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Theory of the Lamb shift in muonic helium ions

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Outline

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CREMA (Charge Radius Experiment with Muonic Atoms) collaboration 2010-2015

Task: to measure fine and hyperfine structure in light muonic atoms (muonic hydrogen, muonic deuterium, ions of muonic helium...); to determine charge radii of the proton, deuteron, helion, alpha-particle with the accuracy 0.0005 fm.



First measurement of the transition $2P_{3/2}^{F=2} - 2S_{1/2}^{F=1}$ in muonic hydrogen

R. Pohl, A. Antognini, F. Nez et al., Nature **466**, 213 (2010). gave the value of proton charge radius $r_p = 0.84184(67)$ fm (CODATA value $r_p=0.8768(69)$ fm)

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The proton radius puzzle

The proton rms charge radius measured with electrons: 0.8770 ± 0.0045 fm muons: 0.8409 ± 0.0004 fm



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Deuteron charge radius

H/D isotope shift: $r_d^2 - r_p^2 = 3.82007(65)$ fm (C.G. Parthey, RP et al., PRL 104, 233001 (2010)) CODATA 2010 $r_d = 2.1424(21)$ fm $r_{p} = 0.84087(39)$ fm from μH gives $r_{d} = 2.12771(22)$ fm Lamb shift in muonic DEUTERIUM $r_d = 2.1289(12)$ fm Preliminary results A. Antognini, R. Pohl for muonic deuterium $\mu d: \Delta E_{LS}^{exp} = 202.8759(34) \text{ meV}$ If the proton radius puzzle is caused by muon-electron universality breakdown (μ *He*)⁺ and (μ *d*) spectroscopy will reveal it!

The transitions in $(\mu^4 He_2)^+$ and $(\mu^3 He_2)^+$ are planned to measure with $\lambda \in [800, 1000]$ nm

One- and two-loop VP corrections in 1γ interaction



$$\begin{aligned} \mathcal{H}_{B} &= \frac{\mathbf{p}^{2}}{2\mu} - \frac{Z\alpha}{r} - \frac{\mathbf{p}^{4}}{8m_{1}^{3}} - \frac{\mathbf{p}^{4}}{8m_{2}^{3}} + \frac{\pi Z\alpha}{2} \left(\frac{1}{m_{1}^{2}} + \frac{\delta_{I}}{m_{2}^{2}}\right) \delta(\mathbf{r}) - \\ &- \frac{Z\alpha}{2m_{1}m_{2}r} \left(\mathbf{p}^{2} + \frac{\mathbf{r}(\mathbf{rp})\mathbf{p}}{r^{2}}\right) + \frac{Z\alpha}{r^{3}} \left(\frac{1}{4m_{1}^{2}} + \frac{1}{2m_{1}m_{2}}\right) (\mathbf{L}\sigma_{1}). \\ \psi_{200}(r) &= \frac{W^{3/2}}{2\sqrt{2\pi}} e^{-\frac{Wr}{2}} \left(1 - \frac{Wr}{2}\right), \psi_{2lm}(r) = \frac{W^{3/2}}{2\sqrt{6}} e^{-\frac{Wr}{2}} WrY_{lm}(\theta, \phi). \\ &W &= \mu Z\alpha. \end{aligned}$$

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One-loop VP correction to the Lamb shift in 1 γ interaction

$$\begin{split} V_{VP}^{C}(r) &= \frac{\alpha}{3\pi} \int_{1}^{\infty} d\xi \rho(\xi) \left(-\frac{Z\alpha}{r} e^{-2m_{\theta}\xi r} \right), \rho(\xi) = \frac{\sqrt{\xi^{2} - 1(2\xi^{2} + 1)}}{\xi^{4}}, \\ \Delta E_{VP}(2S) &= -\frac{\mu(Z\alpha)^{2}\alpha}{6\pi} \int_{1}^{\infty} \rho(\xi) d\xi \int_{0}^{\infty} x dx \left(1 - \frac{x}{2} \right)^{2} e^{-x \left(1 + \frac{2m_{\theta}\xi}{W} \right)} = \\ \frac{1}{12 \left(1 - k_{1}^{2} \right)^{5/2}} \left[\sqrt{1 - k_{1}^{2}} \left(-168k_{1}^{6} + 272k_{1}^{4} - 49k_{1}^{2} + 6\pi \left(k_{1}^{2} - 1 \right)^{2} \left(14k_{1}^{2} + 3 \right) k_{1} - 28 \right) + \\ &+ 3 \left(56k_{1}^{8} - 128k_{1}^{6} + 75k_{1}^{4} + 10k_{1}^{2} - 4 \right) \ln \left(\frac{1 - \sqrt{1 - k_{1}^{2}}}{k_{1}} \right) \right] = \begin{cases} -2041.9990 \ meV}{-2077.2217 \ meV}, \\ \Delta E_{VP}(2P) &= -\frac{\mu(Z\alpha)^{2}\alpha}{72\pi} \int_{1}^{\infty} \rho(\xi) d\xi \int_{0}^{\infty} x^{3} dx e^{-x \left(1 + \frac{2m_{\theta}\xi}{W} \right)} = \\ &= \frac{1}{\left(1 - k_{1}^{2} \right)^{5/2}} \left[\sqrt{1 - k_{1}^{2}} \left(-120k_{1}^{6} + 184k_{1}^{4} - 23k_{1}^{2} + 6\pi \left(k_{1}^{2} - 1 \right)^{2} \left(10k_{1}^{2} + 3 \right) k_{1} - 32 \right) \\ &+ 3 \left(40k_{1}^{8} - 88k_{1}^{6} + 45k_{1}^{4} + 10k_{1}^{2} - 4 \right) \ln \left(\frac{1 - \sqrt{1 - k_{1}^{2}}}{k_{1}} \right) \right] = \begin{cases} -400.1128 \ meV}{-411.4486 \ meV}, \\ \Delta E_{VP}(2P - 2S) = \begin{cases} 1641.8862 \ meV}{1665.7730 \ meV}. \end{cases} \end{split}$$

