

Heavy Quark Physics (at J-PARC) and J-PARC Heavy Ion Project

Makoto Oka

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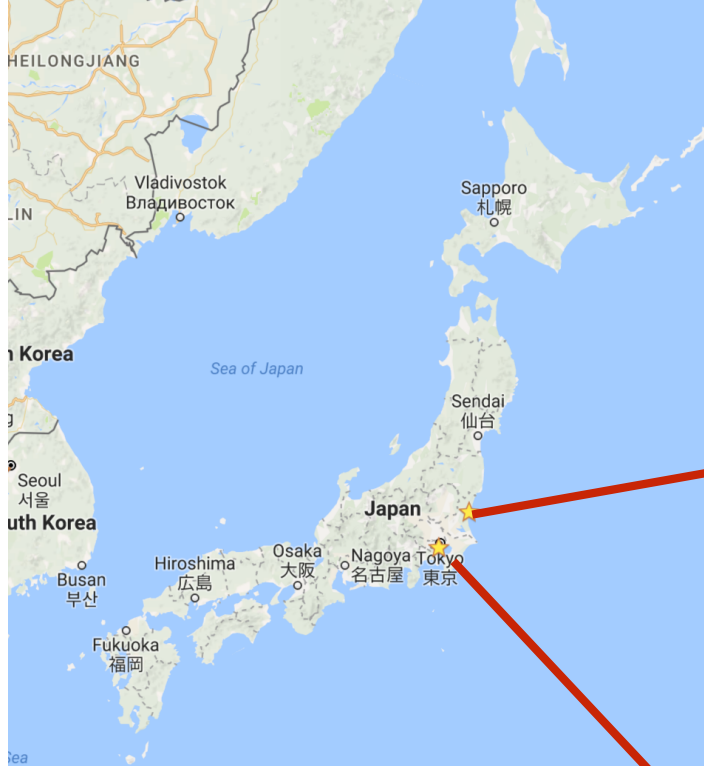
and

Advanced Science Research Center

Japan Atomic Energy Agency (JAEA)

Hadronic Matter under Extreme Conditions

JINR, Dubna, October 31, 2016



JAEA J-PARC

Tokyo Institute of Technology
Science & technology...

Contents

- 1. Introduction: From QCD to Hadron Spectrum**
- 2. Charmed Baryons and Diquark**
- 3. Quarkonium and Hadron Molecules**
- 4. Charmed Dibaryons**
- 5. Conclusion**

+ *J-PARC-HI Experimental Project (by Takao Sakaguchi)*

J-PARC (KEK & JAEA)



**J-PARC
(KEK & JAEA)**

400 MeV H-Linac

**Transmutation Experimental
Facility (TEF)**

**3 GeV Rapid Cycling
Synchrotron (RCS)**

Neutrino experiment (NU)

**50 GeV Main Ring
Synchrotron (MR)
[30 GeV at present]**

**Materials & Life
Science Facility
(MLF)**

**Hadron
Experimental
Hall (HD)**

**J-PARC
(KEK & JAEA)**

400 MeV H⁻ Linac

**Transmutation Experimental
Facility (TEF)**

**3 GeV Rapid Cycling
Synchrotron (RCS)**

HI booster

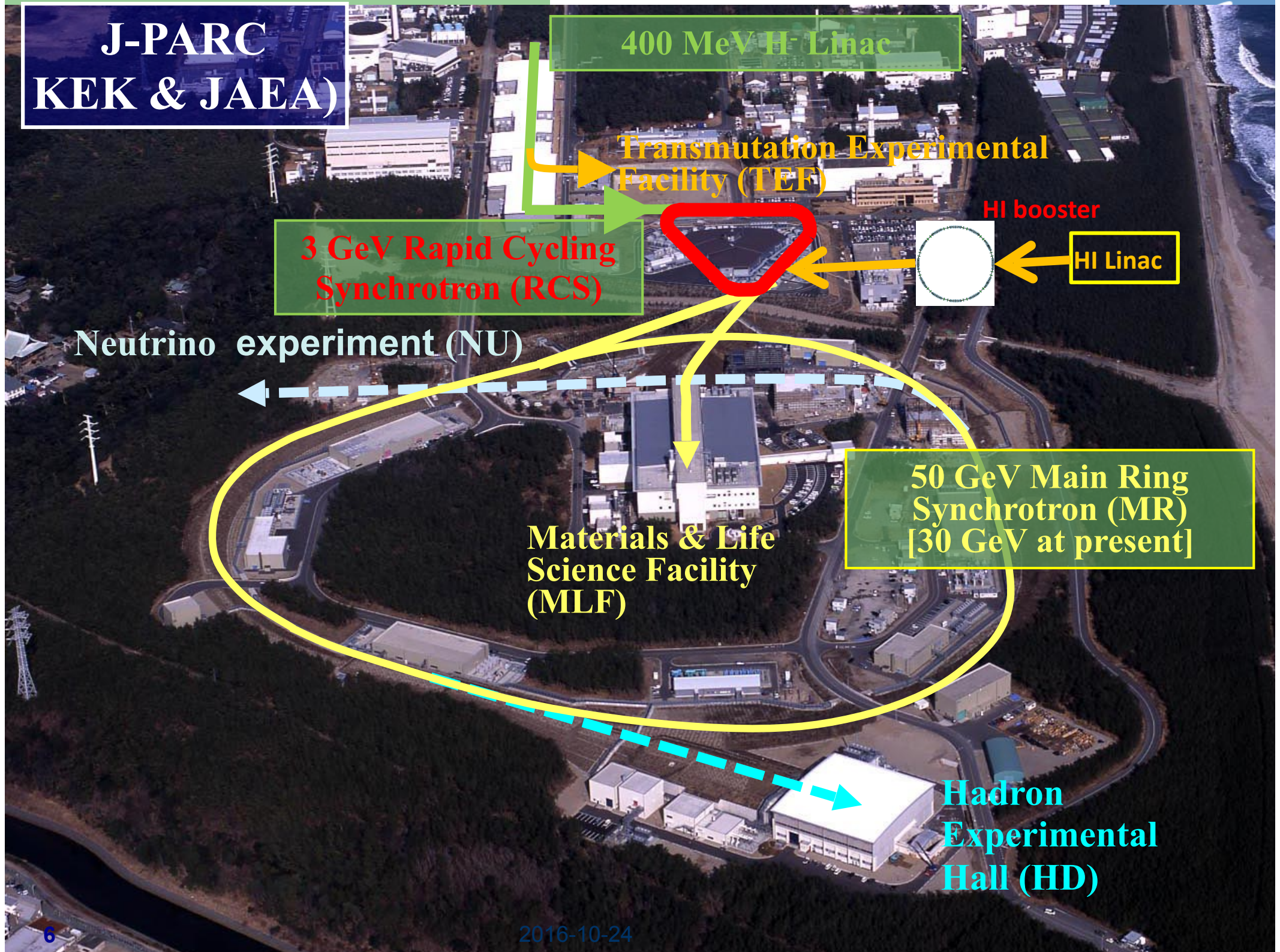
HI Linac

Neutrino experiment (NU)

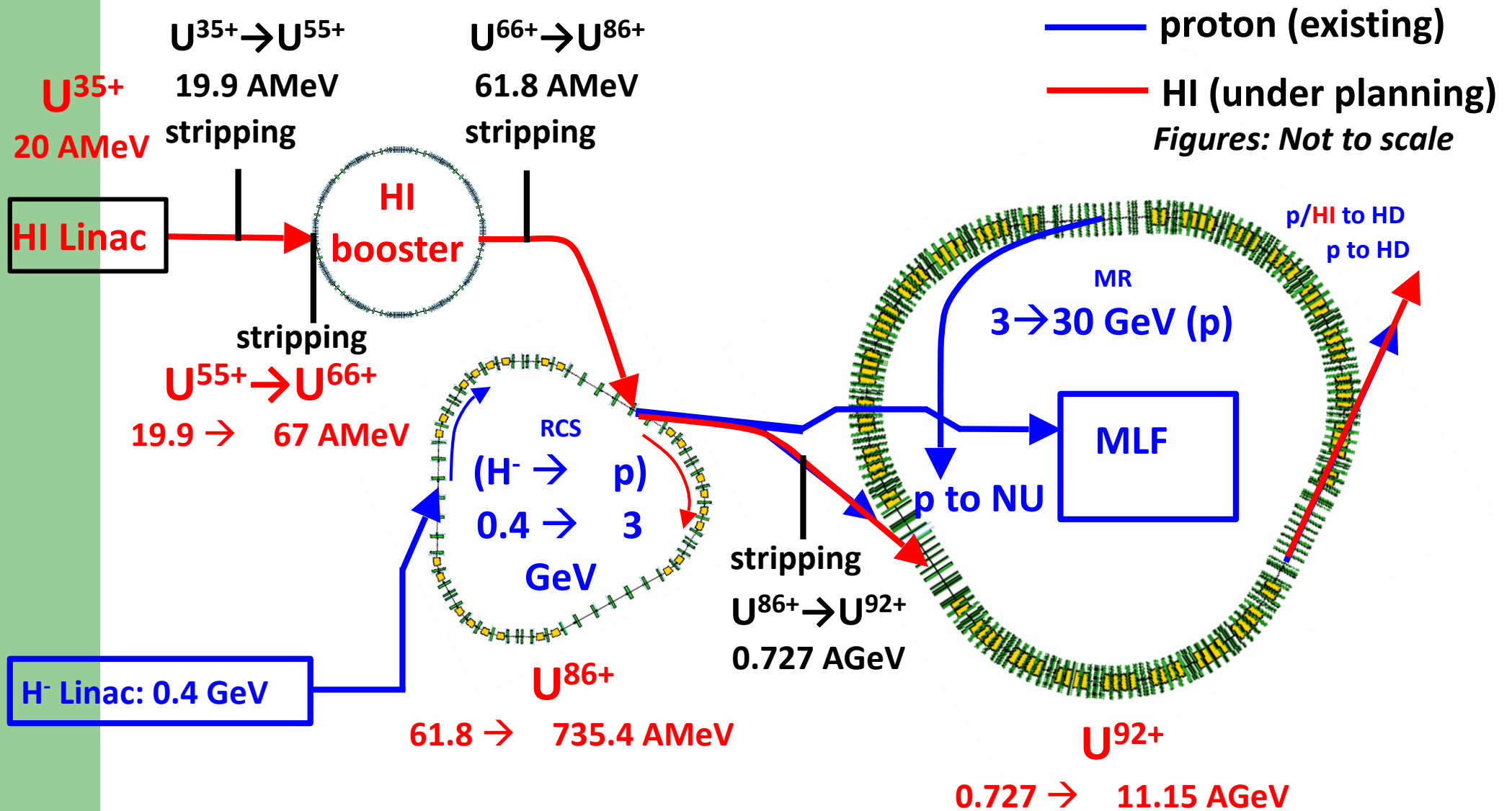
**50 GeV Main Ring
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HI Accelerator scheme in J-PARC

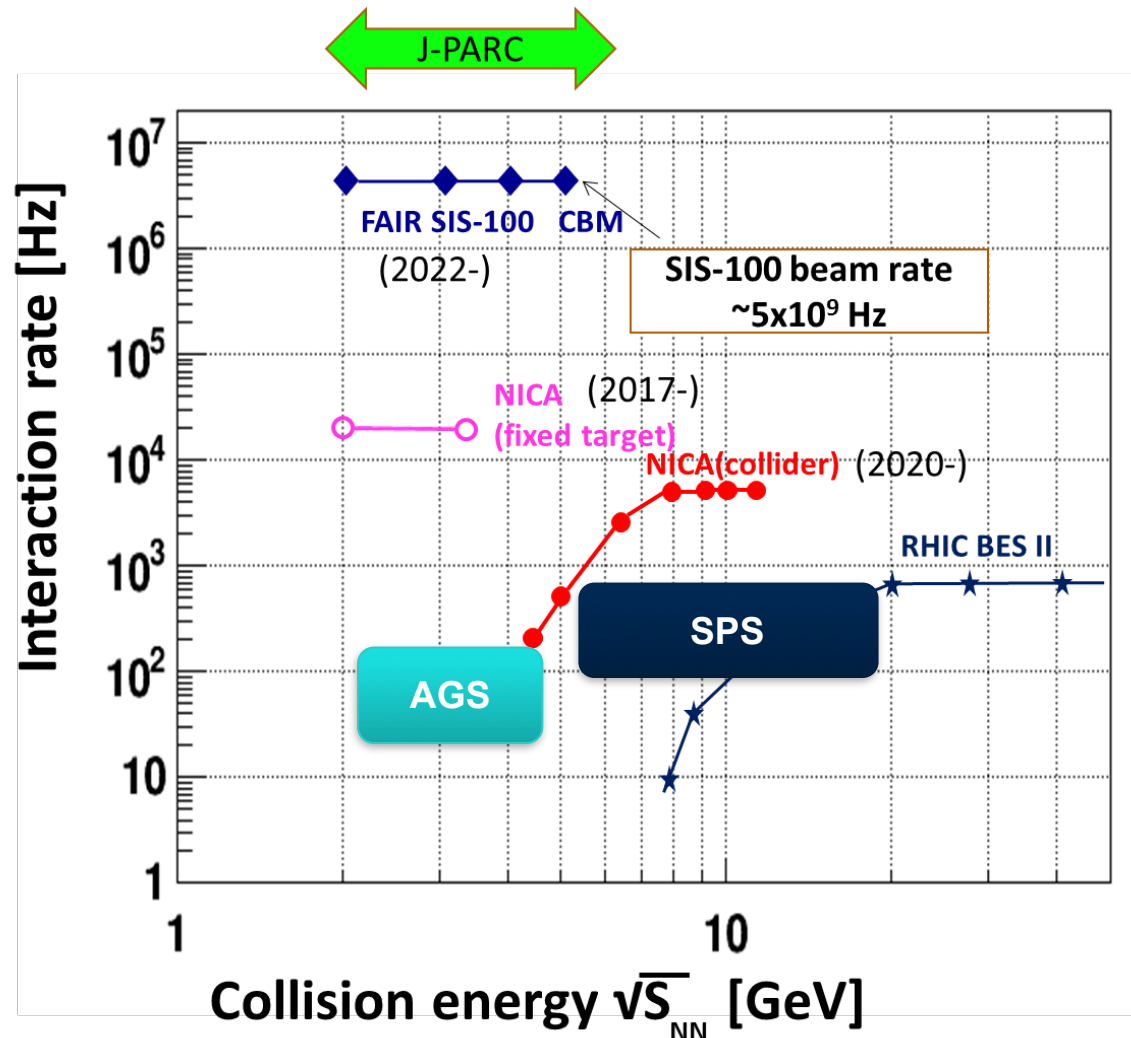


Figures: Not to scale

This HI accelerator scheme has no interference/conflict with proton beam programs

