

# Beauty Physics and $CP$ Violation (II)

Muon/Hadron Detector

## $\sin(2\beta)$ and the Triumph of the Standard Model

Magnet Coil

Electron/Photon Detector

Cherenkov Detector

Tracking Chamber

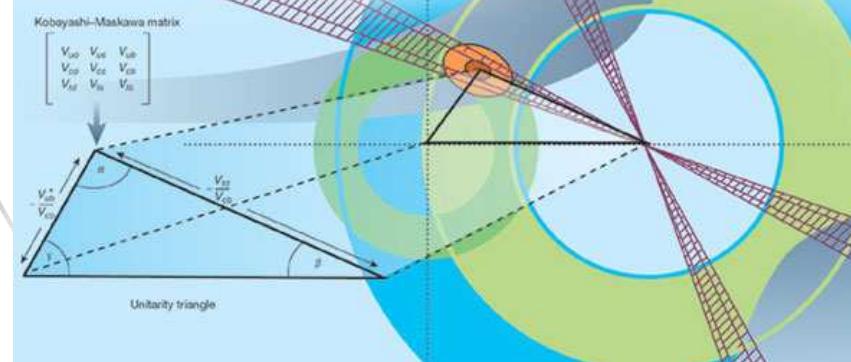
Support Tube

Vertex Detector

$e^-$

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$e^+$

Experimental lecture at the Helmholtz School on Heavy Quark Physics

Dubna, June 6-16, 2005

# Themes

## I. Beauty Physics and $CP$ Violation – the experimental program

- Heavy meson production and decay
- $B$  Physics and  $CP$  Violation
- The  $B$  Factories
- Physics at the  $\Upsilon(4S)$ : time-integrated and time-dependent measurements



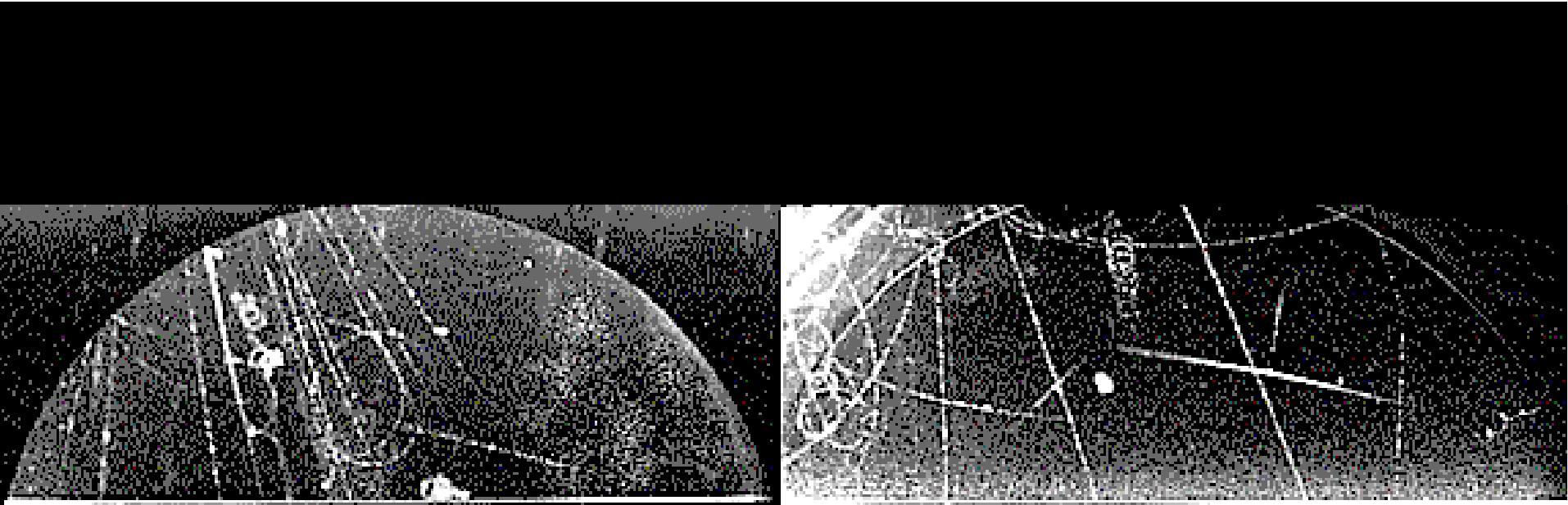
## II. $\sin(2\beta)$ and the triumph of the Standard Model

- $CP$  violation: experimental facts
- $CP$  violation in the  $B$  system
- The measurement of  $\sin(2\beta)$  in tree and loop (penguin) decays
- Briefing on radiative  $B$  decays

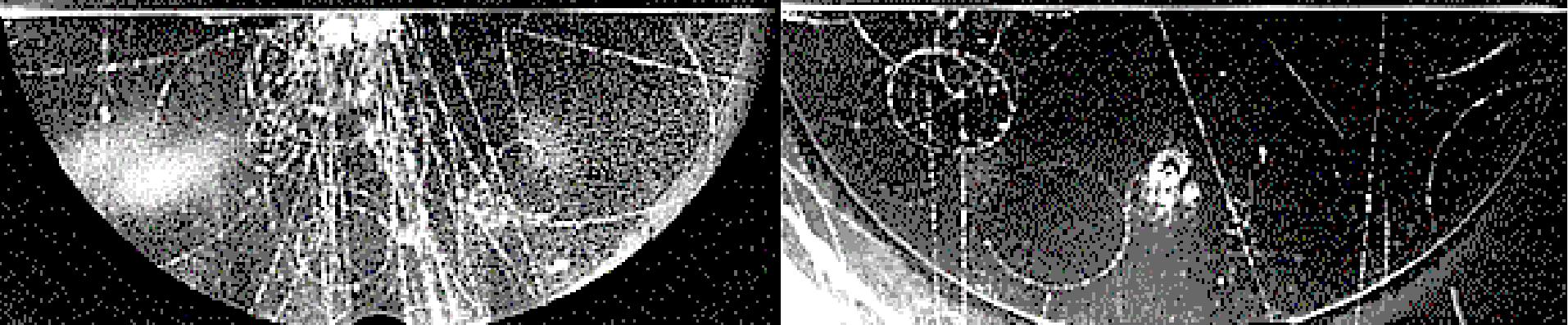
## III. Rare $B$ decays: towards the full unitarity triangle ... and beyond

- Leptonic  $B$  Decays
- Charmless  $B$  decays and the measurement of  $\alpha$
- $B \rightarrow K\pi$  decays (direct  $CP$  violation) and other charmless modes
- Towards  $\gamma$
- Flavor, CPV and CKM: the present picture and the experimental future

# sin(2β) and the Triumph of the Standard Model



## CP Violation: Experimental Facts



# Experimental Facts (I)

- ★ The  $K_L^0$  decays into  $\pi^+\pi^-$  and  $\pi^0\pi^0$   
the  $CP$ -violating parameters are  $\eta_{+-}$  and  $\eta_{00}$

historically, the way  $CP$  violation  
was discovered back in 1964

$$\eta_{+-} = \frac{A(K_L^0 \rightarrow \pi^+\pi^-)}{A(K_S^0 \rightarrow \pi^+\pi^-)} = (2.286 \pm 0.017) \cdot 10^{-3} \cdot e^{i(43.51^\circ \pm 0.06^\circ)}$$

and the measurement of  $\eta_{+-}$  is consistent with that of  $\eta_{00}$

The “super-weak” hypothesis:  $CP$  (time reversal) symmetry violation  
is confined to  $\Delta S = 2$  processes ( $K^0\bar{K}^0$  mixing)

The super-weak model was proposed less than one year after the experimental discovery of  $CP$  violation. A new interaction affecting only kaon mixing seemed a “natural” explanation for the small observed effects of  $CP$  violation. The Kobayashi-Maskawa model came 9 years later, still before the discovery of charm.

→ under the super-weak hypothesis:  $\eta_{+-} = \eta_{00} = \varepsilon$   $|K_L^0\rangle \propto (1 + \varepsilon)|K^0\rangle - (1 - \varepsilon)|\bar{K}^0\rangle$

# Experimental Facts (II)



Direct measurements of  $|\eta_+/\eta_{00}|$  show significant departure from unity

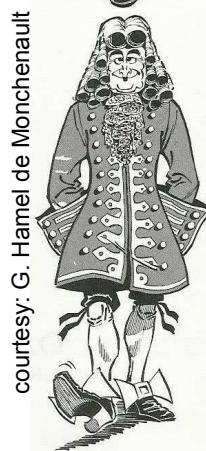
$$\text{Re}(\varepsilon'/\varepsilon) = (1 - |\eta_{00}/\eta_+|)/3 = (16.7 \pm 2.6) \times 10^{-4}$$

→  $\text{Re}(\varepsilon') \propto 10^{-6}$  tiny effect !

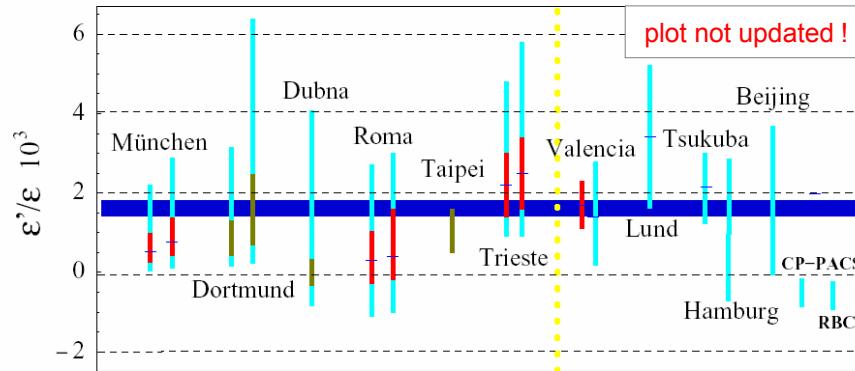
Average: KTeV & NA48

this establishes direct *CP* violation (i.e. in a  $\Delta S=1$  decay process)

...and rules out the super-weak model (~30 years after the discovery of *CP* violation!)



Theoretical pre(post)dictions



...the ball is on the theory side



All observations are consistent with the conservation of the *CPT* symmetry

for instance:

$$m_{\bar{K}^0} - m_{K^0} = (1.7 \pm 4.2) \cdot 10^{-19} \text{ GeV}/c^2$$

including new preliminary result by NA48/1

# Experimental Facts (III)

- ★ 2001: mixing-induced  $CP$  violation has been observed in the  $B$  meson system from time-dependence of  $B^0 \rightarrow J/\psi K_s^0$  and similar decays



$$\sin(2\beta) = 0.726 \pm 0.037$$

BABAR & Belle, 2004

first uncontroversial effect of  $CP$  violation outside the kaon system

The observed  $CP$  violation effect in the  $B$  meson system is in excellent agreement with predictions using the Kobayashi & Maskawa framework in the Standard Model

- ★ 2004: large direct  $CP$  violation has been observed in the  $B$  meson system in  $B^0 \rightarrow K^+ \pi^-$  decays



$$A_{K^+ \pi^-} = -0.109 \pm 0.019$$

BABAR & Belle, 2004

- ★ With rising statistics, many more  $CP$ -violating  $B$  decays should show up  
>  $3\sigma$  candidates are  $B^0 \rightarrow \pi^+ \pi^-$ ,  $\rho^+ \pi^-$

# “Popular Misconceptions”

★ “*CP violation is always a small effect, typically of order  $10^{-3}$  or less*”

not completely exact!

- ➡  $CP$  asymmetry of 10% observed in neutral  $B$  decay to  $K^+\pi^-$
- ➡ Time-dependent  $CP$  asymmetry of 70% in the decay  $B^0, \bar{B}^0 \rightarrow J/\psi K_s^0$
- ➡  $CP$  asymmetry of 100% expected in the rare decay  $K_L^0 \rightarrow \pi^0 \nu \bar{\nu}$

However, it is true that  $CP$  violation is a very rare phenomenon

- ☀ either sizeable after several mean lifetimes
- ☀ or suppressed by very small branching fractions
- ☀ or both

★ Direct  $CP$  “theorem”:

- ☒ don't expect direct  $CP$  violation in copious decays
- ☒ rather look into rare modes

excluded area has CL > 0.95



# The $CP$ Operator

The  $CP$  eigenstates of the  $|P^0\rangle|\bar{P}^0\rangle$  system are:

$$|P_{CP=+1}^0\rangle \equiv \frac{1}{\sqrt{2}} (|P^0\rangle + CP|P^0\rangle)$$
$$|P_{CP=-1}^0\rangle \equiv \frac{1}{\sqrt{2}} (|P^0\rangle - CP|P^0\rangle)$$

with, by construction:

$$CP|P_{CP=+1}^0\rangle = +|P_{CP=+1}^0\rangle$$

$$CP|P_{CP=-1}^0\rangle = -|P_{CP=-1}^0\rangle$$

Note: arbitrary phase  $\xi$  in the definition of the  $CP$  operator

final state of  $P^0$  decay

$$\begin{cases} CP|P^0\rangle = e^{2i\xi_P}|\bar{P}^0\rangle \\ CP|\bar{P}^0\rangle = e^{-2i\xi_P}|P^0\rangle \end{cases} \quad \text{and} \quad CP|f\rangle = e^{2i\xi_f}|\bar{f}\rangle$$

Real eigenvalue for a  $CP$  final state:  $CP|f_{CP}\rangle = \sigma_{f_{CP}}|f_{CP}\rangle$  with  $\sigma_{f_{CP}} = \pm 1$

# Decay Amplitudes

- ★ Decay amplitudes of flavor states into the same final state:

$$\begin{cases} A_f \equiv \mathcal{A}(P^0 \rightarrow f) = \langle f | H | P^0 \rangle \\ \bar{A}_f \equiv \mathcal{A}(\bar{P}^0 \rightarrow f) = \langle f | H | \bar{P}^0 \rangle \end{cases}$$

note: decay amplitudes are phase-convention dependent

define

$$\lambda_f \equiv \frac{q}{p} \frac{\bar{A}_f}{A_f}$$

phase convention independent parameter

note: for flavor eigenstates  
 $A_f = 0$  or  $\bar{A}_f = 0$

- ★ Decay amplitudes of mass states:

$$\begin{cases} A_{Lf} \equiv \langle f | H | P_L \rangle = p A_f + q \bar{A}_f = p A_f (1 + \lambda_f) \\ A_{Hf} \equiv \langle f | H | P_H \rangle = p A_f - q \bar{A}_f = p A_f (1 - \lambda_f) \end{cases}$$

define

$$\eta_f \equiv \frac{A_{Hf}}{A_{Lf}} = \frac{1 - \lambda_f}{1 + \lambda_f}$$

# Why Different Notations ?

☀ In  $B$  physics, the physical (=mass) states cannot be isolated.

One starts with pure  $|B^0\rangle$  or  $|\bar{B}^0\rangle$  initial states,  
identified thanks to flavor tagging



therefore the natural parameter is  $\lambda_f$

☀ In  $K$  physics, the physical (=mass) states are well-isolated,  
thanks to very different lifetimes



therefore the natural parameter is  $\eta_f$

# Time Evolution of Neutral $B$ Meson System



Recall that: the time evolution of the physical states  $|B^0(t)\rangle$  ( $|\bar{B}^0(t)\rangle$ ) is given by

$$|B^0(t)\rangle = g_+(t)|B^0\rangle + \frac{q}{p}g_-(t)|\bar{B}^0\rangle$$

$$|\bar{B}^0(t)\rangle = \frac{p}{q}g_-(t)|B^0\rangle + g_+(t)|\bar{B}^0\rangle$$

where: (assuming  $\Delta\Gamma = 0$ )

$$g_+(t) = e^{-i(\mu_L + \mu_H)t/2} \cdot \cos(\Delta m t / 2)$$

$$g_-(t) = i \cdot e^{-i(\mu_L + \mu_H)t/2} \cdot \sin(\Delta m t / 2)$$

so that one finds for the time-dependent rates of an initially pure flavor state

$$\left| \langle f | H | B^0(t) \rangle \right|^2 = |A_f|^2 \cdot |g_+(t) + \lambda_f g_-(t)|^2 \quad \text{and} \quad \left| \langle f | H | \bar{B}^0(t) \rangle \right|^2 = \left| \frac{p}{q} \right|^2 \cdot |A_f|^2 \cdot |g_-(t) + \lambda_f g_+(t)|^2$$

and after some trigonometry and assuming CPV in mixing is absent ( $|q/p| = 1$ )

$$\left| \langle f | H | B^0(t) \rangle \right|^2 = e^{-t/\tau_B} \cdot |A_f|^2 \frac{1+|\lambda_f|^2}{2} [1 + C_f \cos(\Delta m_d t) - S_f \sin(\Delta m_d t)]$$

$$\left| \langle f | H | \bar{B}^0(t) \rangle \right|^2 = e^{-t/\tau_B} \cdot |A_f|^2 \frac{1+|\lambda_f|^2}{2} [1 - C_f \cos(\Delta m_d t) + S_f \sin(\Delta m_d t)]$$

CP observables

$$C_f = \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2}$$

$$S_f = \frac{2 \operatorname{Im}[\lambda_f]}{1 + |\lambda_f|^2}$$

# Case of $CP$ Eigenstate

- ★ Condition for  $CP$  invariance :  $\left| \langle f_{CP} | H | B^0(t) \rangle \right|^2 = \left| \langle f_{CP} | H | \bar{B}^0(t) \rangle \right|^2, \quad \forall t$
- ★ Definition of “ $CP$  parameter” : 
$$\lambda_{f_{CP}} = \eta_{f_{CP}} \frac{q}{p} \cdot \frac{\bar{A}_{f_{CP}}}{A_{f_{CP}}} \quad \text{← decay amplitude ratio}$$

*CP eigenvalue*

$\approx e^{-2i\beta}$
- ★  $CP$  invariance requires :
 

Note that the  $q/p$  parameter is **not** an **observable**, because its argument depends on phase conventions, but its **modulus**  $|q/p|$  is !