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Two-loop VP correction to the Lamb shift in 1γ interaction

$$\frac{1}{k^2} \rightarrow \frac{\alpha}{3\pi} \int_1^\infty \rho(\xi) d\xi \frac{1}{k^2 + 4m_\theta^2 \xi^2}$$

$$V_{VP-VP}^{C}(r) = \frac{\alpha^2}{9\pi^2} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \left(-\frac{Z\alpha}{r}\right) \frac{1}{(\xi^2 - \eta^2)} \left(\xi^2 e^{-2m_{\theta}\xi r} - \eta^2 e^{-2m_{\theta}\eta r}\right)$$

$$\Delta E_{VP-VP}(2P-2S) = \begin{cases} 3.7207 \ meV \\ 3.7997 \ meV \end{cases}.$$

$$\Delta V_{VP-MVP}(r) = -\frac{4(Z\alpha)\alpha^2}{45\pi^2 m_1^2} \int_1^\infty \rho(\xi) d\xi \left[\pi \delta(\mathbf{r}) - \frac{m_\theta^2 \xi^2}{r} e^{-2m_\theta \xi r}\right]$$

$$\Delta E_{VP-MVP}(2P-2S) = \begin{cases} 0.0022 \ meV \\ 0.0023 \ meV \end{cases}.$$

$$\Delta V_{2-loop VP}^{C} = -\frac{2}{3} \frac{Z\alpha}{r} \left(\frac{\alpha}{\pi}\right)^{2} \int_{0}^{1} \frac{f(v)dv}{(1-v^{2})} e^{-\frac{2m_{\theta}r}{\sqrt{1-v^{2}}}}$$
$$\Delta E_{2-loop VP}(2P-2S) = \begin{cases} 7.6863 \ meV \\ 7.7696 \ meV \end{cases}.$$

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*Precision physics and Fundamental Physical Constants, 1-5 December 2014, JINR, A.A. Krutov et. al. "Theory of the Lamb shift ... **



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$$\begin{split} V_{VP-VP-VP}^{\mathcal{C}}(r) &= -\frac{Z\alpha}{r} \frac{\alpha^3}{(3\pi)^3} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta d\eta \int_1^\infty \rho(\zeta) d\zeta \times \\ &\times \left[e^{-2m_{\theta}\zeta r} \frac{\zeta^4}{(\xi^2 - \zeta^2)(\eta^2 - \zeta^2)} + e^{-2m_{\theta}\xi r} \frac{\xi^4}{(\zeta^2 - \xi^2)(\eta^2 - \xi^2)} + e^{-2m_{\theta}\eta r} \frac{\eta^4}{(\xi^2 - \eta^2)(\zeta^2 - \eta^2)} \right], \\ V_{VP-2-loop}^{\mathcal{C}} V_{P} &= -\frac{4\mu\alpha^3(Z\alpha)}{9\pi^3} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \frac{f(\eta)d\eta}{\eta} \frac{1}{r(\eta^2 - \xi^2)} \left(\eta^2 e^{-2m_{\theta}\eta r} - \xi^2 e^{-2m_{\theta}\xi r} \right), \\ \Delta E_{VP-VP-VP}(2P - 2S) &= \begin{cases} 0.0085 \ meV \\ 0.0088 \ meV \end{cases}, \\ \Delta E_{VP-2-loop} \ VP(2P - 2S) &= \begin{cases} 0.0359 \ meV \\ 0.0366 \ meV \end{cases}. \end{split}$$

There exists also a contribution with three-loop vacuum polarization operator. It was calculated in

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T. Kinoshita and M. Nio, Phys. Rev. Lett. 62, 3240 (1999); Phys. Rev. D60, 053008 (1999).

The Wichmann-Kroll correction to the Lamb shift



$$\Delta V^{WK}(r) = \frac{\alpha(Z\alpha)^3}{\pi r} \int_0^\infty \frac{d\zeta}{\zeta^4} e^{-2m_\theta \zeta r} \left[-\frac{\pi^2}{12} \sqrt{\zeta^2 - 1\theta(\zeta - 1)} + \int_0^\zeta dx \sqrt{\zeta^2 - x^2} t^{WK}(x) \right] dx + \Delta E^{WK}(2P - 2S) = \begin{cases} -0.0197 \ meV \\ -0.0199 \ meV \end{cases}$$