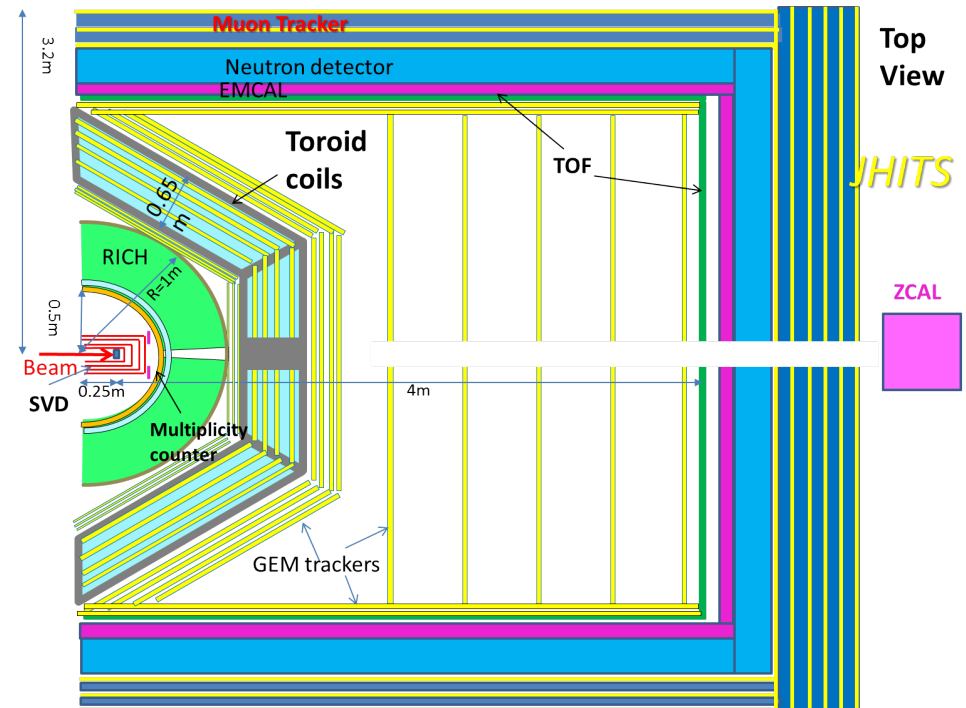
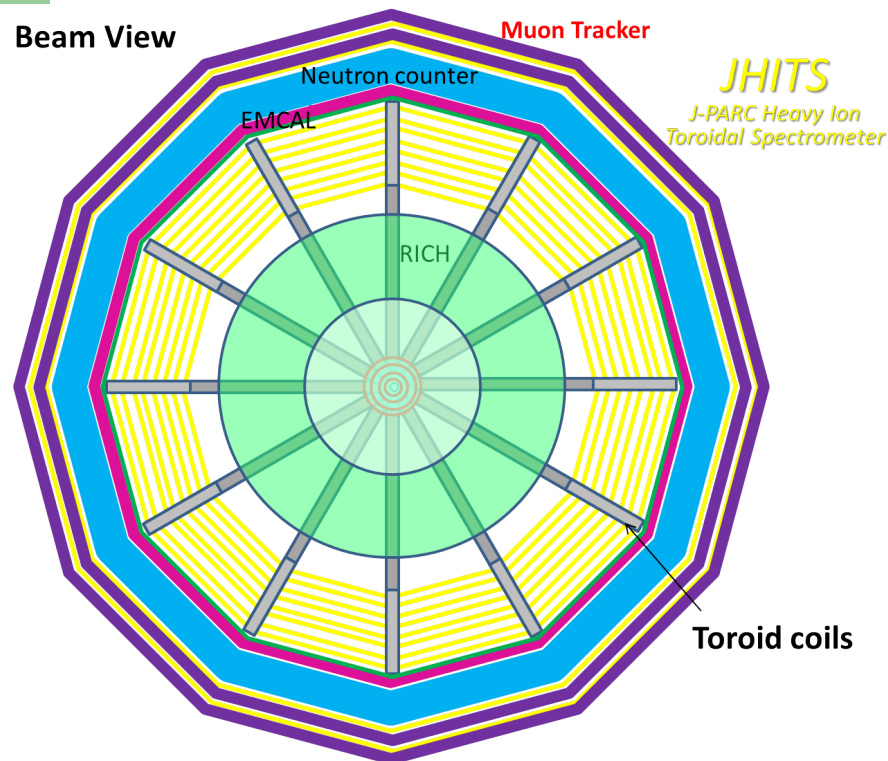
Available beam and rate

- *Very high intensity beam is a feature of J-PARC HI accelerator*
 - $E_{lab}=1-19\text{GeV}/n$, $\sqrt{s_{NN}}=1.9-6.2\text{GeV}$ (\sim AGS), $>10^{11}$ cycle $^{-1}$ (\sim 6s cycle)
- Ion species: p, Si, Ar, Cu, Xe, Au(Pb), U, and also light ions for hypernuclei



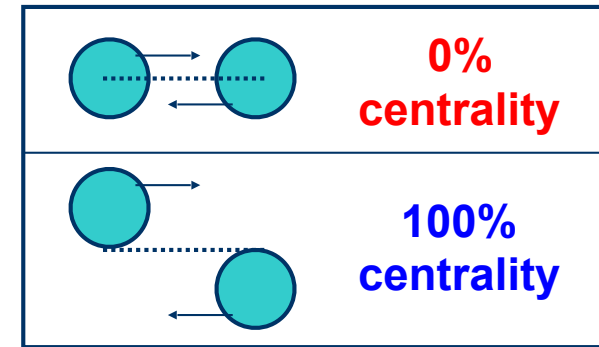
Concept of measurement device

- Detector complex covering wide acceptance
 - High speed tracking, TOF, EM calorimeter, and muon detector
- 0.1% λ_1 target: $\sim 100\text{MHz}$ event rate, 1000 particles/event
- Collect data with minimum bias trigger (Data size: 1TB/s)
 - Continuously take data with no trigger (Import ALICE experience)
 - Select rare events in semi-online, using a high performance computing system



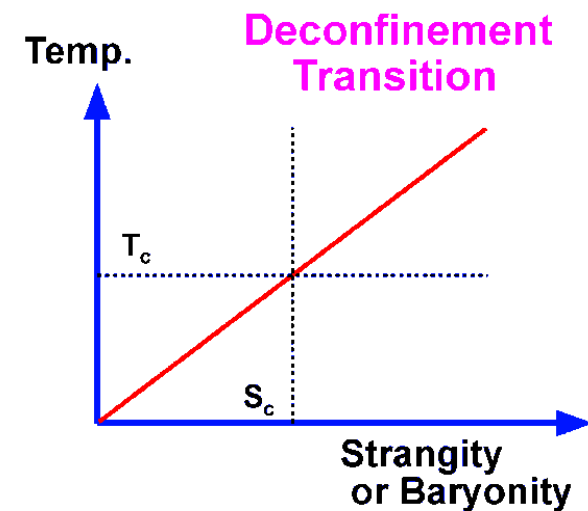
Event selection consideration

- Centrality: Event class variable proportional to impact parameters
 - 0%: $b=0$, **Central collisions**
 - 100%: $b=b_{max}$, **Peripheral collisions**
- Same event selection as we did in the past wouldn't yield new physics
- **We add a new event selection**
 - **After pre-selecting most central collisions**
- **Strangity, Baryonity**
 - Aggressively select interesting events relevant to the new phenomena found by the AGS experiment
 - **Strangeness enhancement, baryon stopping**
- Statistics-starved "very rare event" selection feasible with high luminosity beam at J-PARC-HI



$$\text{Strangity} \equiv \langle N(K^+) \rangle / \langle N_{ch} \rangle$$

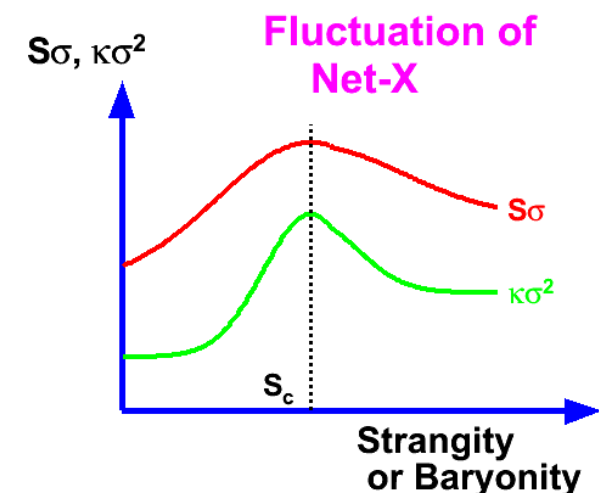
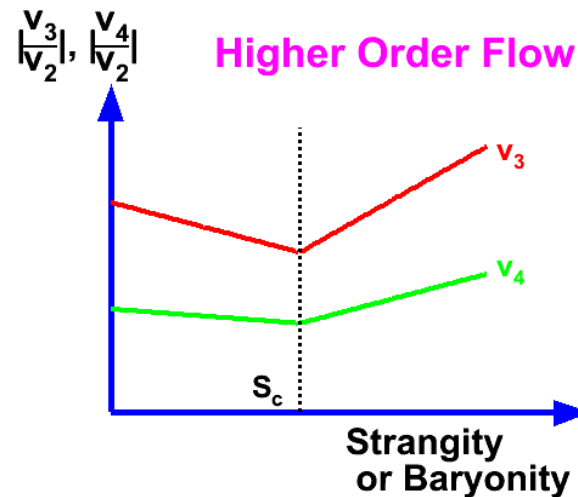
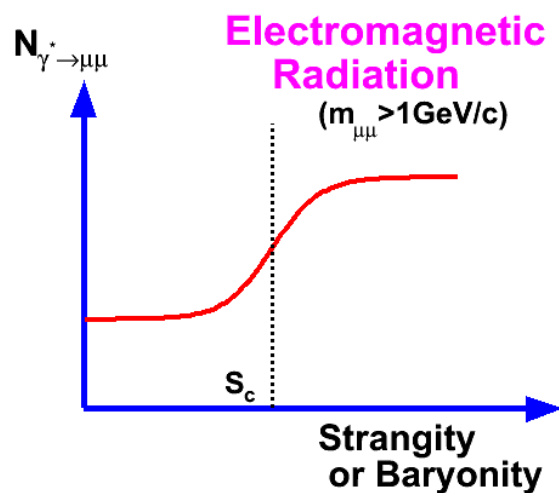
$$\text{Baryonity} \equiv \langle N(p) - N(\bar{p}) \rangle$$



TS, H. Sako and M. Kitazawa, in prep.

Physics observables at J-PARC-HI

- Primarily focus on new observables found at higher energy experiments
 - Based on knowledge gained at RHIC and LHC
- Study characteristics of high density matter
 - Particle emission anisotropy, fluctuation of conserved quantities
 - Lepton pairs, thermal photons



White paper for a Future J-PARC Heavy-Ion Program (J-PARC-HI)

Preliminary plan towards the proposal to J-PARC
by an international research collaboration
(J-PARC-HI Collaboration)

June, 2016

<http://asrc.jaea.go.jp/soshiki/gr/hadron/jparc-hi/>

Letter of Intent for J-PARC Heavy-Ion Program (J-PARC-HI)

August, 2016

J-PARC-HI Collaboration

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Introduction: From QCD to Hadron Spectrum

From QCD to Hadron Spectrum

QCD = quarks + gluons with color $SU(3)_c$ gauge symmetry

$$\mathcal{L} = \bar{q}(i\not{D} - m_q)q - \frac{1}{2}\text{Tr}[G_{\mu\nu}G^{\mu\nu}]$$

expected low energy modes

massless gluons

light quarks ($m_q < 10$ MeV)

But, in reality,

massless gluons \Rightarrow **glueballs** ($m_{GB} \sim 1.4$ GeV or larger)

light quarks \Rightarrow **mesons** (500~800 MeV) except for the pion
baryons (940 MeV ~)

QCD at low energy is strongly correlated.

\Rightarrow color confinement, chiral symmetry breaking, . . .

Nontrivial QCD Vacuum

Properties of the QCD vacuum

– Chiral symmetry breaking

quark condensate $\langle \bar{q}q \rangle \neq 0$

– Scale invariance violation

gluon condensate $\langle G_{\mu\nu}G_{\mu\nu} \rangle \neq 0$

– Topological density

instanton vacuum $\langle G_{\mu\nu}\tilde{G}_{\mu\nu} \rangle \neq 0$

– Color confinement

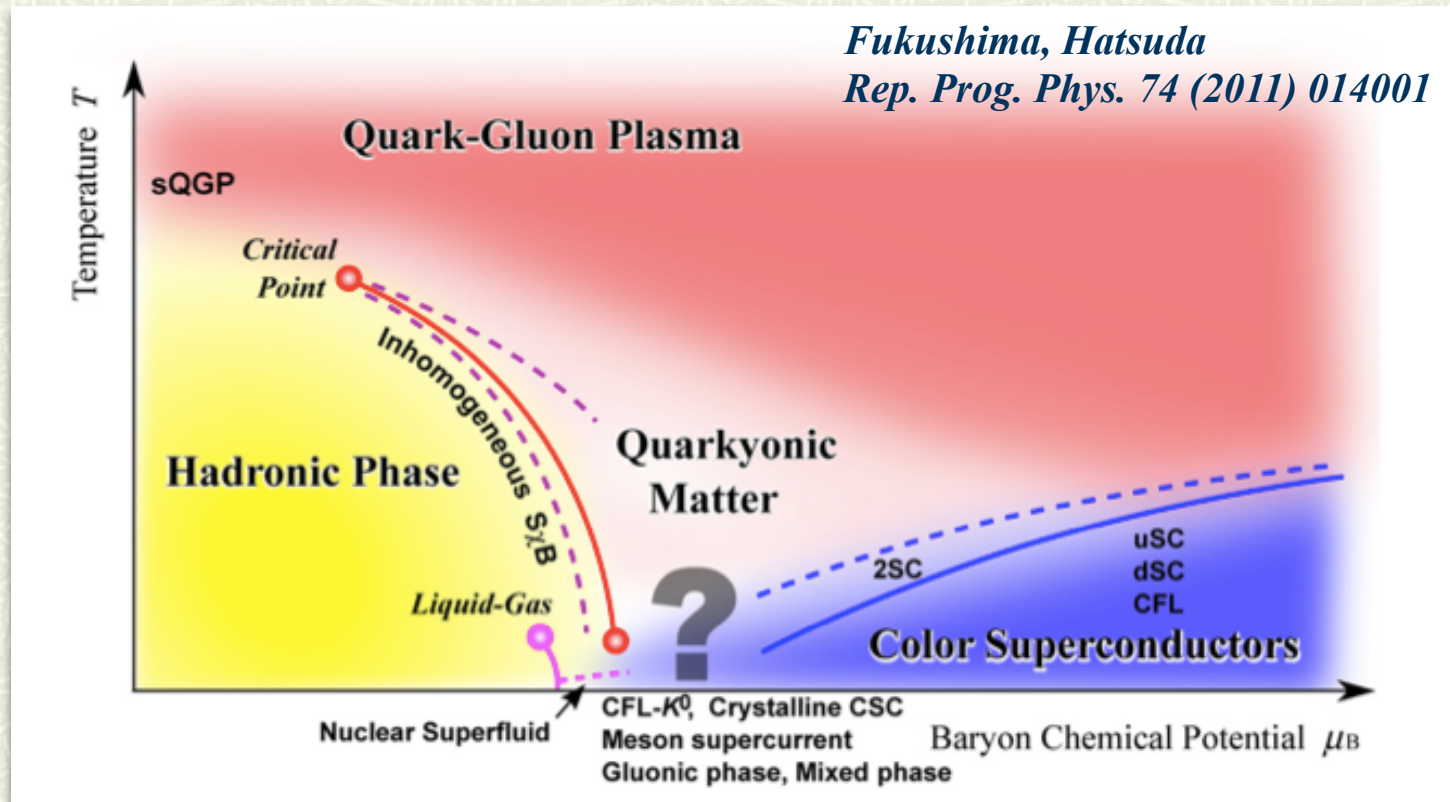
Polyakov loop $\left\langle \mathcal{P} \exp \left(i \int d\tau A_4 \right) \right\rangle = 0$

Hadrons are elementary modes, which probe the QCD vacuum.

