

# Massless QED

$$L = \bar{\psi}_0 i \not{D} \psi_0 - \frac{1}{4} F_{0\mu\nu} F_0^{\mu\nu}$$

$$D_\mu \psi_0 = (\partial_\mu - ie_0 A_{0\mu}) \psi_0 \quad F_{0\mu\nu} = \partial_\mu A_{0\nu} - \partial_\nu A_{0\mu}$$

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$$\begin{array}{ccc} \bullet \xrightarrow[p]{} \bullet & = iS_0(p) & S_0(p) = \frac{1}{p} = \frac{\not{p}}{p^2} \\ \mu \bullet \sim \sim \sim \sim \sim \sim \sim \sim \nu & = -iD_{\mu\nu}^0(p) & \\ & & D_{\mu\nu}^0(p) = \frac{1}{p^2} \left[ g_{\mu\nu} - (1 - a_0) \frac{p_\mu p_\nu}{p^2} \right] \\ \mu \\ \backslash \swarrow \\ \bullet \xrightarrow[\rightarrow]{} \bullet & = ie_0 \gamma^\mu & \end{array}$$

# Renormalization

Renormalized quantities

$$\psi_0 = Z_\psi^{1/2} \psi \quad A_0 = Z_A^{1/2} A \quad a_0 = Z_A a \quad e_0 = Z_\alpha^{1/2} e$$

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Minimal renormalization constants

$$Z_i(\alpha) = 1 + \frac{z_1}{\varepsilon} \frac{\alpha}{4\pi} + \left( \frac{z_{22}}{\varepsilon^2} + \frac{z_{21}}{\varepsilon} \right) \left( \frac{\alpha}{4\pi} \right)^2 + \dots$$

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Dimensionalities:  $[L] = d$ ,  $[A_0] = 1 - \varepsilon$ ,  $[\psi_0] = 3/2 - \varepsilon$ ,  
 $[e_0] = \varepsilon$ . Exactly dimensionless

$$\frac{\alpha(\mu)}{4\pi} = \mu^{-2\varepsilon} \frac{e^2}{(4\pi)^{d/2}} e^{-\gamma\varepsilon} \quad \frac{e_0^2}{(4\pi)^{d/2}} = \mu^{2\varepsilon} \frac{\alpha(\mu)}{4\pi} Z_\alpha(\alpha(\mu)) e^{\gamma\varepsilon}$$

# Photon propagator



$$\begin{aligned} -iD_{\mu\nu}(p) &= -iD_{\mu\nu}^0(p) + (-i)D_{\mu\alpha}^0(p)i\Pi^{\alpha\beta}(p)(-i)D_{\beta\nu}^0(p) \\ &+ (-i)D_{\mu\alpha}^0(p)i\Pi^{\alpha\beta}(p)(-i)D_{\beta\gamma}^0(p)i\Pi^{\gamma\delta}(p)(-i)D_{\gamma\nu}^0(p) + \dots \end{aligned}$$

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$$\left. \begin{aligned} A_{\mu\nu} &= A_{\perp} \left[ g_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right] + A_{||} \frac{p_\mu p_\nu}{p^2} \\ A_{\mu\nu}^{-1} &= A_{\perp}^{-1} \left[ g_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right] + A_{||}^{-1} \frac{p_\mu p_\nu}{p^2} \end{aligned} \right\} \quad A_{\mu\lambda}^{-1} A^{\lambda\nu} = \delta_\mu^\nu$$

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$$D_{\mu\nu}^{-1}(p) = (D^0)_{\mu\nu}^{-1}(p) - \Pi_{\mu\nu}(p)$$

# Photon propagator

Ward identity  $\Pi_{\mu\nu}(p)p^\nu = 0$

$$\Pi_{\mu\nu}(p) = (p^2 g_{\mu\nu} - p_\mu p_\nu) \Pi(p^2)$$

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Renormalized propagator  $D_{\mu\nu}(p) = Z_A(\alpha(\mu)) D_{\mu\nu}^r(p; \mu)$

$$D_{\mu\nu}^r(p; \mu) = D_{\perp}^r(p^2; \mu) \left[ g_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right] + a(\mu) \frac{p_\mu p_\nu}{(p^2)^2}$$

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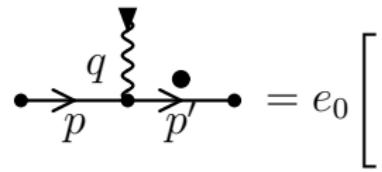
$Z_A(\alpha)$  is constructed to make

$$D_{\perp}^r(p^2; \mu) = Z_A^{-1}(\alpha(\mu)) \frac{1}{p^2(1 - \Pi(p^2))}$$

finite at  $\varepsilon \rightarrow 0$ . But the longitudinal part must be finite too:

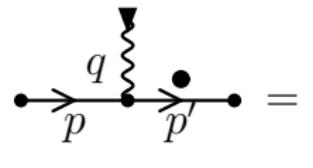
$$a(\mu) = Z_A^{-1}(\alpha(\mu)) a_0$$

# Ward identity


$$\bullet \rightarrow_p \bullet \xrightarrow{q} \bullet \rightarrow_{p'} \bullet = e_0 \left[ \bullet \rightarrow_p \bullet - \bullet \rightarrow_{p'} \bullet \right]$$

$$iS_0(p') ie_0 \not{iS_0(p)} = ie_0 [S_0(p') - S_0(p)]$$

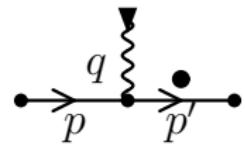
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$$iS_0(p') ie_0 \not{q} iS_0(p) = ie_0 [S_0(p') - S_0(p)]$$

$$\frac{\partial S_0(p)}{\partial p^\mu} = -S_0(p)\gamma_\mu S_0(p)$$

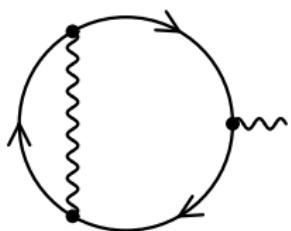
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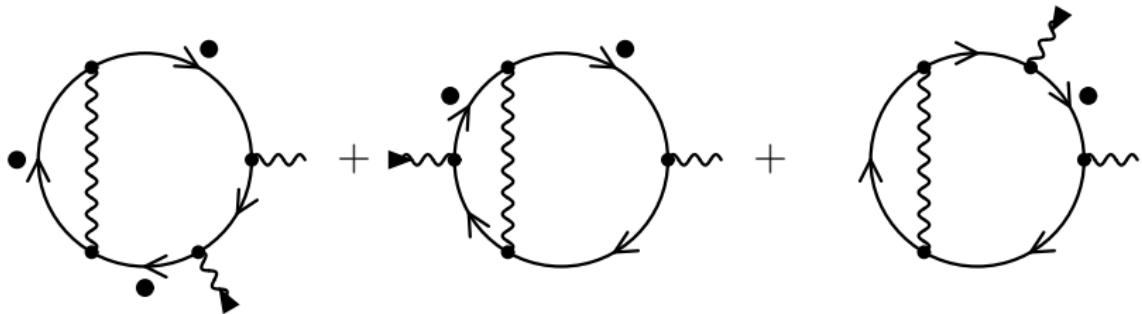

$$\text{Diagram: A quark loop with momentum } q \text{ and external momenta } p \text{ and } p'. \text{ The loop has two vertices connected by a wavy line.}$$
$$= e_0 \left[ \text{Diagram: A quark loop with momentum } p \text{ and external momenta } p \text{ and } p' - \text{Diagram: A quark loop with momentum } p' \text{ and external momenta } p \text{ and } p' \right]$$

$$iS_0(p') ie_0 \not{q} iS_0(p) = ie_0 [S_0(p') - S_0(p)]$$

$$\frac{\partial S_0(p)}{\partial p^\mu} = -S_0(p)\gamma_\mu S_0(p)$$

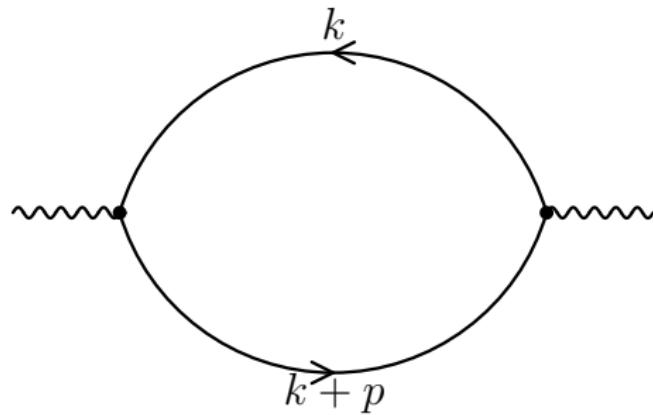

$$\text{Diagram: A quark loop with a self-energy insertion. The loop has two vertices connected by a wavy line. The self-energy insertion is a circle with a dot at the top-left vertex. The loop is oriented clockwise.}$$
$$= e_0 \left[ \text{Diagram: A quark loop with a self-energy insertion and orientation } \rightarrow - \text{Diagram: A quark loop with a self-energy insertion and orientation } \leftarrow \right] = 0$$





