

# History

- ▶ Vanyashin, Terentev (1965)  
Massive vector field

$$\beta_0 = \left( \frac{11}{3} - \frac{1}{6} \right) C_A$$

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- ▶ Gross, Wilczek; Politzer (1973)

# Coulomb gauge

$$\bullet - - - - \bullet = \frac{i}{\vec{q}^2}$$

$$\bullet \text{~~~~~} \bullet = -\frac{i}{q^2 + i0} \left( \delta^{ij} - \frac{q^i q^j}{\vec{q}^2} \right)$$

No ghosts

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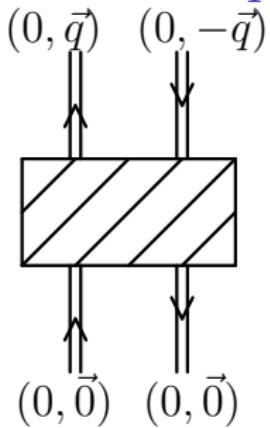
Infinitely heavy quark

$$\begin{array}{c} \text{---} \\ | \quad | \\ \text{---} \end{array} = \frac{i}{p_0 + i0}$$

$$\begin{array}{c} \text{---} \\ | \quad | \\ \text{---} \end{array} - - - = ig_0 t^a$$

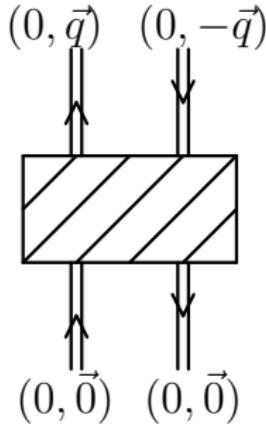
$$\begin{array}{c} \text{---} \\ | \quad | \\ \text{---} \end{array} \text{~~~~~} \text{---} = 0$$

# Quark–antiquark potential

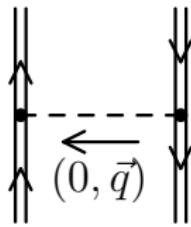


Quantum mechanics:  $iU_{\vec{q}}$

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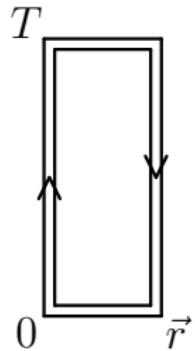


$$= -iC_F g_0^2 D(0, \vec{q}) = iC_F \frac{g_0^2}{\vec{q}^2}$$

$$U_{\vec{q}} = C_F g_0^2 D(0, \vec{q}) = -C_F \frac{g_0^2}{\vec{q}^2}$$

$$U(r) = -C_F \frac{\alpha_s}{r}$$

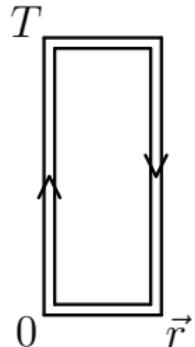
# Wilson loop



$T \gg r$

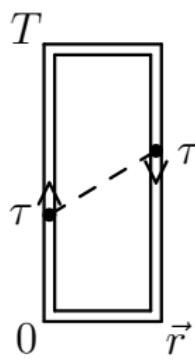
$$= e^{-i U(\vec{r}) T} = 1 - i U(\vec{r}) T$$

# Wilson loop



A diagram of a rectangle representing a Wilson loop. The vertical axis is labeled  $T$  at the top and  $0$  at the bottom. The horizontal axis is labeled  $\vec{r}$  at the right end. A vertical arrow points upwards along the left edge, and a horizontal arrow points to the right along the top edge.

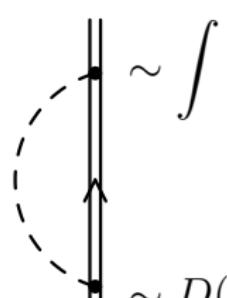
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$$= -i C_F g_0^2 T \int D(t, \vec{r}) dt$$
$$= -i C_F g_0^2 T \int \frac{d^{d-1} \vec{q}}{(2\pi)^{d-1}} D(0, \vec{q}) e^{i \vec{q} \cdot \vec{r}}$$