What are the relevant degrees of freedom in hadron excitations?

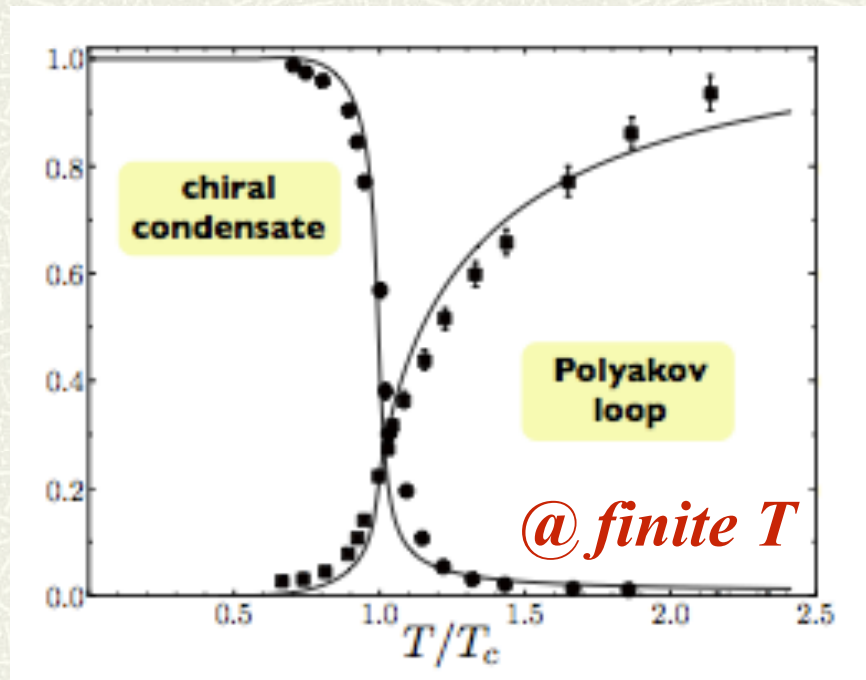
Hadronic Matter @ Finite T and/or ρ

- Hot and Dense hadronic matters show diverse properties.
- Phase transitions to QGP @ high T *heavy ion collisions*
- Color superconductor @ high ρ , low T *neutron star*



Hadrons in Matter

- # Hadrons in matter can probe the hadronic matter at finite T and/or ρ .
- # Interplays of *chiral symmetry* and *color confinement* at finite T and/or ρ are intriguing.

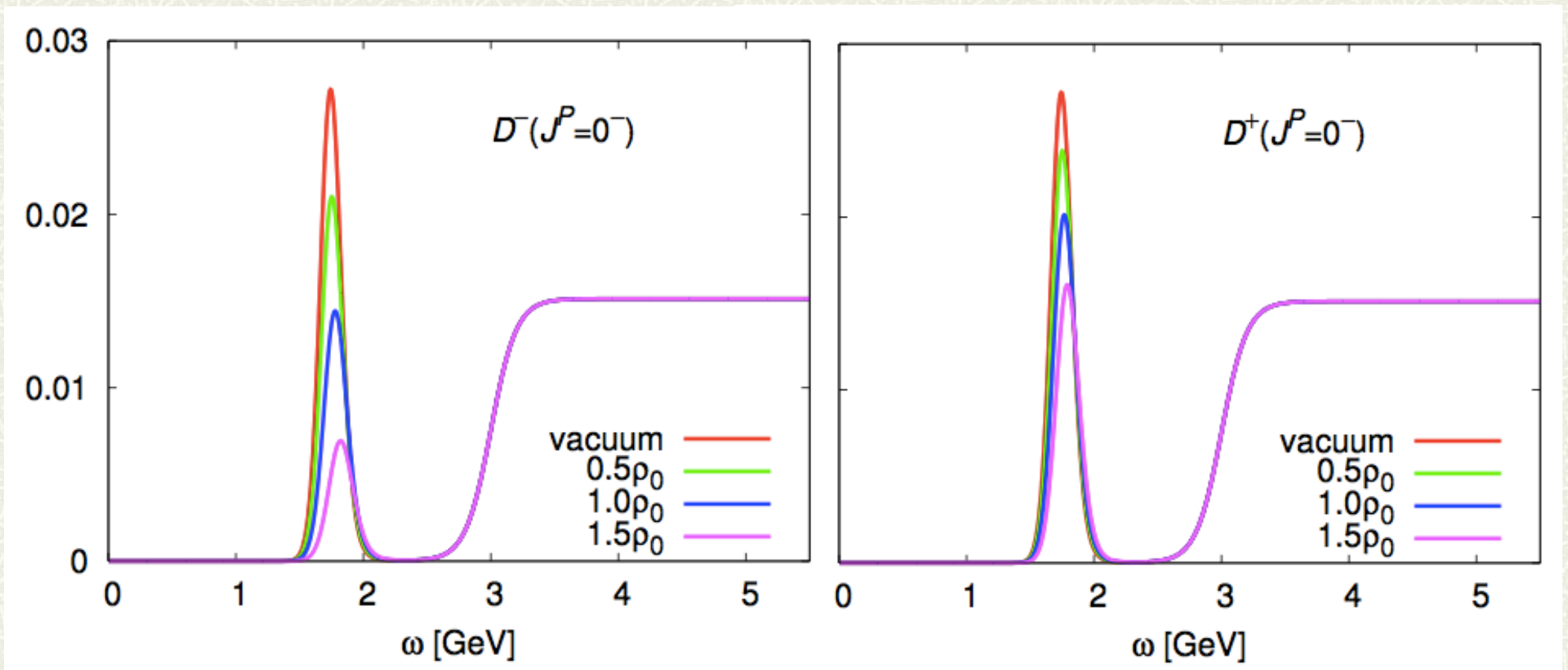


Hadrons in Matter

Masses of D mesons may increase at finite density.

QCD sum rule with the Maximum entropy method:

K. Suzuki, P. Gubler, MO, PR C 93, 045209 (2016)

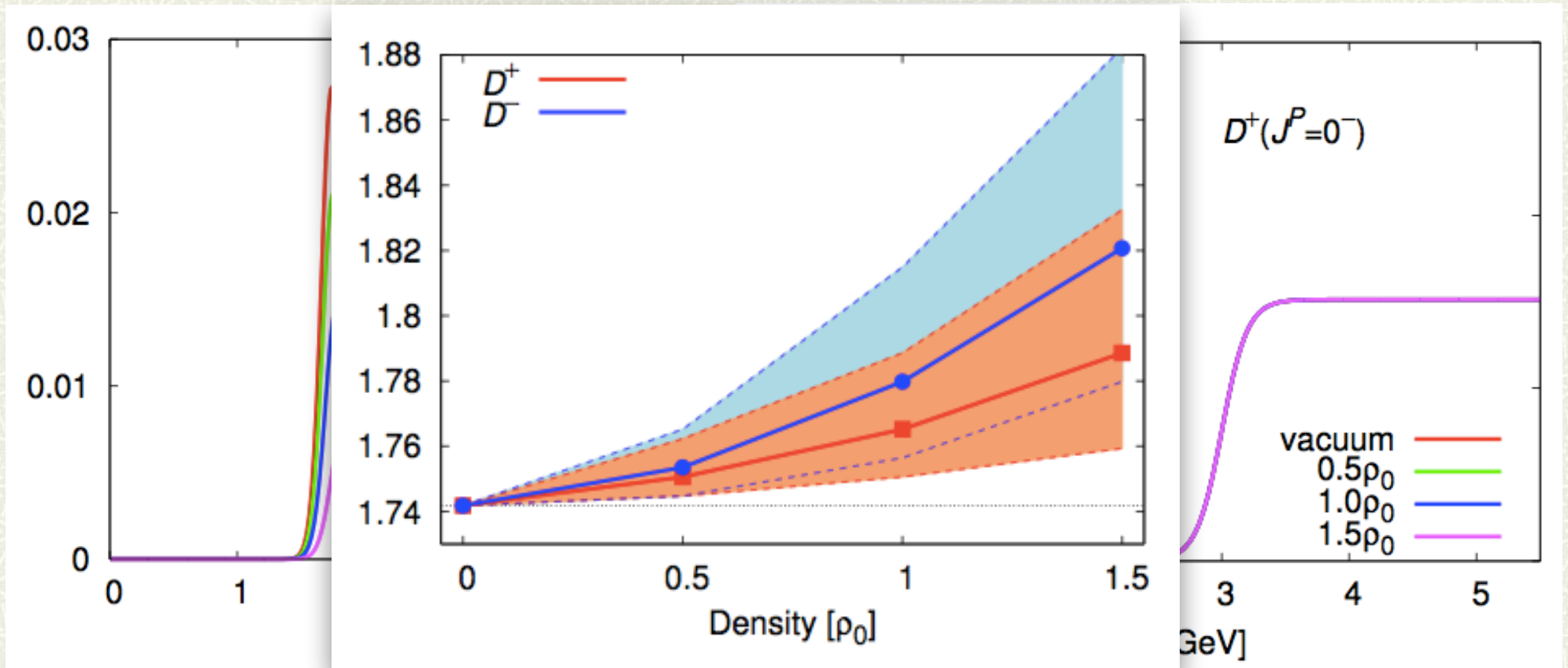


Hadrons in Matter

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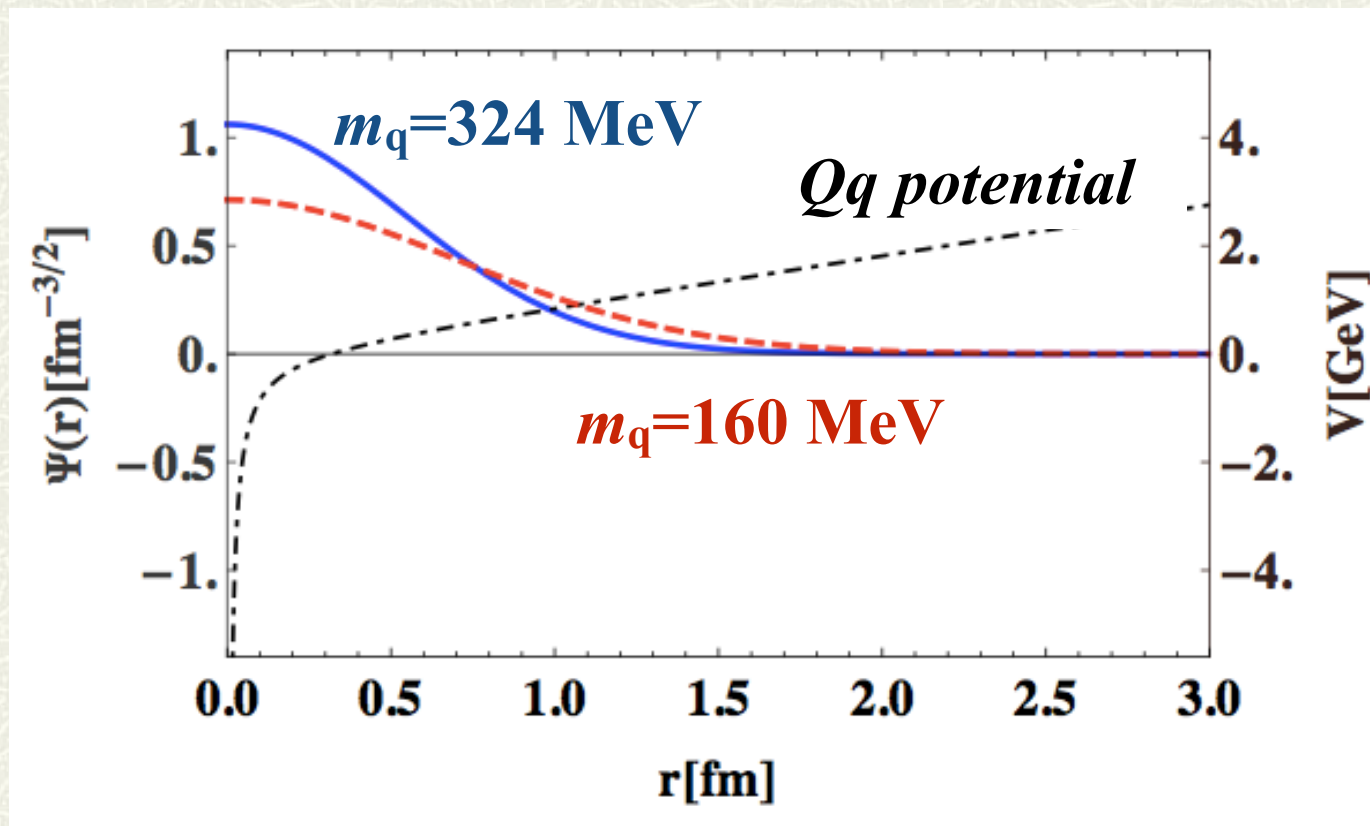
K. Suzuki, P. Gubler, MO, PR C 93, 045209 (2016)



Hadrons in Matter

- ⌘ Interpretation in the constituent quark picture:
The wave function will extend to feel more repulsion.

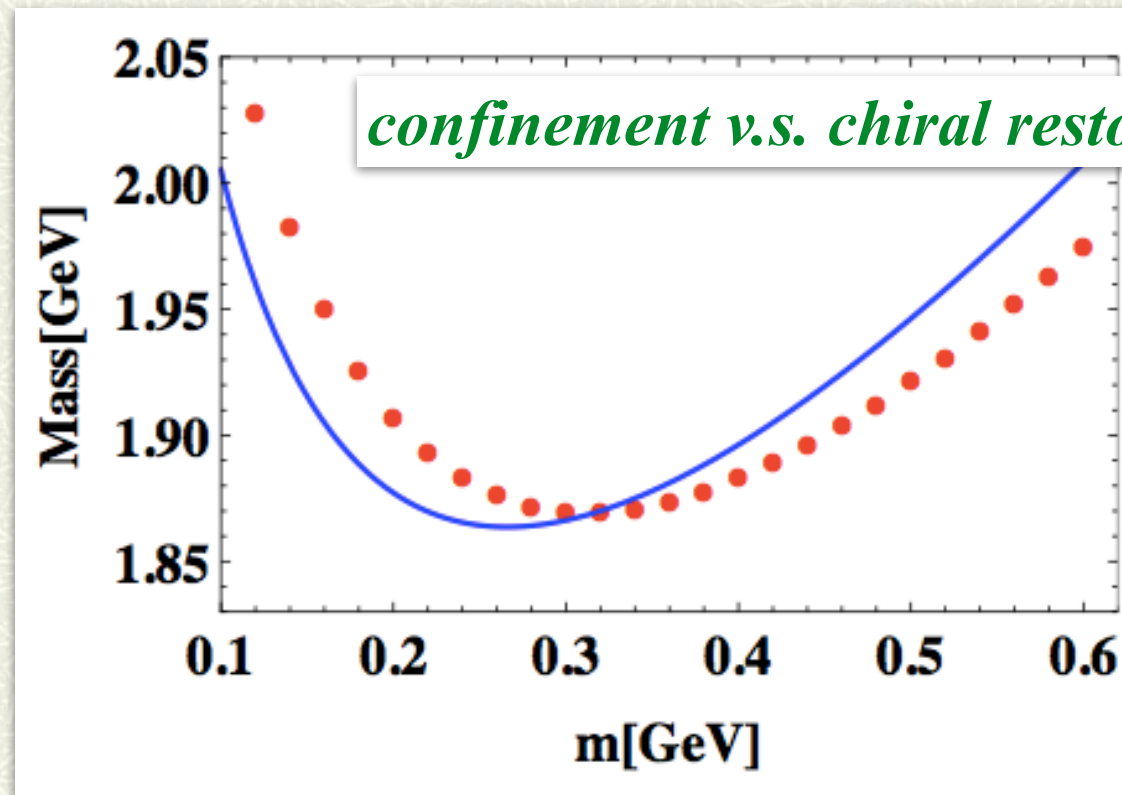
*A. Park, P. Gubler, M. Harada, S.H. Lee, C. Nonaka, W. Park,
PR D93, 054035 (2016)*