One-loop vacuum polarization corrections to the Breit Hamiltonian

$$\begin{split} \Delta V_{VP}^{B}(r) &= \frac{\alpha}{3\pi} \int_{1}^{\infty} \rho(\xi) d\xi \sum_{i=1}^{4} \Delta V_{i,VP}^{B}(r), \\ \Delta V_{1,VP}^{B} &= \frac{Z\alpha}{8} \left(\frac{1}{m_{1}^{2}} + \frac{\delta_{i}}{m_{2}^{2}} \right) \left[4\pi \delta(\mathbf{r}) - \frac{4m_{\theta}^{2}\xi^{2}}{r} e^{-2m_{\theta}\xi r} \right], \\ \Delta V_{2,VP}^{B} &= -\frac{Z\alpha m_{\theta}^{2}\xi^{2}}{m_{1}m_{2}r} e^{-2m_{\theta}\xi r} (1 - m_{\theta}\xi r), \\ \Delta V_{3,VP}^{B} &= -\frac{Z\alpha}{2m_{1}m_{2}} p_{i} \frac{e^{-2m_{\theta}\xi r}}{r} \left[\delta_{ij} + \frac{r_{i}r_{j}}{r^{2}} (1 + 2m_{\theta}\xi r) \right] p_{j}, \\ \Delta V_{4,VP}^{B} &= \frac{Z\alpha}{r^{3}} \left(\frac{1}{4m_{1}^{2}} + \frac{1}{2m_{1}m_{2}} \right) e^{-2m_{\theta}\xi r} (1 + 2m_{\theta}\xi r) (\mathbf{L}\sigma_{1}). \\ \Delta E_{1,VP}^{B}(2P - 2S) &= \begin{cases} -0.8670 \ meV \\ -0.8931 \ meV, \end{cases} \\ \Delta E_{2,VP}^{B}(2P - 2S) &= \begin{cases} 0.0150 \ meV \\ 0.0116 \ meV, \end{cases} \\ \Delta E_{3,VP}^{B}(2P - 2S) &= \begin{cases} 0.0281 \ meV \\ 0.0219 \ meV, \end{cases} \\ \Delta E_{4,VP}^{B}(2P - 2S) &= \begin{cases} -0.0860 \ meV \\ -0.0876 \ meV \end{cases}. \end{split}$$

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Relativistic and VP corrections in second order perturbation theory



$$\begin{split} \Delta E_{SOPT}^{VP} &= \langle \psi | \Delta V_{VP}^{C} \tilde{G} \Delta V_{VP}^{C} | \psi \rangle + 2 \langle \psi | \Delta V^{B} \tilde{G} \Delta V_{VP}^{C} | \psi \rangle . \\ \tilde{G}(2S) &= -\frac{Z \alpha \mu^{2}}{4x_{1}x_{2}} e^{-\frac{x_{1}+x_{2}}{2}} \frac{1}{4\pi} g_{2S}(x_{1}, x_{2}), \\ g_{2S}(x_{1}, x_{2}) &= 8x_{<} - 4x_{<}^{2} + 8x_{>} + 12x_{<}x_{>} - 26x_{<}^{2}x_{>} + 2x_{<}^{3}x_{>} - 4x_{>}^{2} - 26x_{<}x_{>}^{2} + 23x_{<}^{2}x_{>}^{2} - -x_{<}^{3}x_{>}^{2} + 2x_{<}x_{>}^{3} - x_{<}^{2}x_{>}^{3} + 4e^{x}(1 - x_{<})(x_{>} - 2)x_{>} + 4(x_{<} - 2)x_{<}(x_{>} - 2)x_{>} \times \\ &\times [-2C + Ei(x_{<}) - \ln(x_{<}) - \ln(x_{>})], \end{split}$$

$$\tilde{G}(2P) = -\frac{Z\alpha\mu^2}{36x_1^2x_2^2}e^{-\frac{x_1+x_2}{2}}\frac{3}{4\pi}\frac{(\mathbf{x}_1\mathbf{x}_2)}{x_1x_2}g_{2P}(x_1,x_2),$$

$$\begin{split} g_{2P}(x_1,x_2) &= 24x_<^3 + 36x_<^3x_> + 36x_<^3x_>^2 + 24x_>^3 + 36x_< x_>^3 + 36x_<^2x_>^3 + 49x_<^3x_>^3 - 3x_<^4x_>^3 - \\ &- 12e_<^x(2+x_<+x_<^2)x_>^3 - 3x_<^3x_>^4 + 12x_<^3x_>^3[-2C+Ei(x_<) - \ln(x_<) - \ln(x_>)], \end{split}$$

First term $<\psi|\Delta V^{C}_{V\!P}\tilde{G}\Delta V^{C}_{V\!P}|\psi>$

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$$\begin{split} \Delta E_{SOPT}^{VP,VP}(2S) &= -\frac{\mu\alpha^2(Z\alpha)^2}{72\pi^2} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \times \\ \int_0^\infty \left(1 - \frac{x}{2}\right) e^{-x\left(1 - \frac{2m_\theta\xi}{W}\right)} dx \int_0^\infty \left(1 - \frac{x'}{2}\right) e^{-x'\left(1 - \frac{2m_\theta\eta}{W}\right)} dx' g_{2S}(x, x') = \begin{cases} -1.8640 \ meV}{-1.9017 \ meV}, \\ \Delta E_{SOPT}^{VP,VP}(2P) &= -\frac{\mu\alpha^2(Z\alpha)^2}{7776\pi^2} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \times \\ &\times \int_0^\infty e^{-x\left(1 + \frac{2m_\theta\xi}{W}\right)} dx \int_0^\infty e^{-x'\left(1 + \frac{2m_\theta\eta}{W}\right)} dx' g_{2P}(x, x') = \begin{cases} -0.1867 \ meV}{-0.1942 \ meV}, \end{cases} \end{split}$$