$\rightarrow |q/p| = 1 \quad \{|B_L\rangle, |B_H\rangle\} = \{|B_{CP=+1}\rangle, |B_{CP=-1}\rangle\}$   
 $\rightarrow |\bar{A}_{f_{CP}} / A_{f_{CP}}| = 1 \quad \left. \begin{array}{l} \\ \end{array} \right\} \Rightarrow \lambda_{f_{CP}} = \pm 1$   
 $\rightarrow \text{Im } \lambda_{f_{CP}} = 0$
- ★ Classification of  $CP$  violation :
 

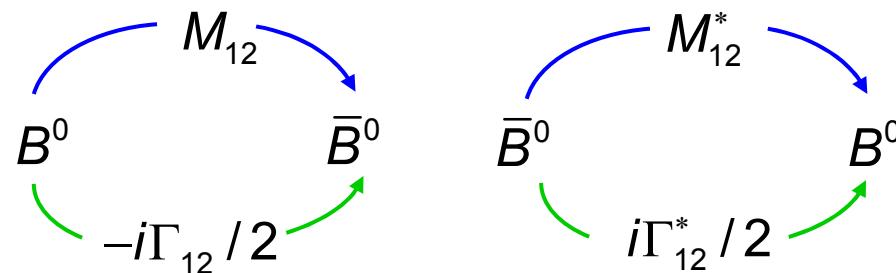
<b>CP-violating phenomena</b>	☀ $CP$ violation in mixing (“indirect”) : $ q/p  \neq 1$
	☀ $CP$ violation in the decay (“direct”) : $ \lambda_{f_{CP}}  \neq 1$
	☀ $CP$ violation in interference between mixing and decay : $\text{Im } \lambda_{f_{CP}} \neq 0$

# CP Violation in Mixing (1<sup>st</sup> type)

The condition for CP conservation in mixing is:

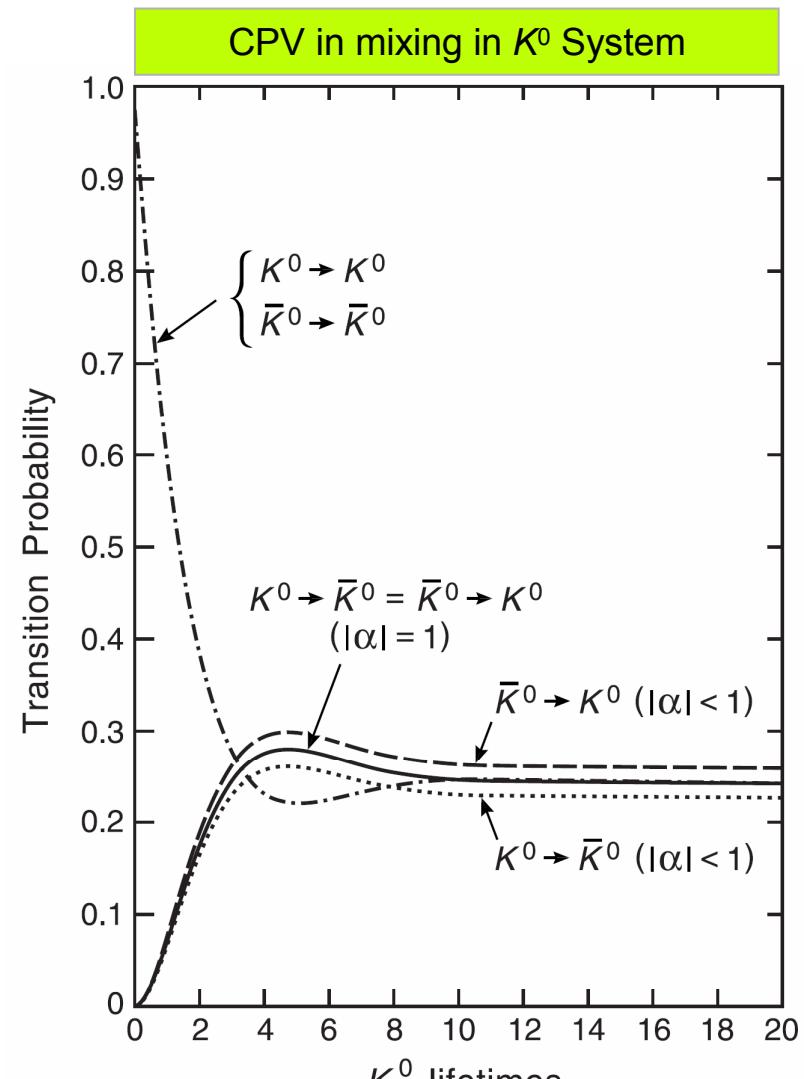
$$\text{Im} [M_{12} \Gamma_{12}^*] = 0$$

CP violation in mixing can be understood as being due to **interference** between the set of amplitudes with **virtual intermediate states** and the set of amplitudes with **on-shell intermediate states**



Indirect CP violation is small if either

- ★ the two amplitudes are **almost relatively real**
- ★ **one** of the two amplitudes **is small**



Small CPV in mixing in  $B^0$  System

# $CP$ Violation in the Decay (2<sup>nd</sup> type)

★  $CP$ -conjugated amplitudes :  $\begin{cases} A_f = A(B \rightarrow f) \\ \bar{A}_f = A(\bar{B} \rightarrow \bar{f}) \end{cases}$  ...  $CP$  invariance implies :  $|A_f| = |\bar{A}_f|$

**define**  $\rightarrow A_{CP} \equiv \frac{\Gamma(\bar{B} \rightarrow \bar{f}) - \Gamma(B \rightarrow f)}{\Gamma(\bar{B} \rightarrow \bar{f}) + \Gamma(B \rightarrow f)} = \frac{|\bar{A}_f|^2 - |A_f|^2}{|\bar{A}_f|^2 + |A_f|^2}$

**CP**  $\rightarrow A_f = \sum_j a_j \cdot e^{i\theta_j} e^{i\phi_j}$        $\bar{A}_f = e^{-2i(\xi_B - \xi_f)} \sum_j a_j \cdot e^{i\theta_j} e^{-i\phi_j}$

$\left\{ \begin{array}{ll} \phi_j & \text{alters sign under } CP \\ & (\text{weak phase}) \\ \theta_j & \text{CP invariant} \\ & (\text{strong phase}) \end{array} \right.$



$$A_{CP} = \frac{\sum_{ij} a_i a_j \cdot \sin(\theta_i - \theta_j) \cdot \sin(\phi_i - \phi_j)}{\sum_{ij} a_i a_j \cdot \cos(\theta_i - \theta_j) \cdot \cos(\phi_i - \phi_j)}$$

Direct  $CP$  violation requires at least two amplitudes with different weak *and* strong phases

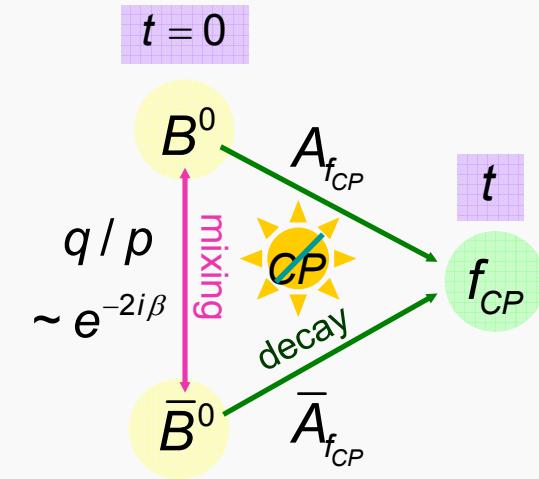
# Mixing-Induced $CP$ Violation (3<sup>rd</sup> type)

- ★  $CP$  Violation due to the interference of decays with and without mixing

$$\lambda_{f_{CP}} \neq \pm 1 \Leftrightarrow \text{Prob}(\bar{B}^0(t) \rightarrow f_{CP}) \neq \text{Prob}(B^0(t) \rightarrow f_{CP})$$

- ★ Time-dependent asymmetry observable

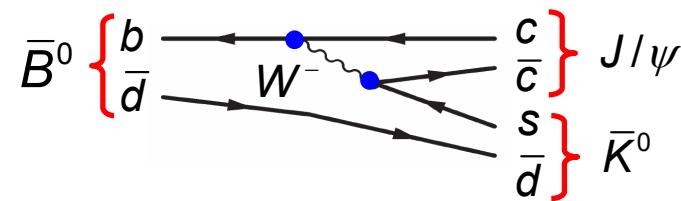
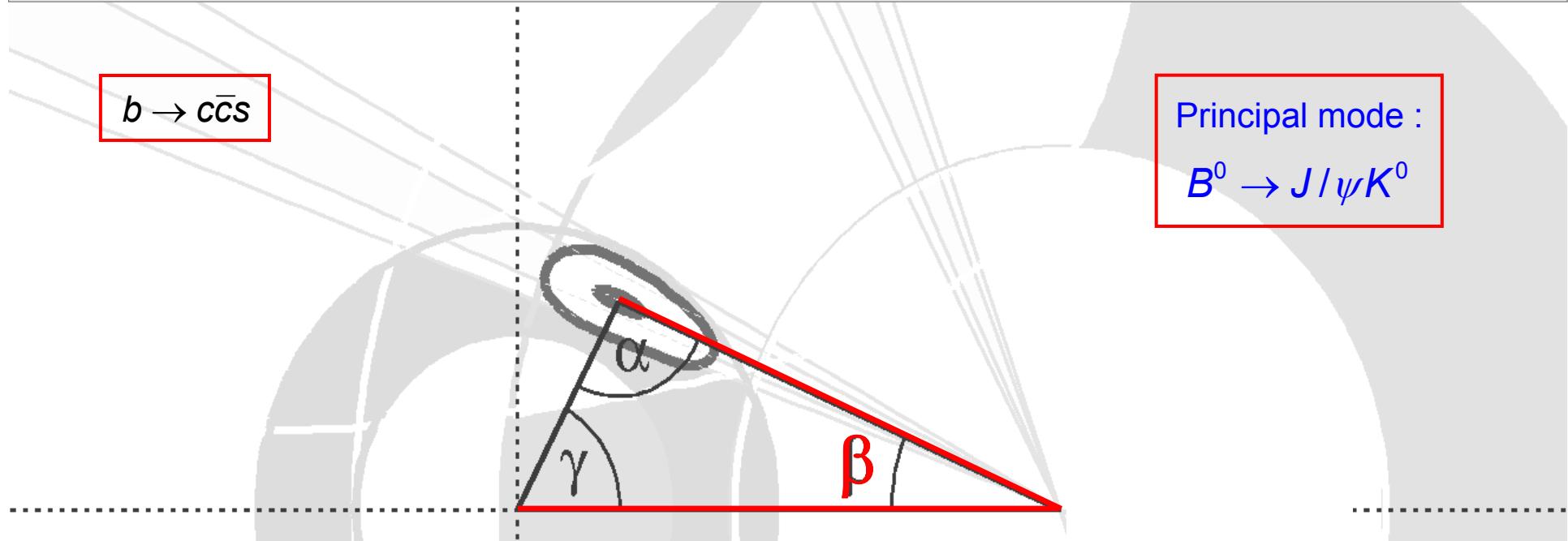
$$\begin{aligned} A_{f_{CP}}(t) &= \frac{\Gamma(\bar{B}^0(t) \rightarrow f_{CP}) - \Gamma(B^0(t) \rightarrow f_{CP})}{\Gamma(\bar{B}^0(t) \rightarrow f_{CP}) + \Gamma(B^0(t) \rightarrow f_{CP})} \\ &= S_{f_{CP}} \sin(\Delta m_d t) - C_{f_{CP}} \cos(\Delta m_d t) \end{aligned}$$



recall:

$$\boxed{\begin{aligned} C_f &= \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2} \\ S_f &= \frac{2 \text{Im}[\lambda_f]}{1 + |\lambda_f|^2} \end{aligned}}$$

