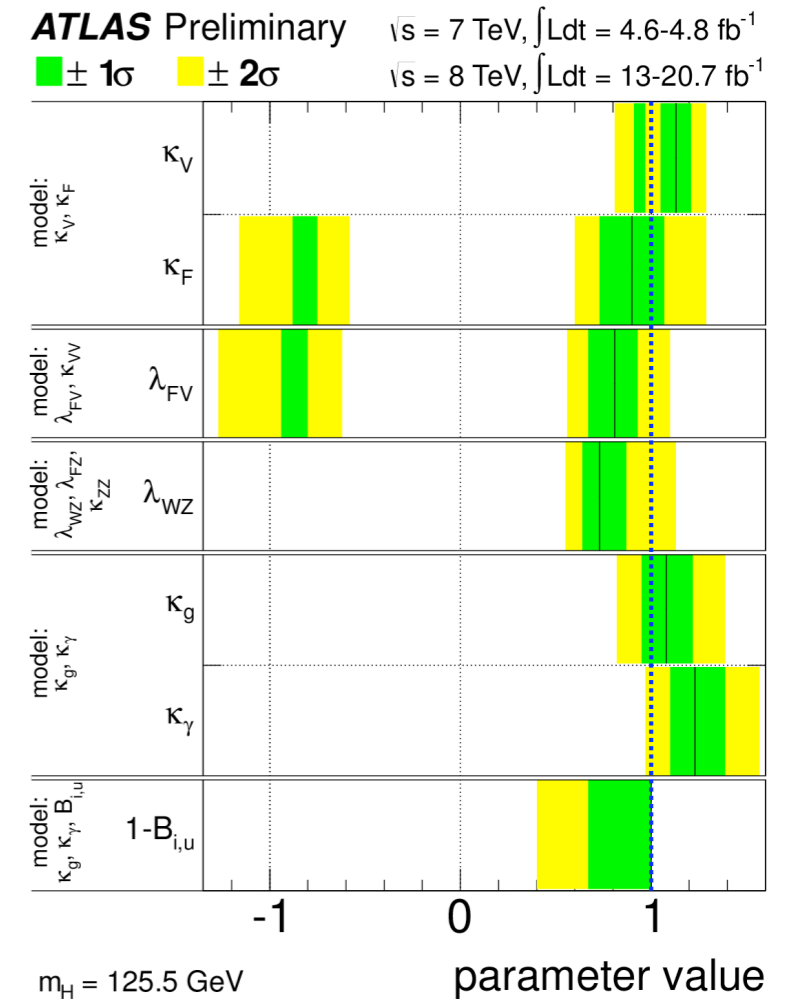
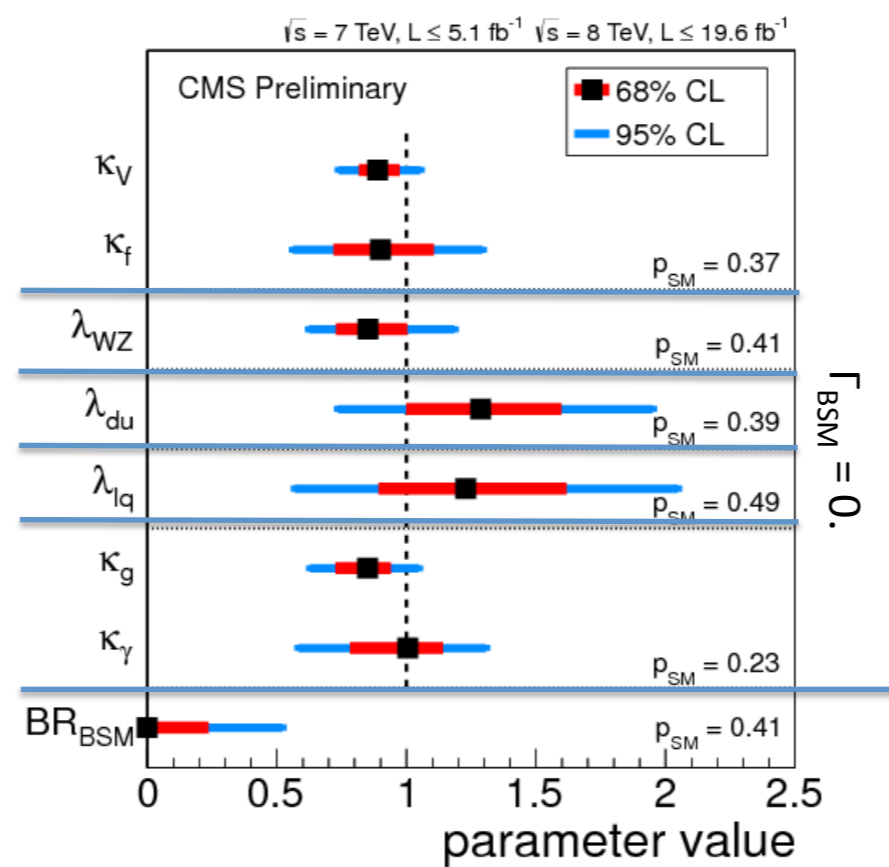
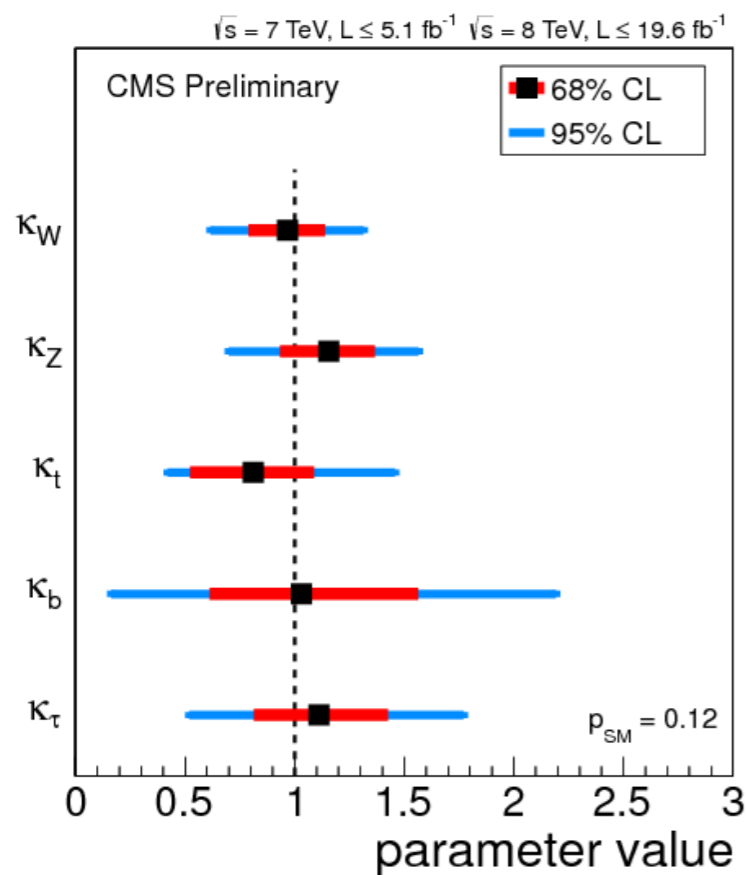
$\sigma/\sigma_{SM} = 0.88 \pm 0.21$

The couplings?

- LHC XS WG benchmark models ([arxiv:1209.0040](https://arxiv.org/abs/1209.0040)):
 - Fermionic vs bosonic couplings: $\kappa_V \kappa_f$
 - Search for asymmetries: λ_{WZ} , λ_{du} , λ_{lq}
 - Search for new physics in loops: $\kappa_g \kappa_\gamma$ BR_{BSM}

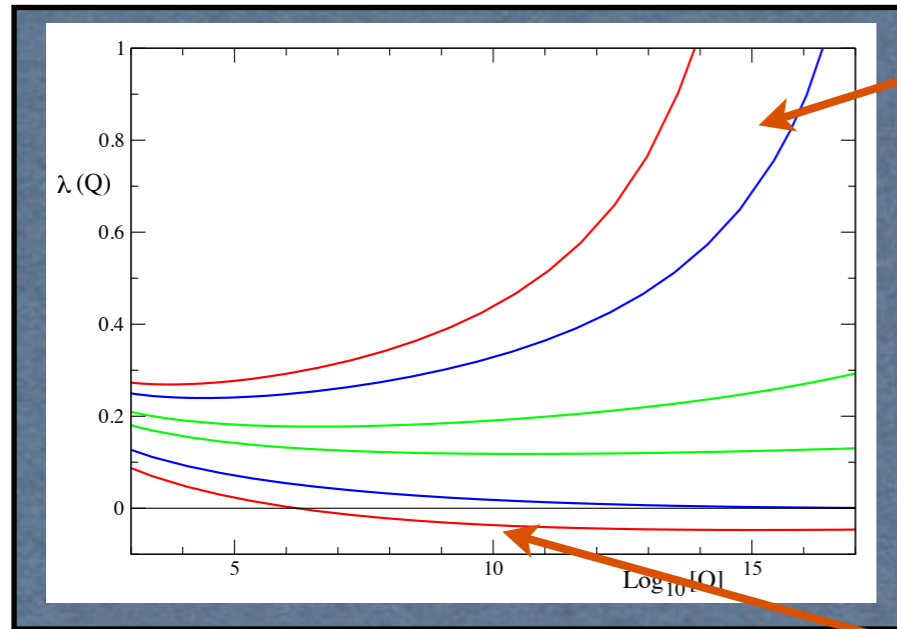


Summary of all searches for coupling deviations



Theoretical Bounds on the Higgs mass in SM

Running of the Higgs coupling



Landau pole

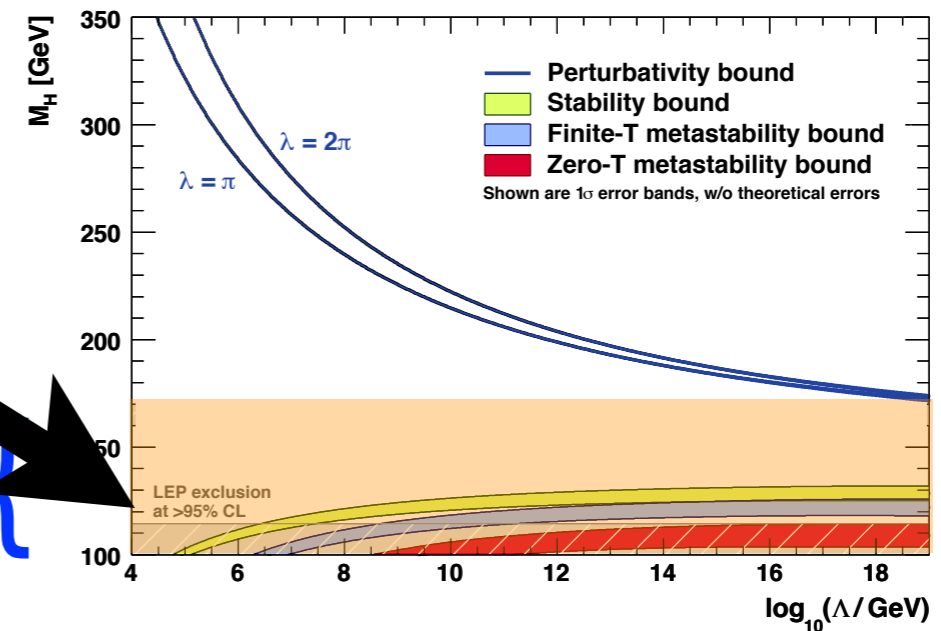
Demanding λ not too large
(keep perturbativity)
not too negative that
destabilizes the Higgs potential

Instability

A 125 GeV Higgs is
in this window!

Only a small window
in the Higgs mass
makes the SM consistent
all the way to the Planck scale

A.Pomarol et al'11

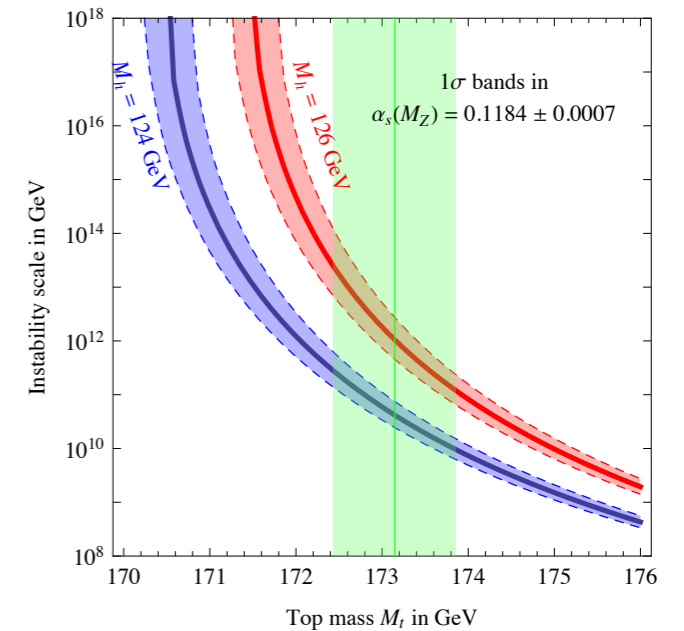
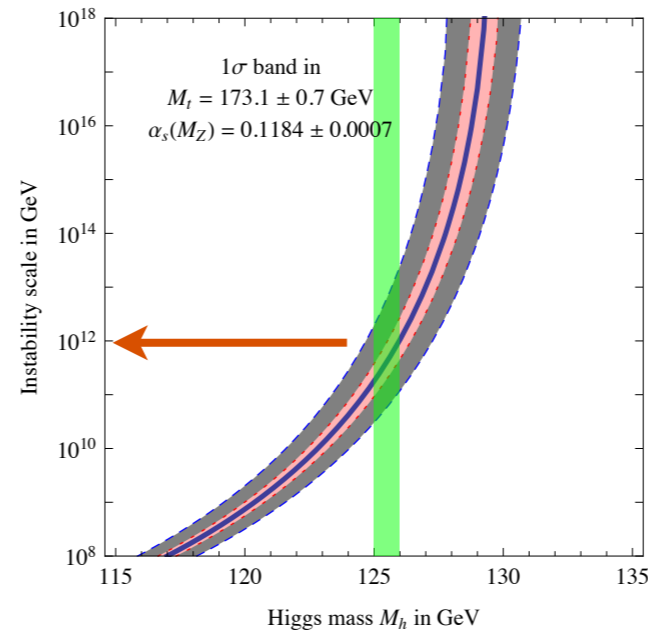
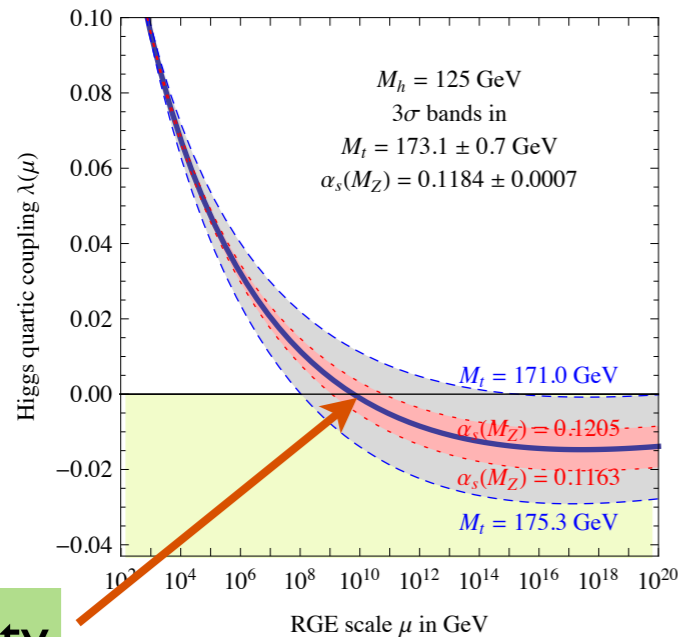