# Self-energy and vertex corrections



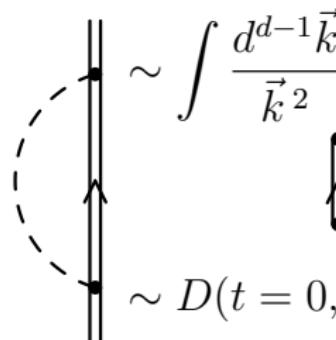
A diagram showing a dashed circular path with two vertical segments. A small black dot is at the top of the left vertical segment, and a horizontal arrow points from right to left along the path.

$$\sim \int \frac{d^{d-1}\vec{k}}{\vec{k}^2} = 0 \quad \text{Self-energy of a classical point charge}$$

  
propagates along time,  along space

$$\sim D(t=0, \vec{r}=0) \sim \int \frac{d^{d-1}\vec{k}}{\vec{k}^2} e^{i\vec{q}\cdot\vec{r}} \Big|_{\vec{r}=0} \sim U(\vec{r}=0) \Rightarrow 0$$

# Self-energy and vertex corrections



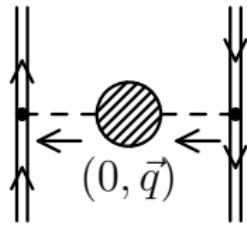
$\sim \int \frac{d^{d-1}\vec{k}}{\vec{k}^2} = 0$     Self-energy of a classical point charge

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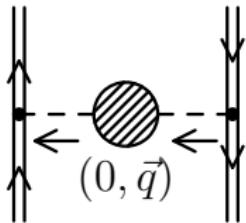

$$= 0$$

# Vacuum polarization



$$U_{\vec{q}} = C_F g_0^2 D(0, \vec{q})$$

# Vacuum polarization

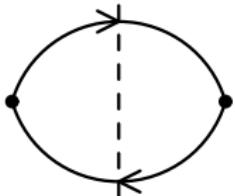


$$U_{\vec{q}} = C_F g_0^2 D(0, \vec{q})$$

$$- \text{---} \circlearrowleft - = -i\vec{q}^2 \Pi(q)$$

$$D(q) = -\frac{1}{\vec{q}^2} \frac{1}{1 - \Pi(q)} = -\frac{1}{\vec{q}^2} (1 + \Pi(q))$$

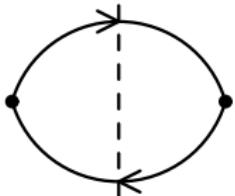
# Quark loop



Lorentz-invariant  $\rho_q(s) \geq 0$

$$\Pi_q(q^2) = \int \frac{\rho_q(s) ds}{q^2 - s + i0}$$

# Quark loop

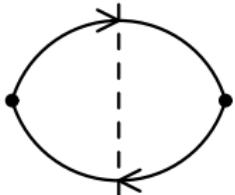


Lorentz-invariant  $\rho_q(s) \geq 0$

$$\Pi_q(q^2) = \int \frac{\rho_q(s) ds}{q^2 - s + i0}$$

$$\begin{aligned} U_{\vec{q}} &= -C_F \frac{g_0^2}{\vec{q}^2} \left[ 1 - \int \frac{\rho_q(s) ds}{\vec{q}^2 + s} + \dots \right] \\ &= -C_F g_0^2 \left[ \left( 1 - \int \frac{\rho_q(s) ds}{s} \right) \frac{1}{\vec{q}^2} + \int \frac{\rho_q(s) ds}{\vec{q}^2 + s} + \dots \right] \end{aligned}$$

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$$U(r) = -C_F \frac{g_0^2}{4\pi r} \left[ 1 - \int \frac{\rho_q(s) ds}{s} + \int \rho_q(s) e^{-\sqrt{s}r} ds + \dots \right]$$

Screening

$$\rho_q(s) = T_F n_f \frac{g_0^2 s^{-\varepsilon}}{(4\pi)^{d/2}} \left( \frac{4}{3} + \dots \right)$$

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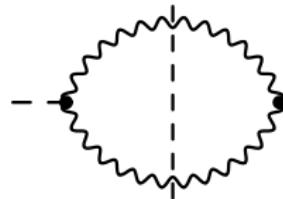
$$\int \frac{\rho_q(s) \, ds}{s + \vec{q}^2} \Big|_{\text{UV}} = \frac{4}{3} T_F n_f \frac{g_0^2}{(4\pi)^{d/2}} \int\limits_{\sim \vec{q}^2}^{\infty} s^{-1-\varepsilon} ds = \frac{4}{3} T_F n_f \frac{\alpha_s}{4\pi \varepsilon}$$