Second term $<\psi|\Delta V^{B}\tilde{G}\Delta V^{C}_{V\!P}|\psi>$

$$<\psi|\frac{\mathbf{p}^{4}}{(2\mu)^{2}}\sum_{m}'\frac{|\psi_{m}><\psi_{m}|}{E_{2}-E_{m}}\Delta V_{VP}^{C}|\psi>=<\psi|(E_{2}+\frac{Z\alpha}{r})(\hat{H}_{0}+\frac{Z\alpha}{r})\sum_{m}'\frac{|\psi_{m}><\psi_{m}|}{E_{2}-E_{m}}\Delta V_{VP}^{C}|\psi>=$$

$$= \langle \psi | \left(E_2 + \frac{Z\alpha}{r} \right)^2 \tilde{G} \Delta V_{VP}^C | \psi \rangle - \langle \psi | \frac{Z\alpha}{r} \Delta V_{VP}^C | \psi \rangle + \langle \psi | \frac{Z\alpha}{r} | \psi \rangle \langle \psi | \Delta V_{VP}^C | \psi \rangle .$$
$$\Delta E_{SOPT}^{B, VP} (2P - 2S) = \begin{cases} 1.4192 \text{ meV} \\ 1.4682 \text{ meV} \end{cases}.$$

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Three-loop vacuum polarization correction in second order perturbation theory



$$\Delta E_{SOPT}^{VP-VP, VP}(2S) = -\frac{\mu\alpha^{3}(Z\alpha)^{2}}{108\pi^{3}} \int_{1}^{\infty} \rho(\xi)d\xi \int_{1}^{\infty} \rho(\eta)d\eta \int_{1}^{\infty} \rho(\zeta)d\zeta \int_{0}^{\infty} dx(1-\frac{x}{2}) \times \int_{0}^{\infty} dx'(1-\frac{x'}{2})e^{-x'(1+\frac{2me\zeta}{W})} \frac{1}{\xi^{2}-\eta^{2}} \left[\xi^{2}e^{-x(1+\frac{2me\zeta}{W})} - \eta^{2}e^{-x(1+\frac{2me\eta}{W})}\right] g_{2S}(x,x') = \begin{cases} -0.0104 \ meV \\ -0.0107 \ meV \end{cases}$$

$$\Delta E_{SOPT}^{2-loop VP, VP}(2S) = -\frac{\mu\alpha^3(Z\alpha)^2}{18\pi^3} \int_0^1 \frac{f(v)dv}{1-v^2} \int_1^\infty \rho(\xi)d\xi \times$$

$$\times \int_0^\infty dx \left(1 - \frac{x}{2}\right) e^{-x(1 + \frac{2m_e}{\sqrt{1 - v^2 W}})} \int_0^\infty dx' \left(1 - \frac{x'}{2}\right) e^{-x'(1 + \frac{2m_e \xi}{W})} g_{2S}(x, x') = \begin{cases} -0.0168 \ meV, \\ -0.0171 \ meV, \end{cases}$$

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Three-loop vacuum polarization correction in third order PT

$$\Delta E = <\psi_2 |\Delta V^C \tilde{G} \Delta V^C \tilde{G} \Delta V^C |\psi_2> - <\psi_2 |\Delta V^C |\psi_2> <\psi_2 |\Delta V^C \tilde{G} \tilde{G} \Delta V^C |\psi_2>.$$

$$\begin{split} \Delta E_{TOPT,1}(2S) &= -\frac{\mu Z^2 \alpha^5}{864 \pi^3} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \int_1^\infty \rho(\zeta) d\zeta \int_0^\infty \left(1 - \frac{x}{2}\right) e^{-x(1+2m_{\theta}\xi/W)} dx \times \\ \int_0^\infty \left(1 - \frac{x''}{2}\right) e^{-x''(1+2m_{\theta}\zeta/W)} dx'' \int_0^\infty \frac{dx'}{x'} e^{-x''(1+2m_{\theta}\zeta/W)} g(x,x') g(x',x'') = \begin{cases} -0.0044 \ meV \\ -0.0045 \ meV \end{cases}, \\ \Delta E_{TOPT,2}(2S) &= \frac{\alpha^2}{288\pi^2} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \int_0^\infty \left(1 - \frac{x}{2}\right) e^{-x(1+2m_{\theta}\xi/W)} dx \times \\ \end{cases}$$

$$\int_0^\infty \left(1 - \frac{x''}{2}\right) e^{-x''(1+2m_\theta\eta/W)} dx'' \int_0^\infty dx' g(x,x')g(x',x'') \begin{cases} 2041.9990 \text{ meV} \\ 2077.2217 \text{ meV} \end{cases} = \begin{cases} 0.0037 \text{ meV} \\ 0.0038 \text{ meV} \end{cases}.$$

$$\text{Replacement } \Delta V_{VP}^{C} \rightarrow \Delta H_{B}, \Delta V_{VP}^{C} \rightarrow \Delta V_{VP, VP}^{C}, \Delta H_{B,1} = (\pi Z \alpha/2)(1/m_{1}^{2} + \delta_{I}/m_{2}^{2})\delta(\mathbf{r}_{A}^{2})$$

$$\Delta E_{SOPT}^{VP-VP,\Delta H_{B,1}}(2S) = \frac{\mu^3(Z\alpha)^4\alpha^2}{144\pi^2} \left(\frac{1}{m_1^2} + \frac{\delta_I}{m_2^2}\right) \int_1^\infty \rho(\xi)d\xi \int_1^\infty \rho(\eta)d\eta \frac{1}{\xi^2 - \eta^2} \times \dots$$

$$\int_0^\infty \left(1 - \frac{x}{2}\right) dx [4x(x-2)(\ln x + C) + x^3 - 13x^2 + 6x + 4] \left[\xi^2 e^{-x(1 + \frac{2me\xi}{W})} - \eta^2 e^{-x(1 + \frac{2me\eta}{W})}\right] =$$