# $\sin(2\beta)$

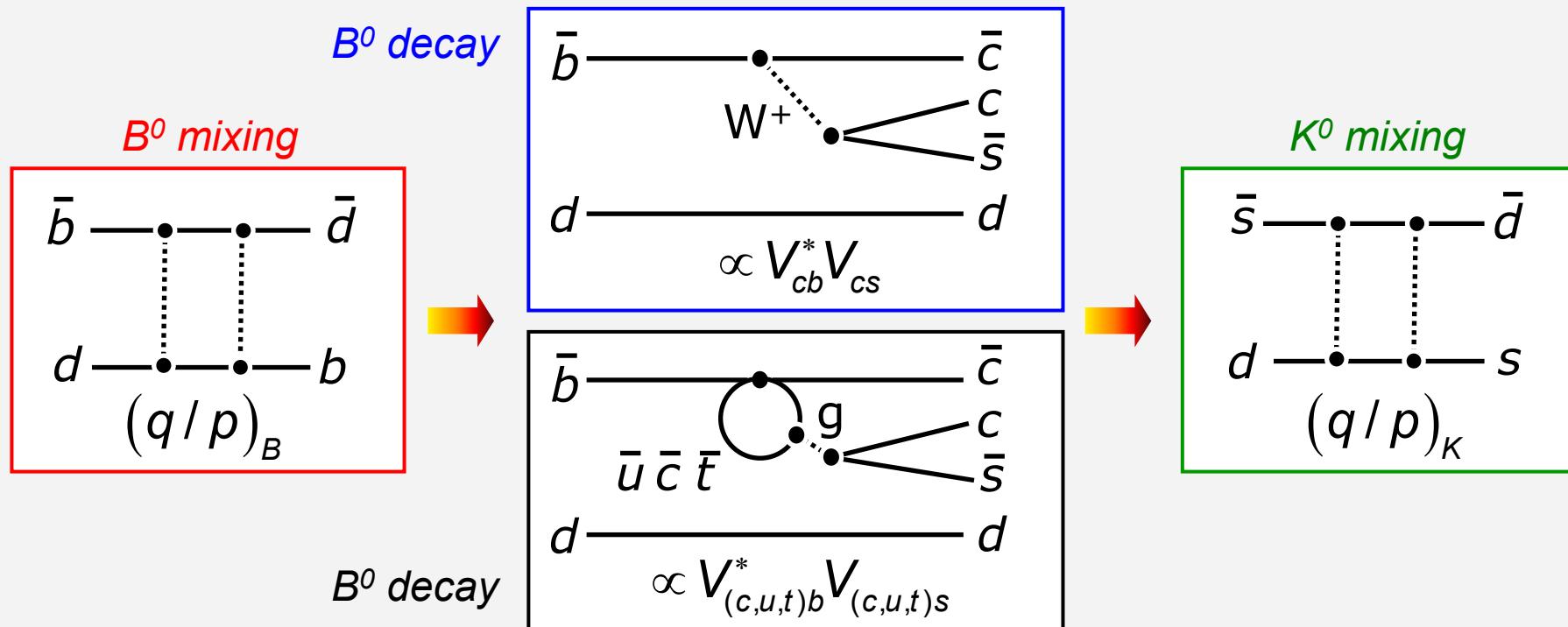


Tree : dominant

$$\begin{aligned} &\propto V_{cb} V_{cs}^* \\ &\propto \lambda^2 \end{aligned}$$

# The Golden Channel: $B^0, \bar{B}^0 \rightarrow J/\psi K^0$

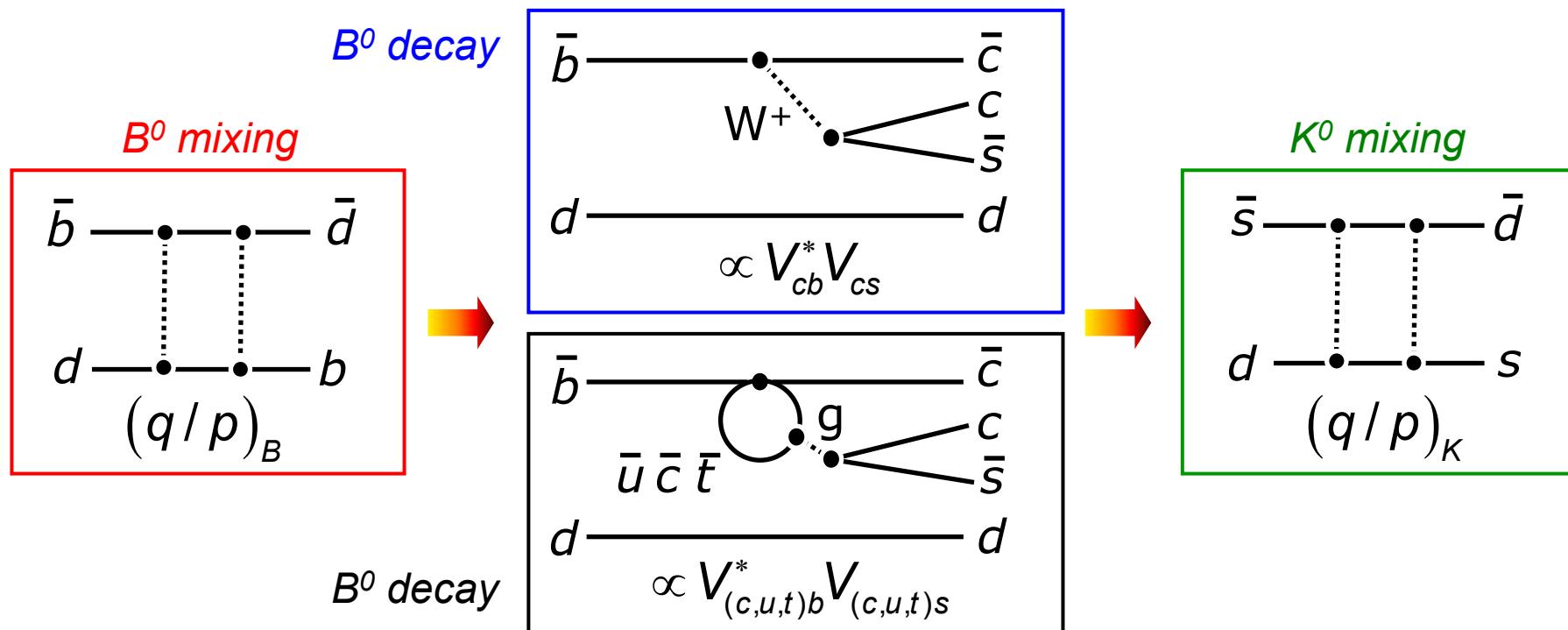
Leading tree :  $\frac{q}{p} \frac{\bar{A}}{A} = \eta_{CP} \left( \frac{q}{p} \right)_B \frac{V_{cb} V_{cs}^*}{V_{cb}^* V_{cs}} \left( \frac{q}{p} \right)_K = -e^{-2i\beta}$



Penguin :  $V_{cb} V_{cs}^*$  real ,  $V_{tb} V_{ts}^*$  real ,  $V_{ub} V_{us}^* \rightarrow \gamma (\propto A \cdot \lambda^4)$

# The Golden Channel: $B^0, \bar{B}^0 \rightarrow J/\psi K^0$

Leading tree :  $\frac{q}{p} \frac{\bar{A}}{A} = \eta_{CP} \left( \frac{q}{p} \right)_B \frac{V_{cb} V_{cs}^*}{V_{cb}^* V_{cs}} \left( \frac{q}{p} \right)_K = -e^{-2i\beta}$



Single weak phase:

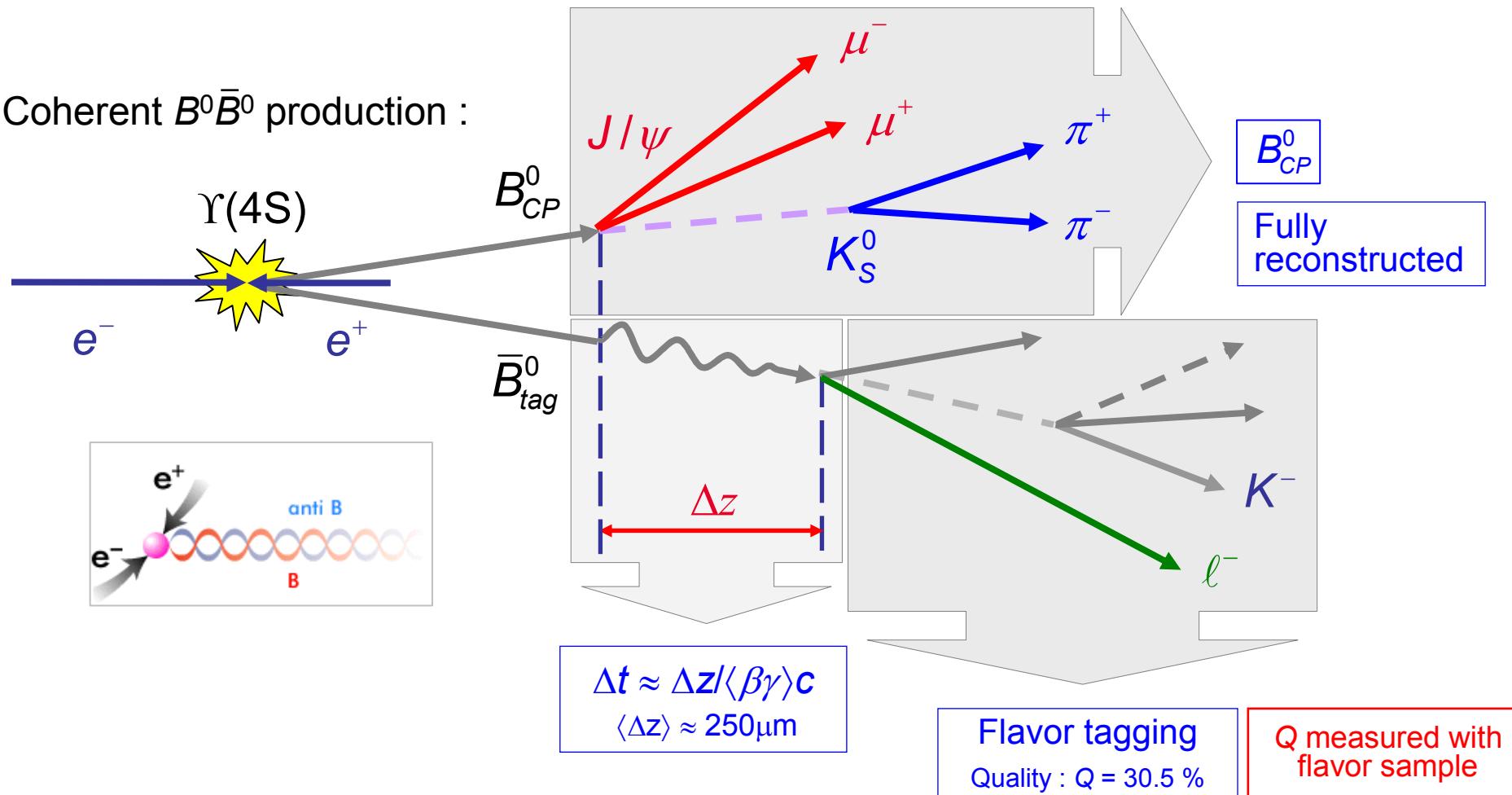
- ☒ clean CP phase
- ☒ no direct CPV:



$$A_{J/\psi K_{S,L}^0}(t) = -\eta_{J/\psi K_{S,L}^0} \cdot \sin(2\beta) \cdot \sin(\Delta m_{B_d} t)$$

# Experimental Technique

Coherent  $B^0\bar{B}^0$  production :



# Data Sample



**Nov 1999- July 2004 data**  
 $213 \text{ fb}^{-1}$  on-peak – 227 million BB pairs

**energy-substituted mass**

$$m_{\text{ES}} \equiv \sqrt{(s/4) - p_B^{*2}}$$

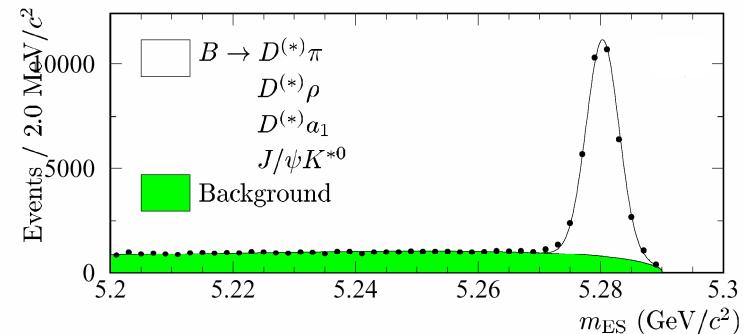
**energy difference**

$$\Delta E \equiv \sqrt{s}/2 - E_B^*$$

Samples of  $B$  decays to flavor-specific final states

$$B^0 \rightarrow D^{(*)-} \pi^+ / \rho^+ / a_1^+$$

$$B^0 \rightarrow J/\psi K^{*0} (\rightarrow K^+ \pi^-)$$



Samples of  $B$  decays to  $CP$  eigenstates with charmonium

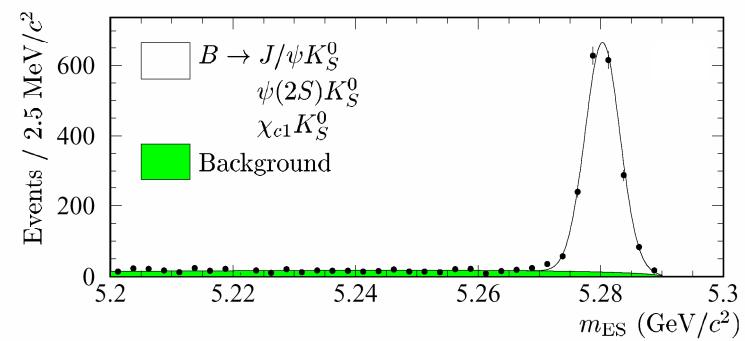
$$B_{CP}^0 \rightarrow \eta_c (\rightarrow K_S^0 K^\pm \pi^\mp) K_S^0 \quad (\text{not shown})$$

$$B_{CP}^0 \rightarrow J/\psi (\rightarrow \ell^+ \ell^-) K_S^0$$

**CP = -1**     $B_{CP}^0 \rightarrow J/\psi K_S^0 (\rightarrow \pi^0 \pi^0)$

$$B_{CP}^0 \rightarrow \psi(2S) (\rightarrow \ell^+ \ell^-, J/\psi \pi\pi) K_S^0$$

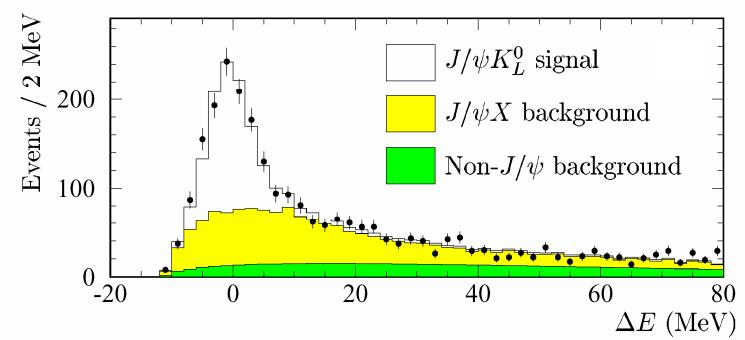
$$B_{CP}^0 \rightarrow \chi_{c1} (\rightarrow J/\psi \gamma) K_S^0$$



**CP = +1**     $B_{CP}^0 \rightarrow J/\psi K_L^0$



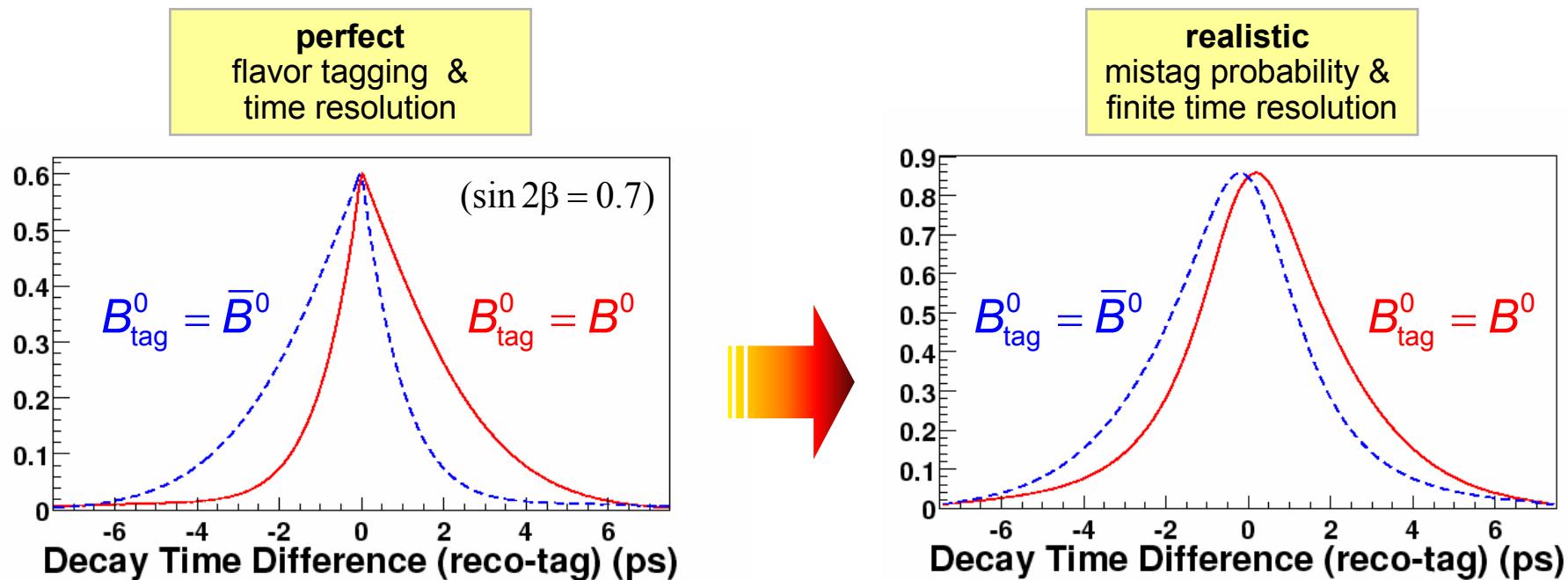
**mixed CP**     $B_{CP}^0 \rightarrow J/\psi K_{CP}^{*0} (\rightarrow K_S^0 \pi^0)$     **requires angular analysis (not shown)**



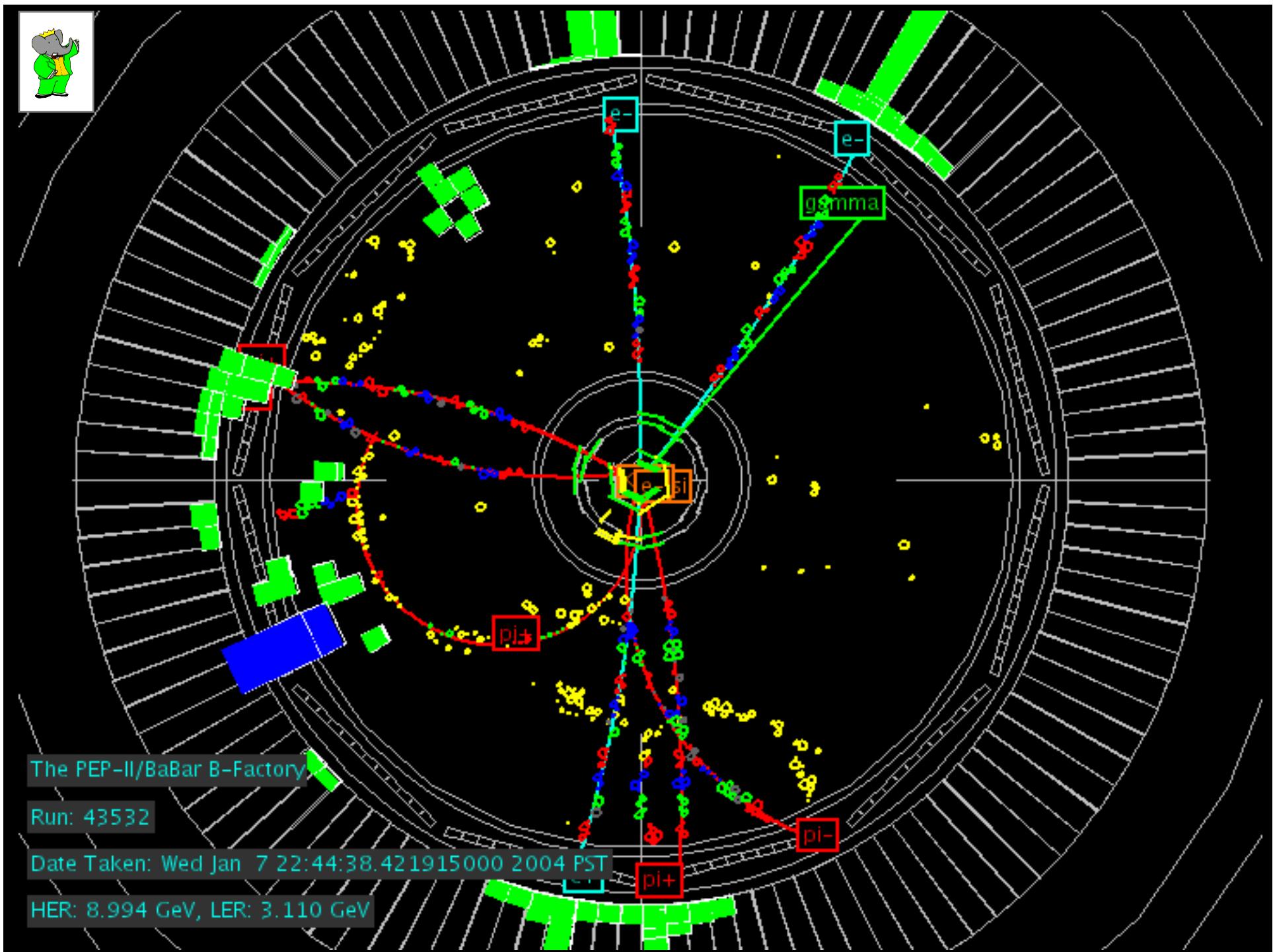
# CPV Experimentally: Mistagging and Resolution