$$= e_0 \left[ \begin{array}{c} \text{Diagram 1: A circle with a wavy line entering from the top-left and exiting to the bottom-right. Arrows on the circle and wavy line indicate clockwise flow.} \\ - \\ \text{Diagram 2: Similar to Diagram 1, but the wavy line enters from the top-right and exits to the bottom-left.} \\ + \\ \text{Diagram 3: Similar to Diagram 1, but the wavy line enters from the bottom-left and exits to the top-right.} \\ - \\ \text{Diagram 4: Similar to Diagram 1, but the wavy line enters from the bottom-right and exits to the top-left.} \\ + \\ \text{Diagram 5: Similar to Diagram 1, but the wavy line enters from the top-left and exits to the bottom-right, with arrows indicating counter-clockwise flow.} \\ - \end{array} \right] = 0$$

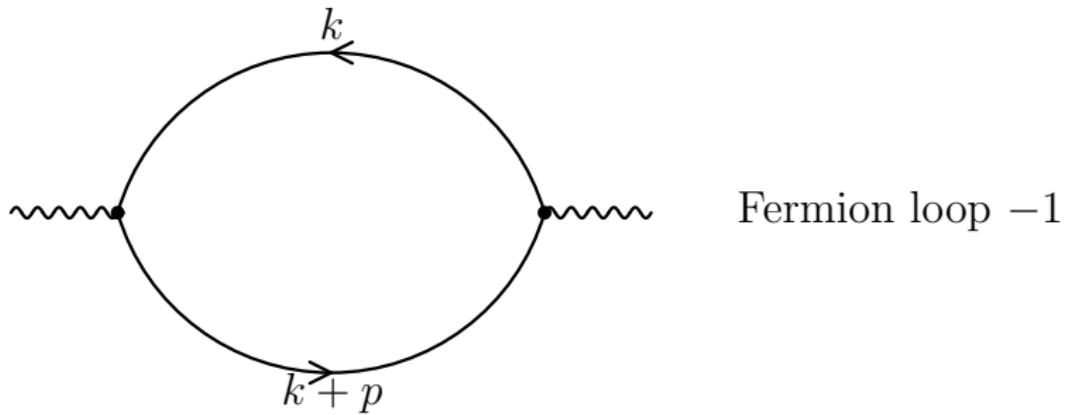
# Photon self-energy



Fermion loop  $-1$

$$\begin{aligned} i(p^2 g_{\mu\nu} - p_\mu p_\nu) \Pi(p^2) = \\ - \int \frac{d^d k}{(2\pi)^d} \text{Tr } ie_0 \gamma_\mu i \frac{\not{k} + \not{p}}{(k + p)^2} ie_0 \gamma_\nu \frac{\not{k}}{k^2} \end{aligned}$$

# Photon self-energy



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$$\Pi(p^2) = \frac{-ie_0^2}{(d-1)(-p^2)} \int \frac{d^d k}{(2\pi)^d} \frac{\text{Tr } \gamma_\mu (\not{k} + \not{p}) \gamma^\mu \not{k}}{[-(k + p)^2] (-k^2)}$$

# Photon self-energy

$$\Pi(p^2) = \frac{d-2}{d-1} \frac{ie_0^2}{-p^2} \int \frac{d^d k}{(2\pi)^d} \frac{4(k+p) \cdot k}{[-(k+p)^2](-k^2)}$$

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Set  $-p^2 = 1$

$$D_1 = -(k+p)^2 \quad D_2 = -k^2$$

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Multiplication table

$$p^2 = -1 \quad k^2 = -D_2 \quad p \cdot k = \frac{1}{2}(1 + D_2 - D_1)$$

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$$\Pi(p^2) = 2 \frac{d-2}{d-1} ie_0^2 \int \frac{d^d k}{(2\pi)^d} \frac{-2D_2 + 1 + D_2 - D_1}{D_1 D_2}$$

# Photon self-energy

Restoring  $-p^2$  by dimensionality

$$\begin{aligned}\Pi(p^2) &= -\frac{e_0^2(-p^2)^{-\varepsilon}}{(4\pi)^{d/2}} 2 \frac{d-2}{d-1} G_1 \\ &= \frac{e_0^2(-p^2)^{-\varepsilon}}{(4\pi)^{d/2}} 4 \frac{d-2}{(d-1)(d-3)(d-4)} g_1\end{aligned}$$