Hadrons in Matter

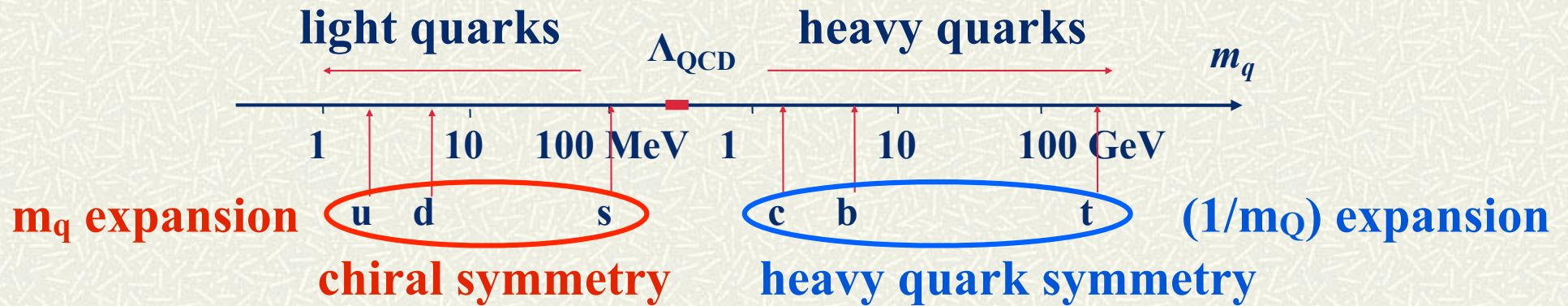
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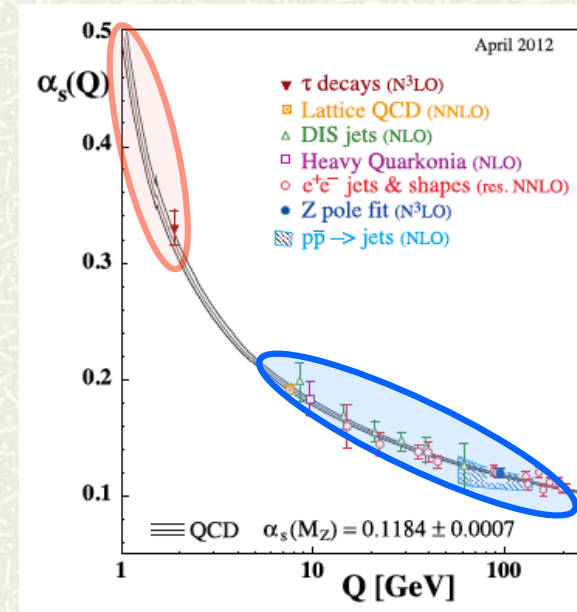


Why Heavy Quarks?

- QCD Lagrangian is flavor independent, but the coupling constant runs.

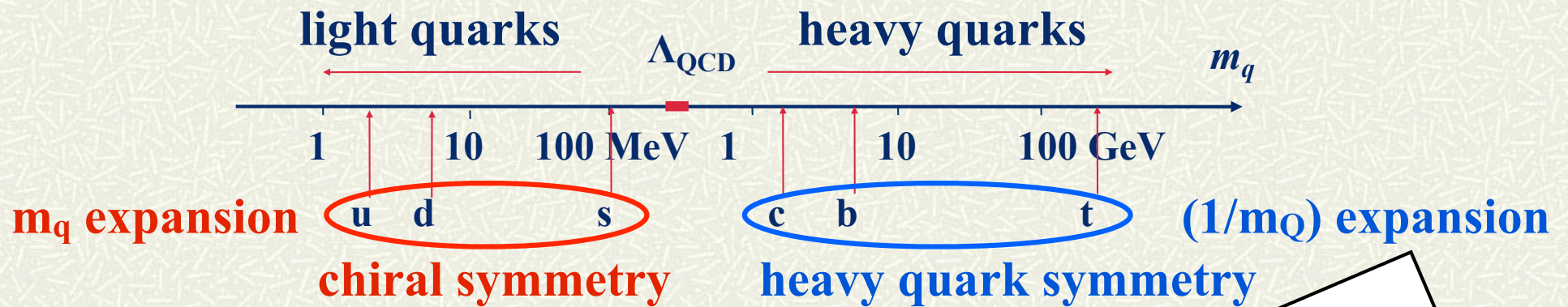


- Light quarks are nonperturbative/ relativistic.
- Heavy quarks are perturbative/ non-relativistic.

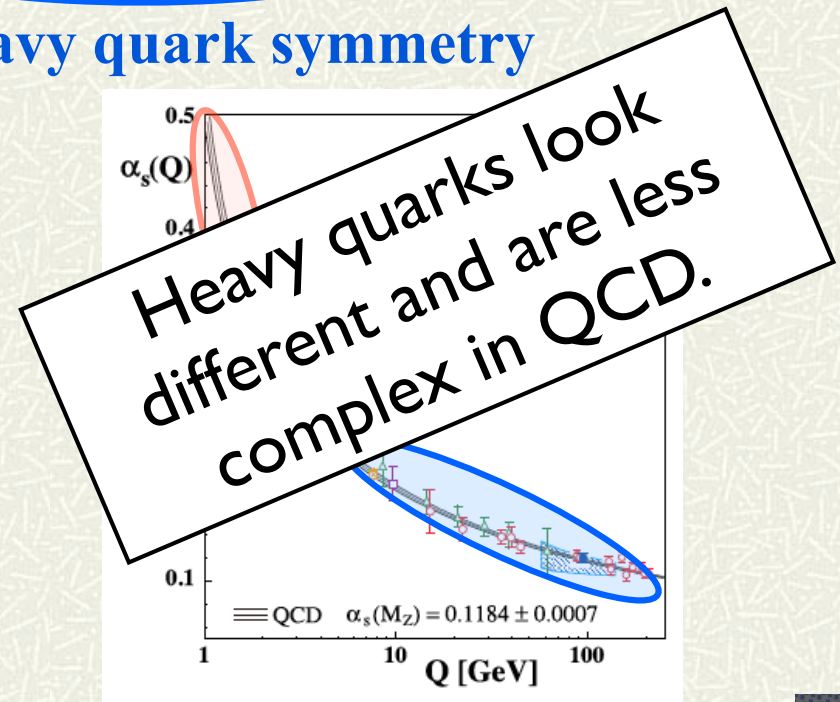


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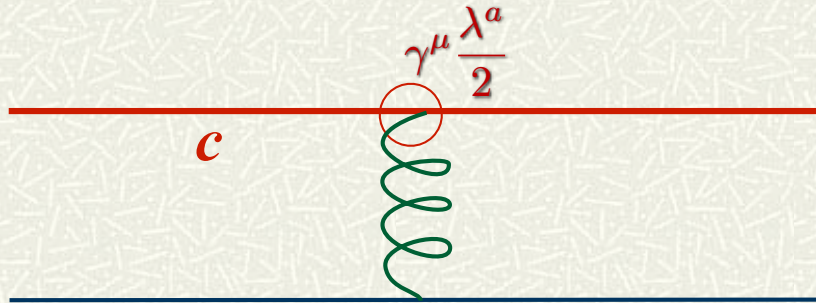




Charmed Baryons and Diquark

Heavy Quark Spin Symmetry

Magnetic gluon coupling is suppressed



$$\bar{\Psi} \gamma^\mu \frac{\lambda^a}{2} \Psi A_\mu^a \sim \underbrace{\Psi^\dagger \frac{\lambda^a}{2} \Psi A_0^a}_{\text{Color Electric coupling}} - \underbrace{\Psi^\dagger \sigma \frac{\lambda^a}{2} \Psi \cdot \frac{1}{m_Q} (\nabla \times A^a)}_{\text{Color Magnetic coupling}}$$

(Color Electric coupling) \gg (Color Magnetic coupling)

HQ spin-flip amplitudes are suppressed by $(1/m_Q)$.

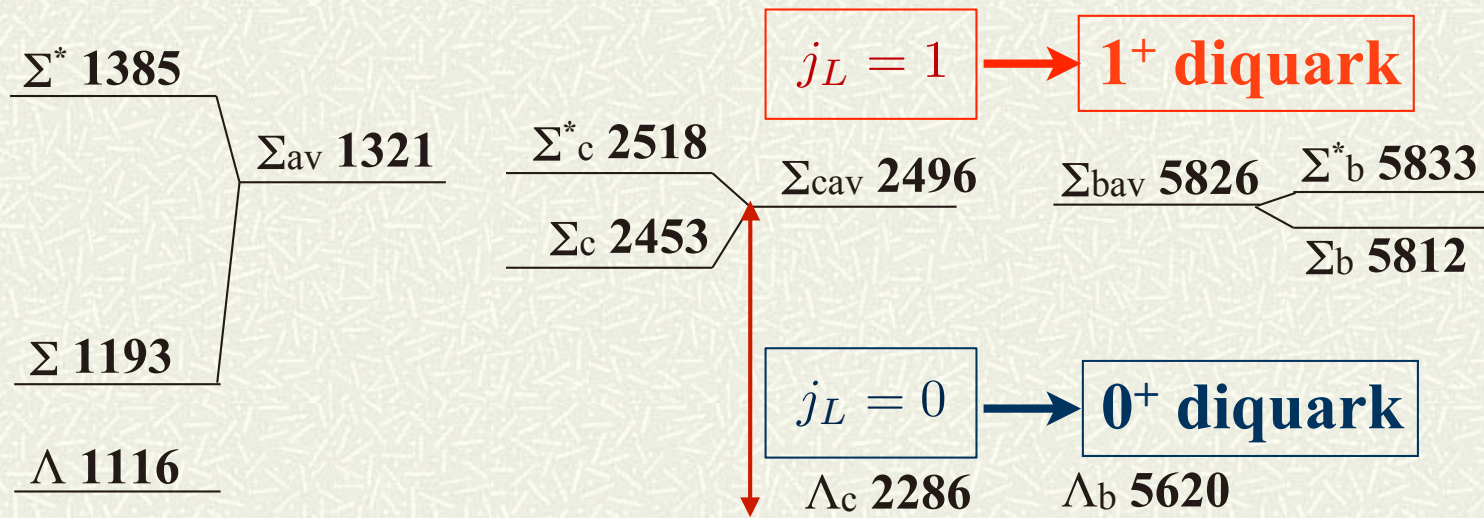
\Rightarrow Heavy Quark Spin Symmetry

Heavy Quark Spin Symmetry

HQ spin symmetry $[S_Q, H] = O\left(\frac{1}{m_Q}\right)$

$$\left. \begin{array}{l} Q \\ \hline qq \end{array} \right\} \vec{J} = \vec{S}_Q + \vec{j}_L \quad \vec{j}_L = \vec{S}_q + \vec{L}_q$$

$J = j_L \pm \frac{1}{2}$ states are degenerate in the HQ limit.



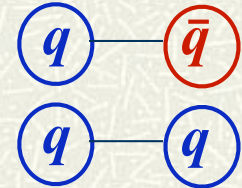
Diquark

The Scalar (0^+) diquark is an analogue of the PS meson:

PS meson $q\bar{q}$: color 1, $J^\pi=0^-$, flavor 1+8

Scalar diquark $[qq]_0$: color 3^{bar} , $J^\pi=0^+$,

flavor SU(3) 3^{bar} : $[ud]_0$, $[ds]_0$, $[sd]_0$



■ Quark model estimate: $S(0^+)$ v.s. $\Lambda(1^+)$

$$M(1^+) - M(0^+) = (2/3) [M(\Delta) - M(N)] \sim 200 \text{ MeV}$$

■ (Quenched) Lattice QCD

Hess, Karsch, Laermann, Wetzorke, PR D58, 111502 (1998)

$$M(1^+) - M(0^+) \sim 120 \text{ MeV (Landau gauge)}$$

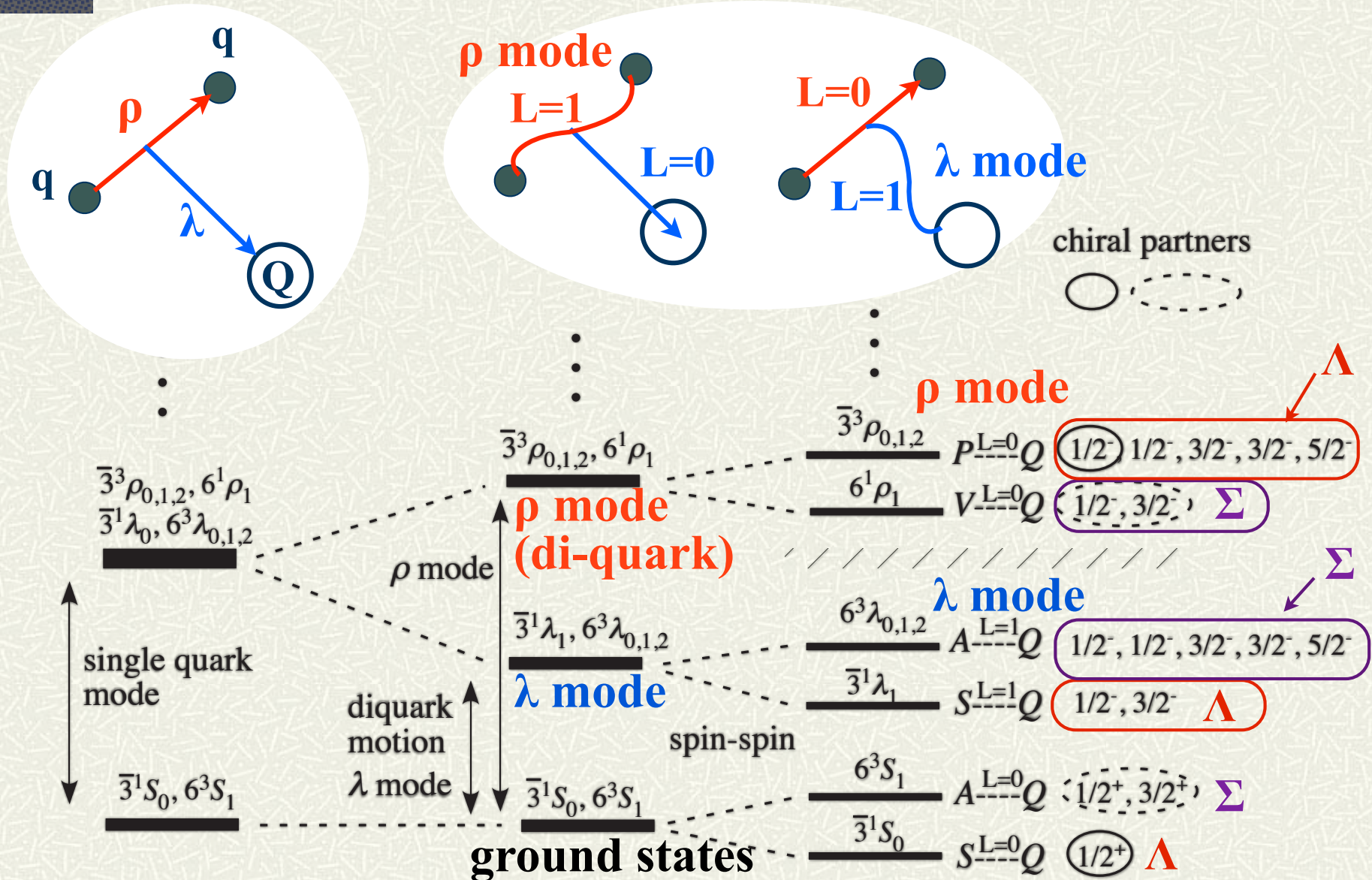
Alexandrou, de Forcrand, Lucini, PRL 97, 222002 (2006)

$$M(1^+) - M(0^+) \sim 100\text{-}150 \text{ MeV (in } Qqq \text{ system)}$$

Babich, et al., PR D76, 074021 (2007)