Stability of the Electroweak Vacuum

The running of the Higgs coupling $\lambda(\mu)$

G.Degrassi et al'12 arXiv:1205.6497

M.Zoller'12



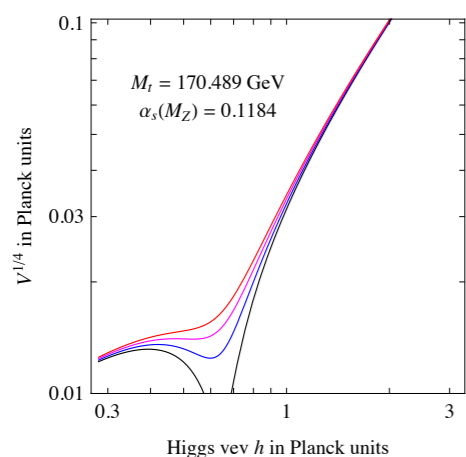
instability

$$M_h \text{ [GeV]} > 129.4 + 1.4 \left(\frac{M_t \text{ [GeV]} - 173.1}{0.7} \right) - 0.5 \left(\frac{\alpha_s(M_Z) - 0.1184}{0.0007} \right) \pm 1.0_{\text{th}} .$$

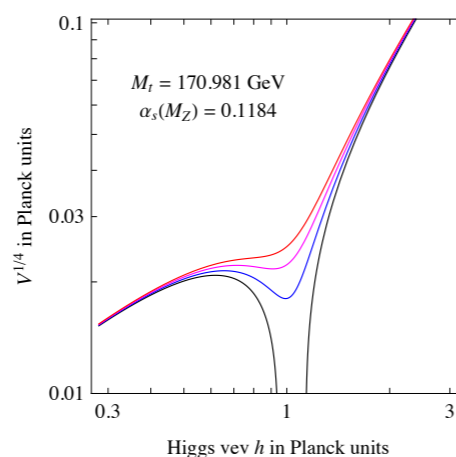
$$M_h > 129.4 \pm 1.8 \text{ GeV} .$$

Effective potential

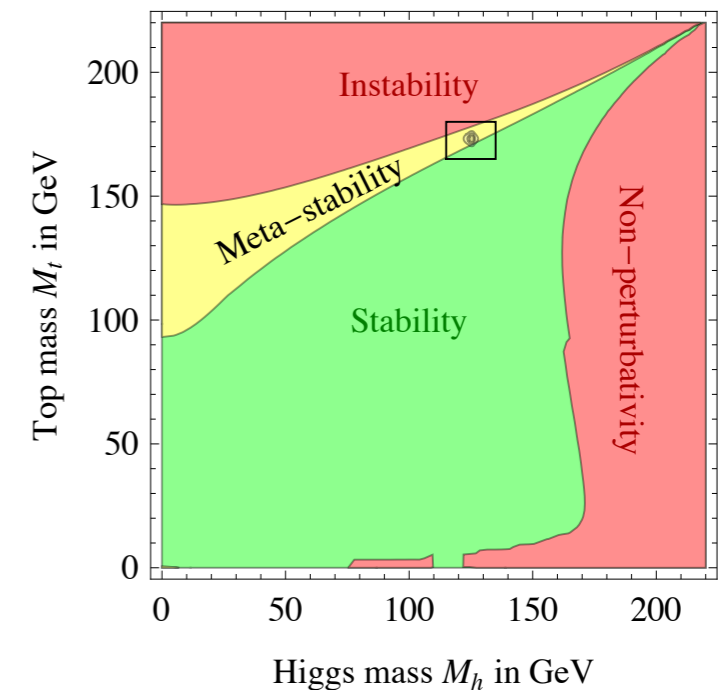
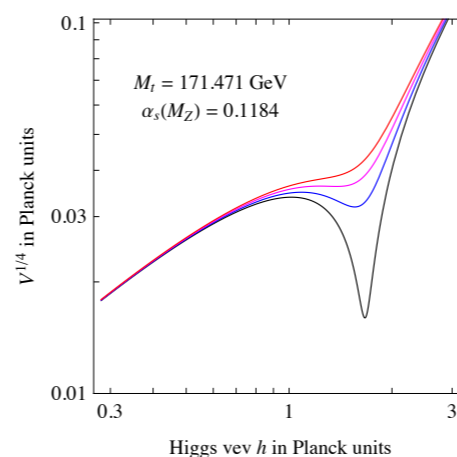
SM Higgs potential, $M_h = 124 \text{ GeV}$



SM Higgs potential, $M_h = 125 \text{ GeV}$



SM Higgs potential, $M_h = 126 \text{ GeV}$



Higgs v SUSY

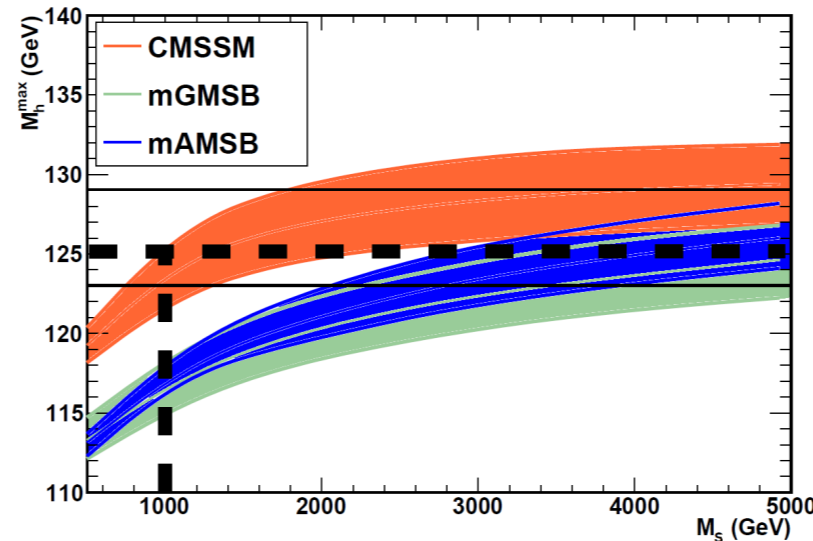
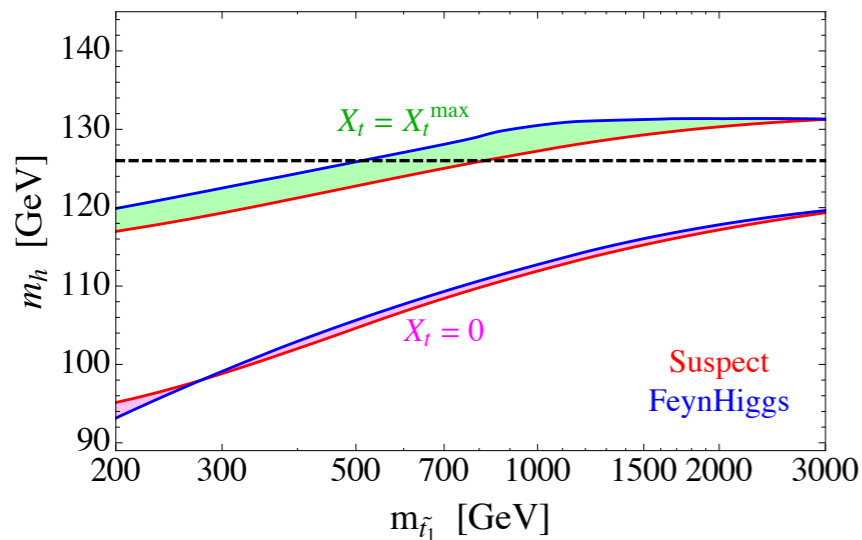
$$m_{Higgs}^2 = M_Z^2 \cos^2 2\beta + \frac{3m_t^4}{4\pi^2 v^2 \sin^2 \beta} \left[\log \frac{M_s^4}{m_t^4} + \frac{X_t^2}{M_S^2} \left(1 - \frac{X_t^2}{6M_S^2} \right) \right] + 2-loop$$

$$M_S^2 = \tilde{m}_{t_1} \tilde{m}_{t_2} \quad X_t = A_t - \mu \cot \beta$$

from JHEP 1204 (2012) 131

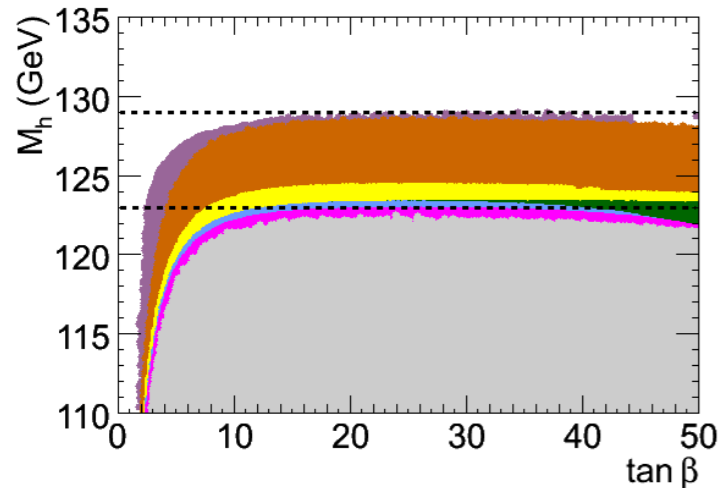
from arXiv:1207.1348

MSSM Higgs Mass

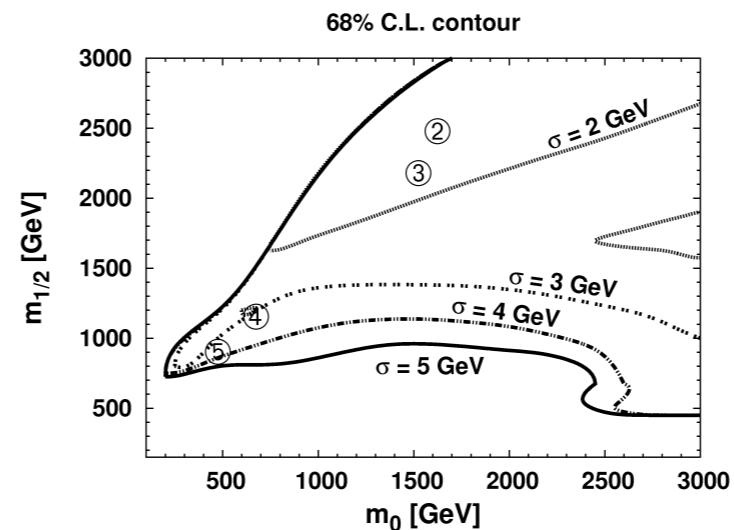


Resume

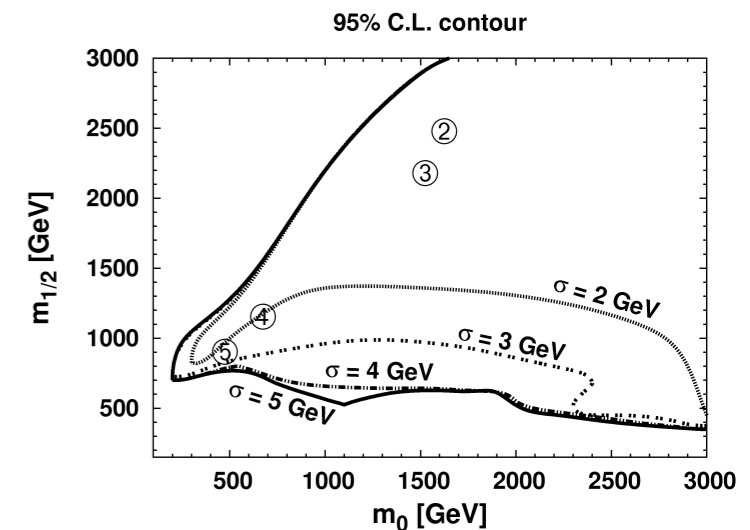
1. MSSM has already troubles to accomodate 126 GeV Higgs
2. Needs $M_S \sim 1\text{TeV}$
3. Large part of the parameter space is closed



A. Arbey et al'12



W. de Boer et al'12



Two Higgses or One Higgs?

Higgs Bosons at 98 and 125 GeV at LEP and LHC

G. Belanger et al'12

M.Drees'12

LEP

$$e^+e^- \rightarrow Zh, \quad h \rightarrow b\bar{b} \quad m_h \approx 98 \text{ GeV}$$

Possible explanation
within NMSSM h, H

LHC

$$pp \rightarrow h \rightarrow \gamma\gamma \quad m_h \approx 126 \text{ GeV}$$

☑ 2σ LEP excess is inconsistent with the SM being only about 10 – 20% of the rate for the SM Higgs but might be consistent within SUSY model

Scenario	m_{h_1}	m_{h_2}	m_{h_3}	m_{a_1}	m_{a_2}	m_{H^\pm}	$m_{\tilde{\chi}_1^0}$	Ωh^2	LSP singlino	LSP Higgsino	$R_{gg}^{h_2}(\gamma\gamma)$
I	99	124	311	140	302	295	76	0.099	18%	75%	1.62
II	97	124	481	217	473	466	92	0.026	20%	74 %	1.53
III	99	126	993	147	991	989	115	0.099	75%	25%	1.14

Two Higgs Bosons in the interval 123-128 GeV

J.Gunion et al'12

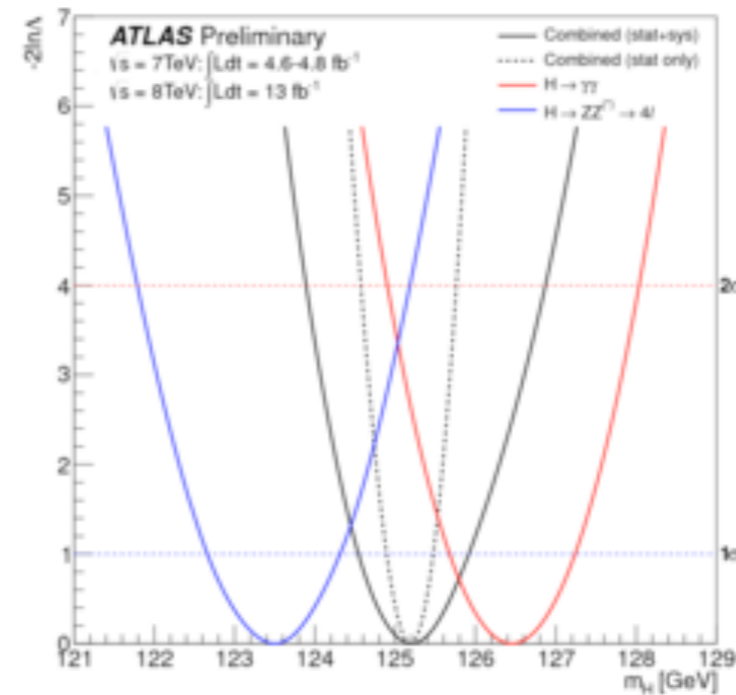
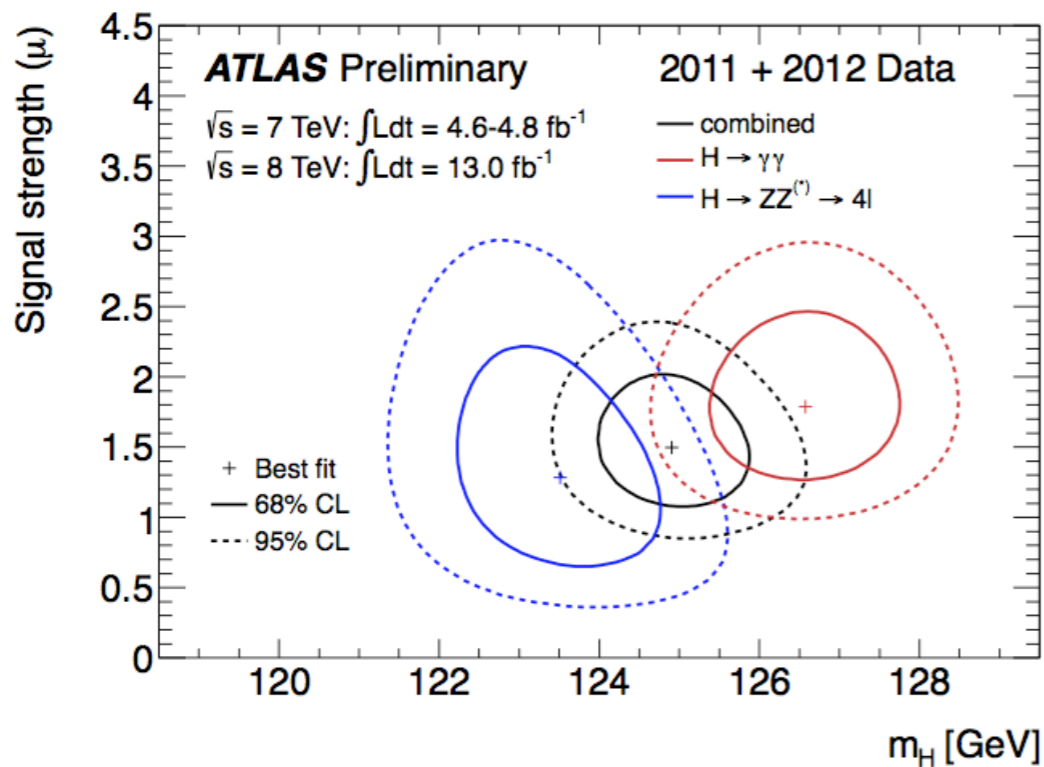
☑ Two CP even Higgses of the NMSSM are degenerate.