# Photon field renormalization

Transverse propagator

$$p^2 D_{\perp}(p^2) = \frac{1}{1 - \Pi(p^2)} = 1 + \frac{e_0^2 (-p^2)^{-\varepsilon}}{(4\pi)^{d/2}} 4 \frac{d-2}{(d-1)(d-3)(d-4)} g_1$$

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Re-expressing via  $\alpha(\mu)$

$$p^2 D_{\perp}(p^2) = 1 + \frac{\alpha(\mu)}{4\pi} e^{-L\varepsilon} e^{\gamma\varepsilon} g_1 4 \frac{d-2}{(d-1)(d-3)(d-4)}$$

$$L = \log \frac{-p^2}{\mu^2}$$

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$$p^2 D_{\perp}(p^2) = 1 - \frac{4}{3} \frac{\alpha(\mu)}{4\pi\varepsilon} \left[ 1 - \left( L - \frac{5}{3} \right) \varepsilon + \dots \right]$$

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This should be  $Z_A(\alpha(\mu)) p^2 D_{\perp}^r(p^2; \mu)$ :

$$Z_A(\alpha) = 1 - \frac{4}{3} \frac{\alpha}{4\pi\varepsilon}$$

$$p^2 D_{\perp}^r(p^2; \mu) = 1 + \frac{4}{3} \frac{\alpha(\mu)}{4\pi} \left( L - \frac{5}{3} \right)$$

# RG equation

$D_\perp(p^2) = Z_A(\alpha(\mu))D_\perp^r(p^2; \mu)$  does not depend on  $\mu$ :

$$\frac{\partial D_\perp^r(p^2; \mu)}{\partial \log \mu} + \gamma_A(\alpha(\mu))D_\perp^r(p^2; \mu) = 0$$

$$\gamma_A(\alpha(\mu)) = \frac{d \log Z_A(\alpha(\mu))}{d \log \mu}$$

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For a minimal renormalization constant using

$$\frac{d \log \alpha(\mu)}{d \log \mu} = -2\varepsilon + \dots$$

we obtain

$$\gamma(\alpha) = \gamma_0 \frac{\alpha}{4\pi} + \dots = -2z_1 \frac{\alpha}{4\pi} + \dots$$

$$Z(\alpha) = 1 - \frac{\gamma_0}{2} \frac{\alpha}{4\pi\varepsilon} + \dots$$

# Anomalous dimension

$$\gamma_A(\alpha) = \frac{8}{3} \frac{\alpha}{4\pi}$$

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$$\frac{\partial p^2 D_{\perp}^r}{\partial L} = \frac{\gamma_A}{2} p^2 D_{\perp}^r$$

with the initial condition

$$p^2 D_{\perp}^r(p^2; \mu^2 = -p^2) = 1 - \frac{20}{9} \frac{\alpha(\mu)}{4\pi}$$

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is enough to reconstruct the result

$a_0 = Z_A(\alpha(\mu))a(\mu)$  does not depend on  $\mu$ :

$$\frac{da(\mu)}{d \log \mu} + \gamma_A(\alpha(\mu))a(\mu) = 0$$

# Electron propagator



$$\begin{aligned} iS(p) = & iS_0(p) + iS_0(p)(-i)\Sigma(p)iS_0(p) \\ & + iS_0(p)(-i)\Sigma(p)iS_0(p)(-i)\Sigma(p)iS_0(p) + \dots \end{aligned}$$

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$$S(p) = S_0(p) + S_0(p)\Sigma(p)S(p)$$

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$$S(p) = S_0(p) + S_0(p)\Sigma(p)S(p)$$

$$S(p) = \frac{1}{S_0^{-1}(p) - \Sigma(p)}$$

# Electron propagator



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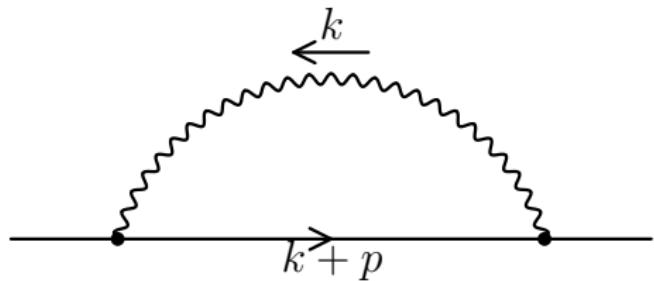
$$S(p) = S_0(p) + S_0(p)\Sigma(p)S(p)$$

$$S(p) = \frac{1}{S_0^{-1}(p) - \Sigma(p)}$$

Massless case:  $\Sigma(p) = \not{p}\Sigma_V(p^2)$  (helicity conservation)