$$= \begin{cases} 0.0050 \text{ meV} \\ 0.0051 \text{ meV} \end{cases}$$

$$\Delta E_{SOPT}^{2-loop VP,\Delta H_{B,1}}(2S) = \frac{\mu^3(Z\alpha)^4\alpha^2}{24\pi^2} \left(\frac{1}{m_1^2} + \frac{\delta_l}{m_2^2}\right) \int_0^1 \frac{f(v)dv}{1-v^2} \times$$

$$\int_0^\infty (1-\frac{x}{2}) dx [4x(x-2)(\ln x+C)+x^3-13x^2+6x+4] e^{-x(1+\frac{2me}{W\sqrt{1-v^2}}} = \begin{cases} 0.0056 \ meV \\ 0.0058 \ meV \end{cases}$$

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Precision physics and Fundamental Physical Constants, 1-5 December 2014, JINR, A.A. Krutov et, al. "Theory of the Lamb shift" Replacement $\Delta V_{VP}^C \rightarrow \Delta H_B, \Delta V_{VP}^C \rightarrow \Delta V_{VP,VP}^C, \Delta H_{B,2} = -\mathbf{p}^4 (1/m_1^3 + 1/m_2^3)$ $\Delta E_{SOPT,1}^{VP-VP,\Delta H_{B,2}}(2S) = -\frac{\mu^4 (Z\alpha)^4 \alpha^2}{72\pi^2} \left(\frac{1}{m^3} + \frac{1}{m^3}\right) \int_{1}^{\infty} \rho(\xi) d\xi \int_{1}^{\infty} \rho(\eta) d\eta \frac{1}{\xi^2 - m^2} \times$ $\int_{0}^{\infty} \left(1 - \frac{x}{2}\right) x dx \left(\frac{1}{x} - \frac{1}{8}\right)^{2} \int_{0}^{\infty} (1 - \frac{x'}{2}) dx' g(x, x') \left[\xi^{2} e^{-x'(1 + \frac{2m_{e}\xi}{W})} - \eta^{2} e^{-x'(1 + \frac{2m_{e}\eta}{W})}\right] = 0$ $=\begin{cases} -0.0029 \ meV \\ -0.0031 \ meV \end{cases}$ $\Delta E_{SOPT,2}^{2-loop VP,\Delta H_{B,2}}(2S) = -\frac{\mu^4 (Z\alpha)^4 \alpha^2}{12\pi^2} \left(\frac{1}{m^3} + \frac{1}{m^3}\right) \int_{1}^{\infty} \rho(\xi) d\xi \int_{0}^{1} \frac{f(v) dv}{1 - v^2} \times$ $\int_{0}^{\infty} (1-\frac{x}{2}) x dx \left(\frac{1}{x}-\frac{1}{8}\right)^{2} \int_{0}^{\infty} (1-\frac{x'}{2}) dx' g(x,x') e^{-x'(1+\frac{2m_{\theta}}{W\sqrt{1-v^{2}}})} = \begin{cases} -0.0045 \text{ meV} \\ -0.0047 \text{ meV} \end{cases}$ $\Delta E_{SOPT,3}^{VP-VP,\Delta H_{B,2}}(2S) = -\frac{\mu^4 (Z\alpha)^4 \alpha^2}{18\pi^2} \left(\frac{1}{m^3} + \frac{1}{m^3}\right) \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \frac{1}{\epsilon^2 - m^2} \times$ $\int_{0}^{\infty} \left(1 - \frac{x}{2}\right)^{2} dx \left[\xi^{2} e^{-x(1 + \frac{2m_{\theta}\xi}{W})} - \eta^{2} e^{-x(1 + \frac{2m_{\theta}\eta}{W})}\right] = \begin{cases} -0.0072 \text{ meV} \\ -0.0075 \text{ meV} \end{cases}$ $\Delta E_{SOPT,4}^{2-loop VP,\Delta H_{B,2}}(2S) = -\frac{\mu^4 (Z\alpha)^4 \alpha^2}{3\pi^2} \left(\frac{1}{m^3} + \frac{1}{m^3}\right) \int_1^\infty \frac{f(v)dv}{1 - v^2} \times \frac{f(v)dv}{1 - v^2}$ $\int_{0}^{\infty} \left(1 - \frac{x}{2}\right)^{2} dx e^{-x(1 + \frac{2m_{\theta}}{W\sqrt{1 - v^{2}}})} = \begin{cases} -0.0083 \ meV \\ -0.0086 \ meV \end{cases}$ $\Delta E_{SOPT}^{VP, VP; \Delta V^{B}}(2P-2S) = \begin{cases} 0.0120 \text{ meV} \\ 0.0127 \text{ meV} \end{cases}, \Delta E_{SOPT}^{VP, \Delta V^{B}}(2P-2S) = \begin{cases} -0.0066 \text{ meV} \\ -0.0069 \text{ meV} \end{cases}.$