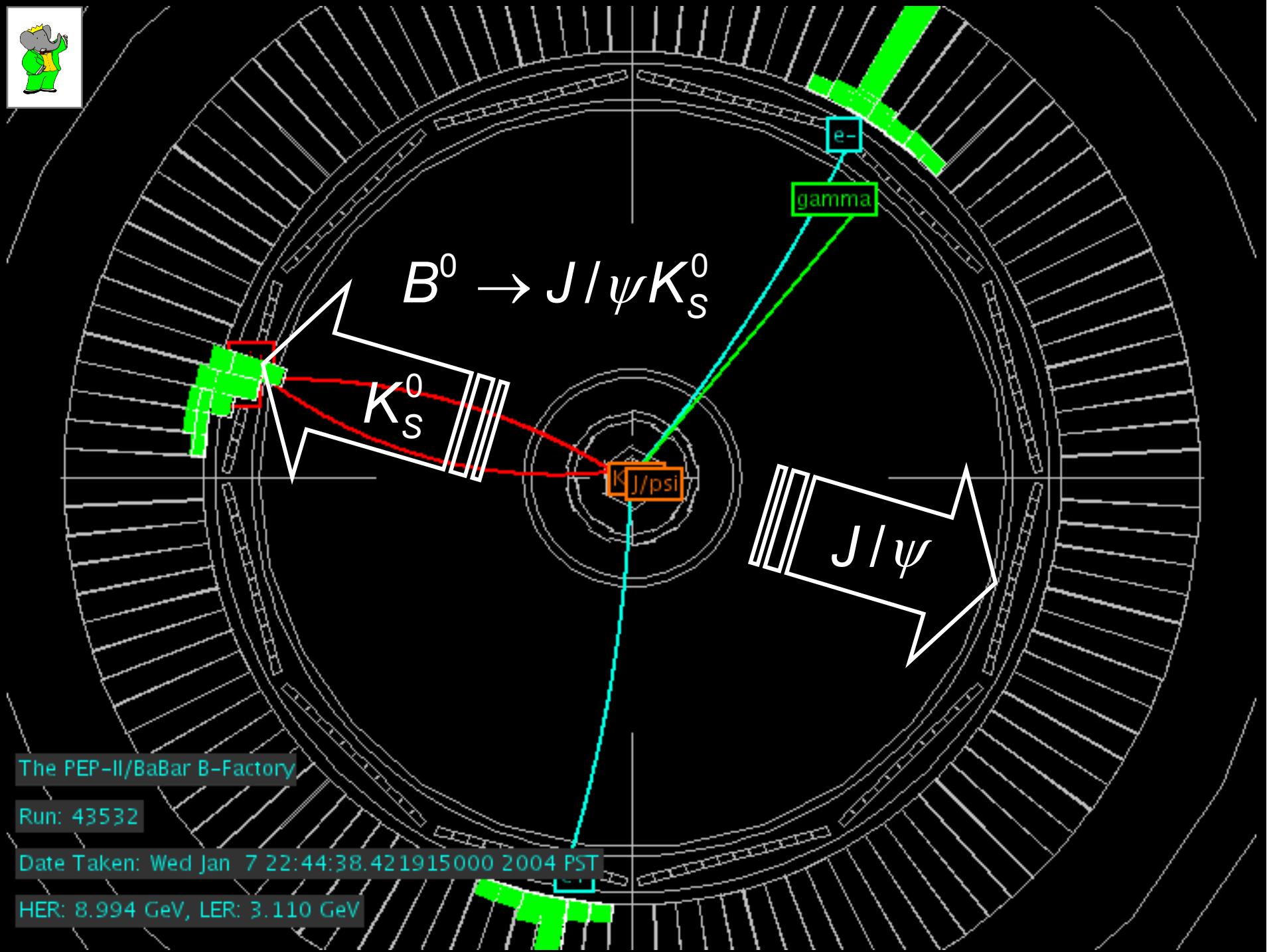
Two categories of events:

- $B_{\text{tag}}^0 = B^0$  (+) : the “other” (=tag)  $B$  is tagged as a  $B^0$
- $B_{\text{tag}}^0 = \bar{B}^0$  (-) : the “other” (=tag)  $B$  is tagged as a  $\bar{B}^0$



$$A_{J/\psi K_{S,L}^0}(t) = -\eta_{J/\psi K_{S,L}^0} (1 - 2\omega) \cdot \sin(2\beta) \cdot \sin(\Delta m_{B_d} t') \otimes R(t' - t)$$







Recoil *B*

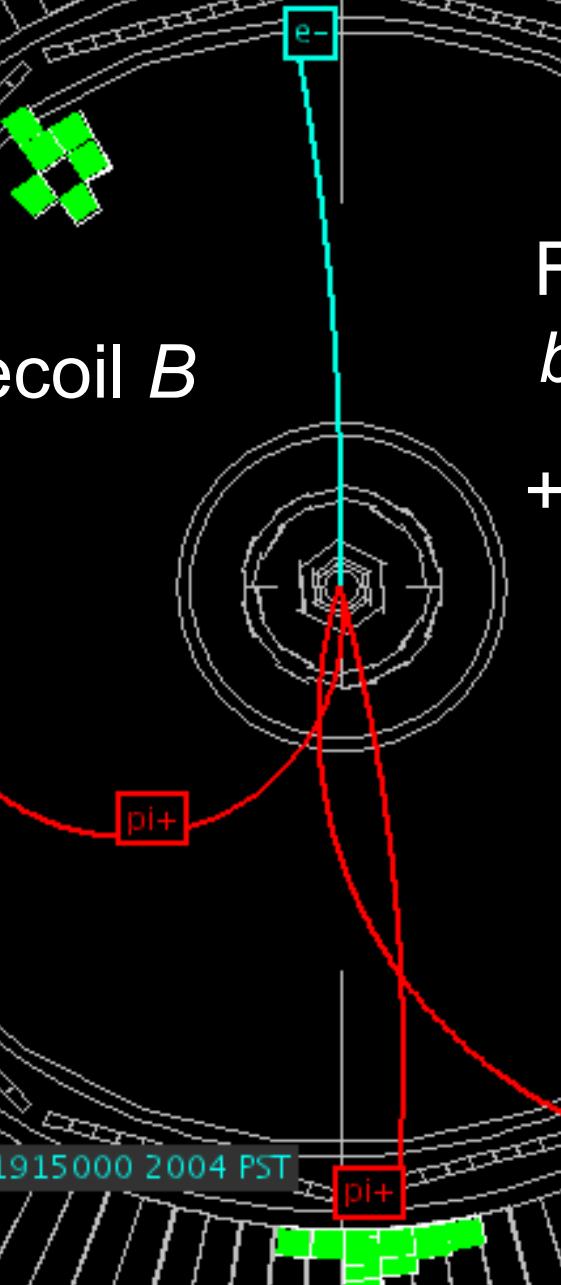
Flavor tag :  
 $b \rightarrow c e^- \bar{\nu}_e$   
+ vertex separation  
= time difference