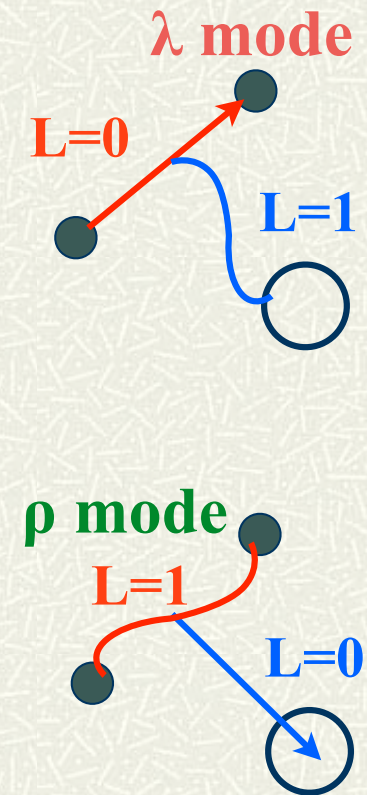
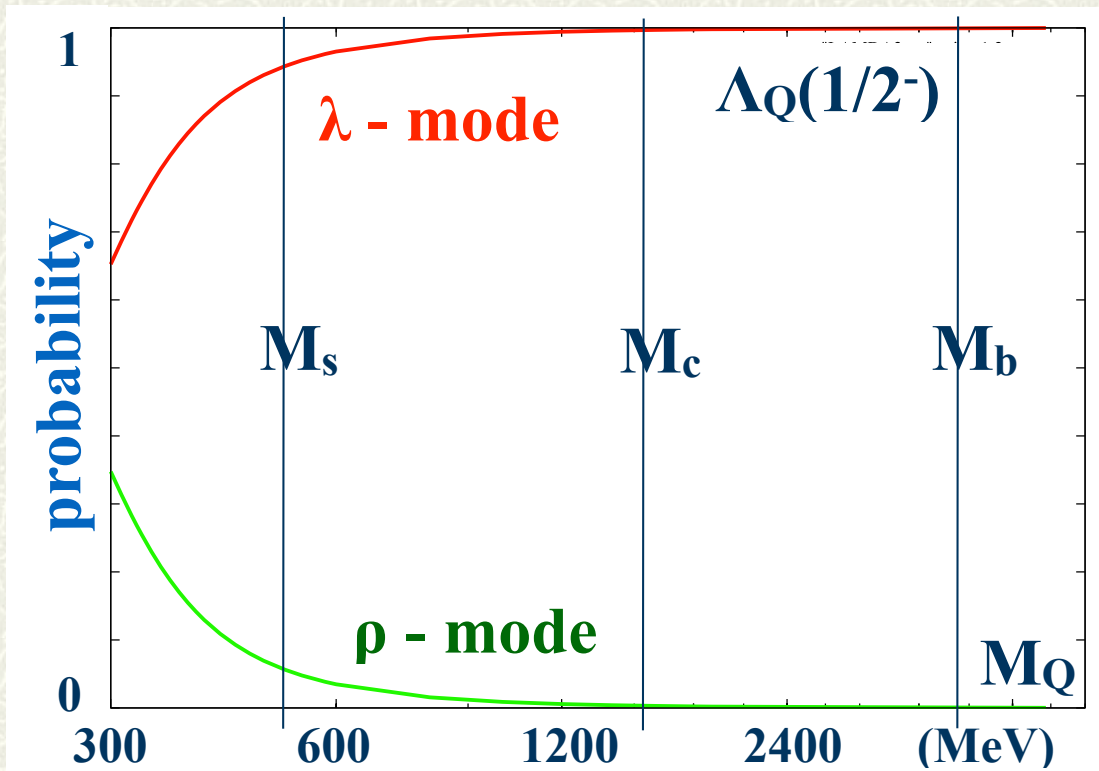
$$M(1^+) - M(0^+) \sim 162 \text{ MeV (Landau gauge)}$$

Diquarks in P-wave Baryons



Diquarks in P-wave Baryons

- ▣ Probabilities of λ and ρ modes v.s. heavy quark mass by a Hamiltonian quark model with spin-spin, spin-orbit and tensor forces

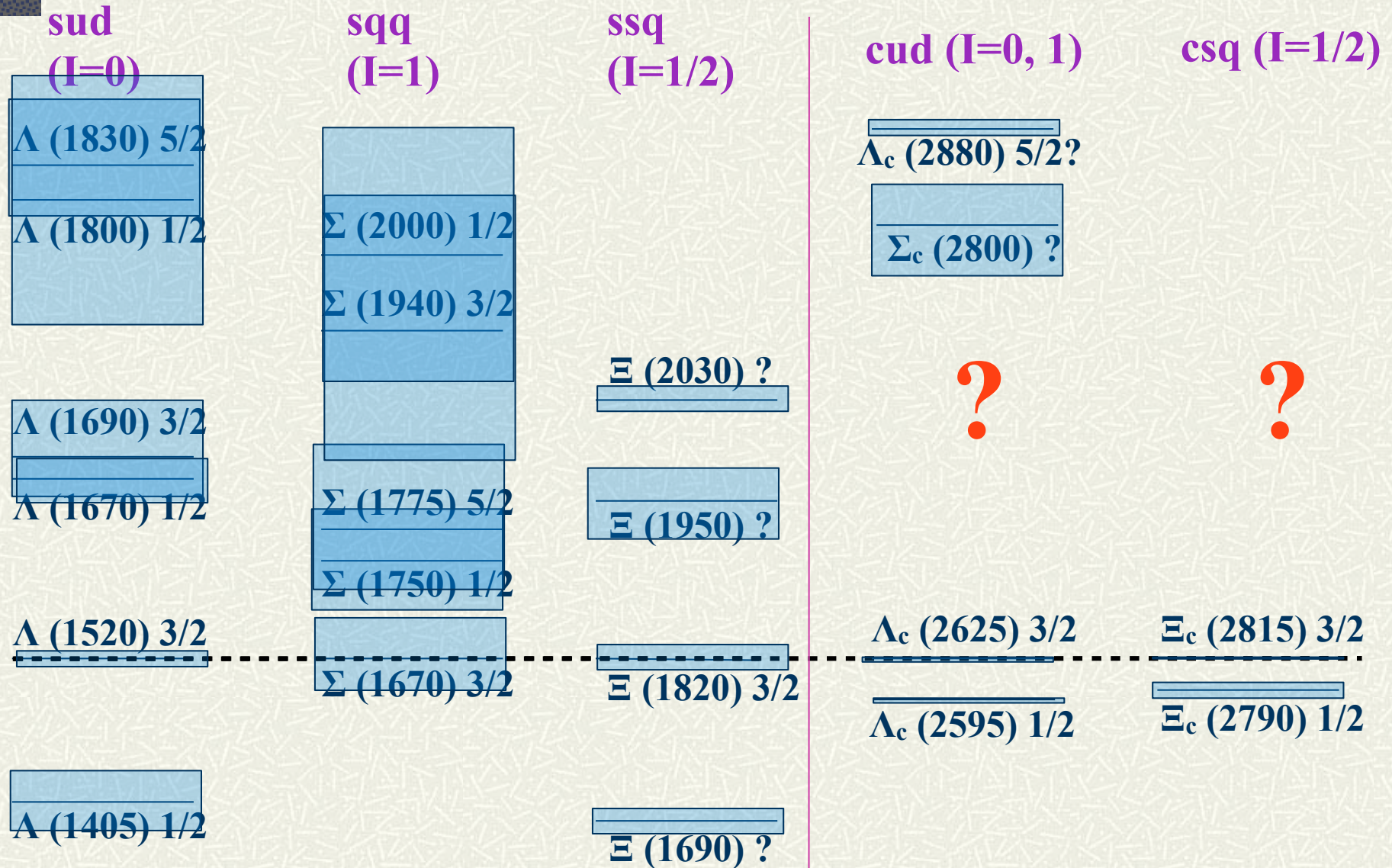


*T. Yoshida, E. Hiyama, A. Hosaka, M. Oka, K. Sabato
PRD 92, 114029-1-19 (2015)*

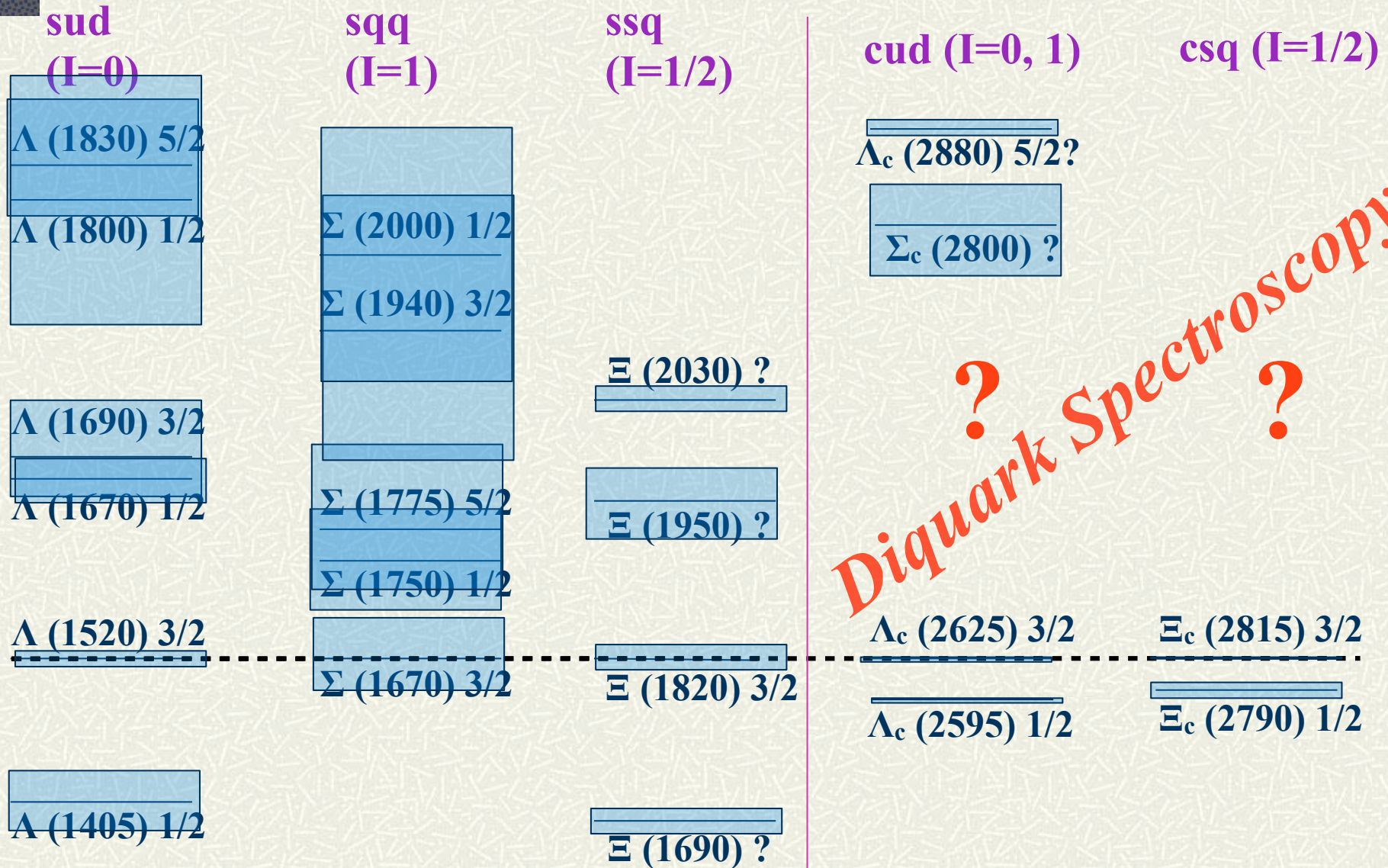
Diquarks in P-wave Baryons

sud ($I=0$)	sqq ($I=1$)	ssq ($I=1/2$)	cud ($I=0, 1$)	csq ($I=1/2$)
$\underline{\Lambda (1830) 5/2}$	$(S=1/2)_\rho$		$\underline{\Lambda_c (2880) 5/2?}$	
$\underline{\Lambda (1800) 1/2}$	$\underline{\Sigma (2000) 1/2}$		$\underline{\Sigma_c (2800) ?}$	
$(S=3/2)_\rho$	$\underline{\Sigma (1940) 3/2}$?	?
$\underline{\Lambda (1690) 3/2}$		$\underline{\Xi (2030) ?}$		
$\underline{\Lambda (1670) 1/2}$		$\underline{\Xi (1950) ?}$		
$\underline{\Lambda (1520) 3/2}$	$\underline{\Sigma (1775) 5/2}$		$\underline{\Lambda_c (2625) 3/2}$	$\underline{\Xi_c (2815) 3/2}$
	$\underline{\Sigma (1750) 1/2}$		$\underline{\Lambda_c (2595) 1/2}$	$\underline{\Xi_c (2790) 1/2}$
	$\underline{\Sigma (1670) 3/2}$	$\underline{\Xi (1820) 3/2}$		
$(S=1/2)_\lambda$	$(S=3/2)_\lambda$		$(S=1/2)_\lambda$	
$\underline{\Lambda (1405) 1/2}$		$\underline{\Xi (1690) ?}$		