Nuclear structure correction in 1γ and 2γ interaction



Nuclear structure and one-loop VP correction in second order PT

$$\Delta V_{str}^{VP}(r) = \frac{2}{3}\pi Z\alpha < r_N^2 > \frac{\alpha}{3\pi} \int_1^\infty \rho(\xi) d\xi \left[\delta(\mathbf{r}) - \frac{m_\theta^2 \xi^2}{\pi r} e^{-2m_\theta \xi r} \right]$$

$$\begin{split} \Delta E_{str}^{VP}(2S) &= \frac{\alpha (Z\alpha)^4 < r_N^2 > \mu^3}{36\pi} \int_1^\infty \rho(\xi) d\xi \left[1 - \frac{4m_\theta^2 \xi^2}{W^2} \int_0^\infty x dx (1 - \frac{x}{2})^2 e^{-x(1 + \frac{2m_\theta \xi}{W})} \right] = \\ &= \begin{cases} 1.2493 \ meV \\ 0.9365 \ meV \end{cases}, \end{split}$$

$$\Delta E_{str}^{VP}(2P) = -\frac{\alpha(Z\alpha)^4 \mu^3 < r_N^2 >}{108\pi} \frac{m_e^2}{W^2} \int_1^\infty \xi^2 \rho(\xi) d\xi \int_0^\infty x^3 e^{-x(1+\frac{2m_e\xi}{W})} dx = \begin{cases} -0.0300 \ meV \\ -0.0225 \ meV \end{cases},$$

$$\Delta E_{str}^{VP}(2P-2S) = \begin{cases} -1.2793 \pm 0.0130 \text{ meV} \\ -0.9590 \pm 0.0092 \text{ meV} \end{cases}$$
$$\Delta E_{str,SOPT}^{VP}(2P-2S) = -\frac{\alpha(Z\alpha)^4 \mu^3 < r_N^2 >}{36\pi} \int_1^\infty \rho(\xi) d\xi \times$$
$$\times \int_0^\infty dx e^{-x(1+\frac{2m_e\xi}{W})} (1-\frac{x}{2}) \left[4x(x-2)(\ln x+C) + x^3 - 13x^2 + 6x + 4 \right] = \begin{cases} -2.0083 \text{ meV} \\ -1.5063 \text{ meV} \end{cases}$$

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Nuclear structure and two-loop VP correction



$$\Delta E_{str,SOPT}^{VP, VP(1)}(2S) = \frac{\alpha^2 (Z\alpha)^4 \mu^3 r_N^2}{108\pi^2} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \times \int_0^\infty \left(1 - \frac{x}{2}\right) dx e^{-x\left(1 + \frac{2me\eta}{W}\right)} \left[4x(x-2)(\ln x + C) + x^3 - 13x^2 + 6x + 4\right],$$

$$\Delta E_{str,SOPT}^{VP, VP(2)}(2S) = -\frac{\alpha^2 (Z\alpha)^4 \mu^3 r_N^2 m_e^2}{54\pi^2 W^2} \int_1^\infty \xi^2 \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \times \int_0^\infty \left(1 - \frac{x}{2}\right) dx e^{-x\left(1 + \frac{2me\eta}{W}\right)} \int_0^\infty \left(1 - \frac{x'}{2}\right) dx' e^{-x'\left(1 + \frac{2me\eta}{W}\right)} g_{2S}(x, x').$$

$$\Delta E_{str,SOPT}^{VP, VP}(2P - 2S) = \begin{cases} -0.0086 \ meV \\ -0.0065 \ meV. \end{cases}$$

$$\begin{split} \Delta E_{str, VP}^{2\gamma}(nS) &= -\frac{2\mu^3 \alpha (Z\alpha)^5}{\pi^2 n^3} \int_0^\infty kV(k) dk \int_0^1 \frac{v^2 (1 - \frac{v^2}{2}) dv}{k^2 (1 - v^2) + 4m_e^2}, \\ \Delta E_{str, VP}^{2\gamma}(2P - 2S) &= \begin{cases} 0.2214 \pm 0.0022 \ meV}{0.1270 \pm 0.0013 \ meV}. \end{split}$$

Recoil correction of order $(Z\alpha)^4$

$$\Delta E_{rec}(2P-2S) = \begin{cases} \frac{\mu^3(Z\alpha)^4}{48m_2^2}, & \delta_I = 1\\ \frac{\mu^3(Z\alpha)^4}{12m_2^2}, & \delta_I = 0 \end{cases} = \begin{cases} 0.1265 \ meV\\ 0.2952 \ meV \end{cases}.$$

Recoil correction of order $(Z\alpha)^5$

$$\Delta E_{rec}^{(Z\alpha)^5} = \frac{\mu^3 (Z\alpha)^5}{m_1 m_2 \pi n^3} \Big[\frac{2}{3} \delta_{l0} \ln \frac{1}{Z\alpha} - \frac{8}{3} \ln k_0(n,l) - \frac{1}{9} \delta_{l0} - \frac{7}{3} a_n - \frac{2}{m_2^2 - m_1^2} \delta_{l0}(m_2^2 \ln \frac{m_1}{\mu} - m_1^2 \ln \frac{m_2}{\mu}) \Big],$$

 $\ln k_0(2S) = 2.811769893120563$, $\ln k_0(2P) = -0.030016708630213$,

$$a_n = -2\left[\ln\frac{2}{n} + \left(1 + \frac{1}{2} + \dots + \frac{1}{n} + 1 - \frac{1}{2n}\right]\delta_{l0} + \frac{(1 - \delta_{l0})}{l(l+1)(2l+1)}$$
$$\Delta E_{lec}^{(Z\alpha)^5}(2P - 2S) = \begin{cases} -0.5581 \text{ meV} \\ -0.4330 \text{ meV} \end{cases}.$$

Recoil correction of order $(Z\alpha)^6$

$$\Delta E_{rec}^{(Z\alpha)^6}(2P-2S) = \frac{(Z\alpha)^6 m_1^2}{8m_2} \left(\frac{23}{6} - 4\ln 2\right) = \begin{cases} 0.0051 \ meV \\ 0.0038 \ meV \end{cases}.$$