☑ Large rates (relative to $gg \rightarrow h_{SM} \rightarrow \gamma\gamma$ or $gg \rightarrow h_{SM} \rightarrow ZZ \rightarrow 4l$) for $gg \rightarrow h_{1,2} \rightarrow \gamma\gamma$ and $gg \rightarrow h_{1,2} \rightarrow ZZ \rightarrow 4l$ are possible when either one rate is large or the sum is large

$m_0 \in [0.9, 1.3] \text{ TeV}$, $m_{1/2} \in [500, 700] \text{ GeV}$, $A_0 \in [-1.8, -1.0] \text{ TeV}$, $A_\kappa \in [-400, -250] \text{ GeV}$, $A_\lambda \in [-600, -400] \text{ GeV}$, $m_S \text{ (GUT)} \in [1.4, 2.2] \text{ TeV}$, $m_{Hu} \text{ (GUT)} \in [2, 2.2] \text{ TeV}$ and $m_{Hd} \text{ (GUT)} \in [0.7, 1.2] \text{ TeV}$

Two Higgs Bosons at ATLAS?

Two Higgs Bosons in the interval 123-128 GeV



$$m_H = 123.5 \pm 0.9 \text{ (stat)} \begin{matrix} +0.4 \\ -0.2 \end{matrix} \text{ (syst)} \text{ GeV} \quad m_H = 126.6 \pm 0.3 \text{ (stat)} \pm 0.7 \text{ (syst)} \text{ GeV}$$

$$m_H = 125.2 \pm 0.3 \text{ (stat)} \pm 0.6 \text{ (syst)} \text{ GeV}$$

NO the measurements of the mass in the two channels are each other compatible (2.3-2.7 σ) and everything is compatible with what CMS observes.



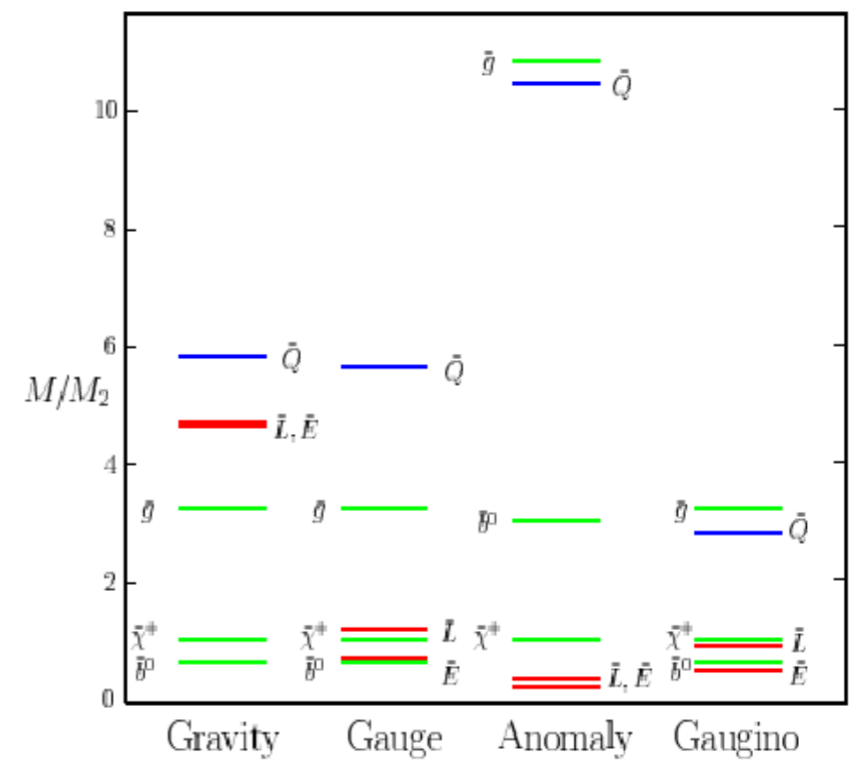
Supersymmetry



What
SUSY?

MSSM
 CMSSM
 mSUGRA
 mGMSB
 mAMSB
 NUHM
 NMSSM
 No Scale
 ...

Superparticle spectrum



SUSY at TeV scale:

- Unification of the gauge couplings
- Solution of the hierarchy problem
- Explains the electroweak symmetry breaking

SUSY:

- Provides unification with gravity
- Provides the Dark matter candidate

Search for SUSY Manifestation

Particle Phys

- Direct production at colliders at high energies
- Indirect manifestation at low energies
 - Rare decays ($B_s \rightarrow s\gamma$, $B_s \rightarrow \mu^+\mu^-$, $B_s \rightarrow \tau\nu$)
 - g-2 of the muon
- Search for long-lived SUSY particles

**Astro Phys
(if SUSY DM)**

- Relic abundance of Dark Matter in the Universe
- DM annihilation signal in cosmic rays
- Direct DM interaction with nucleons

Nothing so far ...

Exp and Theor Framework

Two ways to present and analyse data:

1. High energy input:

introduce universal parameters at high energy scale (GUT)

Example $m_0, m_{1/2}, A_0, \tan \beta$ of MSSM

Advantage: small number of universal parameters for all masses

Disadvantage: strictly model dependent (MSSM, NMSSM, etc)

2. Low energy input:

use low energy parameters like masses of superpartners

Example $\tilde{m}_g, \tilde{m}_q, \tilde{m}_\chi$ or $m_A, \tan \beta$

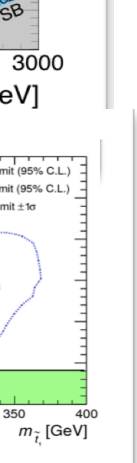
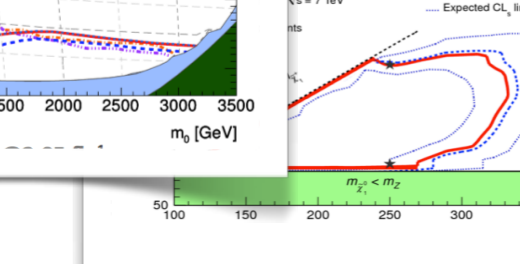
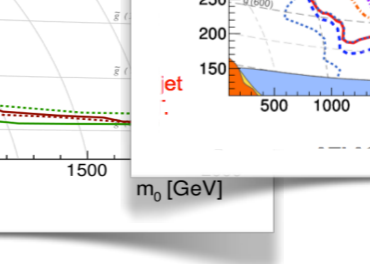
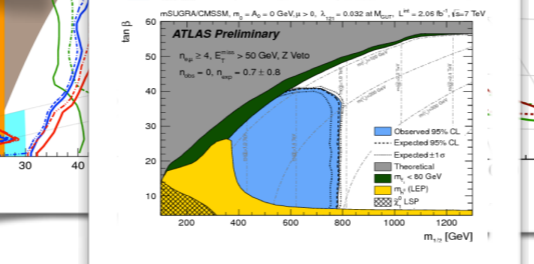
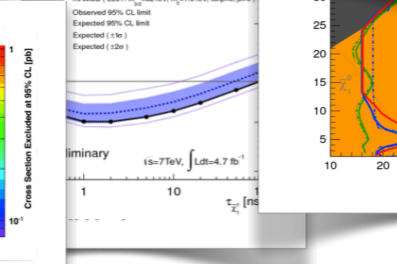
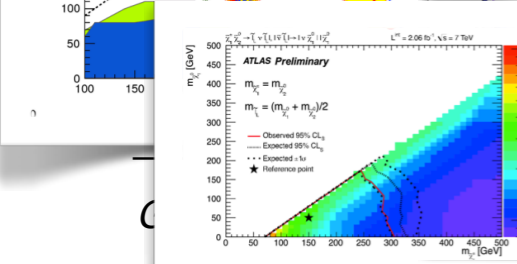
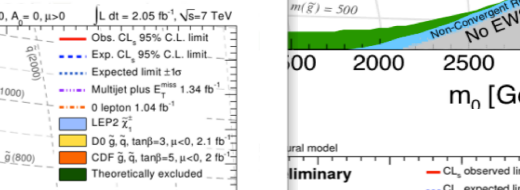
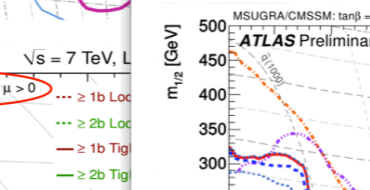
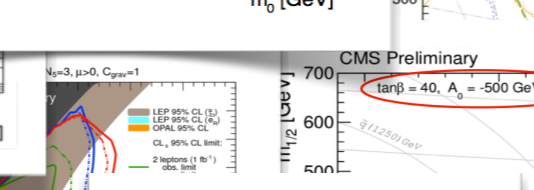
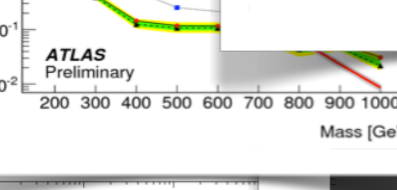
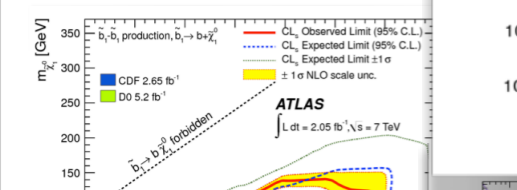
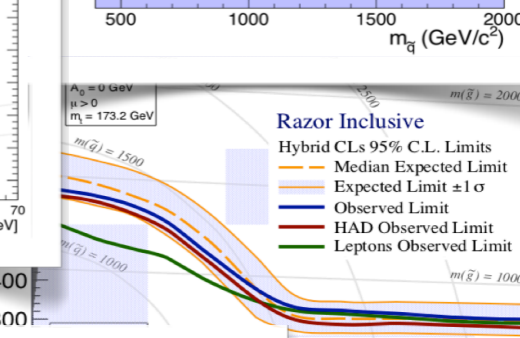
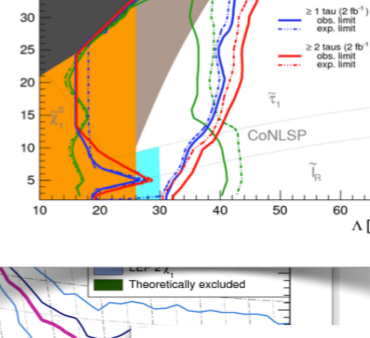
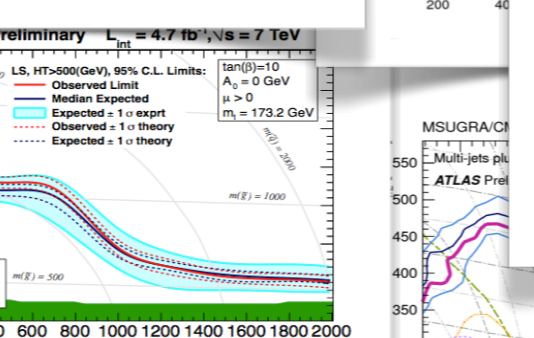
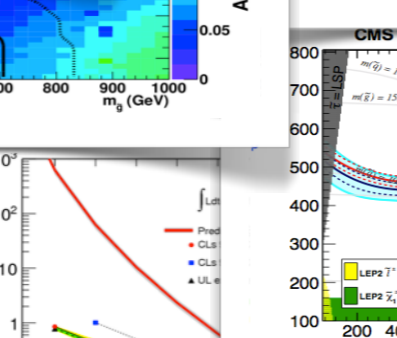
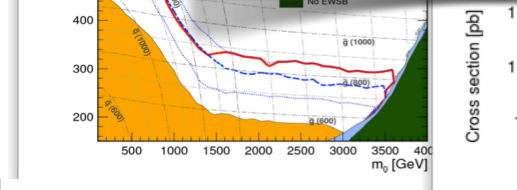
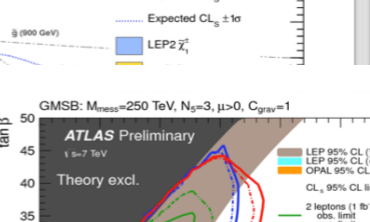
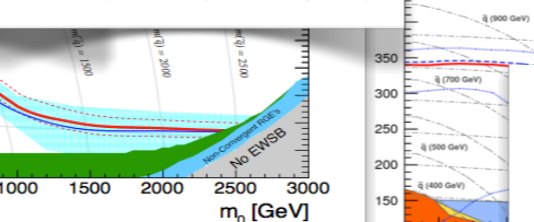
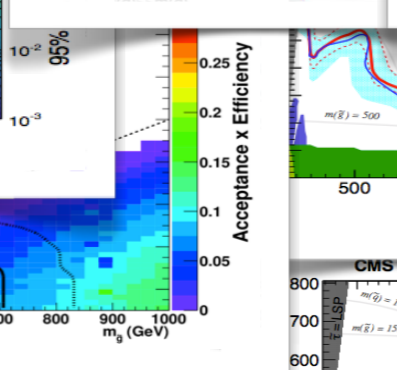
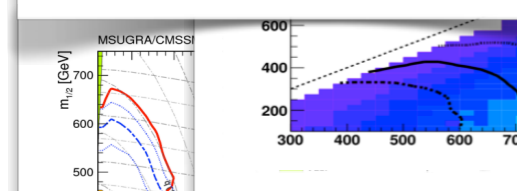
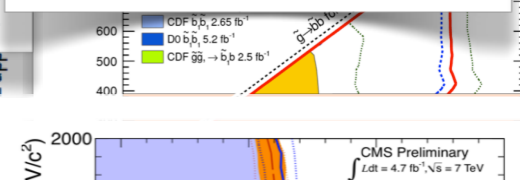
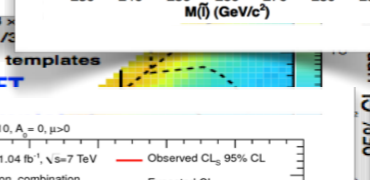
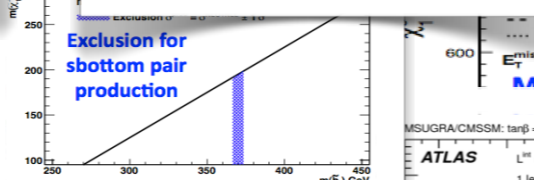
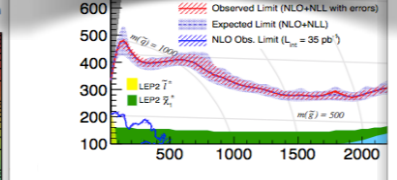
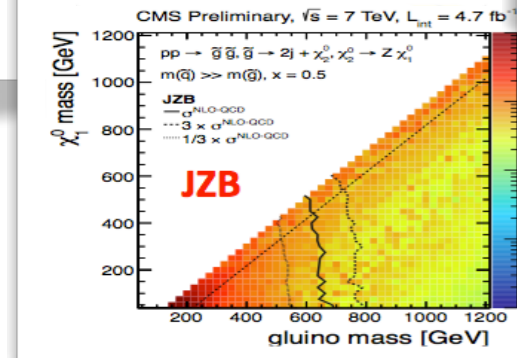
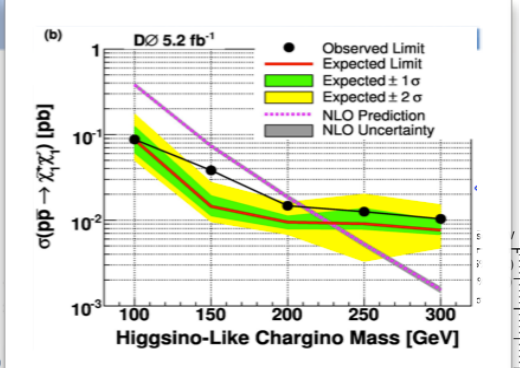
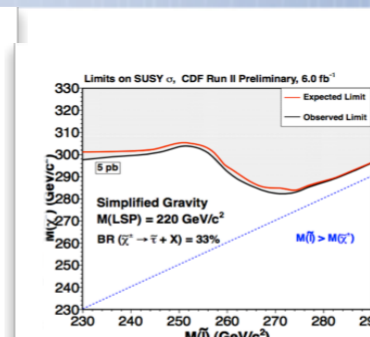
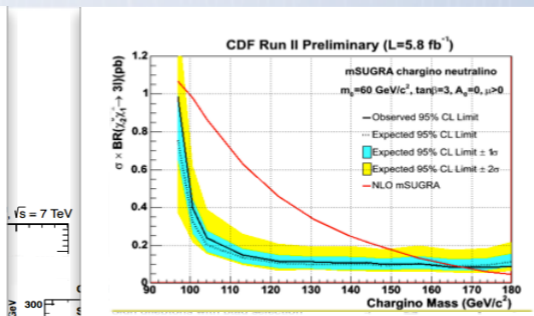
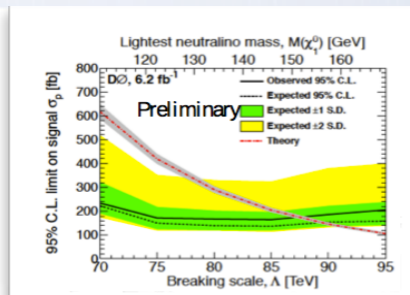
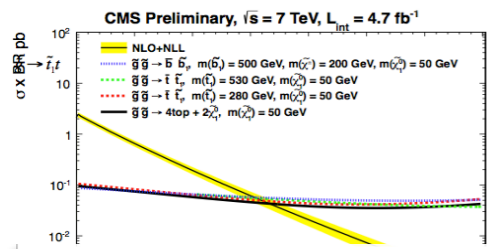
Advantage: less model dependent