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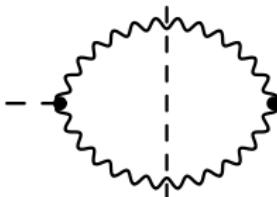
$$U_{\vec{q}} = -C_F \frac{g_0^2}{\vec{q}^{\,2}} \left[ 1 + \frac{4}{3}T_F n_f \frac{\alpha_s}{4\pi\varepsilon} + \dots \right]$$

# Transverse gluons



$$- \text{ - } \Pi_t(q_0^2, \vec{q}^2) = \int \frac{\rho_t(s, \vec{q}^2) ds}{q^2 - s + i0}$$

# Transverse gluons



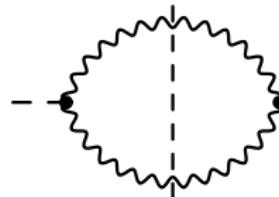
A Feynman diagram showing a loop of gluons (wavy lines) attached to a vertical dashed line. The loop has two external gluon lines meeting at a central vertex. The vertical dashed line extends downwards from this vertex.

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$s \gg \vec{q}^2$ :

$$\rho_t(s, \vec{q}^2) = C_A \frac{g_0^2 s^{-\varepsilon}}{(4\pi)^{d/2}} \left( \frac{1}{3} + \dots \right)$$

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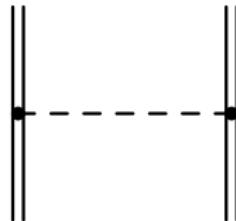
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$$U_{\vec{q}} = -C_F \frac{g_0^2}{\vec{q}^2} \left[ 1 + \frac{1}{3} C_A \frac{\alpha_s}{4\pi\varepsilon} + \dots \right]$$

# Coulomb gluon

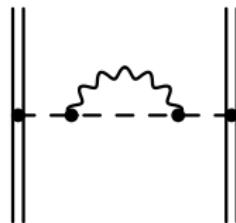


# Coulomb gluon



+ transverse-gluon  
vacuum

$$E_0 = U(r)$$



Second order of perturbation theory  
Energy decreases — **antiscreening**



Depends on  $\vec{q}^2$  but not on  $q^0$   
Is not given by a spectral representation



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$$\Pi_c(\vec{q}^2) = \int \frac{d^d k}{(2\pi)^d} \frac{f(\vec{k}, \vec{q})}{k^2 + i0}$$



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$$\Pi_c(\vec{q}^2) = C_A \frac{g_0^2 (\vec{q}^2)^{-\varepsilon}}{(4\pi)^{d/2}} \left( \frac{4}{\varepsilon} + \dots \right)$$

# Results

$$U_{\vec{q}} = -C_F \frac{g_0^2}{\vec{q}^2} \left\{ 1 + \frac{g_0^2 (\vec{q}^2)^{-\varepsilon}}{(4\pi)^{d/2}} \left[ \left( \left( \textcolor{red}{4} - \frac{1}{3} \right) C_A - \frac{4}{3} T_F n_f \right) \frac{1}{\varepsilon} + \dots \right] \right\}$$

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Renormalization

$$\frac{g_0^2}{(4\pi)^{d/2}} = \mu^{2\varepsilon} \frac{\alpha_s(\mu)}{4\pi} Z_\alpha e^{\gamma\varepsilon} \quad Z_\alpha = 1 - \beta_0 \frac{\alpha_s}{4\pi\varepsilon}$$

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$$\beta_0 = \left( \textcolor{red}{4} - \frac{1}{3} \right) C_A - \frac{4}{3} T_F n_f$$

Asymptotic freedom

# Ward identity

External Coulomb line with incoming momentum  $q = \omega v$   
(a dot means shifting the momentum by  $q$ )

$$\begin{aligned} \omega \bullet \rightarrow \bullet \bullet \rightarrow \bullet &= g_0 \rightarrow \bullet \rightarrow \otimes \left[ \bullet \rightarrow \bullet - \bullet \rightarrow \bullet \right] \\ \omega \text{---} \bullet \bullet \text{---} \bullet &= g_0 \text{---} \bullet \text{---} \otimes \left[ \text{---} \bullet \text{---} - \text{---} \bullet \text{---} \right] \\ \text{---} \bullet \bullet \text{---} \bullet &= \bullet \text{---} \bullet \text{---} \bullet = 0 \end{aligned}$$