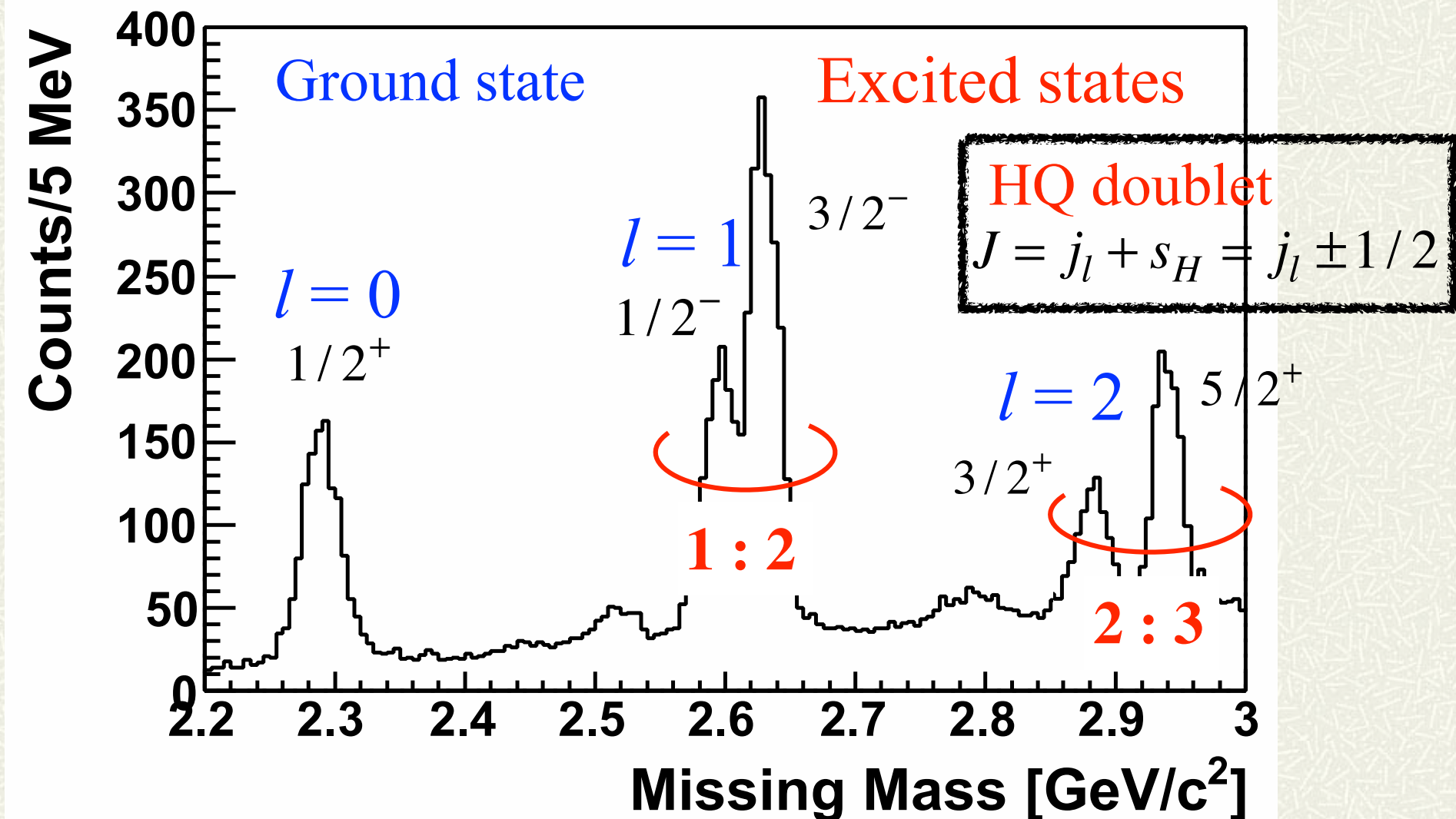
Diquarks in P-wave Baryons



Diquarks in P-wave Baryons



Charm production spectrum



Simulation by K. Shirotori for the J-PARC high momentum BL
S.H. Kim, et al.,
PTEP 2014 (2014) 103D01, and PRD92 (2015) 094021



Quarkonium and Hadron Molecule

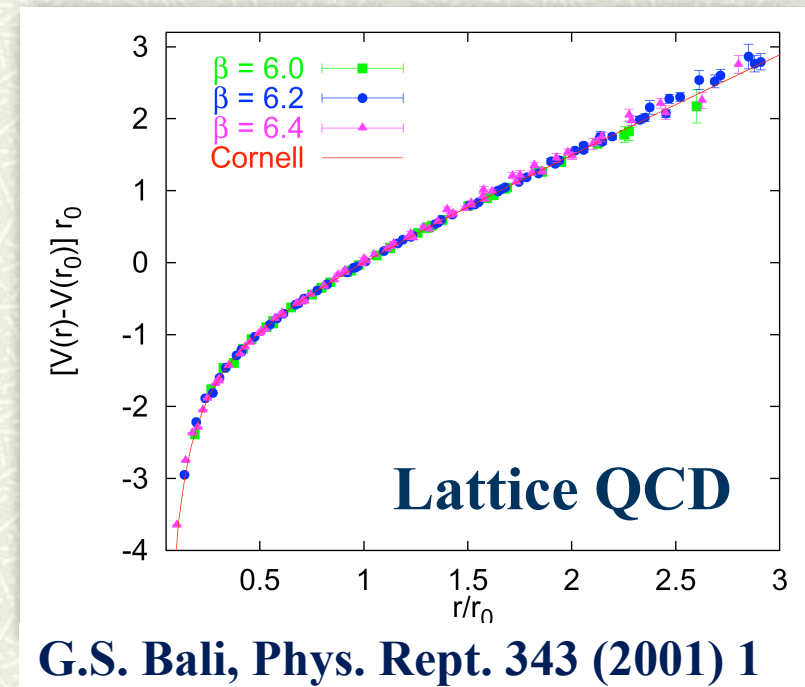
Quarkonium

- # After 50 years since it was born, the quark model gives very good guidelines to classify and interpret the hadron spectrum.
- # The charmonium spectrum is a textbook example.
“hydrogen atom” in QCD

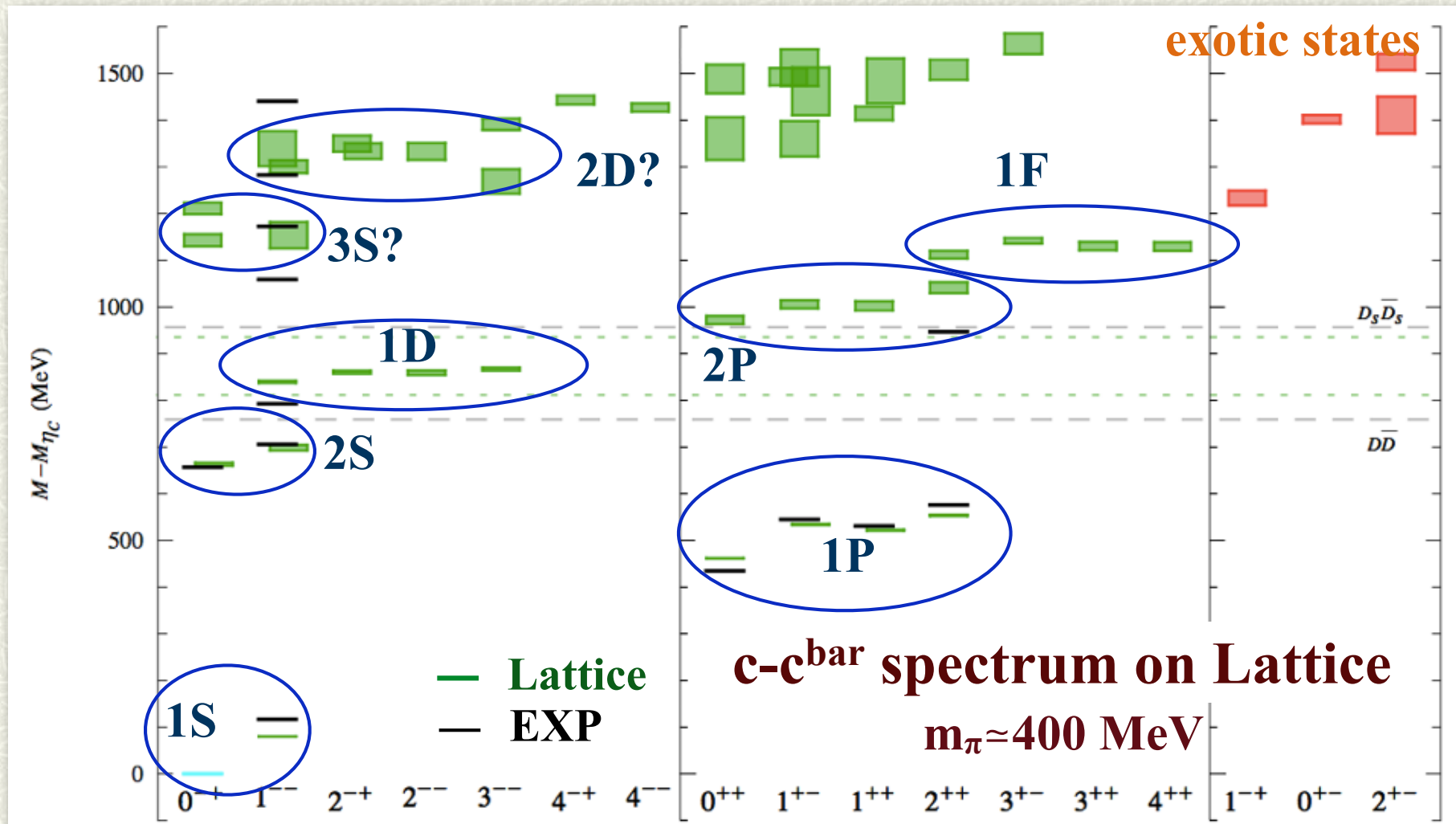
- # The Hamiltonian with a Linear + Coulomb potential

$$V(r) = -\frac{e}{r} + \sigma r$$

E. Eichten, et al., PRL 34 (1975) 369
gives a good fit to the 1S, 1P, 2S, . . .
charmonium (and bottomonium)
states.



Charmonium spectra on Lattice



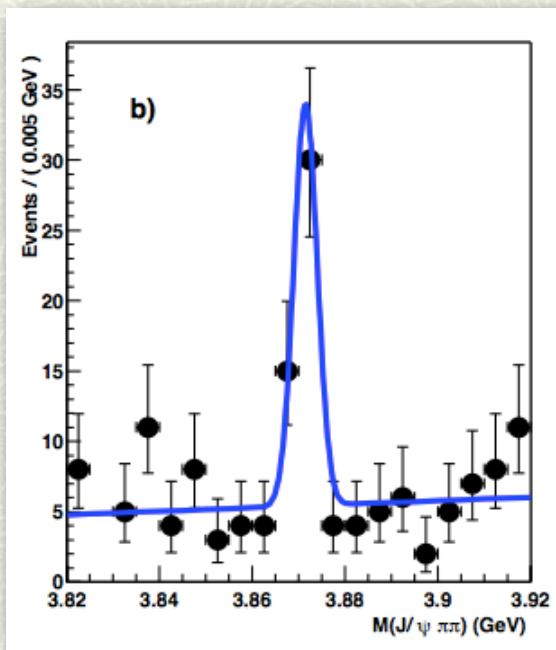
L. Liu, et al. (Hadron Spectrum Collaboration), JHEP 07, 126 (2012)

Charmonium

- # X(3872) found in 2003 by Belle (KEK)
→ *not reproduced by lattice QCD using only $q-q^{bar}$ operators.*
- # Z(3900), Z(4430) etc. : charged hidden charm states

X(3872)

Belle

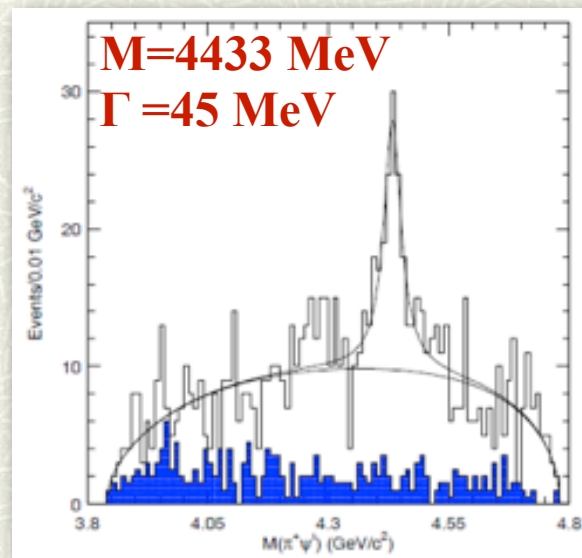


PRL 91 (2003) 262001

M.Oka (Tokyo Tech. and JAEA)

Z_c⁺(4430)

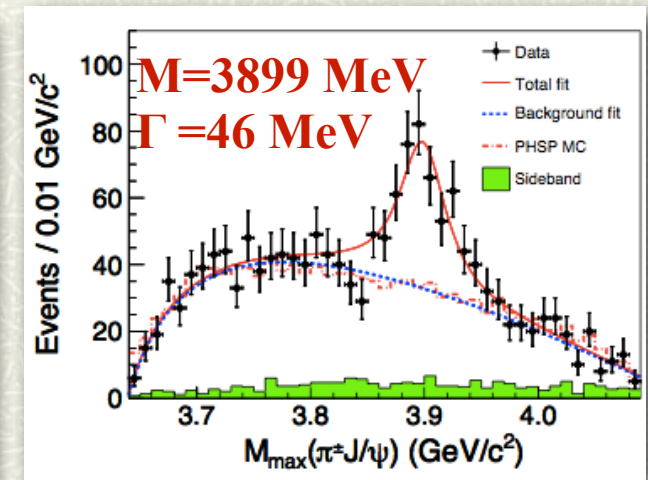
Belle



PRL 100 (2008) 142001

Z_c⁺(3900)

BES III

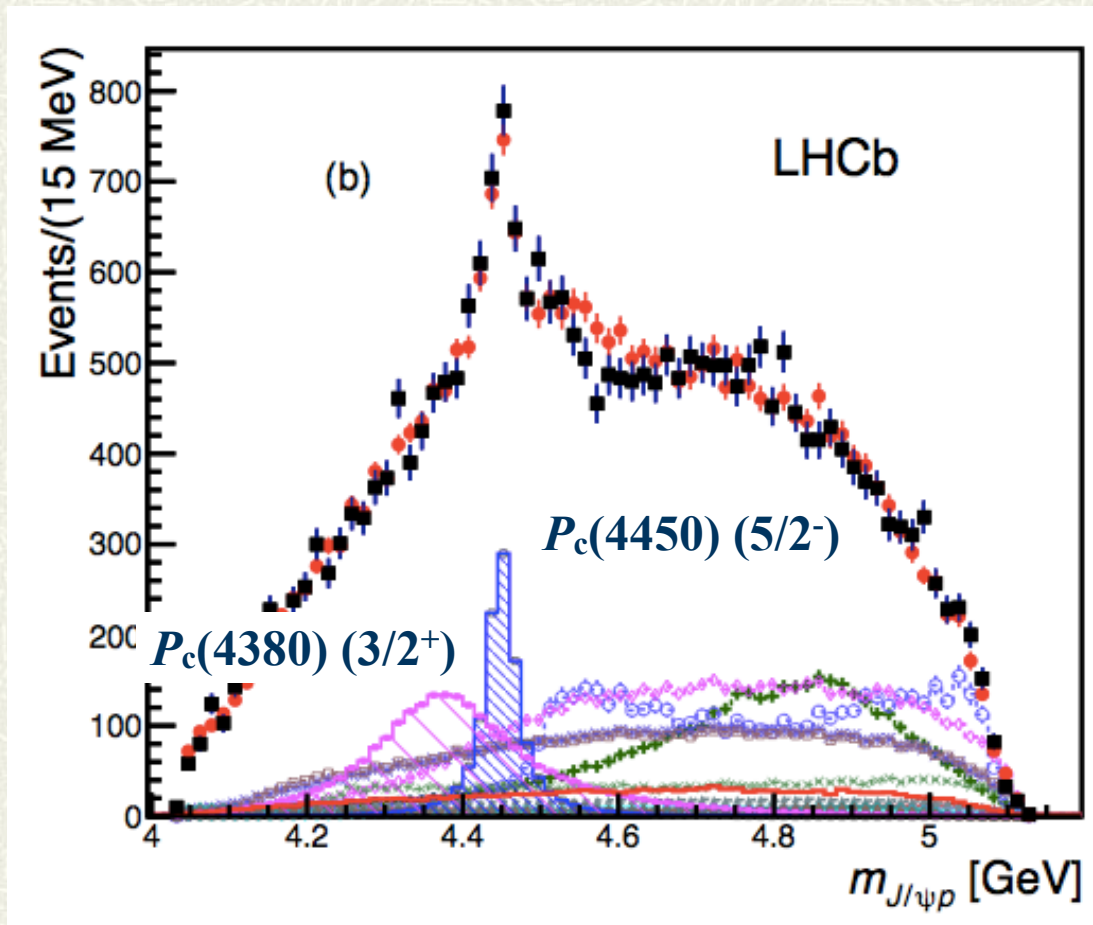


PRL 110 (2013) 252001

P_c Pentaquark@LHCb

$P_c \rightarrow J/\psi + p$ ($cc^{\text{bar}}uud$)