Disadvantage: many parameters, process dependent

Both approaches are used

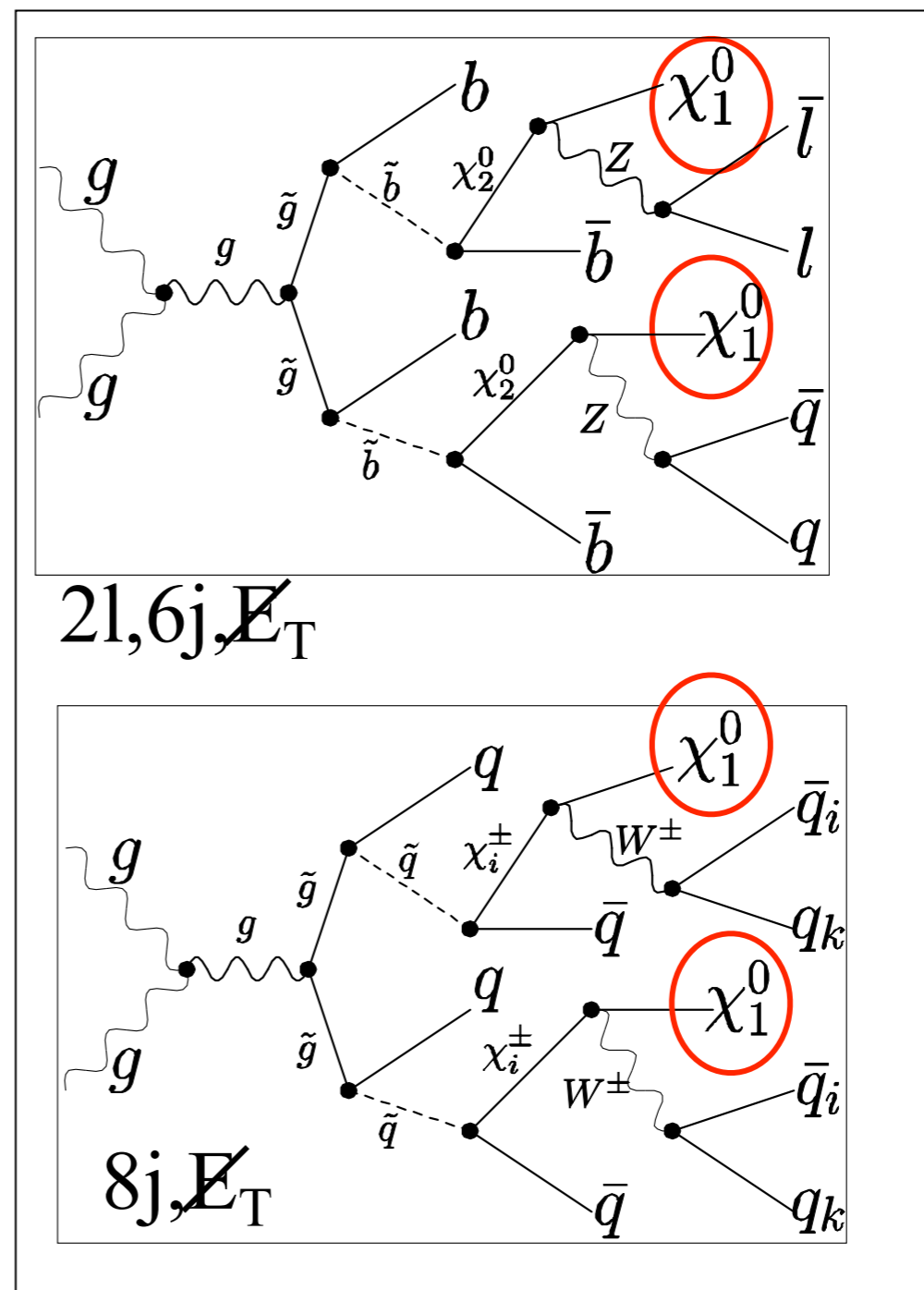
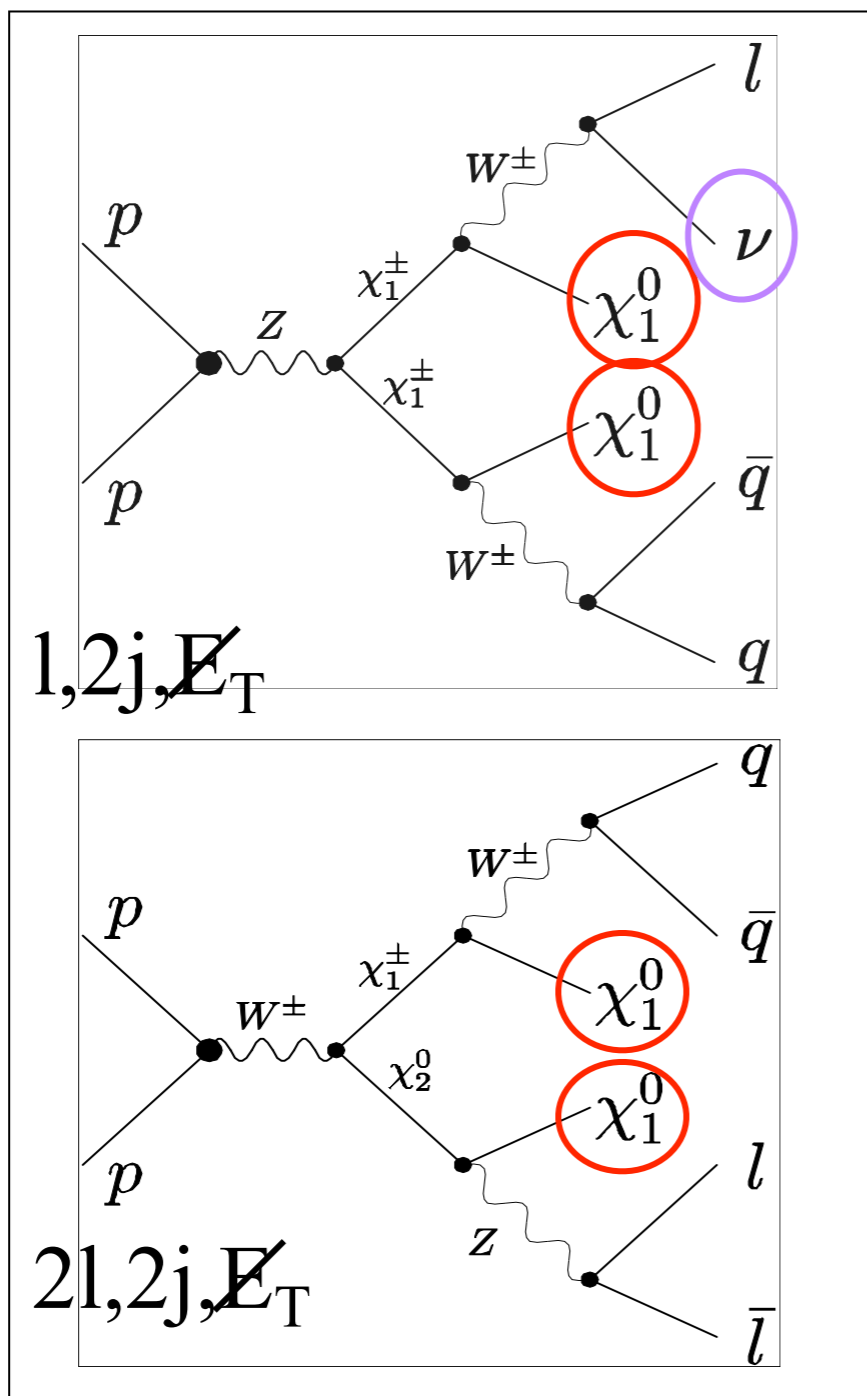


SUSY searches



Creation and Decay of Superpartners in Cascade Processes @ LHC

weak int's

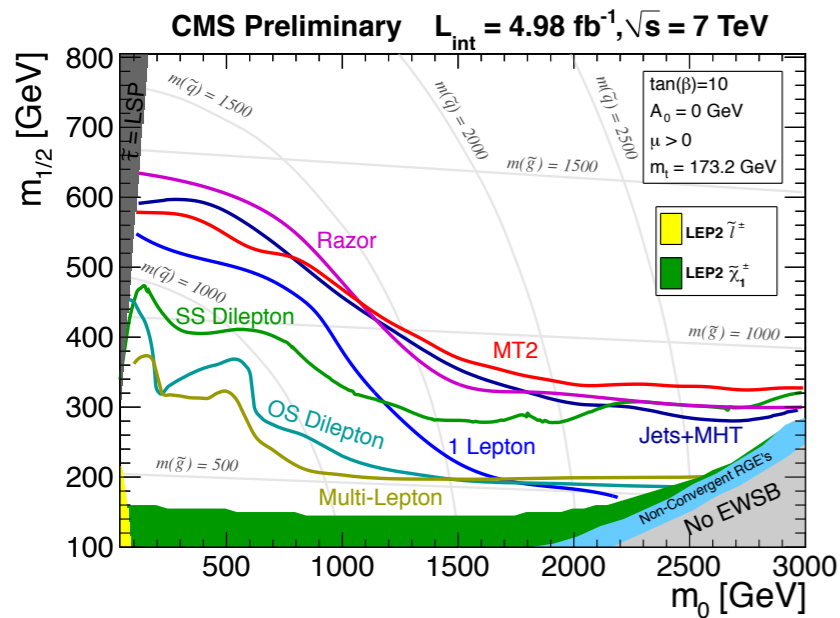


Strong int's

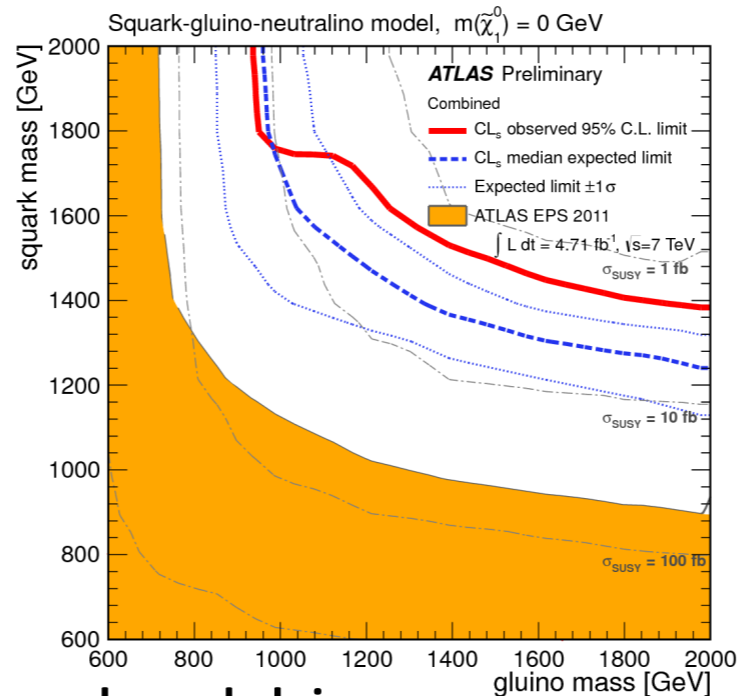
Typical SUSY signature: Missing Energy and Transverse Momentum

The Progress of LHC

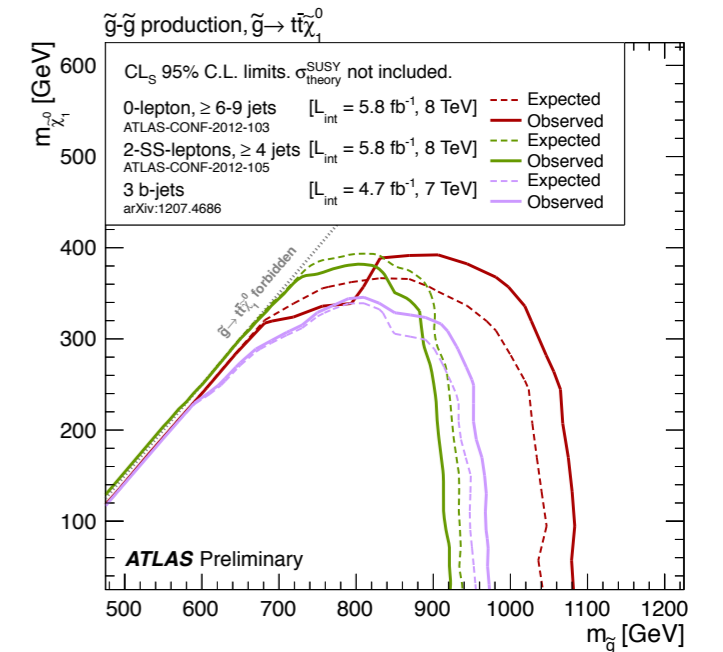
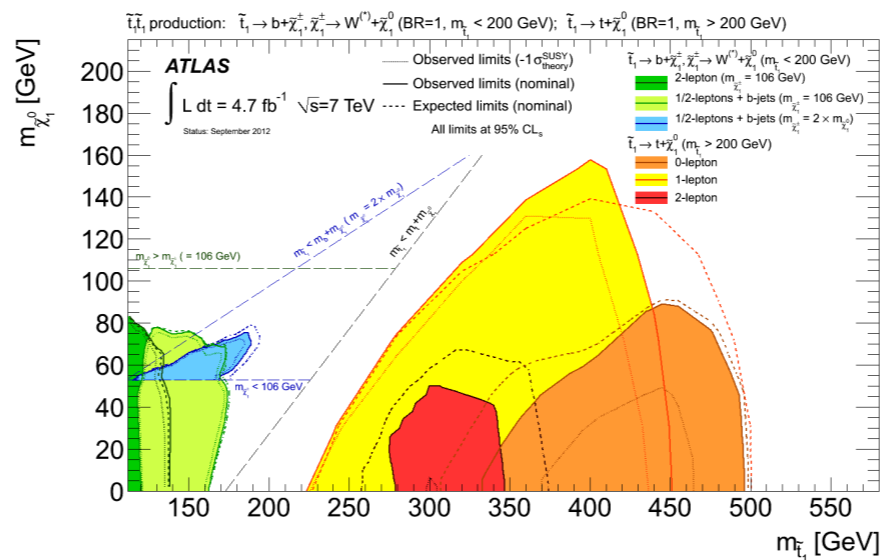
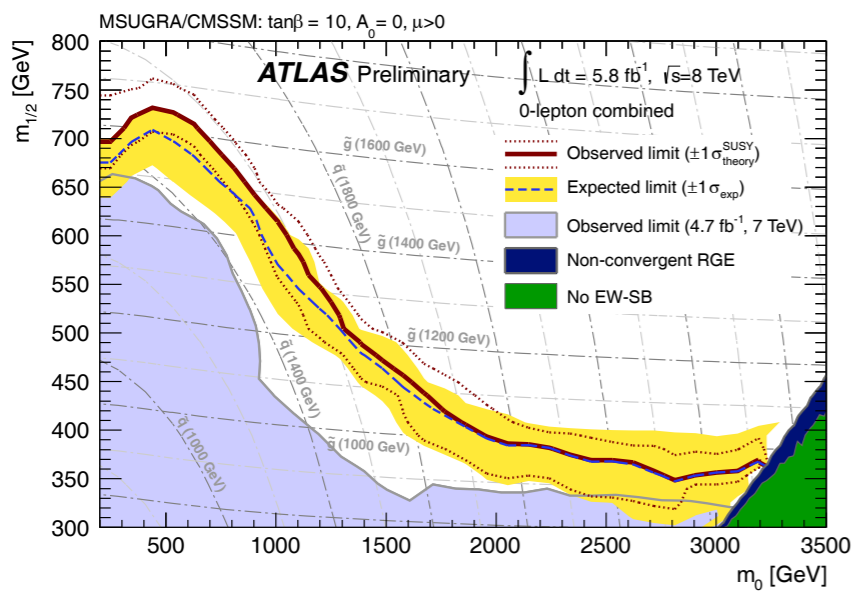
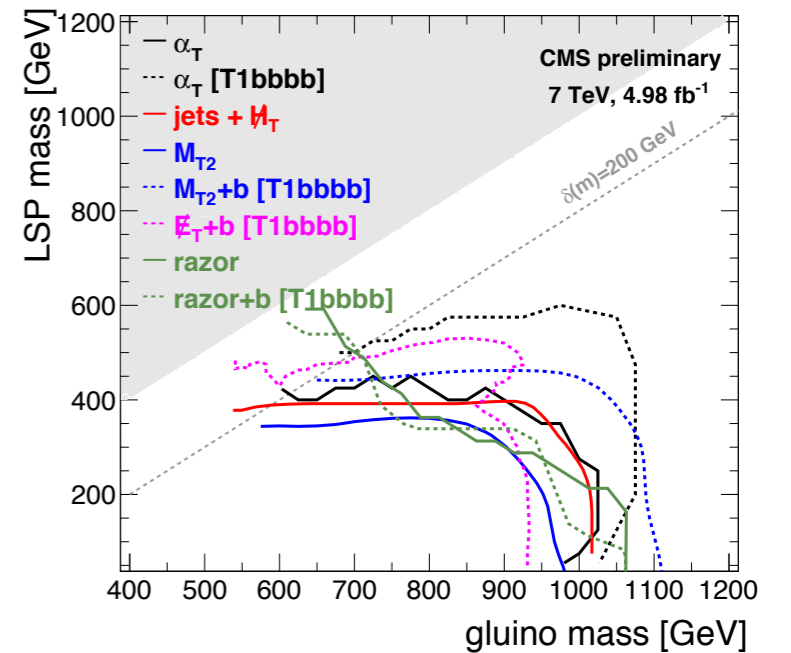
CMS CMSSM summary