The PEP-II/BaBar B-Factory

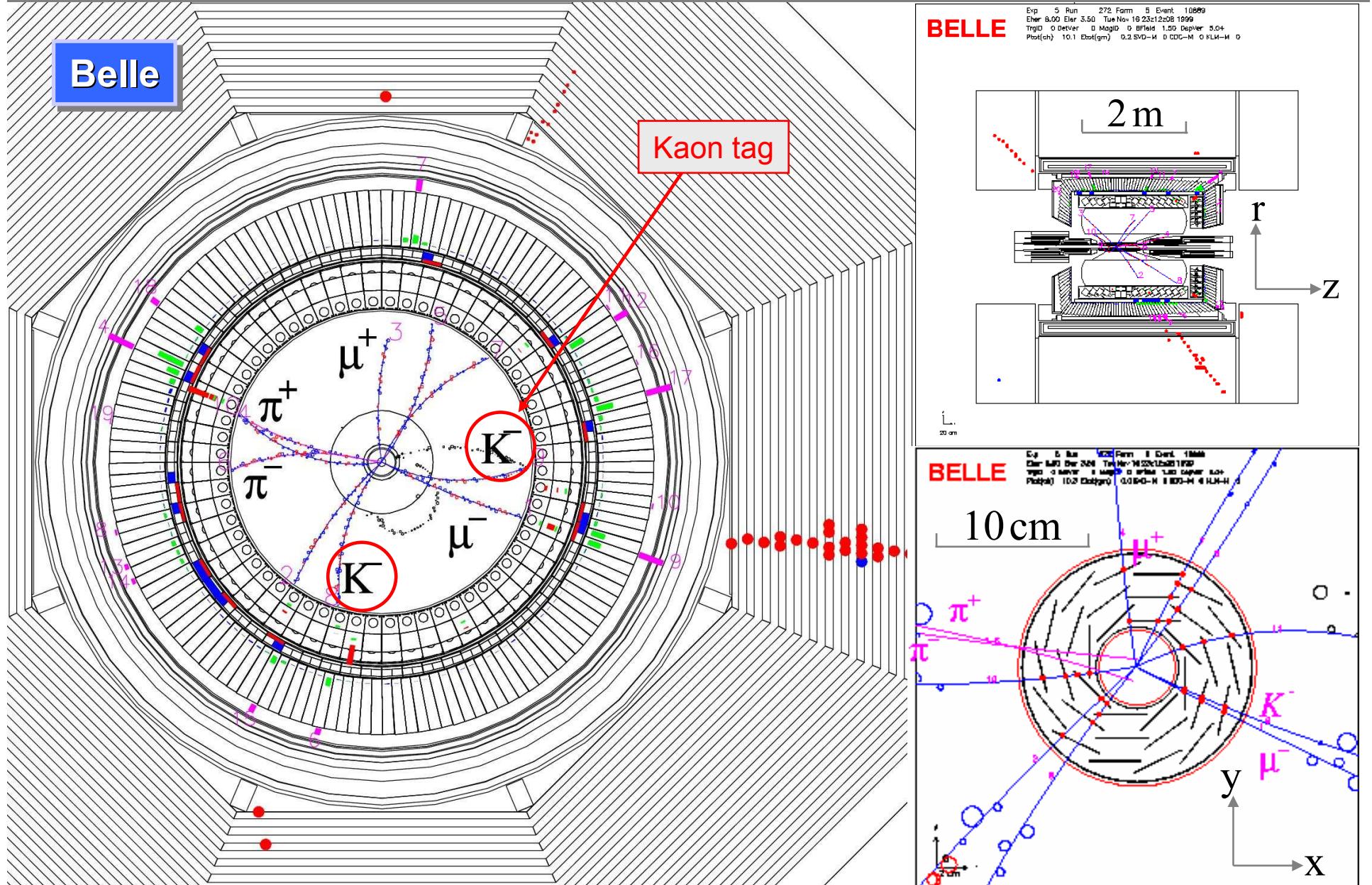
Run: 43532

Date Taken: Wed Jan 7 22:44:38.421915000 2004 PST

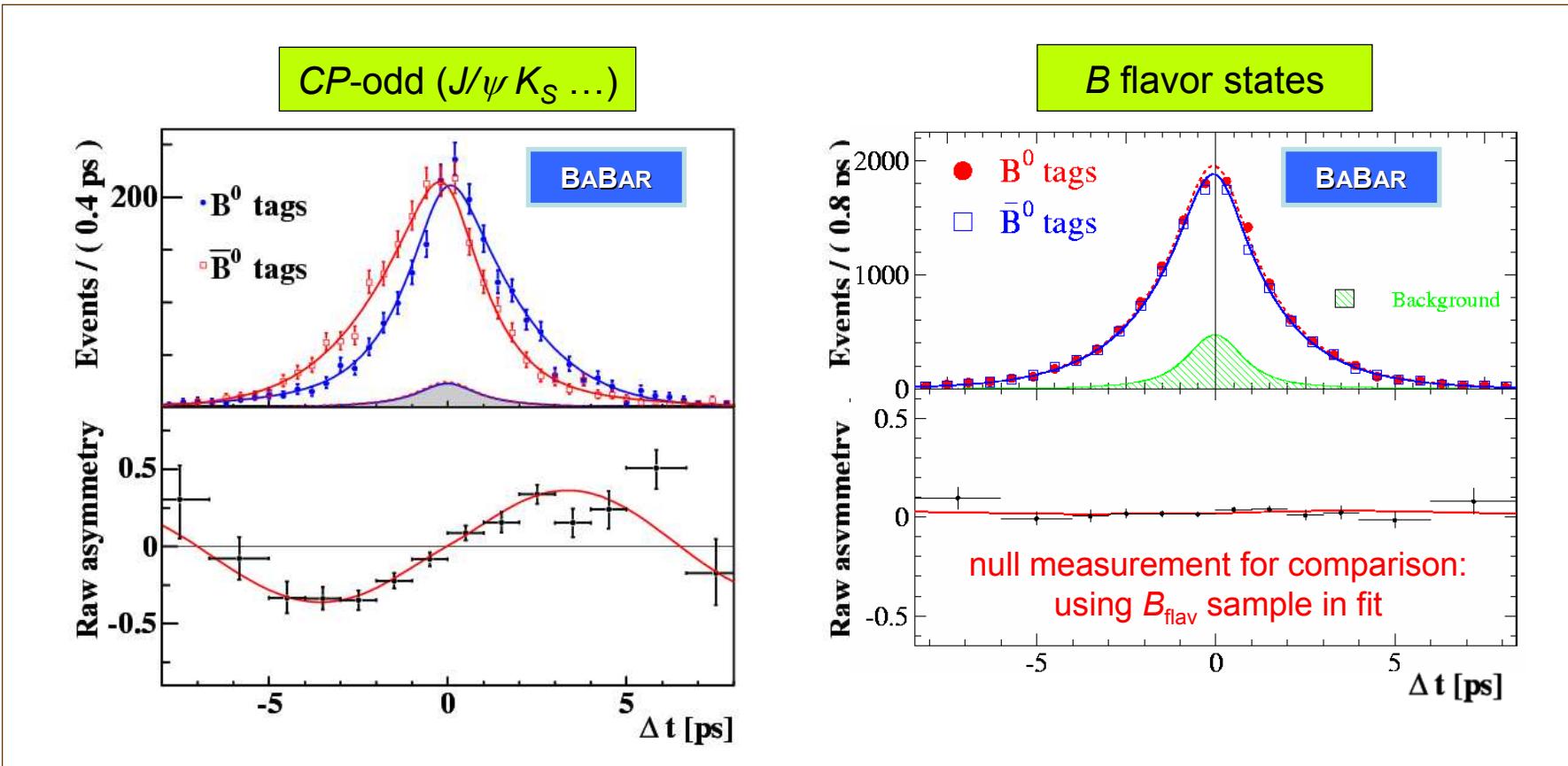
HER: 8.994 GeV, LER: 3.110 GeV



# A $B \rightarrow \psi(\rightarrow \mu^+ \mu^-) K_S$ Event as seen by Belle



# $\sin(2\beta)$ with $B^0$ to charmonium



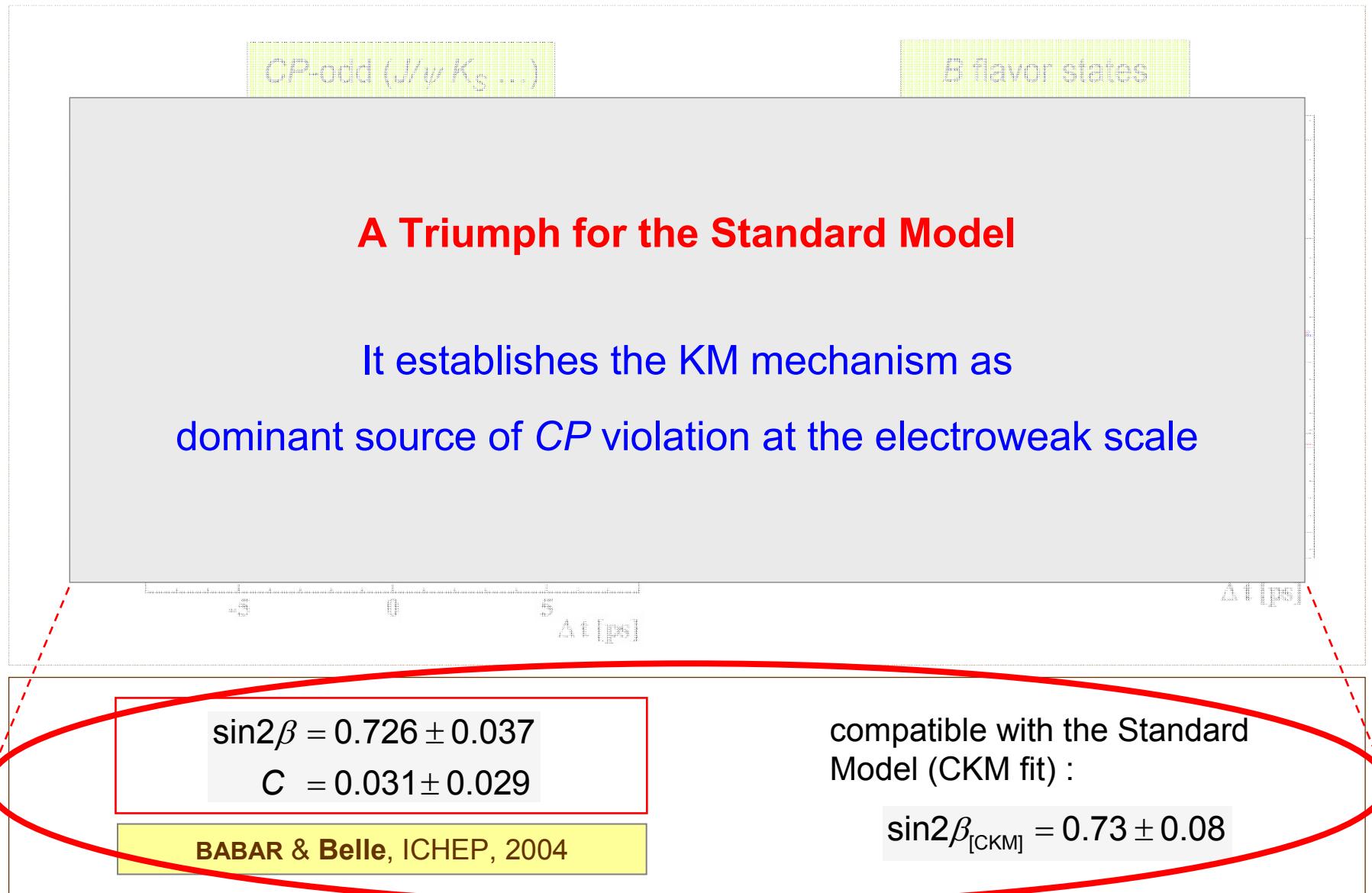
$$\begin{aligned}\sin 2\beta &= 0.726 \pm 0.037 \\ C &= 0.031 \pm 0.029\end{aligned}$$

BABAR & Belle, ICHEP, 2004

compatible with the Standard Model (CKM fit) :

$$\sin 2\beta_{[\text{CKM}]} = 0.73 \pm 0.08$$

# $\sin(2\beta)$ with $B^0$ to charmonium

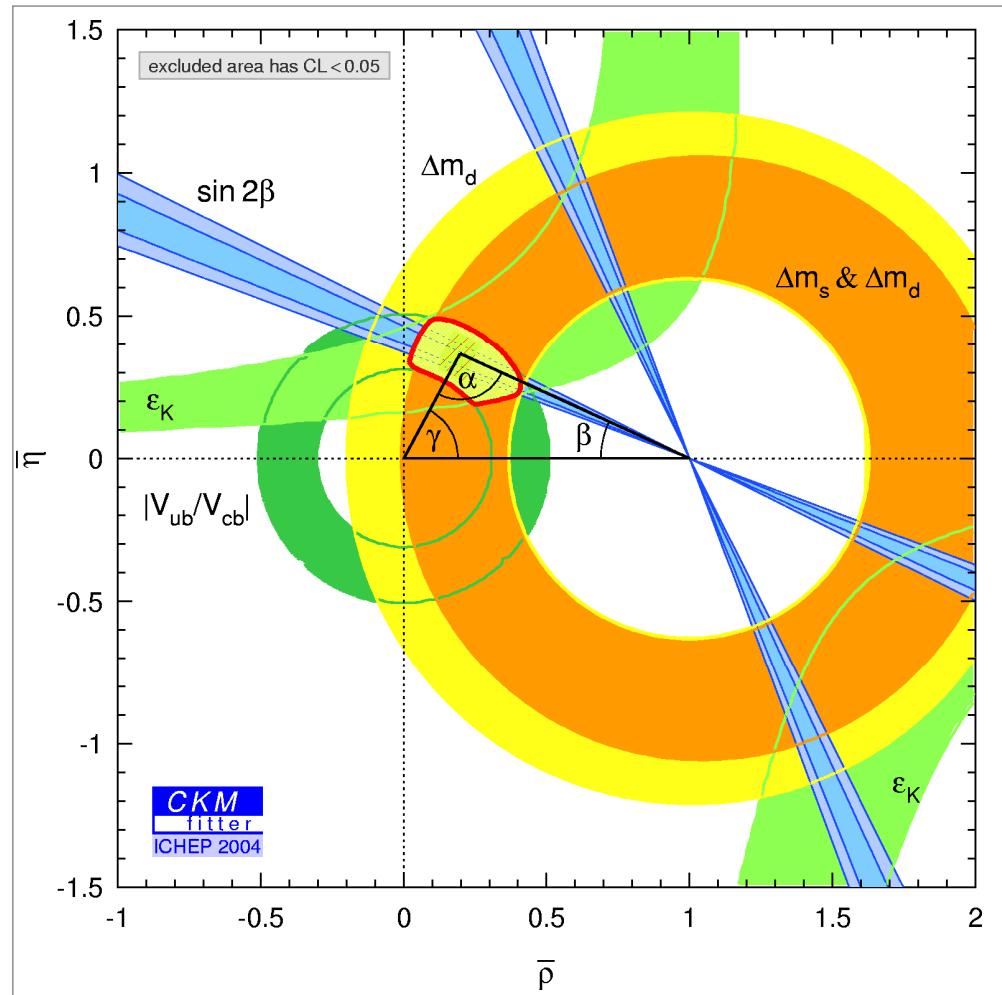


# Comparison with CKM Fit



Precision measurement :

$$\beta = [23.3 \begin{array}{l} +1.6 \\ -1.5 \end{array}]^\circ$$



**CKMfitter**, EPJ C41, 1-131,2005 [hep-ph/0406184 (2004)]  
**UTfit**, hep-ph/0408079 (2004)  
 and others !

# Comparison with CKM Fit



Precision measurement :

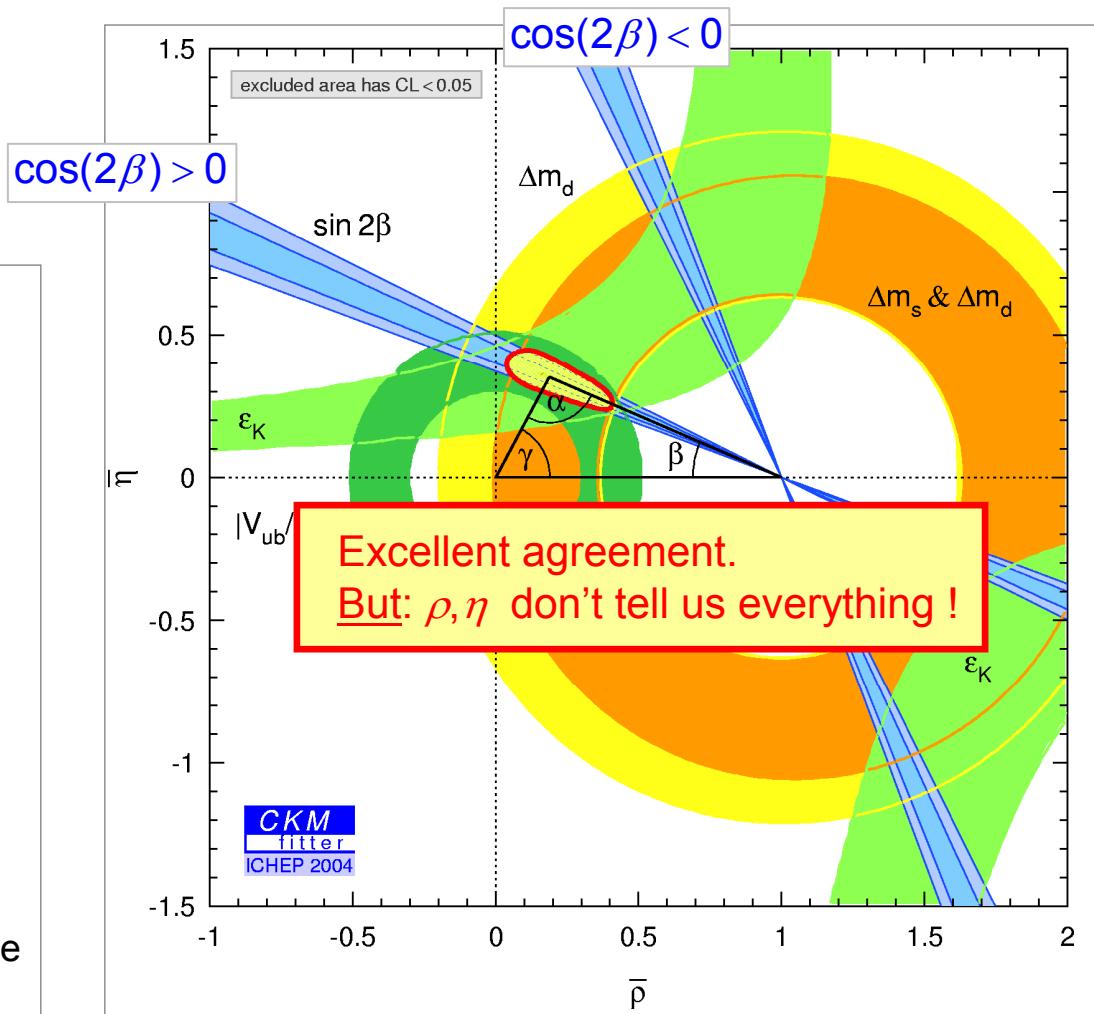
$$\beta = [23.3^{+1.6}_{-1.5}]^\circ$$

- Measurement of the sign of  $\cos 2\beta$  is a direct test of the SM : if SM  $\Rightarrow \cos 2\beta > 0$
- $\cos 2\beta$  accessible through time-dependent angular correlations in

$$B^0 \rightarrow J/\psi K^{*0} \left( K^{*0} \rightarrow K_S^0 \pi^0 \right)$$

- But :
  - with one ambiguity on sign ... ☺
  - induced by ambiguity on strong phase
- Ambiguity can be lifted with analysis of interference pattern between S and P wave in  $(K\pi)$  system.

→  $\cos(2\beta) > 0$  at 89% C.L.



CKMfitter, EPJ C41, 1-131,2005 [hep-ph/0406184 (2004)]  
 UTfit, hep-ph/0408079 (2004)  
 and others !

# Measurement of $\sin 2\beta$ : Trivial Assumptions ???

- ★ All  $B \rightarrow$  charmonium  $K_S$  modes measure the same  $\sin 2\beta$
- ★  $J/\psi K_S$  and  $J/\psi K_L$  measure the same  $\sin 2\beta$
- ★ Direct  $CP$  violation in  $B^0 \rightarrow J/\psi K^0$  is negligible
- ★ Physical  $B$  mesons have equal lifetimes
- ★  $CP$  violation in  $B$  mixing is negligible
- ★  $CPT$  is an exact symmetry

All conserved OR  
 $T, CPT$  violated ( $z=0, |q/p|=1$ )

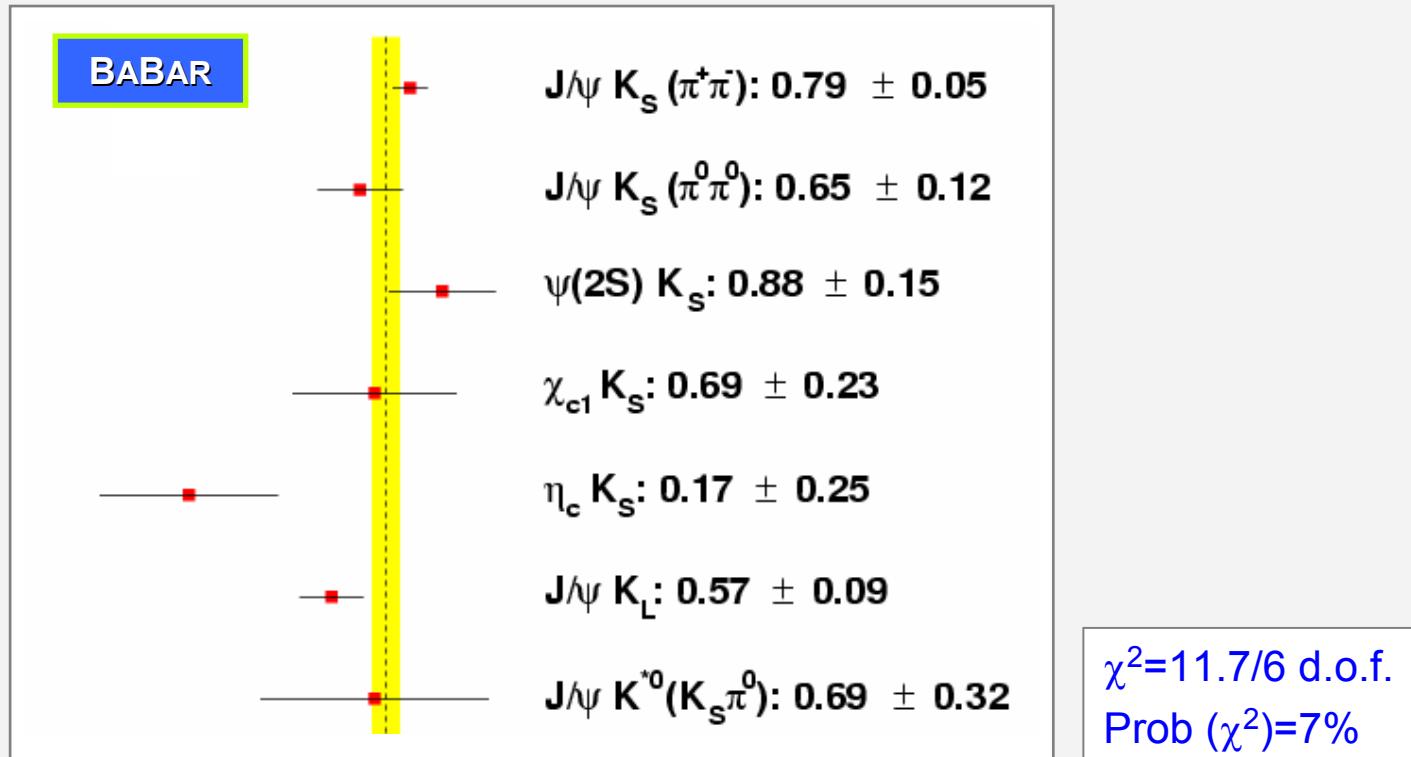
Standard Model

( $z=0, |q/p|=1 = (2.5-6.5) \times 10^{-4}$ )

# Trivial Assumptions ?

Example: all  $B \rightarrow$  charmonium  $K_S$  modes measure the same  $\sin 2\beta$

Compare  $\sin(2\beta)$  among charmonium modes



→ reasonable consistency among the various charmonium modes

# Trivial Assumptions ?

Example: all  $B \rightarrow$  charmonium  $K_s$  modes measure the same  $\sin 2\beta$