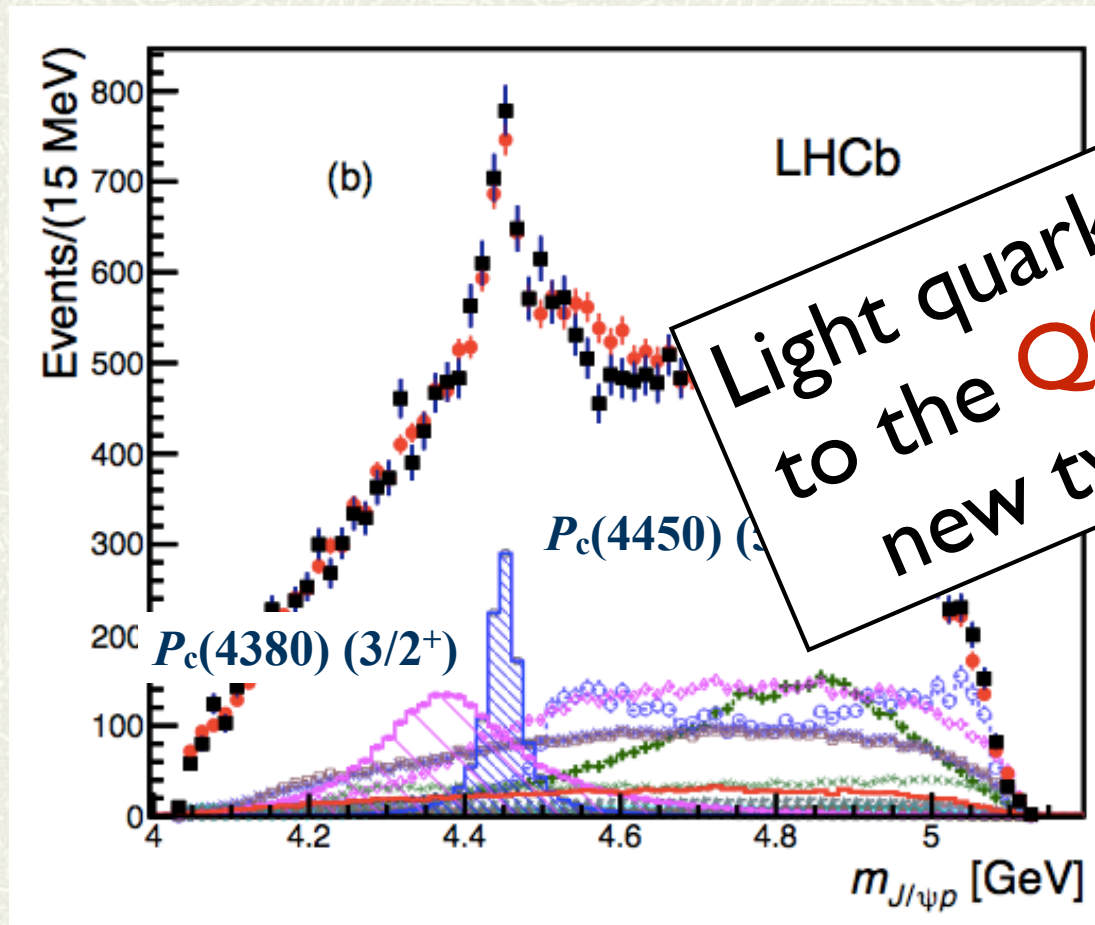
LHCb (*PRL* 115 (2015) 07201) found two penta-quark states with hidden cc^{bar} .



P_c Pentaquark@LHCb

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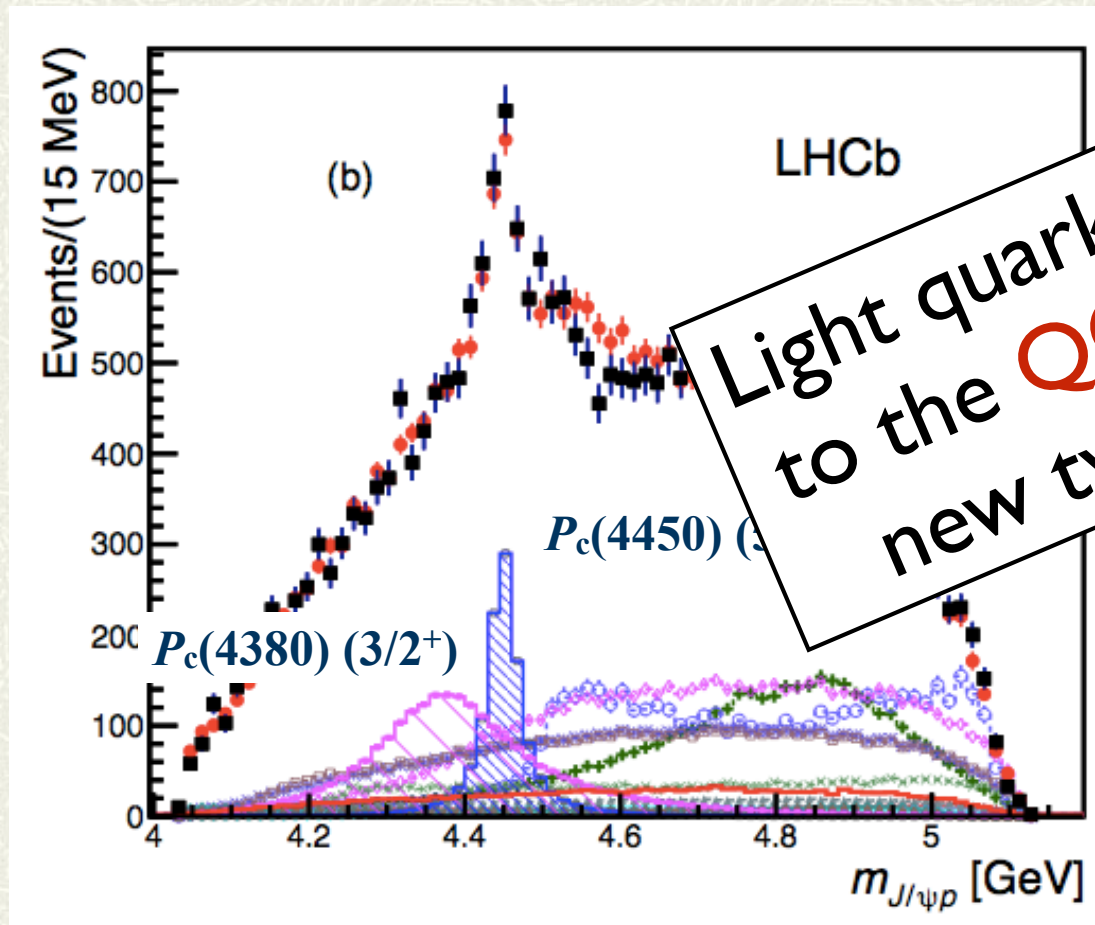


Light quarks tend to stick to the QQ^{bar} and form a new type of hadrons.

P_c Pentaquark@LHCb

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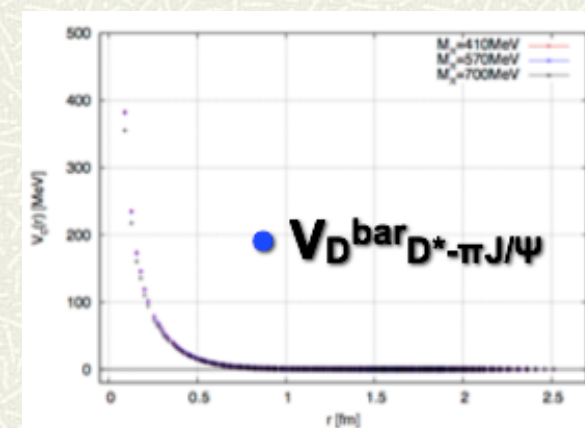
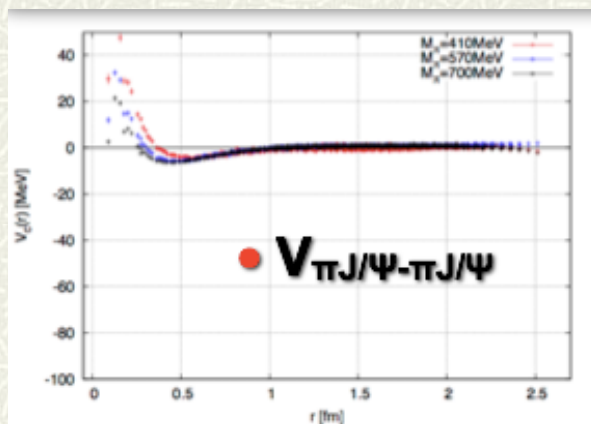
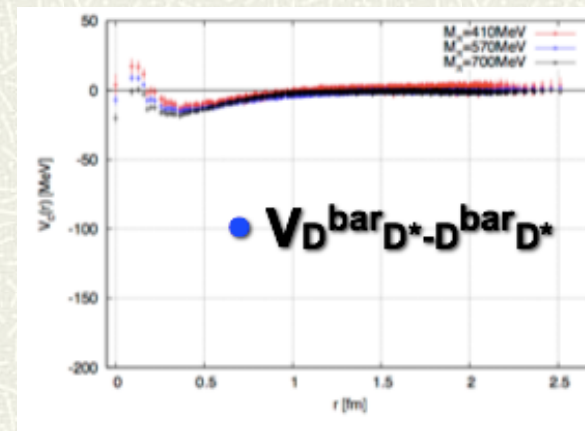
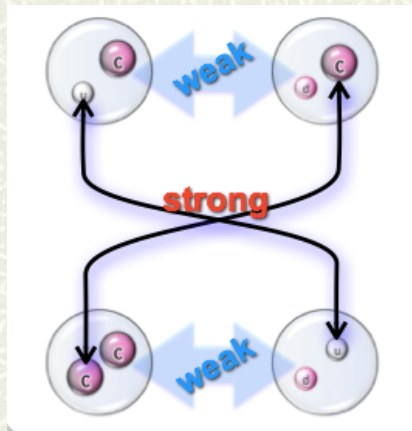


Light quarks tend to stick to the QQ^{bar} and form a new type of hadrons.

Why?

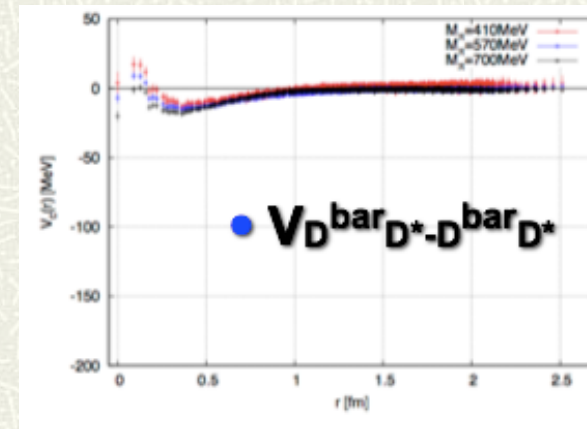
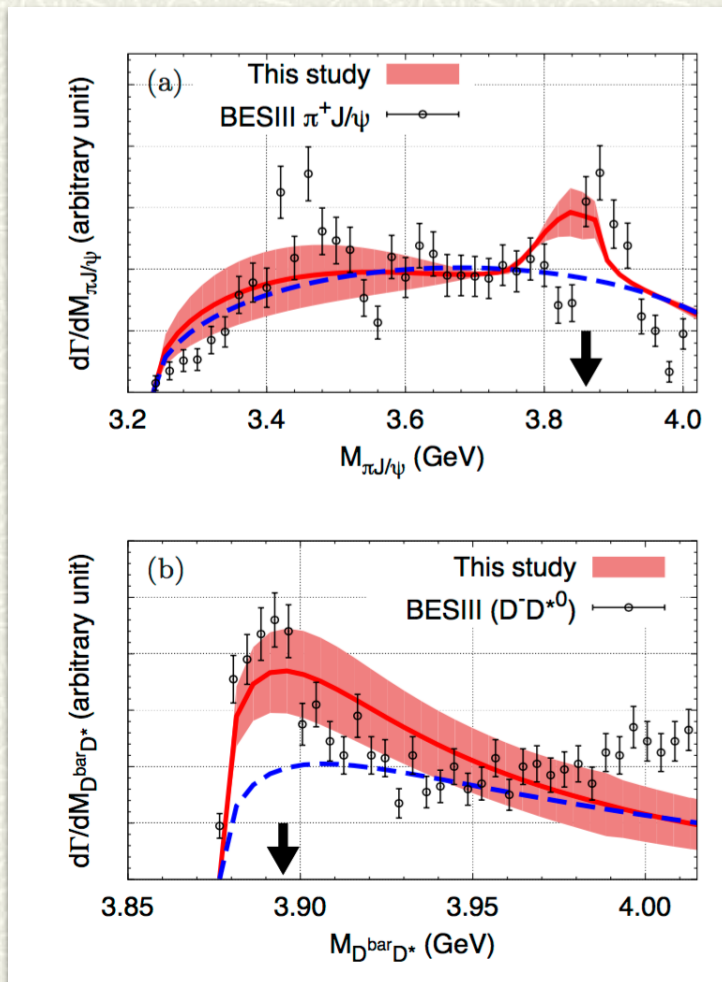
Exotic States on Lattice

- # $Z_c(3900)$ v.s. $(D^{\text{bar}}D^*) + (\pi J/\psi)$ using the HAL QCD method
Y. Ikeda (HALQCD), arXiv:1602.03465



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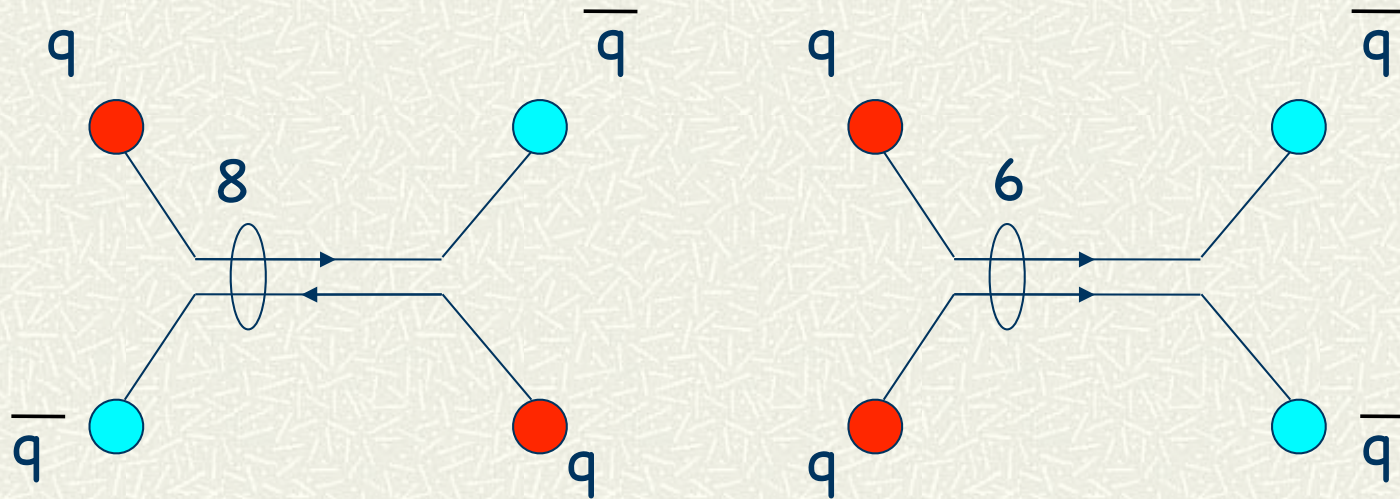


Strong (cusp) enhancement above the $D^{\text{bar}}D^*$ threshold

Exotic Hadrons

Exotics are “Colorful” ! (Lipkin@YKIS06)

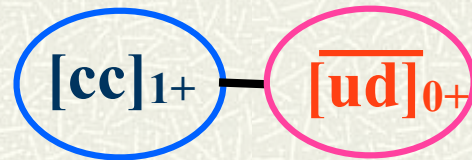
$(qq)_8$ or $(qq)_6$ are allowed only in the multi-quarks.