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Muon vacuum polarization, muon self-energy correction

$$\begin{split} \Delta E_{MVP,MSE}(2S) &= \frac{\alpha(Z\alpha)^4}{8\pi} \frac{\mu^3}{m_1^2} \left[\frac{4}{3} \ln \frac{m_1}{\mu(Z\alpha)^2} - \frac{4}{3} \ln k_0(2S) + \frac{38}{45} + \right. \\ &+ \frac{\alpha}{\pi} \left(-\frac{9}{4} \zeta(3) + \frac{3}{2} \pi^2 \ln 2 - \frac{10}{27} \pi^2 - \frac{2179}{648} \right) + 4\pi Z\alpha \left(\frac{427}{384} - \frac{\ln 2}{2} \right) \right] = \begin{cases} 10.6633 \ meV}{10.9392 \ meV} \\ \Delta E_{MVP,MSE}(2P) &= \frac{\alpha(Z\alpha)^4}{8\pi} \frac{\mu^3}{m_1^2} \left[-\frac{4}{3} \ln k_0(2P) - \frac{m_1}{6\mu} - \right. \\ &- \frac{\alpha}{3\pi} \frac{m_1}{\mu} \left(\frac{3}{4} \zeta(3) - \frac{\pi^2}{2} \ln 2 + \frac{\pi^2}{12} + \frac{197}{144} \right) \right] = \begin{cases} -0.1653 \ meV}{-0.1678 \ meV}. \end{cases}$$

Radiative-recoil corrections of orders $\alpha(Z\alpha)^5$, $(Z^2\alpha)(Z\alpha)^4$

$$\Delta E_{rad-rec}(2S) = -1.324 \frac{\alpha(Z\alpha)^5 m_1^2}{8m_2} + \left(\frac{2\pi^2}{9} - \frac{70}{27}\right) \frac{\alpha(Z\alpha)^5 m_1^2}{8\pi^2 m_2} +$$

$$+\left[\frac{1}{3}\ln\frac{\Lambda(Z\alpha)^{-2}}{\mu}+\frac{11}{72}-\frac{1}{24}-\frac{7\pi}{32}\frac{\Lambda^2}{4m_2^2}+\frac{2}{3}\left(\frac{\Lambda^2}{4m_2^2}\right)^2-\frac{1}{3}\ln k_0(2S)\right]\frac{4(Z^2\alpha)(Z\alpha)^4\mu^3}{8\pi m_2^2},$$

$$\Delta E_{rad-rec}(2P) = -\frac{1}{3} \ln k_0(2P) \frac{4(Z^2 \alpha (Z\alpha)^4 \mu^3)}{8\pi m_2^2}.$$

$$\Delta E_{rad-rec}(2P-2S) = \begin{cases} -0.0656 \ meV \\ -0.0377 \ meV \end{cases}.$$

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Nuclear structure corrections of orders $(Z\alpha)^6$, $\alpha(Z\alpha)^5$

$$\begin{split} \Delta E_{str}^{(Z\alpha)^6}(2P-2S) &= \frac{(Z\alpha)^6}{12} \mu^3 \Big\{ r_N^2 \left[\langle \ln \mu Z \alpha r \rangle + C - \frac{3}{2} \right] - \frac{1}{2} r_N^2 + \frac{1}{3} \langle r^3 \rangle \langle \frac{1}{r} \rangle - \\ - l_2^{rel} - l_3^{rel} - \mu^2 F_{NR} + \frac{1}{40} \mu^2 \langle r^4 \rangle \Big\} = \begin{cases} -0.005064 \cdot r_h^2 + 0.11445 = -0.3882 \text{ meV} \\ -0.00533 \cdot r_\alpha^2 + 0.07846 = -0.3063 \text{ meV} \end{cases} \\ \Delta E_{str}^{\alpha(Z\alpha)^5}(2P-2S) = \begin{cases} 0.0940 \text{ meV} \\ 0.0702 \text{ meV} \end{cases} \end{split}$$



$$\begin{split} \Delta E_{rad+VP}(nS) &= \frac{\mu^3}{m_1^2} \frac{(Z\alpha)^4}{n^3} \left[4m_1^2 F_1'(0)\delta_{l0} + F_2(0)\frac{C_{jl}}{2l+1} \right], \ C_{jl} = \delta_{l0} + (1-\delta_{l0})\frac{j(j+1) - l(l+1) - \frac{3}{4}}{l(l+1)}. \\ m_1^2 F_1'(0) &= \left(\frac{\alpha}{\pi}\right)^2 \left[\frac{1}{9} \ln^2 \frac{m_1}{m_e} - \frac{29}{108} \ln \frac{m_1}{m_e} + \frac{1}{9}\zeta(2) + \frac{395}{1296} \right], \\ F_2(0) &= \left(\frac{\alpha}{\pi}\right)^2 \left[\frac{1}{3} \ln \frac{m_1}{m_e} - \frac{25}{36} + \frac{\pi^2}{4} \frac{m_e}{m_1} - 4 \frac{m_e^2}{m_1^2} \ln \frac{m_1}{m_e} + 3 \frac{m_e^2}{m_1^2} \right]. \\ \Delta E_{rad+VP}(2P - 2S) &= \begin{cases} -0.0299 \ meV \\ -0.0307 \ meV \end{cases}. \end{split}$$

$$\begin{split} \Delta E_{MSE}^{VP} &= \frac{\alpha}{3\pi m_1^2} \ln \frac{m_1}{\mu(Z\alpha)^2} \left[<\psi_n | \Delta \cdot \Delta V_{VP}^C | \psi_n > +2 <\psi_n | \Delta V_{VP}^C \tilde{G} \Delta \left(-\frac{Z\alpha}{r}\right) | \psi_n > \right]. \\ \Delta E_{MSE}^{VP} (2P-2S) &= \begin{cases} -0.1008 \ meV \\ -0.1074 \ meV \end{cases}. \end{split}$$

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K. Pachucki, Phys. Rev. A 54, 1994 (1996)

U.D. Jentschura and B.J. Wundt, Eur. Phys. Jour. D 65, 357 (2011).