Compare  $\sin(2\beta)$  among charmonium modes



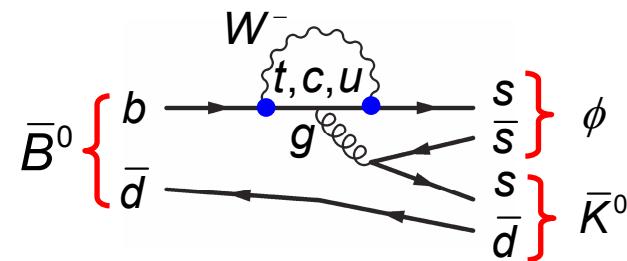
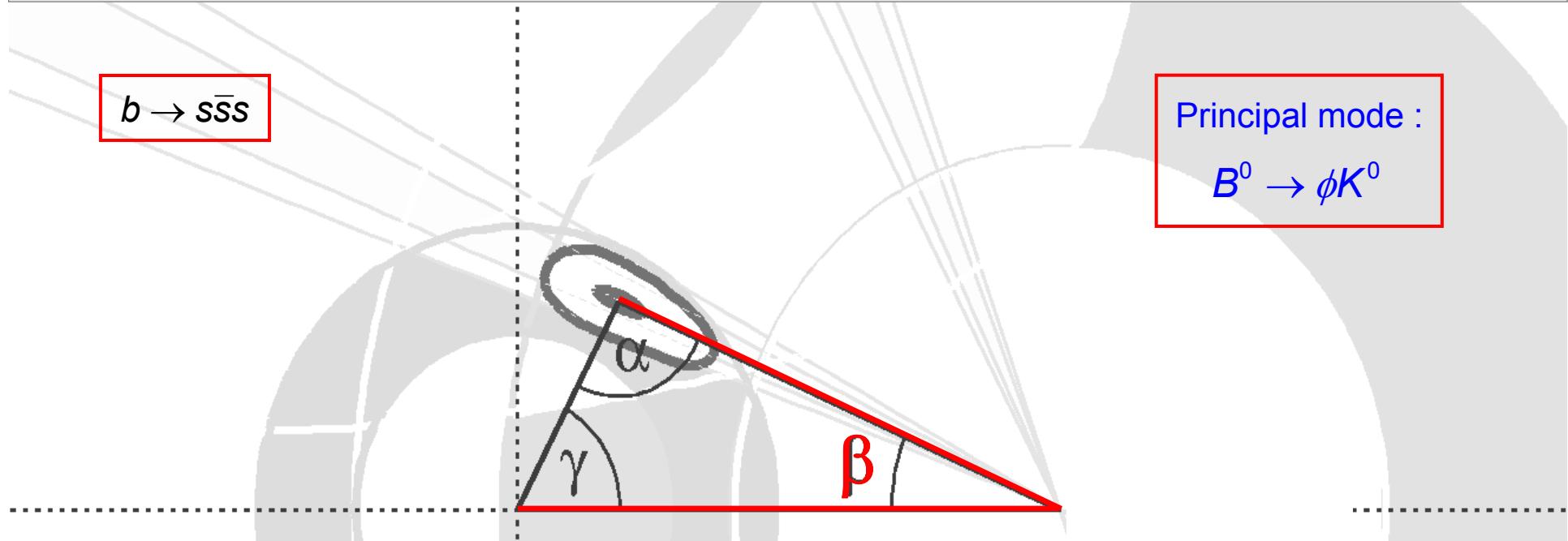
## How can we Further Challenge the Standard Model ?



$\chi^2 = 11.7 / 6$  d.o.f.  
Prob ( $\chi^2$ ) = 7%

→ reasonable consistency among the various charmonium modes

# Use the *Penguins* to Measure $\sin(2\beta_{\text{eff}})$



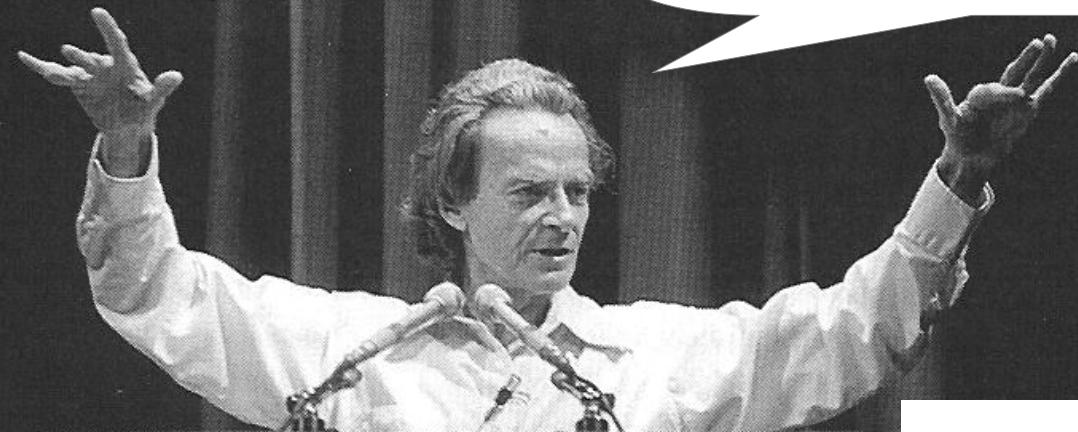
Penguin : dominant

$$\begin{aligned} &\propto V_{tb} V_{ts}^* \\ &\propto \lambda^2 \end{aligned}$$

BUT ...

a controversy...

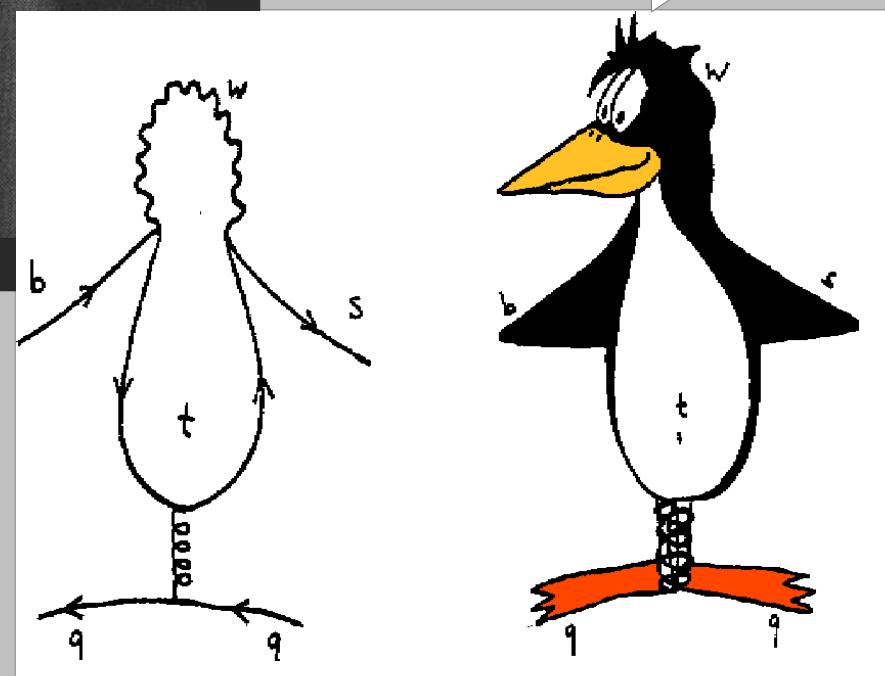
why the hell do you call these  
**Penguin diagrams?**  
They don't look like penguins!



mirror image of Richard Feynman

courtesy: G. Hamel de Monchenault

I've never seen a  
**Feynman diagram**  
that looks like you ☺

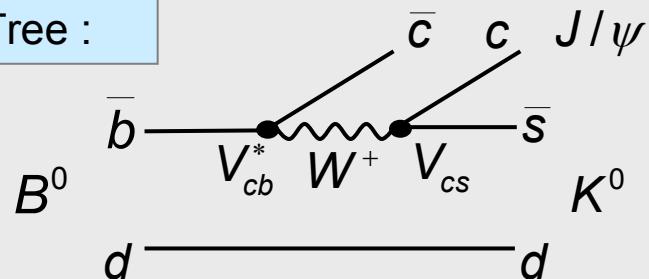


# Confronting Loop and Tree Decays

- ▶  $b \rightarrow c\bar{c}s$  decays are tree and penguin diagrams, with equal dominant weak phases
  - ▶  $b \rightarrow s\bar{s}s$  decays are pure “internal” and “flavor-singlet” penguin diagrams
  - ▶ High virtual mass scales involved: believed to be sensitive to New Physics

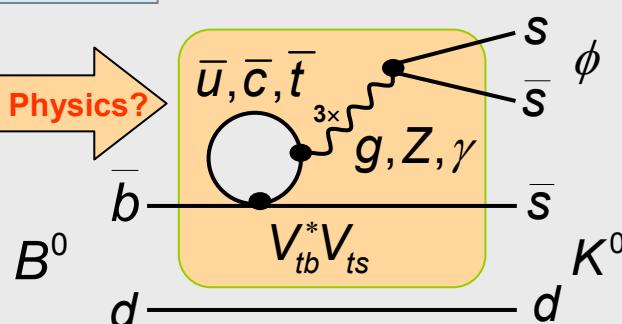
Both decays dominated by single weak phase

## Tree :



# Penguin :

New Physics?



$b \rightarrow c\bar{c}s$

$$\lambda_{J/\psi K^0_{S,L}} = \eta_{J/\psi K^0_{S,L}} \left( \frac{q}{p} \right)_B \cdot \left( \frac{V_{cb} V_{cs}^*}{V_{cb}^* V_{cs}} \right) \cdot \left( \frac{q}{p} \right)_K = \eta_{J/\psi K^0_{S,L}} e^{-2i\beta}$$

$b \rightarrow s\bar{s}s$

$$\lambda_{\phi K_{S,L}^0} = \eta_{\phi K_{S,L}^0} \left( \frac{q}{p} \right)_B \cdot \left( \frac{V_{cb} V_{cs}^*}{V_{cb}^* V_{cs}} \right) \cdot \left( \frac{q}{p} \right)_K = \eta_{\phi K_{S,L}^0} e^{-2i\beta}$$

$$\sin 2\beta \text{ [charmonium]} = \sin 2\beta \text{ [s-pingouin]}$$

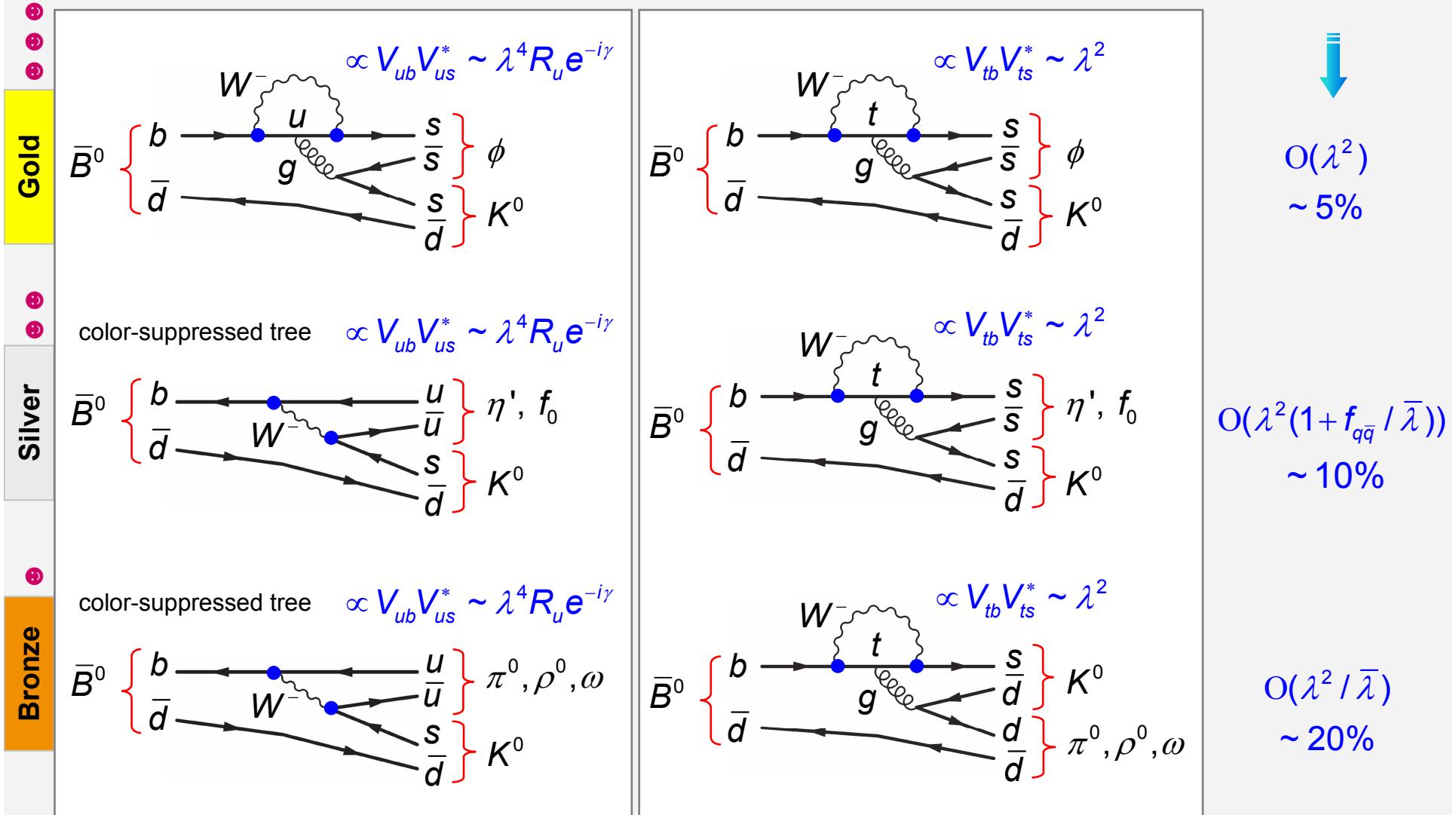
?

(7000 tagged events)	(1500 tagged events)
----------------------	----------------------

# Naïve Classification of the penguins

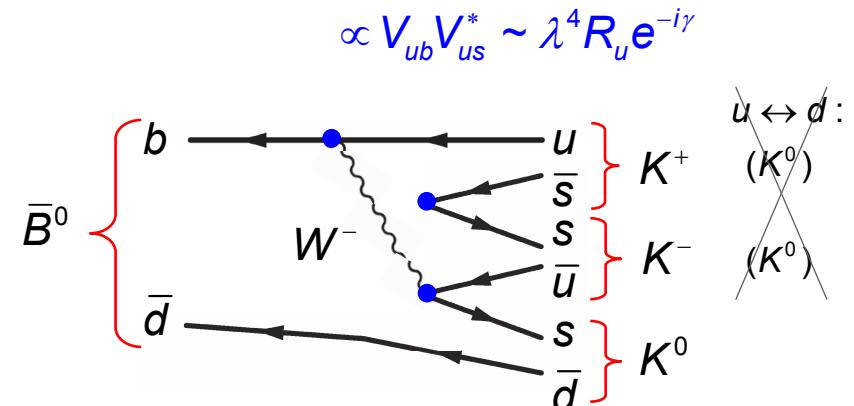
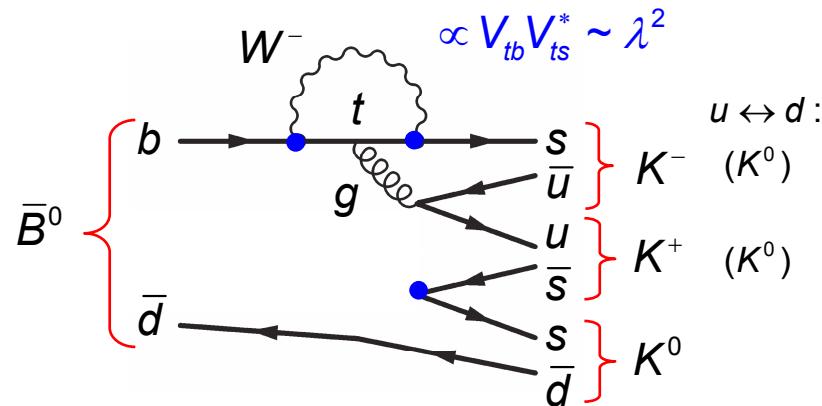
- Can we identify golden, silver and bronze-plated s-penguin modes ?

Naïve (dimensional) uncertainties on  $\sin 2\beta$



# $KKK^0$ Modes

- How about  $B^0 \rightarrow K^+ K^- K^0$  and  $B^0 \rightarrow K_S K_S K_S$  ?



$B^0 \rightarrow K^+ K^- K^0$  bronze and  $B^0 \rightarrow K_S K_S K_S$  seems to be a new golden mode !

- but : « *popup* » ss-bar dynamically disfafored with respect to *popup uu-bar*  
 → is  $B^0 \rightarrow K^+ K^- K^0$  silver ?