(colour structure)  $\otimes$  (Lorentz structure)

(In covariant gauges, the second identity contains extra ghost terms)

# 1 loop

$$\begin{aligned}
 & \omega \text{---} \bullet \text{---} \bullet \text{---} \omega = g_0 \text{---} \bullet \text{---} \bullet \text{---} \omega \otimes \left[ \text{---} \bullet \text{---} \text{---} - \text{---} \bullet \text{---} \text{---} \right] \\
 & \omega \text{---} \bullet \text{---} \bullet \text{---} \omega = g_0 \text{---} \bullet \text{---} \bullet \text{---} \omega \otimes \left[ \text{---} \bullet \text{---} \text{---} - \text{---} \bullet \text{---} \text{---} \right] \\
 & = g_0 \text{---} \bullet \text{---} \bullet \text{---} \omega \otimes \left[ \text{---} \bullet \text{---} \text{---} - \text{---} \bullet \text{---} \text{---} \right] = 0
 \end{aligned}$$

Diagrammatic representation of a 1-loop Feynman diagram. The top row shows a solid horizontal line with arrows pointing right, labeled  $\omega$  at both ends. A vertical dashed line with a downward arrow connects the first two vertices. A wavy line connects the second vertex to a black dot. The result is equated to  $g_0$  times a similar diagram where the wavy line is connected to the third vertex, followed by a tensor product symbol ( $\otimes$ ) and a bracket containing two terms: one with a wavy line to the first vertex and another with a wavy line to the third vertex. The bottom row shows a similar diagram but with a dashed horizontal line instead of a solid one. The final result is zero.

$$\begin{aligned}
 & \text{Diagram 1:} \\
 & \omega \xrightarrow{\text{---}} \text{---} \xrightarrow{\text{---}} \text{---} \quad = g_0 \left[ \text{---} \xrightarrow{\text{---}} \text{---} \xrightarrow{\text{---}} \text{---} - \text{---} \xrightarrow{\text{---}} \text{---} \xrightarrow{\text{---}} \text{---} \right] \\
 & \otimes \left[ \text{---} \xrightarrow{\text{---}} \text{---} \xrightarrow{\text{---}} \text{---} - \text{---} \xrightarrow{\text{---}} \text{---} \xrightarrow{\text{---}} \text{---} \right] \\
 & \text{Diagram 2:} \\
 & \omega \xrightarrow{\text{---}} \text{---} \xrightarrow{\text{---}} \text{---} = \text{---} \xrightarrow{\text{---}} \text{---} \xrightarrow{\text{---}} \text{---} = 0
 \end{aligned}$$

# Ward identity

$$\overrightarrow{p} \circlearrowleft \text{---} \overrightarrow{\phantom{p}} = -i\Sigma(p)$$

$$S(p) = \frac{1}{p - m_0 - \Sigma(p)}$$

$$\overrightarrow{p} \circlearrowleft \text{---} \overrightarrow{\phantom{p}} = ig_0 t^a \Gamma(p, q)$$

$$\Gamma = \gamma_0 + \Lambda$$

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$$\begin{array}{c} \rightarrow \\ p \end{array} \circlearrowleft \begin{array}{c} \downarrow \\ q \end{array} = ig_0 t^a \Gamma(p, q) \quad \Gamma = \gamma_0 + \Lambda$$

$$\omega\Lambda(p, \omega v) = \Sigma(p) - \Sigma(p + \omega v)$$

$$\omega\Gamma(p, \omega v) = S^{-1}(p + \omega v) - S^{-1}(p)$$

# Coupling constant renormalization

$$g_0 = Z_\alpha^{1/2} g \quad \Gamma = Z_\Gamma \Gamma_r \quad S = Z_\psi S_r$$

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$$Z_\alpha = (Z_\Gamma Z_\psi)^{-2} Z_A^{-1}$$

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Ward identity

$$Z_\Gamma Z_\psi = 1 \quad Z_\alpha = Z_A^{-1}$$