Look for squarks and gluinos with direct decays to SM+LSP



95% exclusion limits for $\tilde{g} \rightarrow q q \tilde{\chi}^0$; $m(\tilde{q}) \gg m(\tilde{g})$



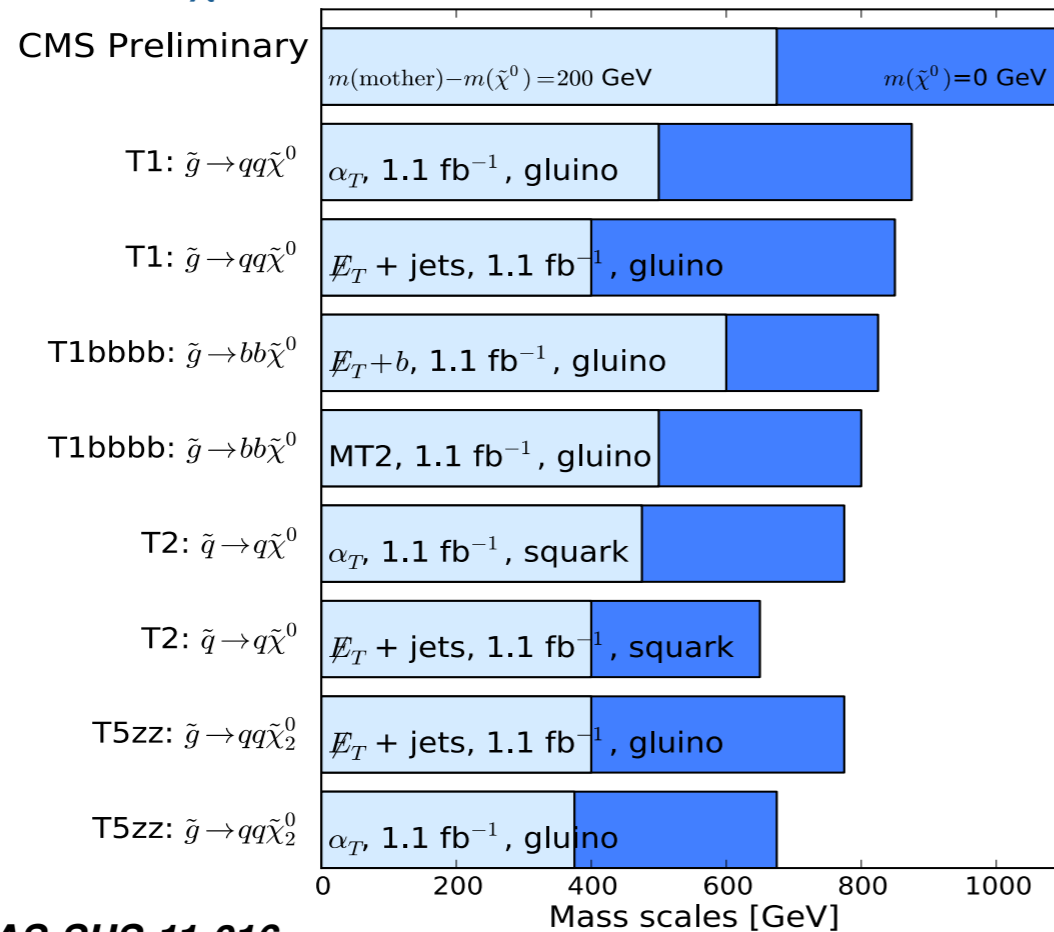
$$\tilde{m}_g > 1000 \text{ GeV}$$

$$\tilde{m}_q > 1400 \text{ GeV}$$

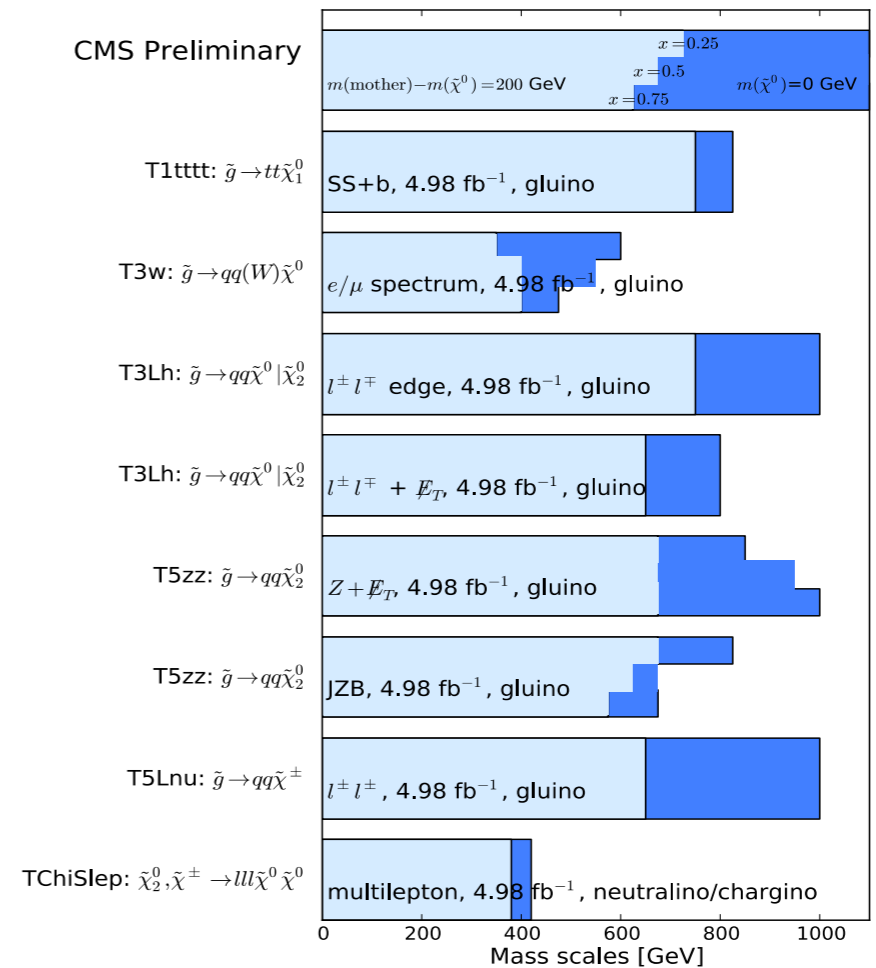


SUSY in simplified models

Hadronic (left) and leptonic (right) SUSY searches in simplified SUSY models. Exclusion limits for gluino and squark masses, for $m_{\tilde{\chi}^0} = 0$ GeV (dark blue) and $m_{\text{mother}} - m_{\tilde{\chi}^0} = 200$ GeV (light blue).

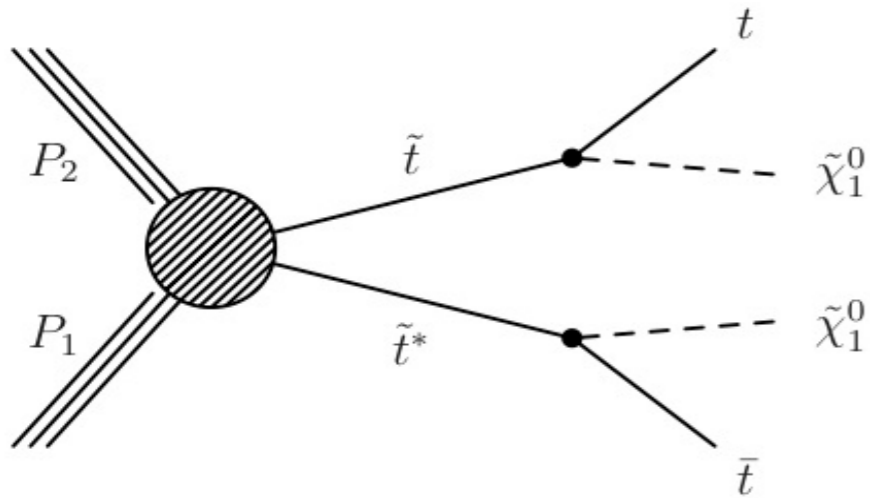


CMS-PAS-SUS-11-016

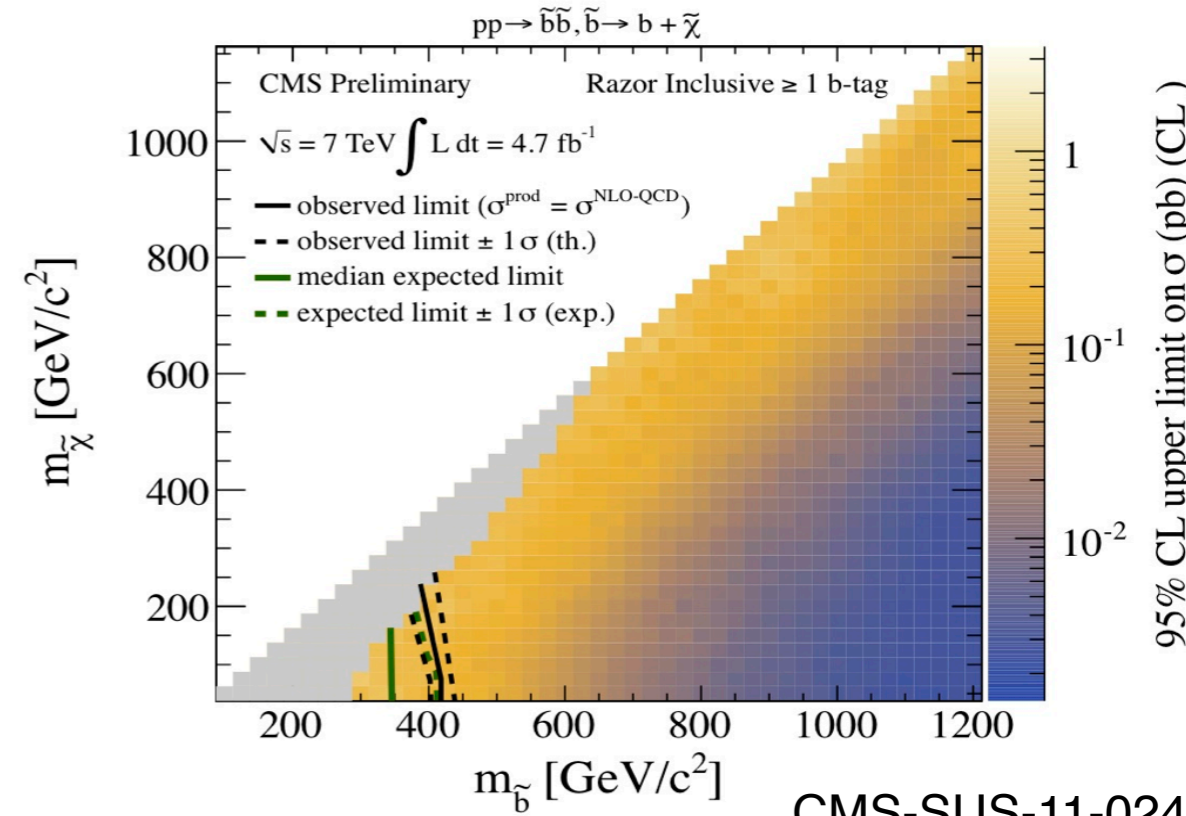
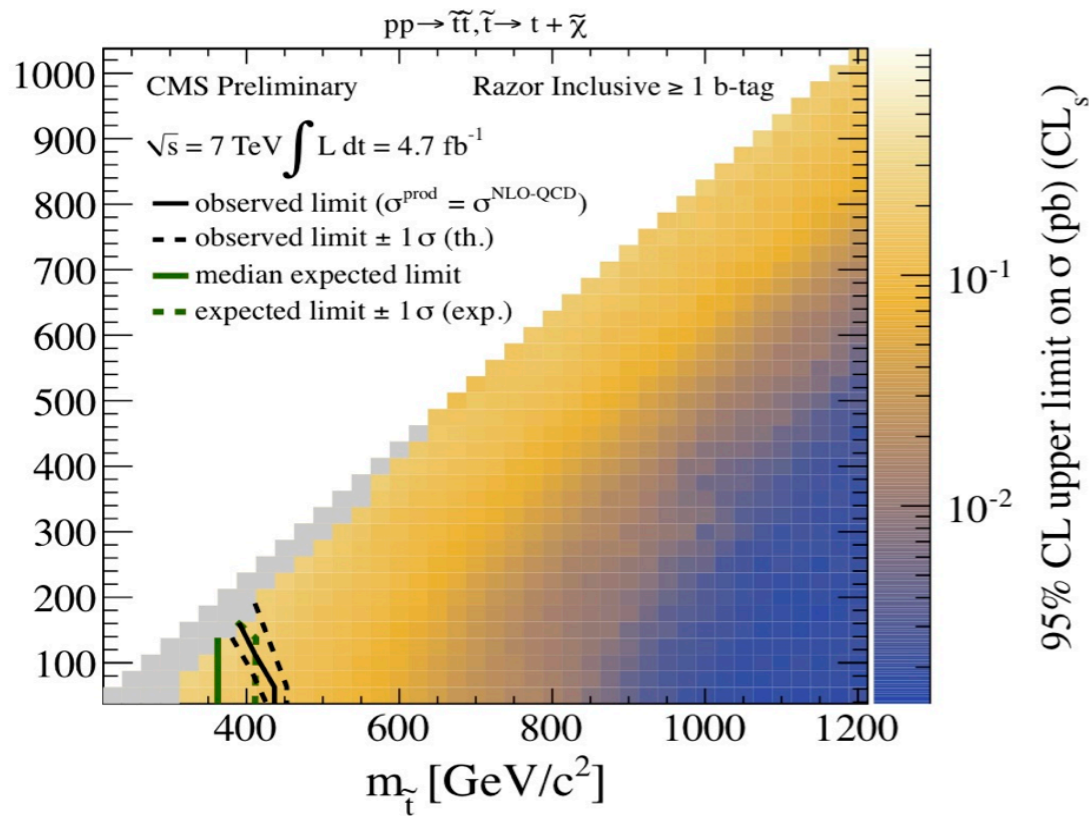
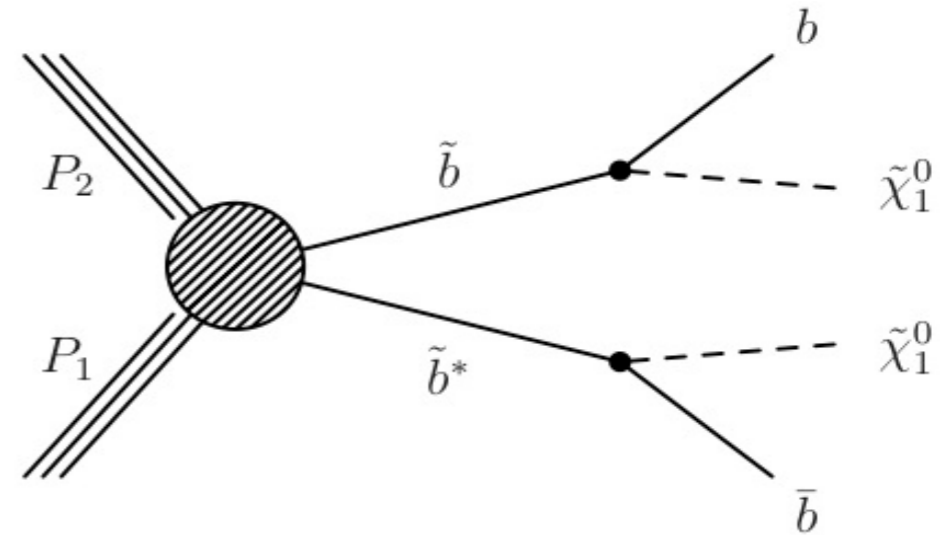


SUSY is not dead (yet). It might still hide in low MET/low HT events. More complicated models are under investigation → more challenging searches. For some it is hard to even get the data on tape.