- Note that

if NP contributes significantly to CPV in loop decays, we naturally expect it to be different among the modes

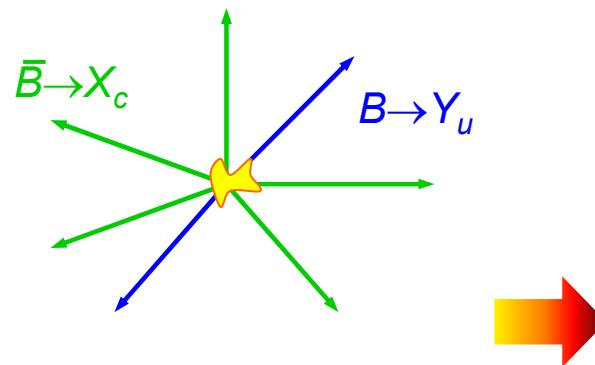
→ averaging only meaningful (if at all) in case of SM

# “Rare Modes” → Dedicated Background Fighting (I)

- ★ Suppress continuum, background using **event shape** variables

$$e^+e^- \rightarrow Y(4s) \rightarrow B\bar{B}$$

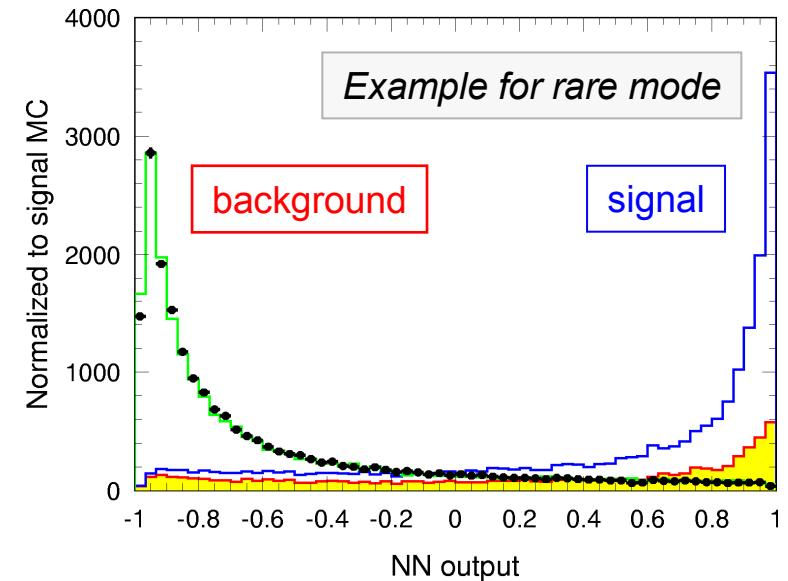
$B$  decay  $\sim$  in rest  $\rightarrow$  event shape spherical



Dominant  $e^+e^- \rightarrow q\bar{q}$   
background is jet-like:



Apply multivariate analyzer techniques:  
**Neural Network** or **Fisher Discriminants**



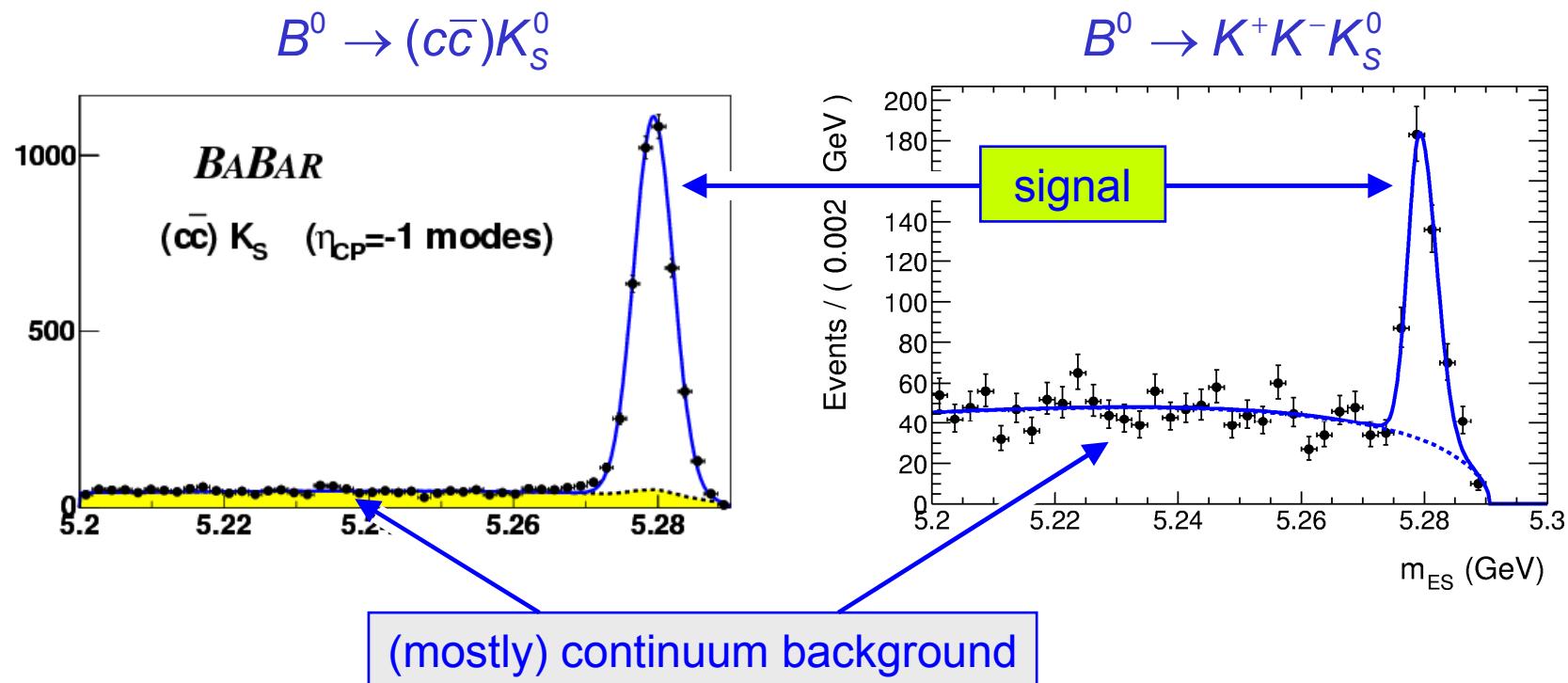
# “Rare Modes” → Dedicated Background Fighting (II)

- ★ Also: kinematic variables:  $m_{\text{ES}}$  and  $\Delta E$

Typical resolution:

$$\sigma(m_{\text{ES}}) \approx 2.5 \text{ MeV}/c^2$$

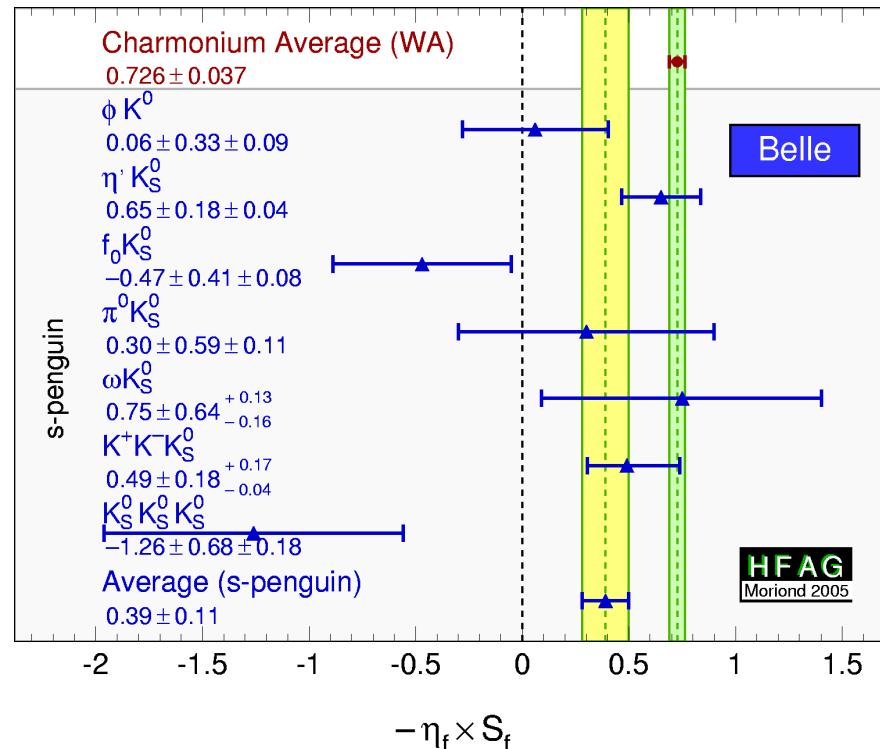
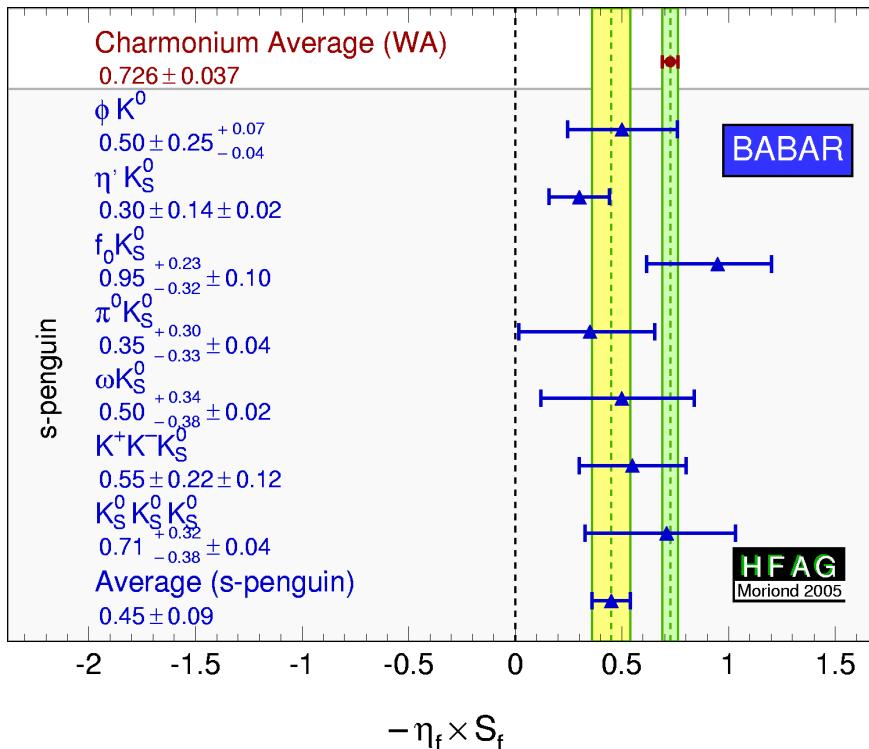
$$\sigma(\Delta E) \approx 25 - 40 \text{ MeV}$$



# s-penguin : comparison BABAR vs. Belle

Moriond 2005

BABAR : s-penguin average at  $2.9\sigma$  from  $\sin 2\beta[c\bar{c}]$  (WA)



Belle : s-penguin average at  $2.9\sigma$  from  $\sin 2\beta[c\bar{c}]$  (WA)

# Confronting Loop and Tree Decays

- Conflict with  $\sin 2\beta_{\text{eff}}$  from s-penguin modes ?

$$\langle \sin 2\beta_{[\text{s-peng}]} \rangle - \sin 2\beta_{[c\bar{c}]} = \underbrace{-0.30 \pm 0.08}_{3.7\sigma}$$

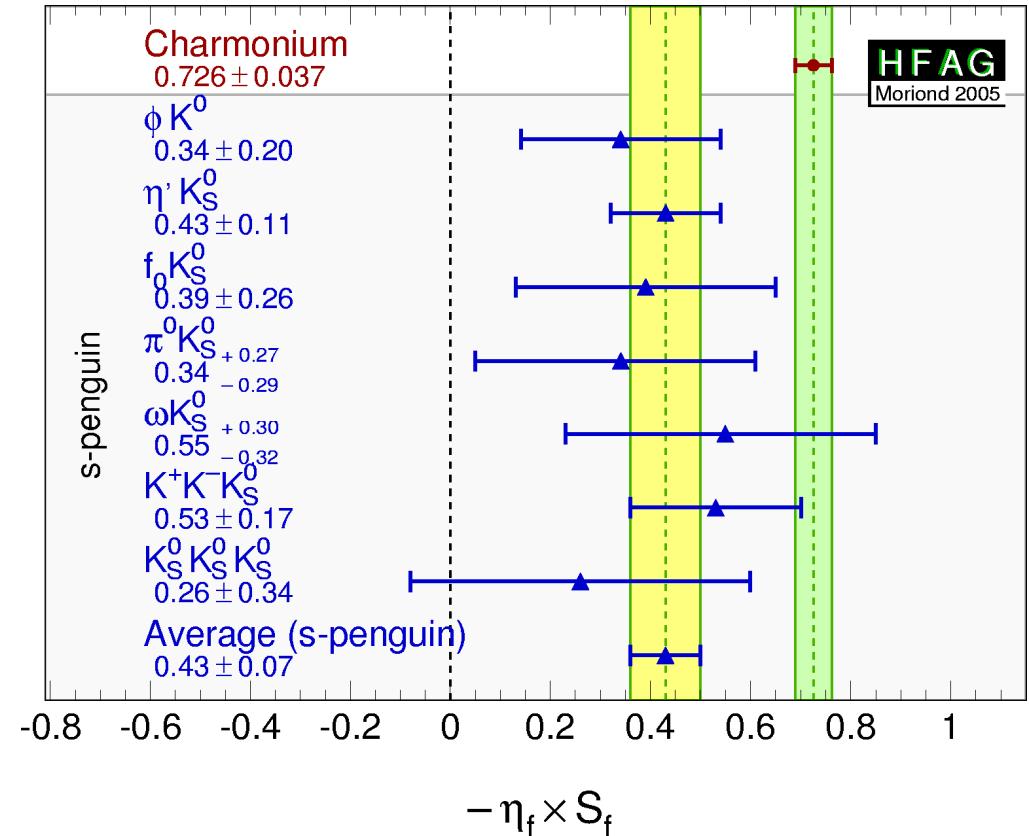
$$\langle \sin 2\beta_{[\phi/\eta'/2K_S]} \rangle - \sin 2\beta_{[c\bar{c}]} = \underbrace{-0.33 \pm 0.10}_{3.3\sigma}$$

Theory uncertainty ?

what is  $\Delta S_{[\text{s-peng}]}$  ? positive ?



see theory lectures



# Confronting Loop and Tree Decays

- Conflict with  $\sin 2\beta_{\text{eff}}$  from s-penguin modes ?