Double Charm Tetraquark

Double charm meson

$$T_{cc} (ccu^{\text{bar}}d^{\text{bar}}, 1^+, I=0) = [cc]_{1+} [u^{\text{bar}}d^{\text{bar}}]_{0+}$$



- The lowest strong-decay threshold is $D(0^-) - D^*(1^-)$ ($L=0$).
- If the scalar diquark is light enough to make T_{cc} bound below DD^* threshold, T_{cc} will be a stable tetra-quark resonance.

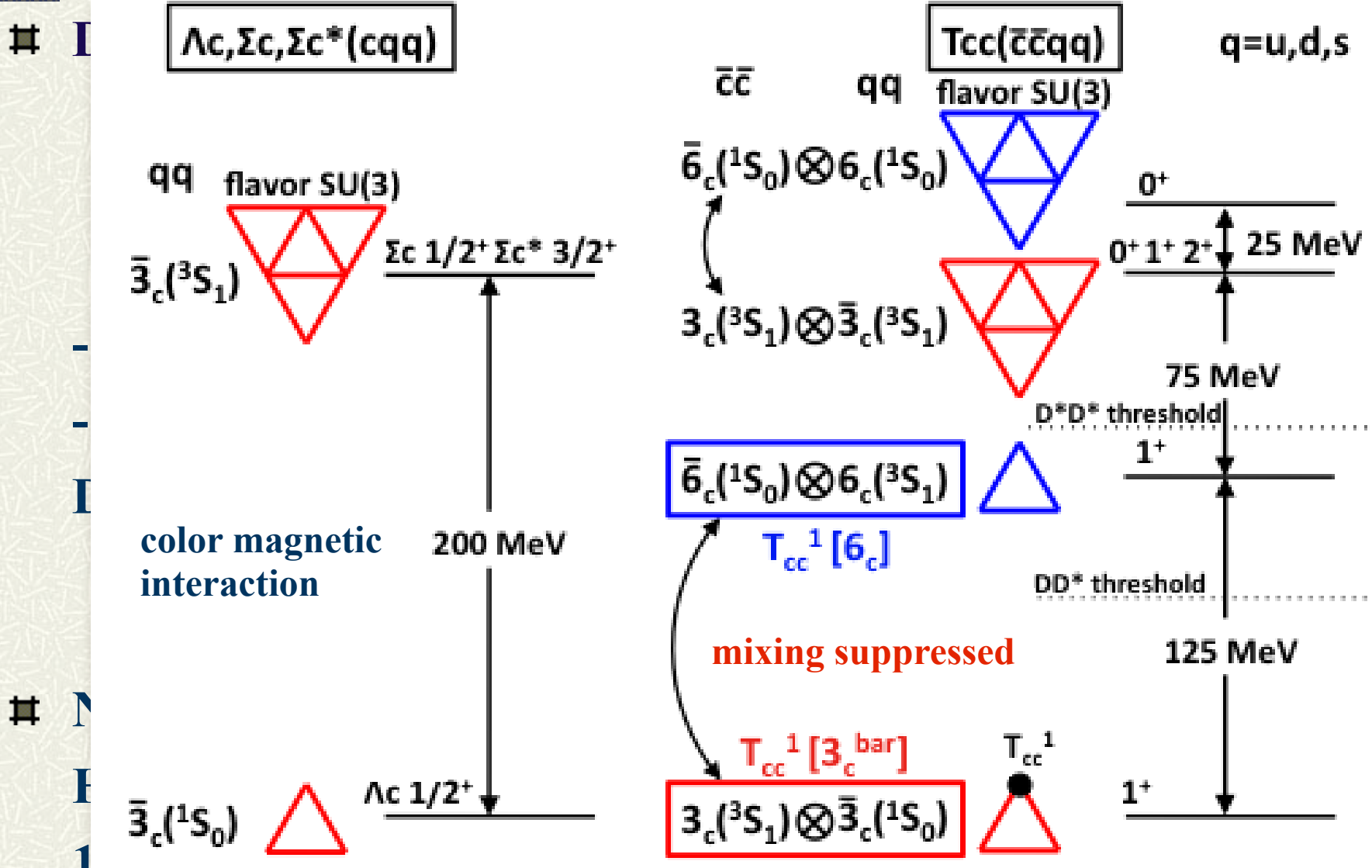
S. Zouzou, et al., Z. Phys. C30 (1986)457

H.J. Lipkin, Phys. Lett. B172 (1986) 242

New possible color correlations

Hyodo, Liu, Oka, Sudoh, Yasui, PLB721 (2013) 56-60, ArXiv
1209.6207

Double Charm Tetraquark



Hyodo, Liu, Oka, Sudoh, Yasui, *PL B721 (2013) 56-60, ArXiv 1209.6207*



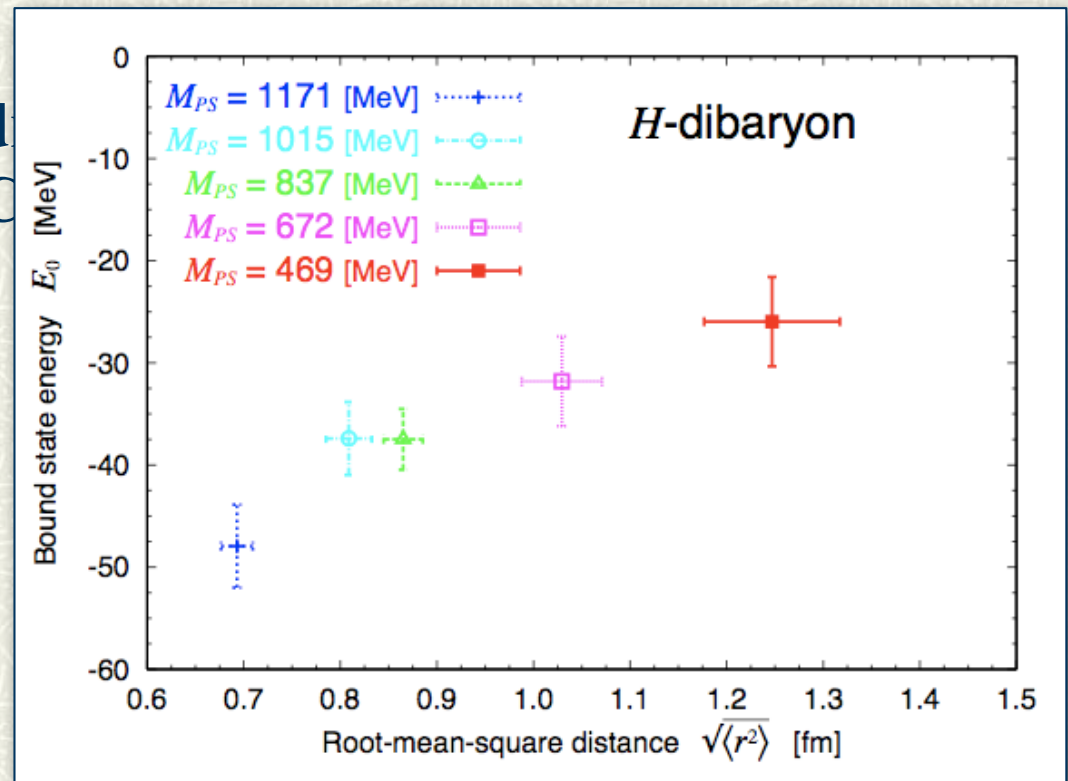
Charmed Dibaryons

Heavy dibaryons

- # **H dibaryon ($= u^2d^2s^2$) predicted by Jaffe (1977)**
New Lattice QCD calculations of H dibaryon
- **Bound H di-baryon in Flavor SU(3) Limit of Lattice QCD**
Takashi Inoue (HAL QCD Collaboration)
PRL 106, 162002 (2011)
- **Evidence for a Bound H di-baryon from Lattice QCD**
S. R. Beane et al. (NPLQCD Collaboration)
PRL 106, 162001 (2011)

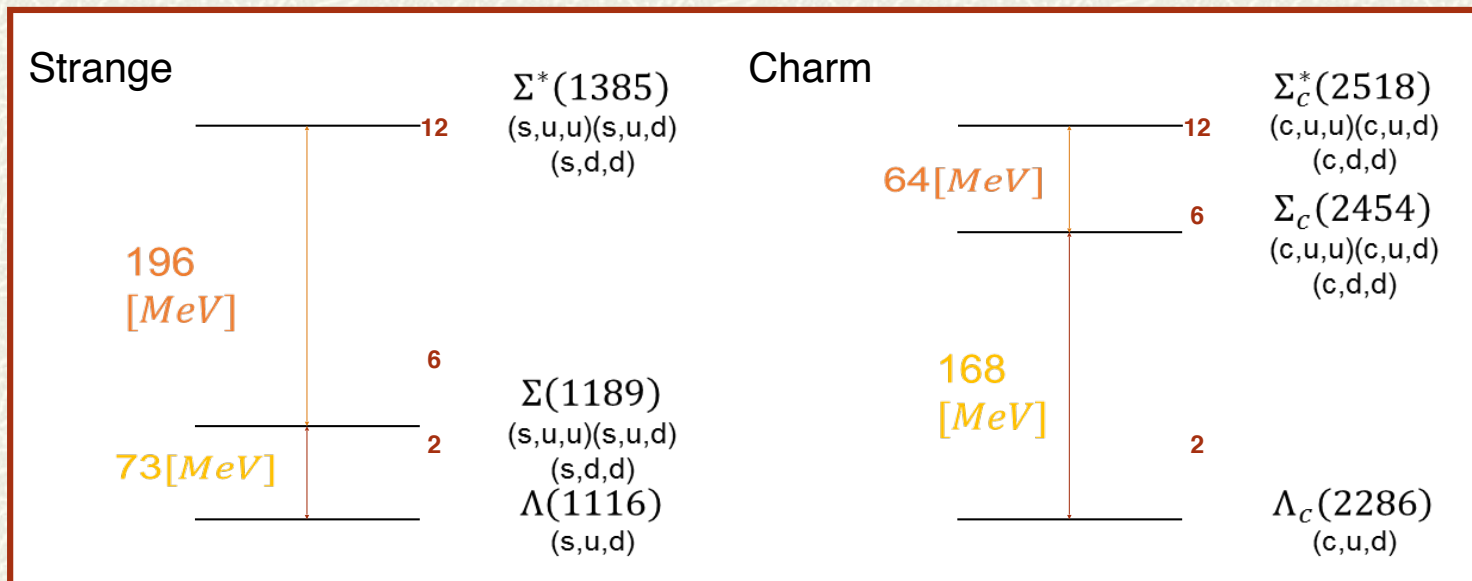
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Heavy dibaryons

- # Coupling between Σ_c and Σ_c^* is enhanced.



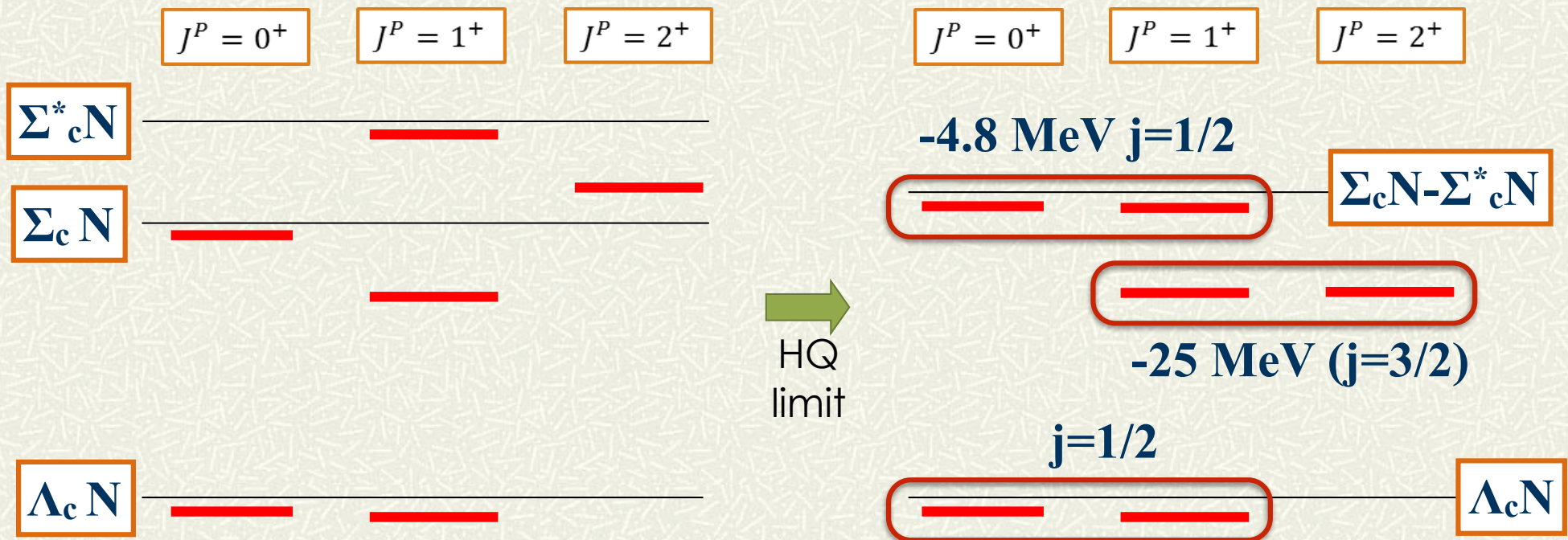
- # $\Lambda_c N$ - $\Sigma_c N$ - $\Sigma_c^* N$ bound/resonance states are considered.
S. Maeda et al., Prog. Theor. Exp. Phys. (2016) 023D02
and in preparation.

Heavy dibaryons

HQ doublets

A shallow bound state of $\Lambda_c N$ with $j=1/2$

A shallow ($j=1/2$) and deep bound ($j=3/2$) state of $\Sigma_c^{(*)} N$.



S. Maeda et al.,

Prog. Theor. Exp. Phys. (2016) 023D02 and in preparation.

Conclusion

- # **Hadron Physics is developing from**
Up-Down \rightarrow Strangeness \rightarrow Charm/ Bottom
Chiral symmetry \rightarrow Confinement and HQ symmetry
- # **It is important to identify effective degrees of freedom in QCD resonances?**
constituent quarks \rightarrow diquarks, glue/string vibration, hadrons
- # **Experiment:**
High statistics data of productions, decays, transitions will reveal the quantum numbers and structures of resonances
- # **Theory:**
Define effective quasi-particles in QCD and make systematic predictions of QCD resonances.