Di-stop production resulting in 2 top quarks +MET final states



Di-sbottom production resulting in 2 b quarks +MET final states



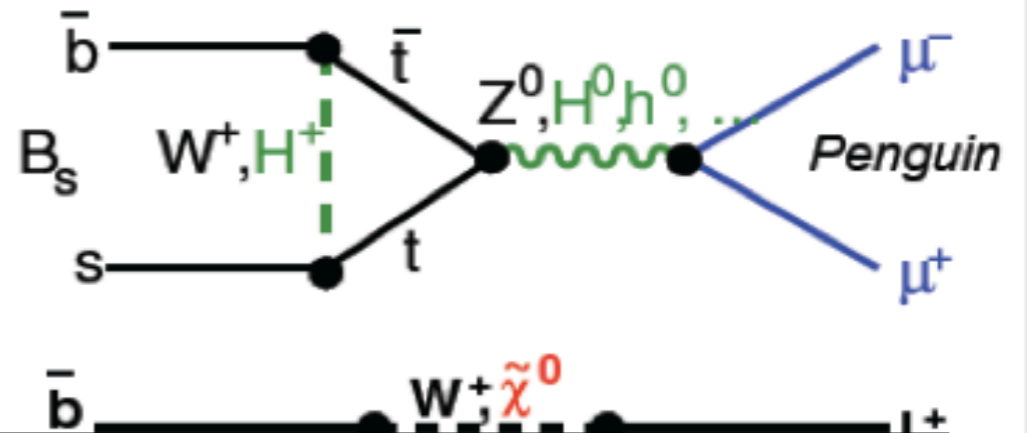
CMS-SUS-11-024

Probing SUSY with

$$B_{s,d} \rightarrow \mu^+ \mu^-$$

- **Decays highly suppressed in SM**

- Forbidden at tree level
- $b \rightarrow s(d)$ FCNC transition only through penguin and box diagrams
- Helicity suppressed by factors of (m_μ/m_B)



$$Br[B_s \rightarrow \mu\mu] = \frac{2\tau_B m_B^5}{64\pi} f_{B_s}^2 \sqrt{1 - \frac{4m_l^2}{m_B^2}} \left[\left(1 - \frac{4m_l^2}{m_B^2}\right) \left| \frac{(C_S - C'_S)}{(m_b + m_s)} \right|^2 + \left| \frac{(C_P - C'_P)}{(m_b + m_s)} + 2 \frac{m_\mu}{m_{B_s}^2} (C_A - C'_A) \right|^2 \right]$$

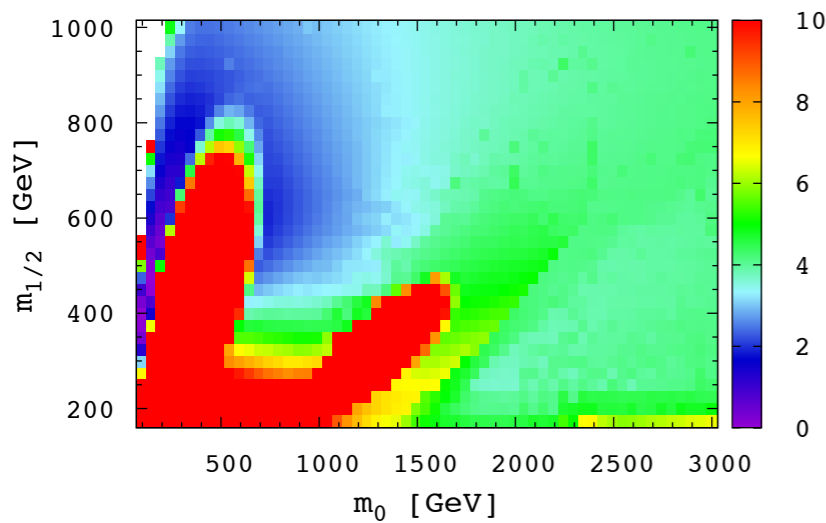
$$C_S \simeq \frac{G_F \alpha}{\sqrt{2}\pi} V_{tb} V_{ts}^* \left(\frac{\tan^3 \beta}{\sin^2 \theta_W} \right) \left(\frac{m_b m_\mu m_t \mu}{M_W^2 M_A^2} \right) \frac{\sin 2\theta_{\tilde{t}}}{2} \left(\frac{m_{\tilde{t}_1}^2 \log \left[\frac{m_{\tilde{t}_1}^2}{\mu^2} \right]}{\mu^2 - m_{\tilde{t}_1}^2} - \frac{m_{\tilde{t}_2}^2 \log \left[\frac{m_{\tilde{t}_2}^2}{\mu^2} \right]}{\mu^2 - m_{\tilde{t}_2}^2} \right)$$

Enhancement
Suppression

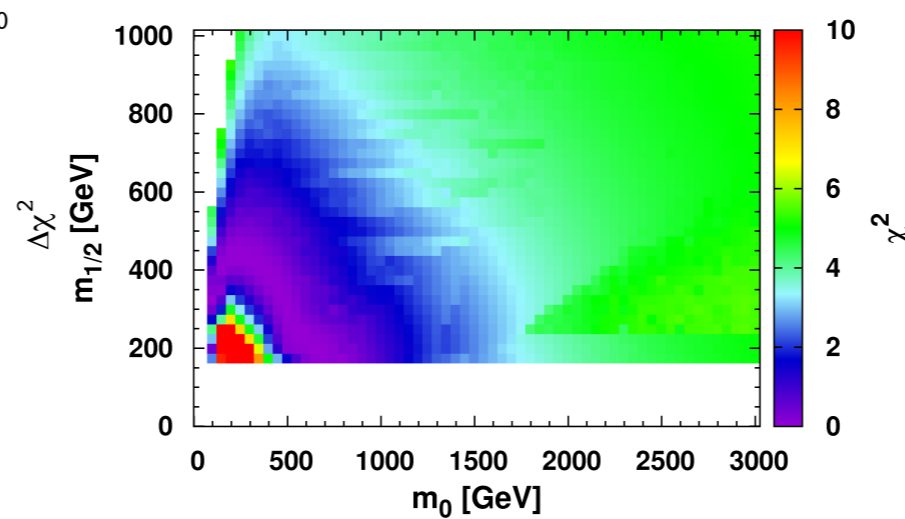
Electroweak, g-2, and Dark Matter constraints

W.de Boer, C.Beskidt, D.K.'11'12

95% CL exclusion

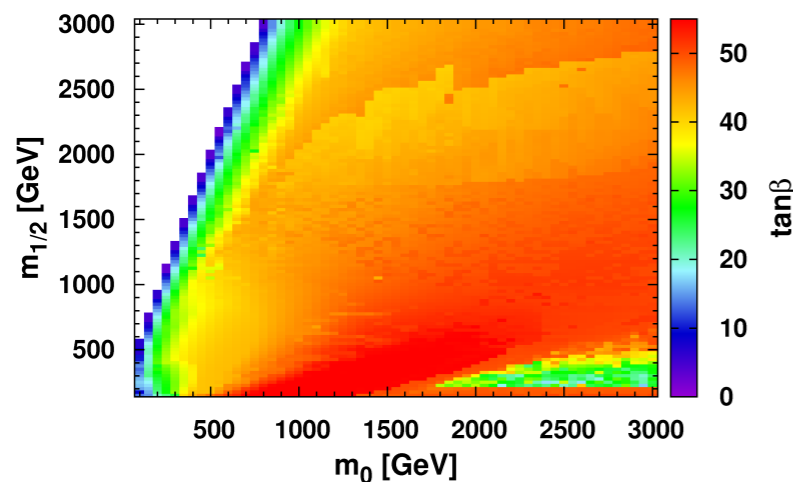


$B_s \rightarrow s\gamma, B_s \rightarrow \mu^+\mu^-, B_s \rightarrow \tau\nu$

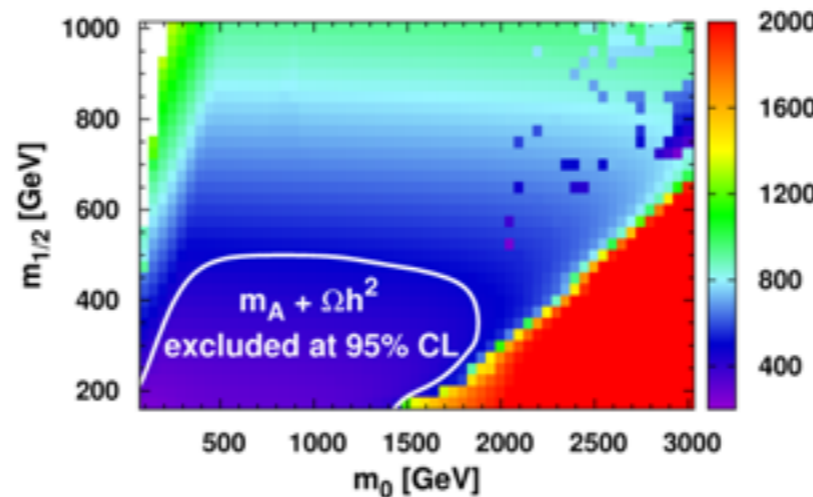


muon $g - 2$

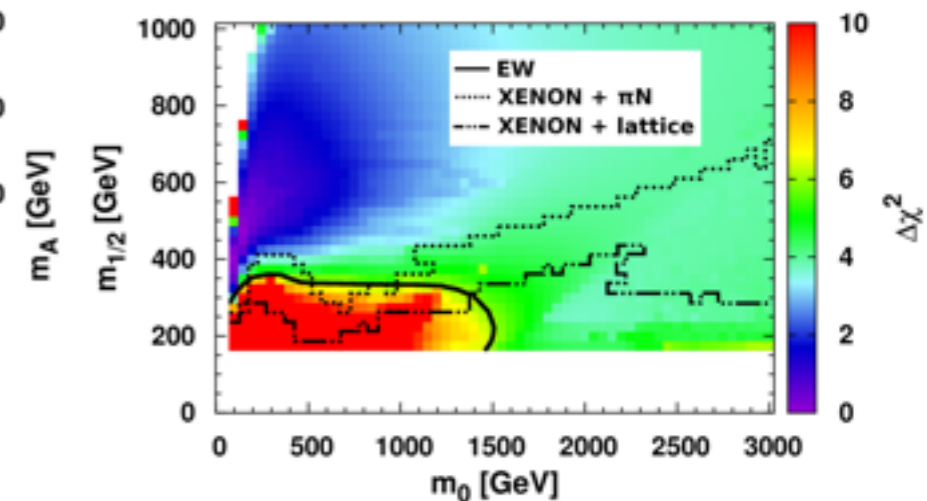
Constraint	Data
Ωh^2	0.113 ± 0.004
$b \rightarrow X_s \gamma$	$(3.55 \pm 0.24) \cdot 10^{-4}$
$B_u \rightarrow \tau\nu$	$(1.68 \pm 0.31) \cdot 10^{-4}$
Δa_μ	$(302 \pm 63(exp) \pm 61(theo)) \cdot 10^{-11}$
$B_s^0 \rightarrow \mu^+\mu^-$	$B_s^0 \rightarrow \mu^+\mu^- < 4.5 \cdot 10^{-9}$
m_h	$m_h > 114.4 \text{ GeV}$
m_A	$m_A > 480 \text{ GeV for } \tan\beta \approx 50$
ATLAS	$\sigma_{had}^{SUSY} < 0.003 - 0.03 \text{ pb}$
CMS	$\sigma_{had}^{SUSY} < 0.005 - 0.03 \text{ pb}$
XENON100	$\sigma_{\chi N} < 8 \cdot 10^{-45} - 2 \cdot 10^{-44} \text{ cm}^2$