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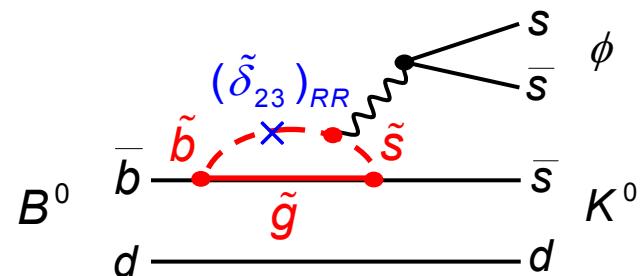
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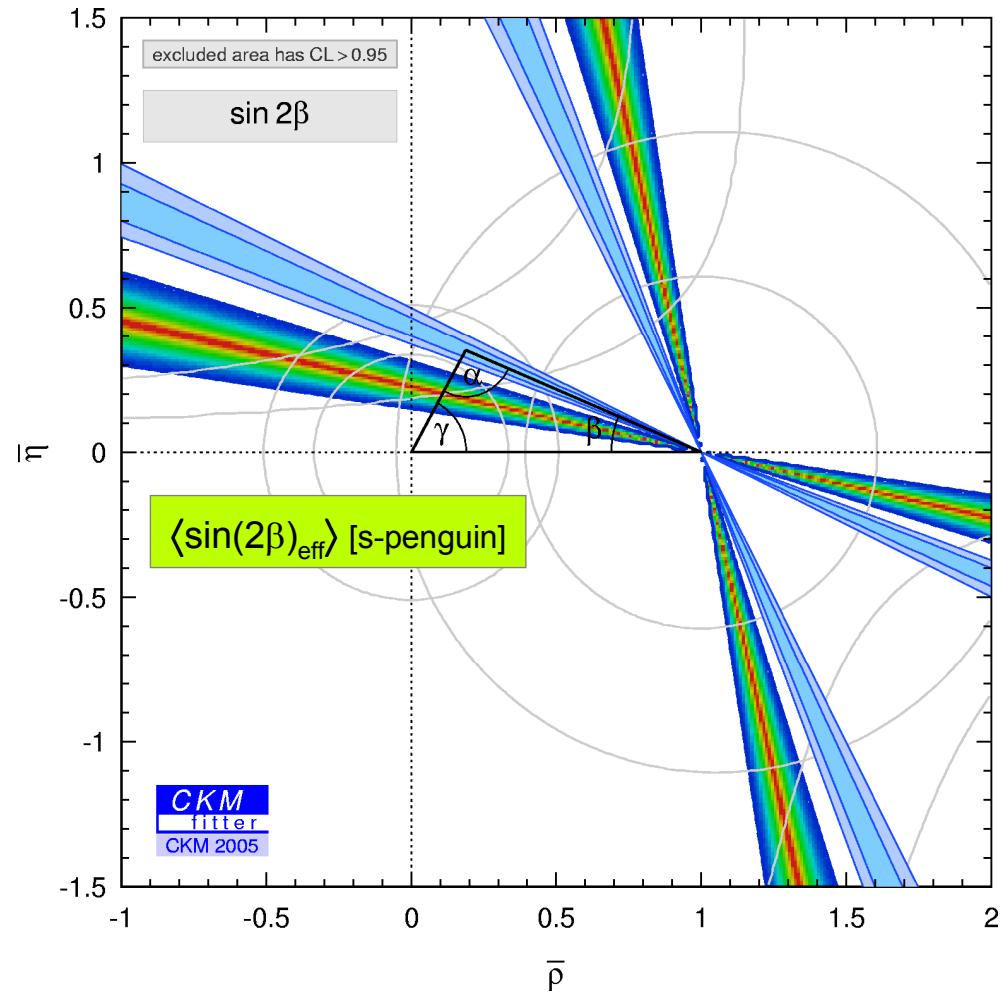
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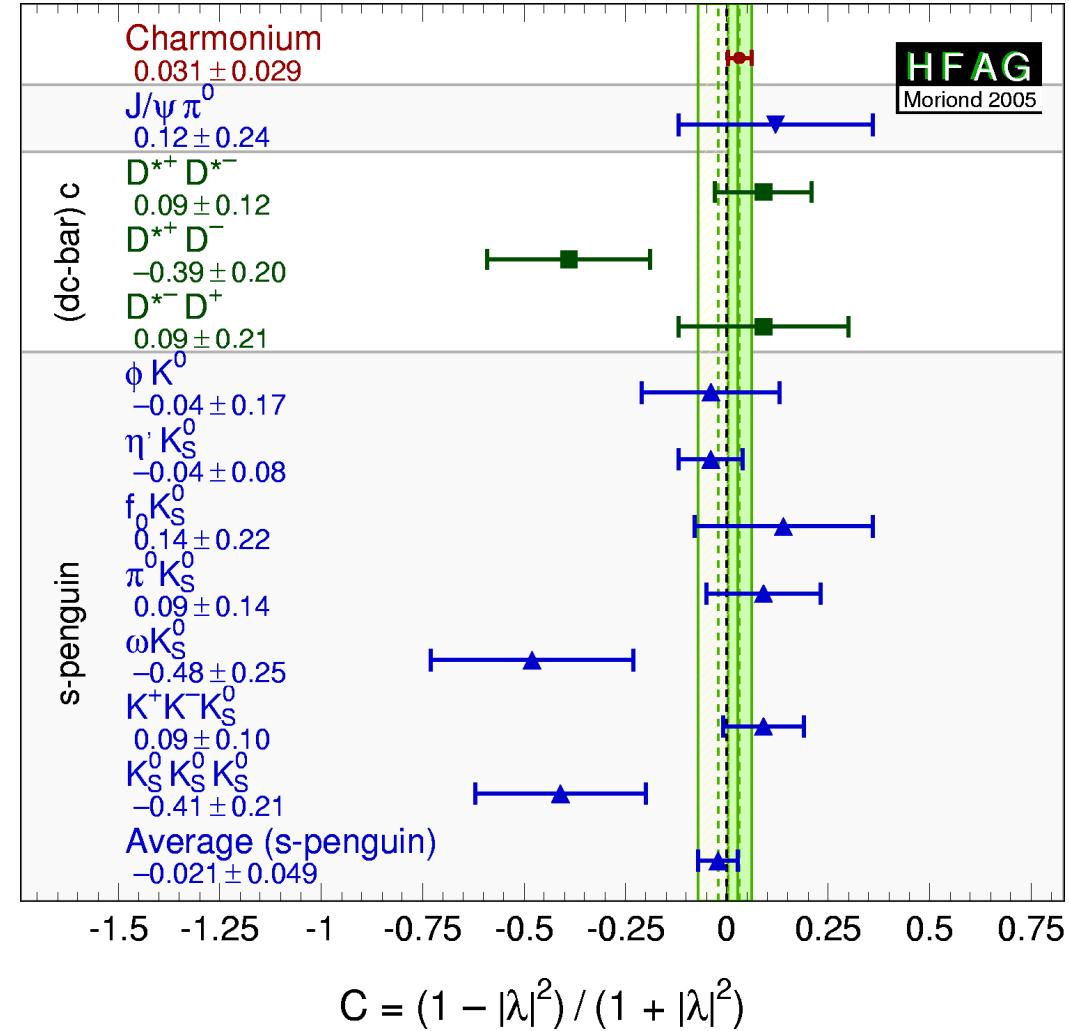
- New Physics in  $b \rightarrow s$  transitions?



Masiero-Murayama,  
PRL 83, 907 (1999)

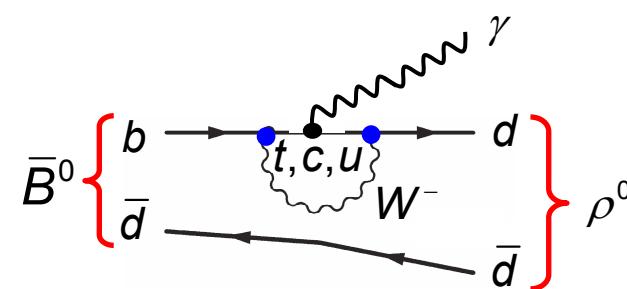
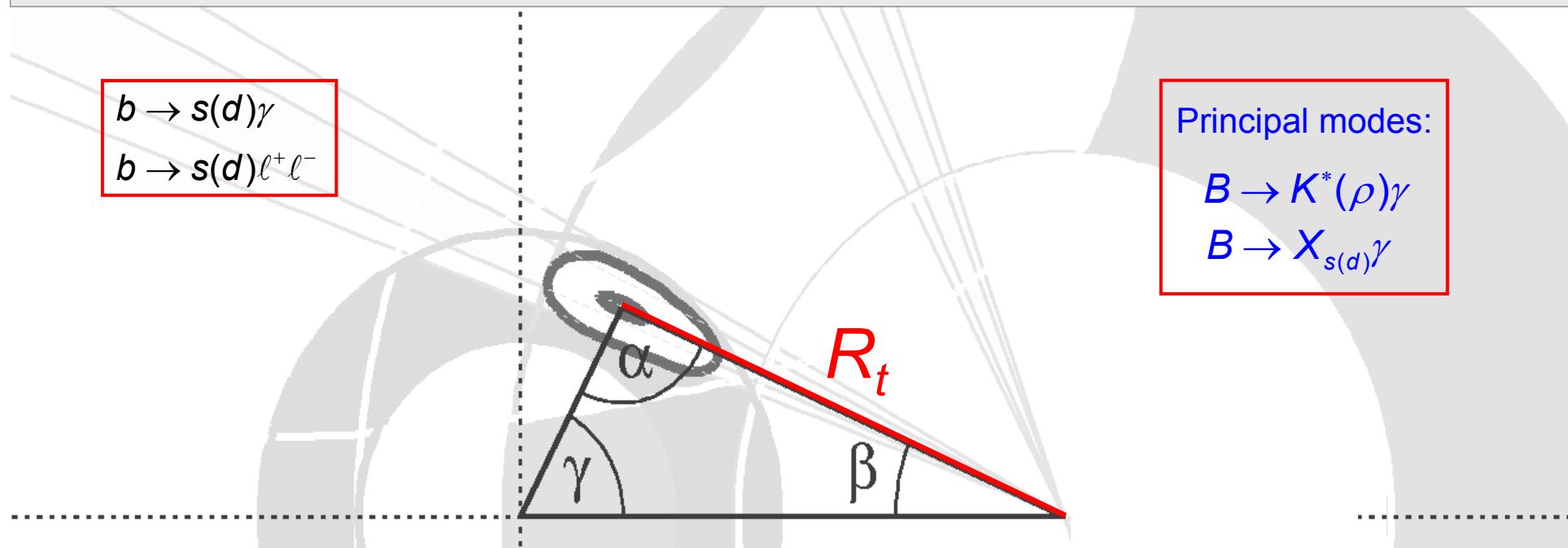


# If there were New Physics ... how about direct CPV ?



... none.

# Radiative Penguin Decays



Penguin : dominant

$$\begin{aligned} &\propto V_{tb} V_{td}^* \\ &\propto \lambda^3 \end{aligned}$$

# Radiative Penguin Decays

- Radiative penguin decays  $B \rightarrow \rho\gamma$  ( $\propto |V_{td}|^2$ ) and  $B \rightarrow K^*\gamma$  ( $\propto |V_{ts}|^2$ ) sensitive to New Physics
- Ratio of BRs predicted more cleanly than the individual rates: SU(3) breaking correction

$$\frac{\text{BR}(B^0 \rightarrow \rho^0\gamma)}{\text{BR}(B^0 \rightarrow K^{*0}\gamma)} \propto \frac{1}{2} \left| \frac{V_{td}}{V_{ts}} \right|^2 \xi^{-2} (1 + \Delta)$$

$$\underbrace{\xi = 1.2 \pm 0.1}_{\text{SU(3) breaking}}, \quad \underbrace{\Delta < 0.04}_{\text{NP contrib.}}$$

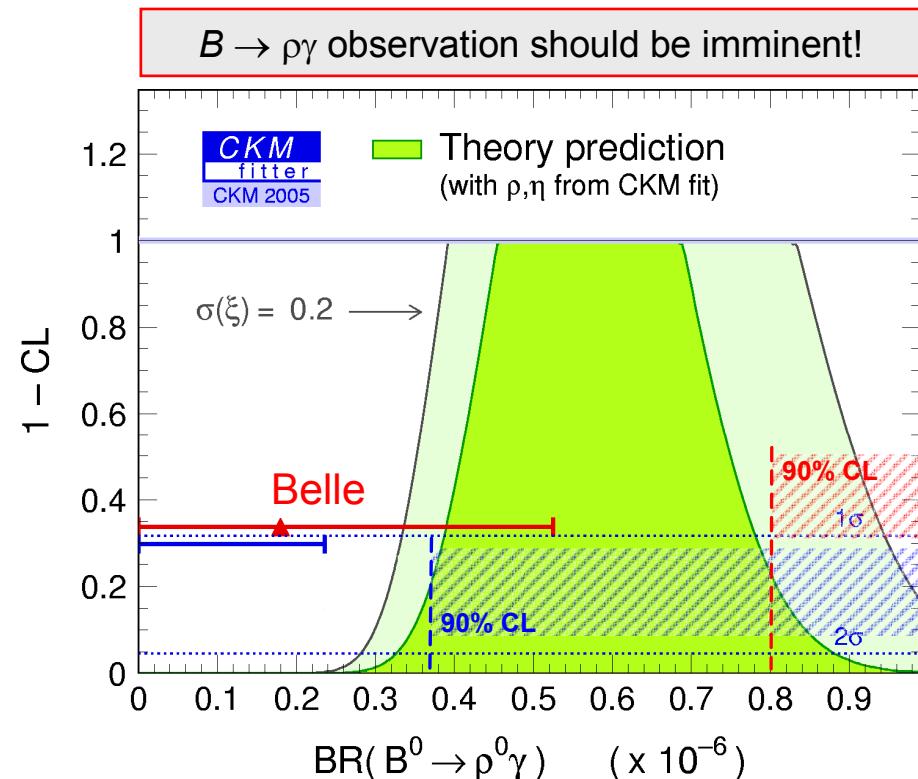
Ali, Parkhomenko, EPJ C23 (2002) 89  
 Bosch, Buchalla, NP B621 (2002) 459  
 (and later papers); errors from CKM 05

- So far only upper limit for  $B \rightarrow \rho\gamma$

$$\text{BR}(B^0 \rightarrow \rho^0\gamma) = (0.06^{+0.19}_{-0.14}) \times 10^{-6}$$

$$\text{BR}(B^0 \rightarrow K^{*0}\gamma) = (40.1 \pm 2.0) \times 10^{-6}$$

BABAR, PRL 94, 011801 (2005)  
 Belle, hep-ex/0408137 (prelim.)



Charged modes larger limit:  $\text{BR}(B^+ \rightarrow \rho^+\gamma) = (1.0 \pm 0.4) \times 10^{-6}$ , but less theoretically clean

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SU(3) breaking      NP contrib.

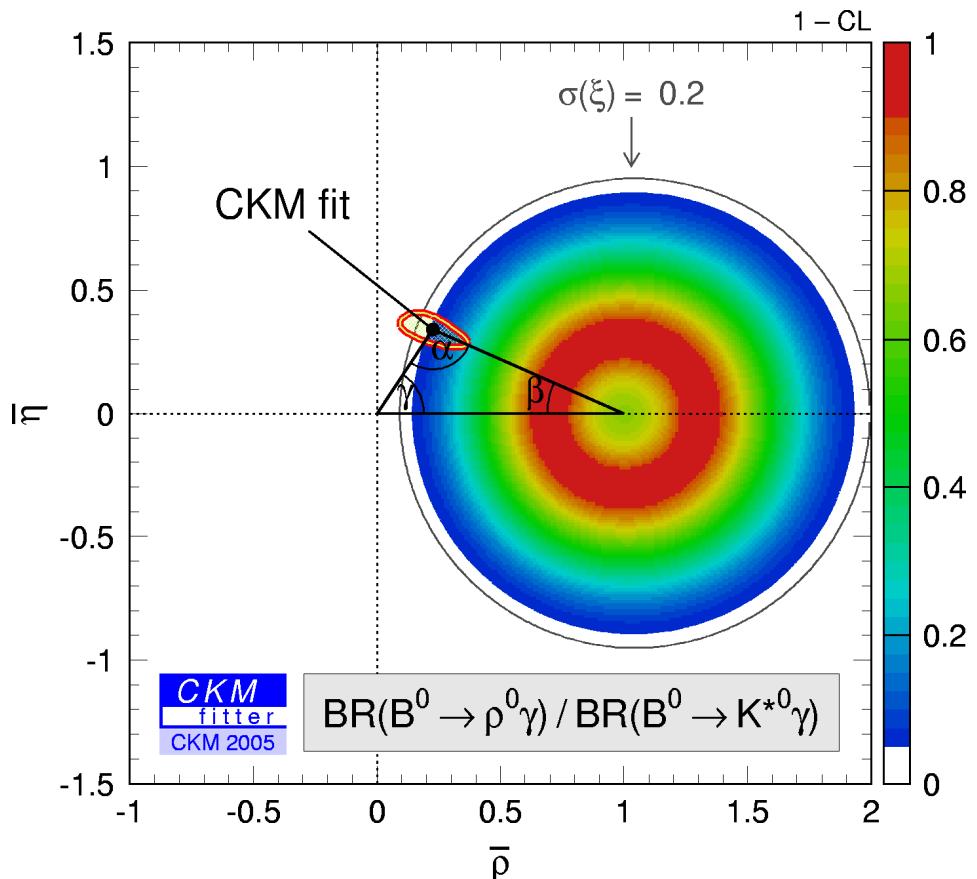
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# New Physics in Penguins ?

copyright: J. Berryhill (SLAC)



Penguins being inspected carefully for new physics (J.B.)

# appendix

## ■ penguin outlook

# The Experimental Program for $\sin 2\beta_{\text{eff}}$

Mode	$CP$	Tot. error Belle $\mathcal{L} \sim 253 \text{ fb}^{-1}$	Tot. error BABAR $\mathcal{L} \sim 195\text{-}212 \text{ fb}^{-1}$	$\langle \Delta(\text{SM}) \rangle$ [in $\sigma$ ]	Error estimate at $2 \text{ ab}^{-1}$	System- atics	Max. central value for $5\sigma$ deviation at $2 \text{ ab}^{-1}$	Quality [naïve theoretical cleanliness]
$\phi K^0$	-1	0.34	0.26	-1.9	0.10	small	0.22	😊 😊 😊
$\eta' K^0$	-1	0.18	0.14	-2.6	< 0.05	small	0.45	😊 😊 (😊)
$f_0(980)K^0$	+1	0.42	0.29	-1.3	< 0.12	Q2B	0.12	😊 😊
$K_S K_S K^0$	±1	0.71	0.36	-1.4	< 0.16	vertex	-0.08	😊 😊 😊
$K^+ K^- K^0$	~+1	0.25	0.25	-1.1	< 0.08	$CP$	0.31	(😊)
$\pi^0 K_S$	-1	0.60	0.32	-1.4	0.13	vertex	0.07	😊
$\omega K^0$	-1	0.66	0.36	-0.6	< 0.15	small	-0.03	(😊)
$\rho^0 K^0$	-1	-	-	?	?	Q2B	?	(😊)
$\eta K_S$	+1	-	-	?	?	vertex	?	-
Average	-	$0.39 \pm 0.11$	$0.45 \pm 0.09$	-3.7	< 0.034	ok	0.53	😊 😊