HVP and nuclear polarizability contributions

 $\Delta E^{HVP} = \begin{cases} 0.2170 \text{ meV} \\ 0.2229 \text{ meV} \end{cases}.$

E. Borie, Z. Phys. A 302, 187 (1981)

J.L. Friar, J. Martorell and D.W.L. Sprung, PRA 59, 4061 (1999).

R.N. Faustov and A.P. Martynenko, EPJC 6, 1 (1999)

 $\Delta E^{NP} = \begin{cases} 4.9 \pm 1.0 \text{ meV} \\ 2.47 \pm 0.15 \text{ meV} \end{cases}.$

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J. Bernabeu and C. Jarlskog, Nucl. Phys. B 75, 59 (1974)

C. Ji, N.N. Dinur, S. Bacca and N. Barnea PRL 111, 143402 (2013).

Numerical results, comparison with other calculations

- E. Borie, Ann. Phys. (NY) **72**, 052511 (2012).
- E.Yu. Korzinin, V.G. Ivanov and S.G. Karshenboim, PRD 88, 125019 (2013); S.G. Karshenboim, V.G. Ivanov, E.Yu. Korzinin, and V.A. Shelyuto, PRA 81, 060501 (2010).
- U.D. Jentschura, Ann.Phys. **326**, 500 (2011); U.D. Jentschura, PRA **84**, 012505 (2011); U.D. Jentschura, EPJD **61**, 7 (2011).

Our one-loop VP result coincides with the calculation KKIS.

KKIS, meV

First order VP: 1665.7729

Our result, meV

 VP contribution of order α(Zα)² in 1γ interaction: 1665.7730

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Total two-loop contribution from KKIS is equal to

► 13.2769 meV, (µ₂⁴He)⁺

This agrees with our results

▶ 13.2789 meV, (µ₂⁴He)⁺

with the accuracy 0.002 meV (a number of two-loop corrections to the Breit Hamiltonian were estimated approximately).

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Our three-loop VP result is also in agreement with the calculation KKIS.

KKIS, meV

▶ 0.074 (µ₂⁴He)⁺

Our result, meV

▶ 0.0703 (µ₂⁴He)⁺

Relativistic corrections with vacuum polarization effects (FOPT, SOPT) in our work coincide with the results of Jentschura.

Jentschura, meV

• $\delta E_{vp} = 0.521$

Our results, meV

 Relativistic-VP correction of order α(Zα)⁴ in FOPT: -0.9472

► Relativistic-VP correction of order $\alpha(Z\alpha)^4$ in SOPT: 1.4682 Total: 0.521

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There exists the only calculation of E. Borie where total results for the Lamb shift in muonic helium ions were obtained. In the case of $(\mu_2^4 He)^+$:

Borie, meV

Uehling: 1666.305

Our results, meV

- VP contribution of order α(Zα)² in 1γ interaction: 1665.7730
- Relativistic-VP contribution of order α(Zα)⁴ in FOPT: -0.9472
- Relativistic-VP contribution of order α(Zα)⁴ in SOPT: 1.4682

Total: 1666.2940

Borie, meV

Kallen-Sabry: 11.573

Our results, meV

- 2-loop VP contribution of order α²(Zα)² in 1γ interaction: 11.5693
- Relativistic-2loop VP contribution of order α²(Zα)⁴ in FOPT: -0.0037
- Relativistic-2loop VP contribution of order α²(Zα)⁴ in SOPT: 0.0058

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Total: 11.5714

The small difference may be related with recoil terms accounted in our calculation.

We can easily compare our results for nuclear structure corrections with Borie's results.

We used the same value for charge radius of α -particle

 $r_{He} = 1.676 \ fm$

I. Sick Phys. Lett. B **116**, 212 (1982)

We also use the same Gaussian parametrization for the formfactors.

Our results, meV

- Nuclear structure of order $(Z\alpha)^4$: -295.848± 2.83
- Nuclear structure of order (Zα)⁵ in 2γ interaction: 6.605±
 0.07
- ► Nuclear structure-VP of order α(Zα)⁴ (FOPT): -0.960± 0.0092
- Nuclear structure and VP correction of order α(Zα)⁴ (SOPT): -1.5063± 0.0092
- Nuclear structure-2-loop VP correction of order α²(Zα)⁴ in 1γ interaction: -0.0076
- Nuclear structure and 2-loop VP correction of order α²(Zα)⁴ (SOPT): -0.0182
- Nuclear structure-VP contribution in 2γ interaction: 0.1279± 0.0013

Total: -291.844, Borie, meV:-292.045

Comparison between total results for $(\mu_2^4 He)^+$ of <u>Borie</u>: $\Delta E = 1379.2479 \ meV$ <u>Our result</u>: $\Delta E = 1379.1107 \ meV$ Discrepancy is equal 0.1 meV.

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Thank you for your attention.