$\tan\beta$ fit of DM abundance

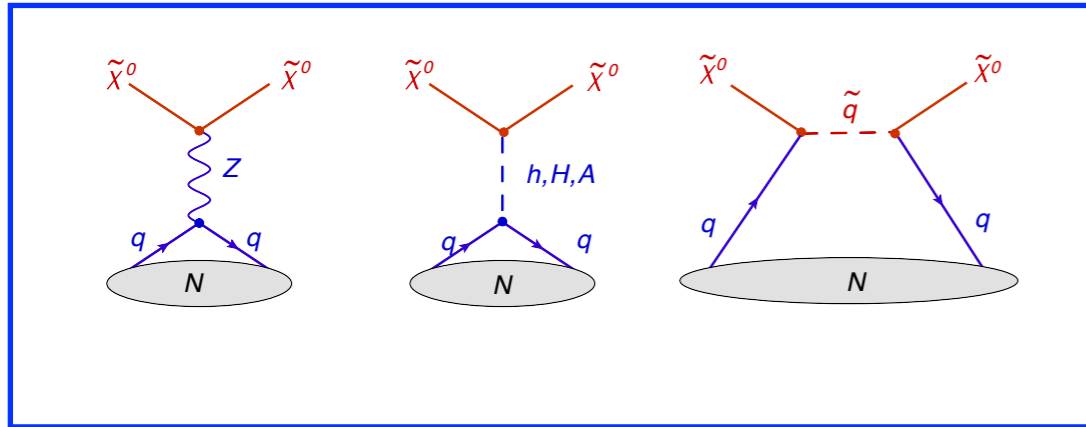


Dark matter abundance

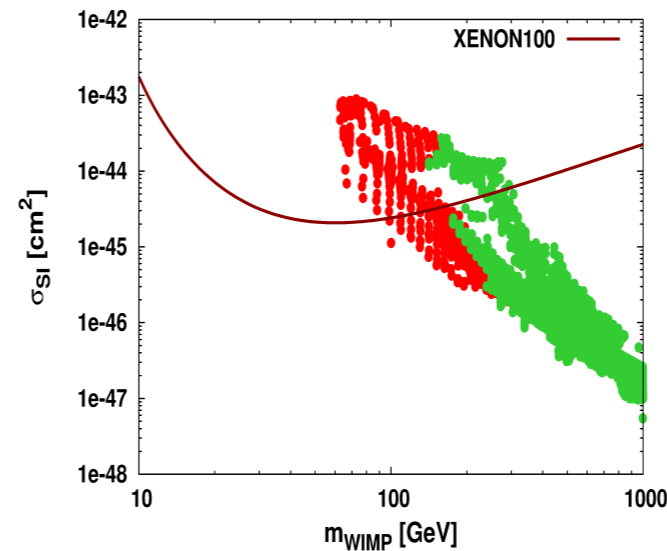


Direct DM search

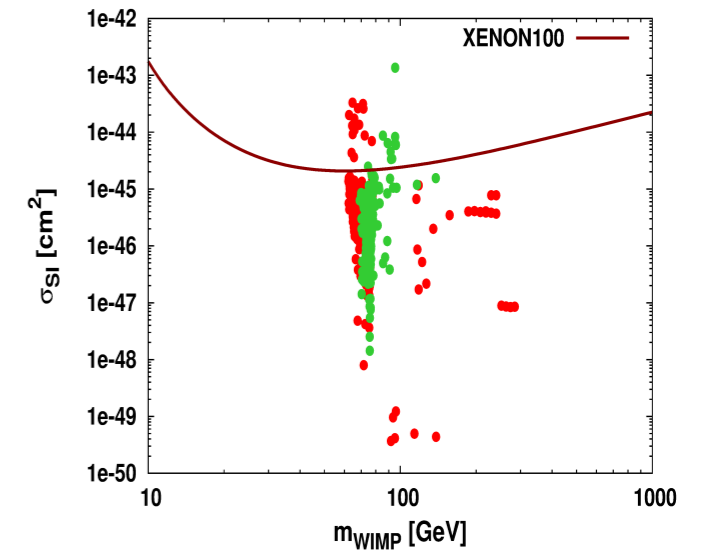
DM Nucleons Interaction



MSSM

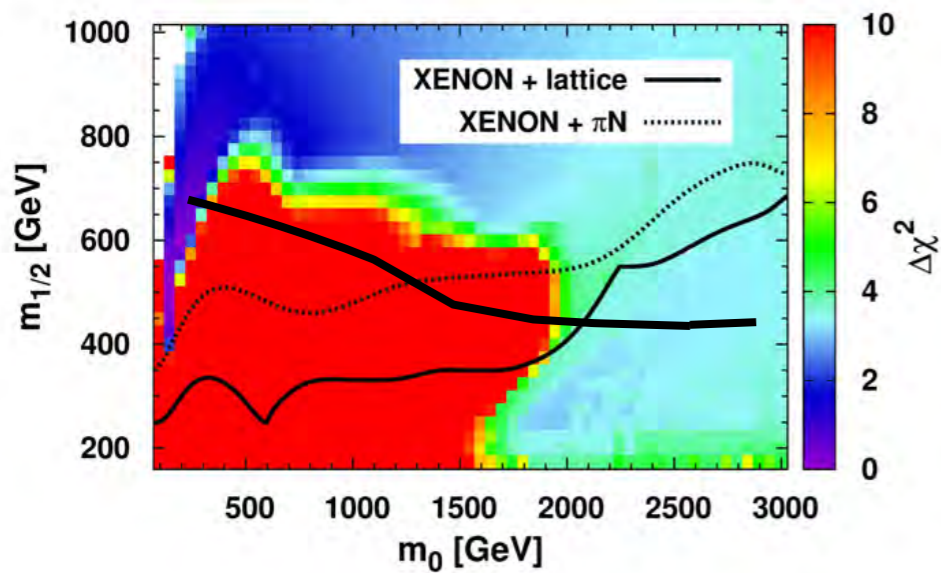


NMSSM



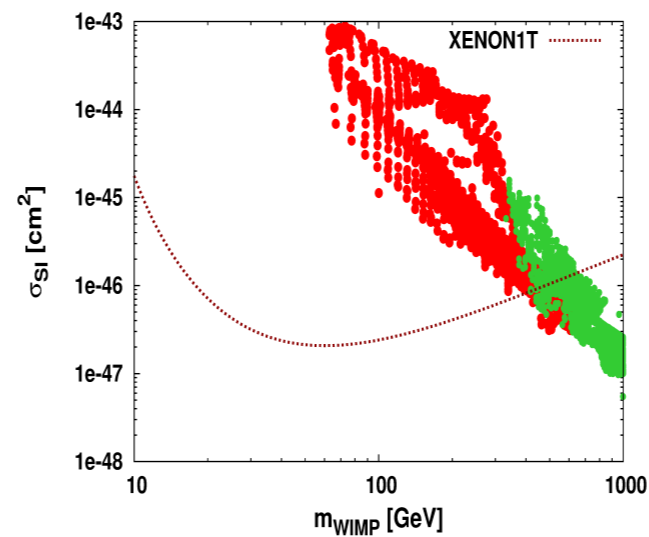
Red points - excluded by LHC

XENON100+LHC 8

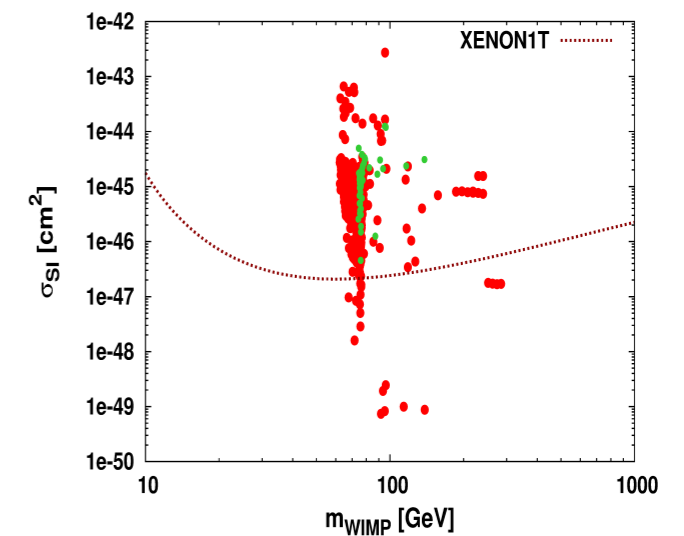


MSSM

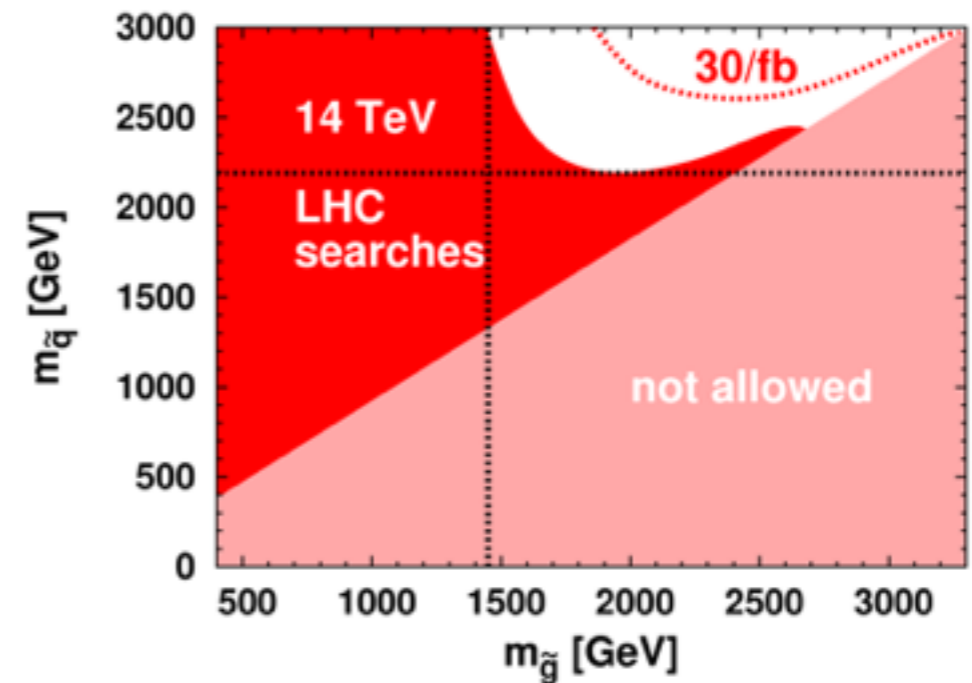
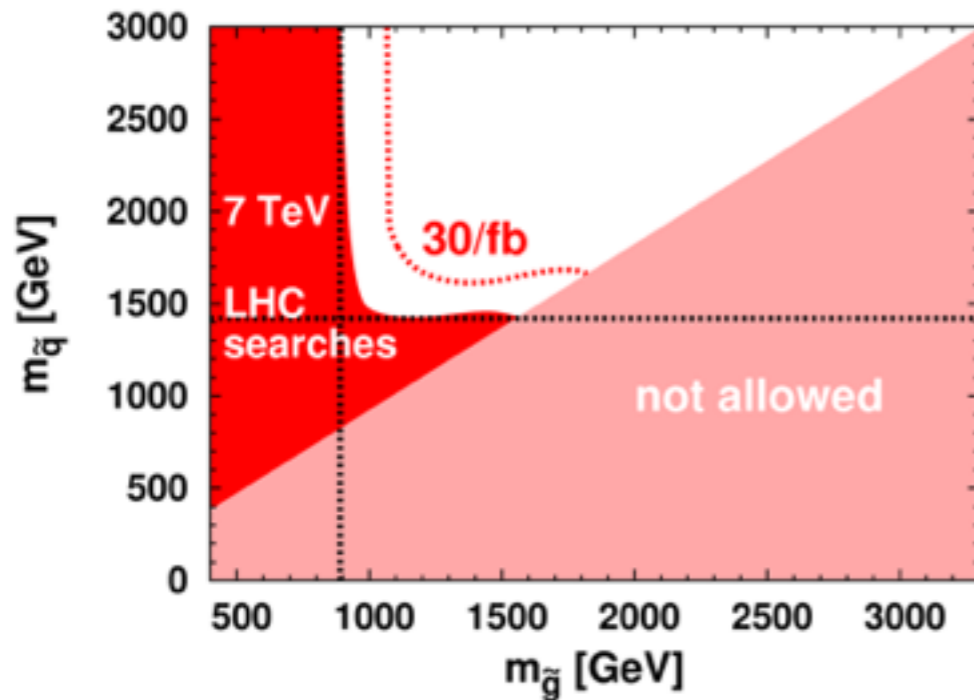
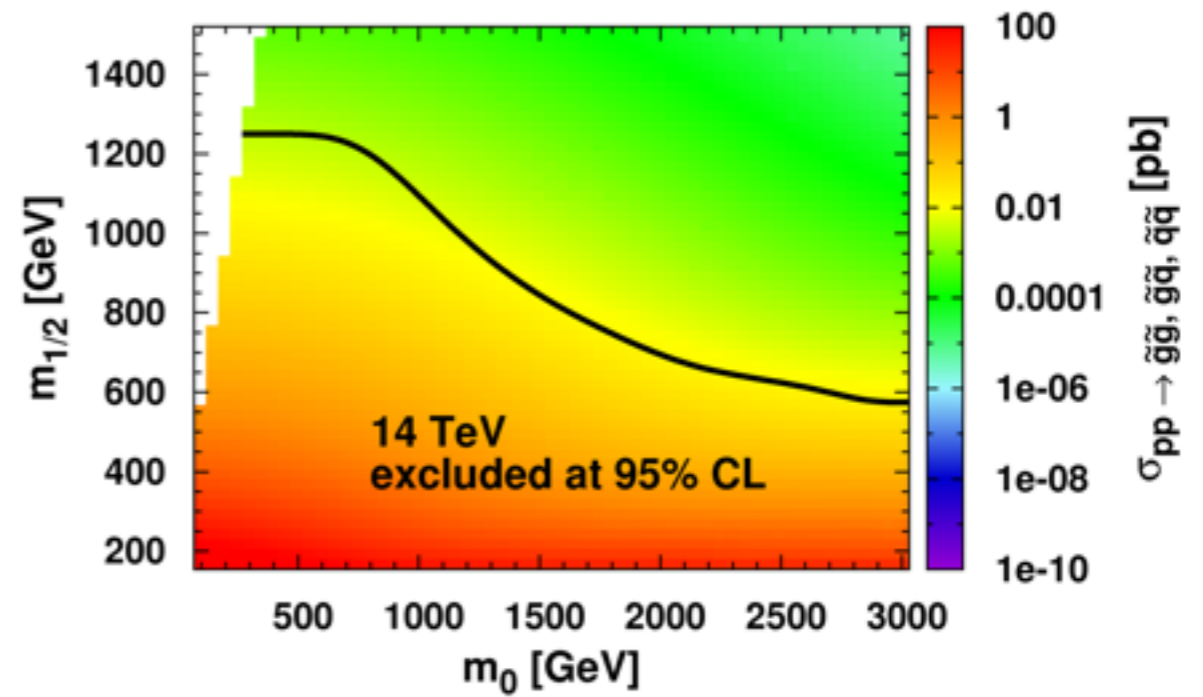
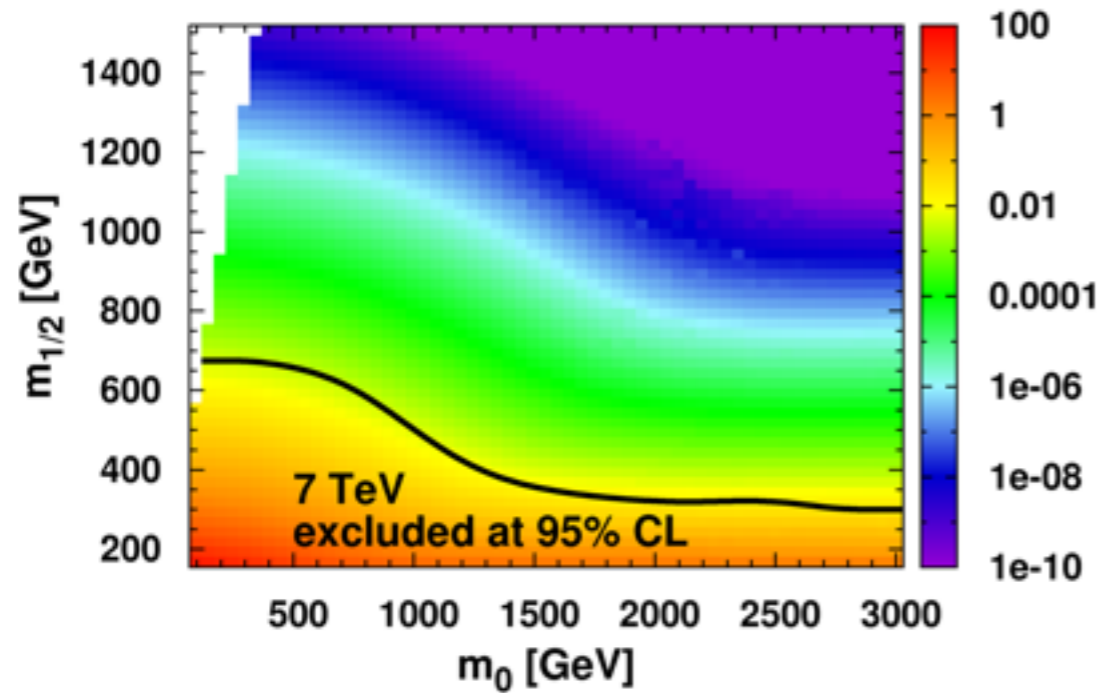
Direct DM search



XENON1T+LHC 14



LHC Reach at 7 and 14 TeV

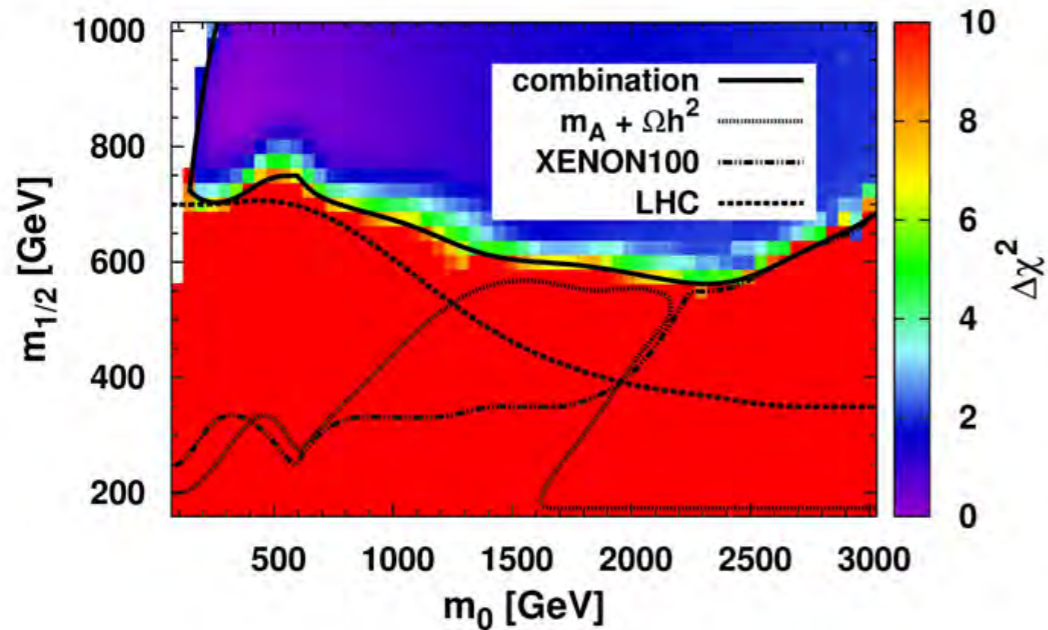


Energy is more important than luminosity

Combined Fit to all Data

W.de Boer, C.Beskidt, D.K.'11'12

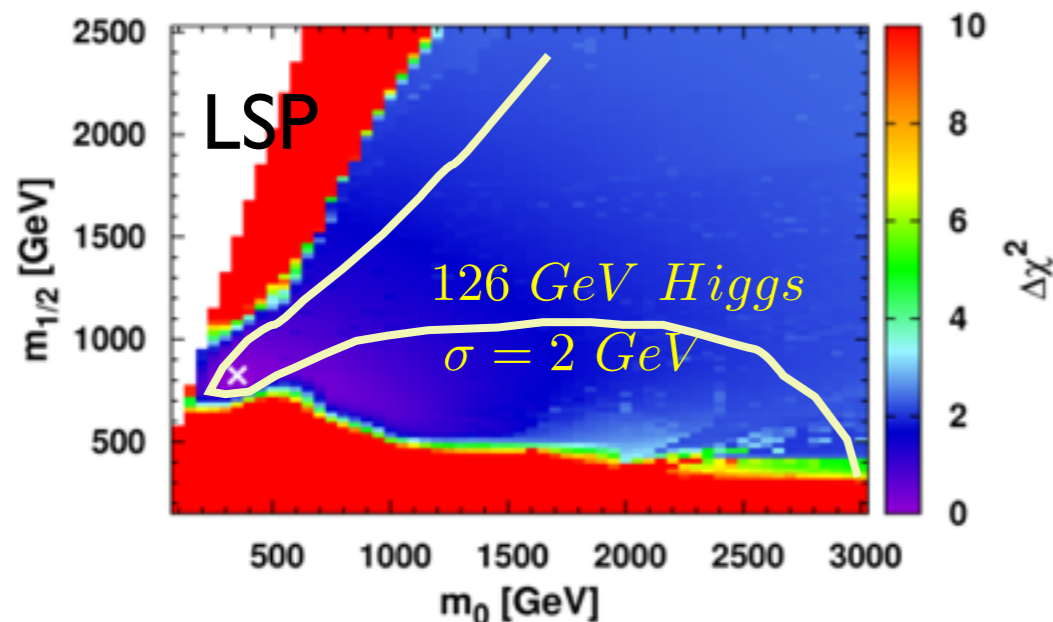
MSSM



without
125 GeV
Scalar

Constraint	Data
Ωh^2	0.113 ± 0.004
$b \rightarrow X_s \gamma$	$(3.55 \pm 0.24) \cdot 10^{-4}$
$B_u \rightarrow \tau \nu$	$(1.68 \pm 0.31) \cdot 10^{-4}$
Δa_μ	$(302 \pm 63(\text{exp}) \pm 61(\text{theo})) \cdot 10^{-11}$
$B_s^0 \rightarrow \mu^+ \mu^-$	$B_s^0 \rightarrow \mu^+ \mu^- < 4.5 \cdot 10^{-9}$
m_h	$m_h > 114.4 \text{ GeV}$
m_A	$m_A > 480 \text{ GeV for } \tan \beta \approx 50$
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XENON100	$\sigma_{\chi N} < 8 \cdot 10^{-45} - 2 \cdot 10^{-44} \text{ cm}^2$

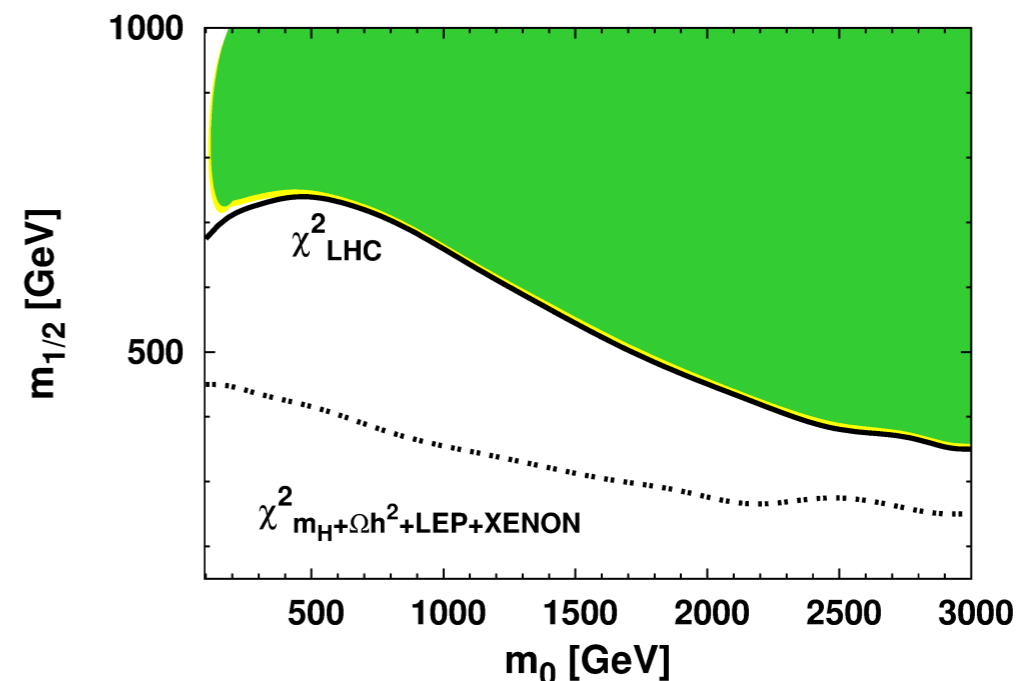
The values of $\tan \beta$ and A_0 are adjusted



with 125
GeV
Scalar

Larger scale for $m_{1/2}$

NMSSM



Conclusion



SUSY today:

- No signal so far, but do not give up
- There is still plenty of room for SUSY
- Interpretations of searches are model dependent
- LHC run at 14 TeV might be crucial for low energy SUSY



- Give me something better and I will stick to it

Concluding Remarks



The Higgs Boson:

- Everything is developing as expected
- Wait till the end of the run
- It looks like the Higgs boson so far



SUSY:

- No signal so far, but do not give up
- There is still plenty of room for SUSY



LHC:

- Excellent performance, new results are ahead



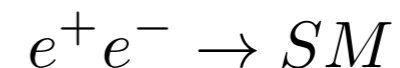
Linear Collider:

- It is time to think of it again as a precision physics and discovery machine

Discovery Machine



Precision Machine



Precision physics of Higgs bosons

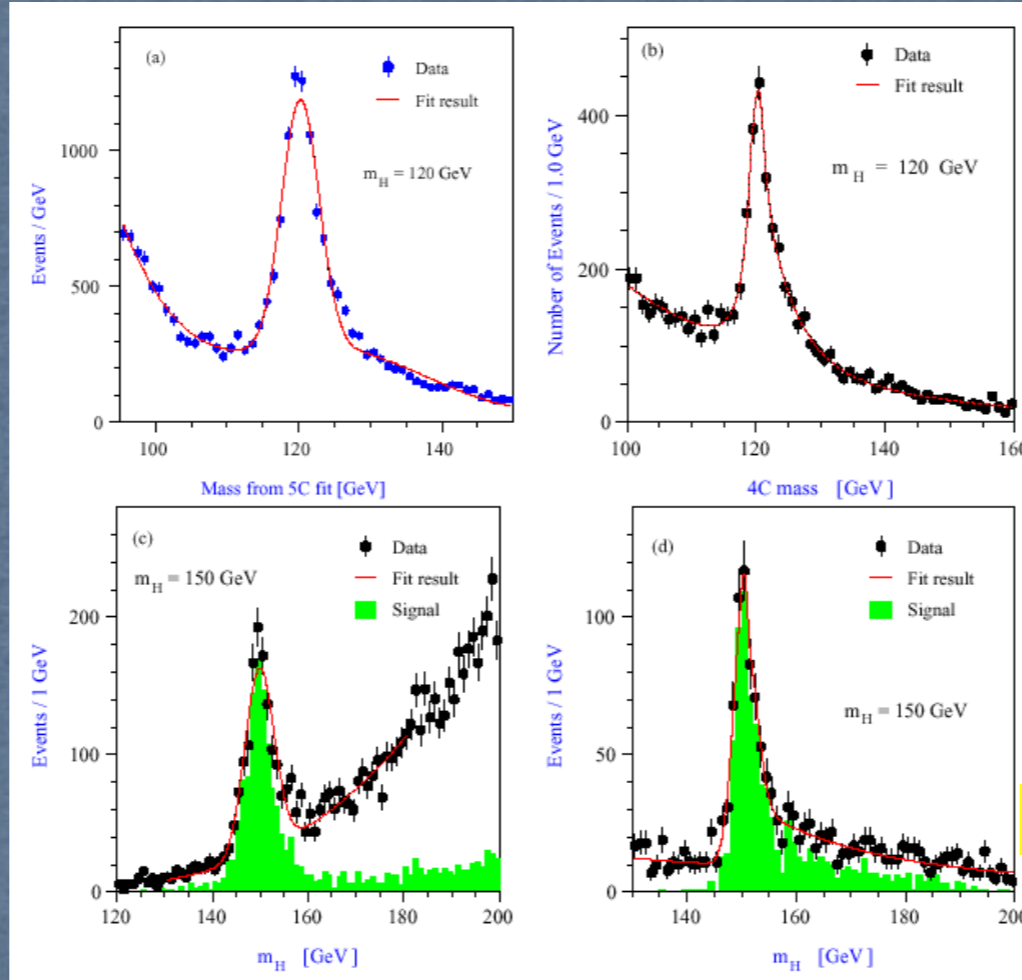
$ee \rightarrow HZ$ diff. decay channels

$m_H = 120 \text{ GeV}$

$\rightarrow b\bar{b}q\bar{q}$

$m_H = 150 \text{ GeV}$

$\rightarrow W^+W^-q\bar{q}$



$\rightarrow q\bar{q}l^+l^-$

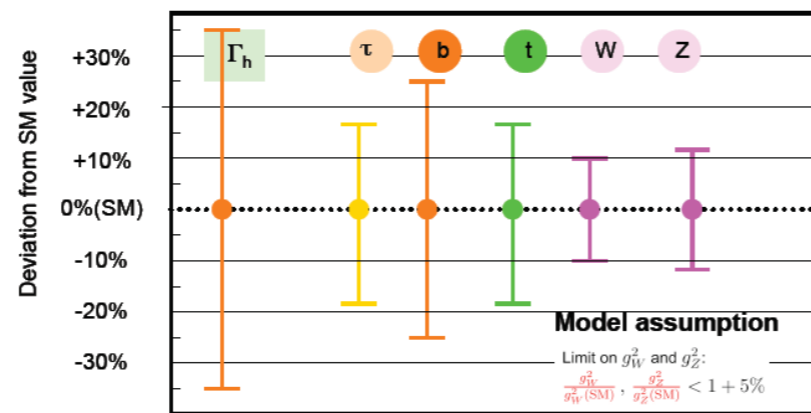
$\Delta m_H = 40 \text{ MeV}$

$\rightarrow W^+W^-l^+l^-$

$\Delta m_H = 70 \text{ MeV}$

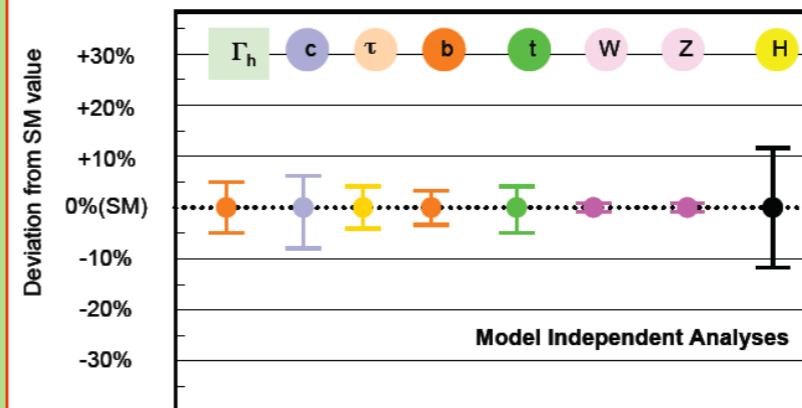
Coupling Precision

LHC $300 \text{ fb}^{-1} \times 2$

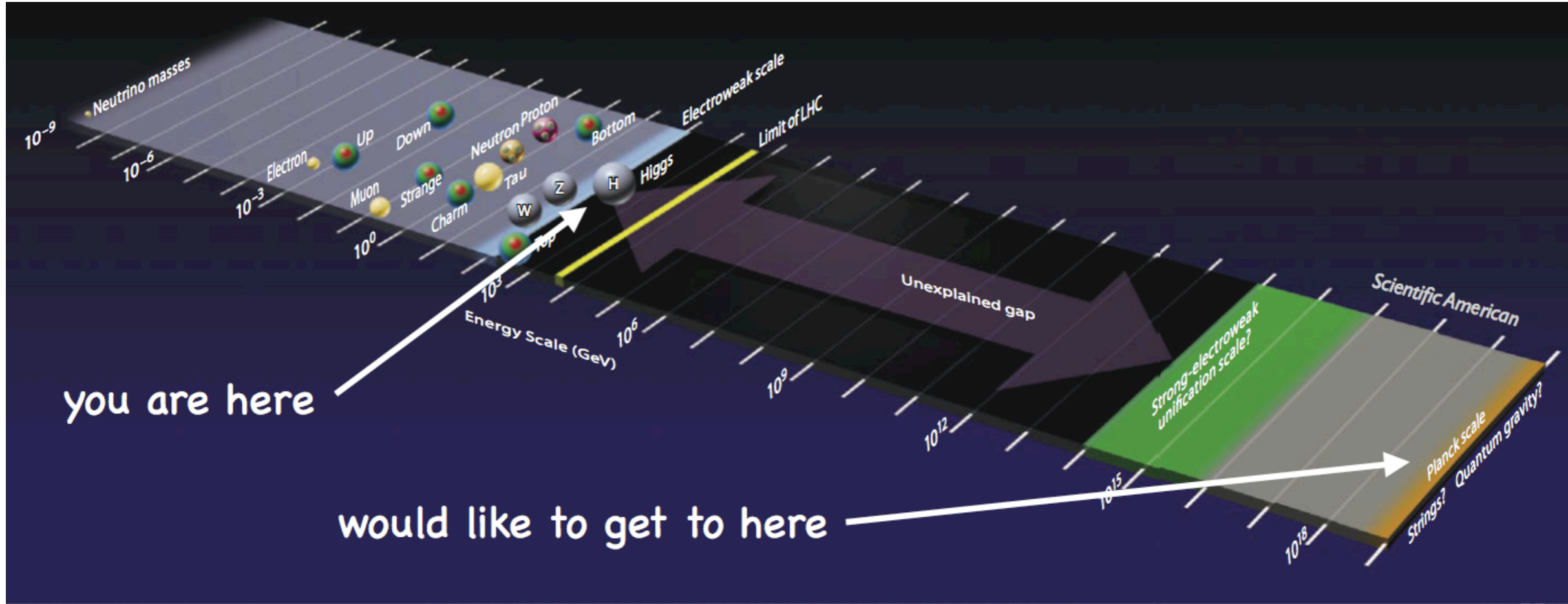


Coupling Precision

ILC



HEP Scale



The End