$$S(p) = \frac{1}{1 - \Sigma_V(p^2)} \frac{1}{\not{p}}$$

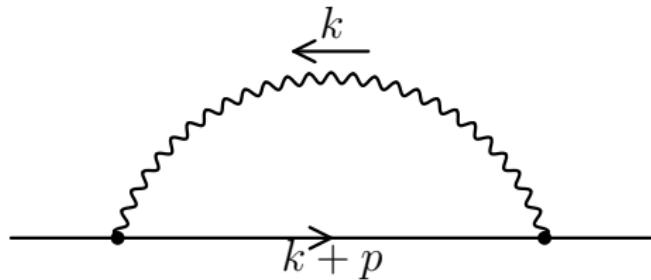
# Electron self-energy



$$-i\cancel{p}\Sigma_V(p^2) = \int \frac{d^d k}{(2\pi)^d} ie_0 \gamma^\mu i \frac{\cancel{k} + \cancel{p}}{(k + p)^2} ie_0 \gamma^\nu \frac{-i}{k^2} \left( g_{\mu\nu} - \xi \frac{k_\mu k_\nu}{k^2} \right)$$

where  $\xi = 1 - a_0$ .

# Electron self-energy



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where  $\xi = 1 - a_0$ . Taking  $\frac{1}{4} \text{Tr } \cancel{p}$

$$\Sigma_V(p^2) = \frac{ie_0^2}{-p^2} \int \frac{d^d k}{(2\pi)^d} \frac{N}{D_1 D_2}$$

$$N = \frac{1}{4} \text{Tr } \cancel{p} \gamma^\mu (\cancel{k} + \cancel{p}) \gamma^\nu \left( g_{\mu\nu} + \xi \frac{k_\mu k_\nu}{D_2} \right)$$

# Electron self-energy

Using the multiplication table

$$\begin{aligned} N &= \frac{1}{4} \text{Tr } \not{p} \gamma_\mu (\not{k} + \not{p}) \gamma^\mu + \frac{\xi}{D_2} \frac{1}{4} \text{Tr } \not{p} \not{k} (\not{k} + \not{p}) \not{k} \\ &= -(d-2)(p^2 + p \cdot k) + \frac{\xi}{D_2} [k^2 p \cdot k + 2(p \cdot k)^2 - p^2 k^2] \\ &= \frac{1}{2} \left[ d - 2 + \xi \left( \frac{1}{D_2} - 1 \right) \right] \end{aligned}$$

# Electron self-energy

Using the multiplication table

$$\begin{aligned} N &= \frac{1}{4} \text{Tr } \not{p} \gamma_\mu (\not{k} + \not{p}) \gamma^\mu + \frac{\xi}{D_2} \frac{1}{4} \text{Tr } \not{p} \not{k} (\not{k} + \not{p}) \not{k} \\ &= -(d-2)(p^2 + p \cdot k) + \frac{\xi}{D_2} [k^2 p \cdot k + 2(p \cdot k)^2 - p^2 k^2] \\ &= \frac{1}{2} \left[ d - 2 + \xi \left( \frac{1}{D_2} - 1 \right) \right] \end{aligned}$$

$$\Sigma_V(p^2) = -\frac{e_0^2(-p^2)^{-\varepsilon}}{(4\pi)^{d/2}} \frac{1}{2} [(d-2-\xi)G(1,1) + \xi G(1,2)]$$

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# Electron field renormalization

$$\not{p}S(p) = \frac{1}{1 - \Sigma_V(p^2)}$$

expressed via the renormalized quantities

$$\begin{aligned}\not{p}S(p) &= 1 + \frac{\alpha(\mu)}{4\pi} e^{-L\varepsilon} e^{\gamma\varepsilon} g_1 a(\mu) \frac{d-2}{(d-3)(d-4)} \\ &= 1 - \frac{\alpha(\mu)}{4\pi\varepsilon} a(\mu) e^{-L\varepsilon} (1 + \varepsilon + \dots)\end{aligned}$$

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should be  $Z_\psi(\alpha(\mu), a(\mu))\not{p}S_r(p; \mu)$ :

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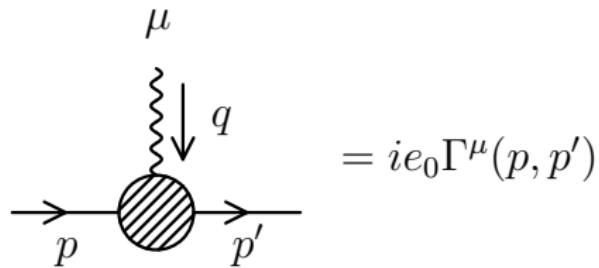
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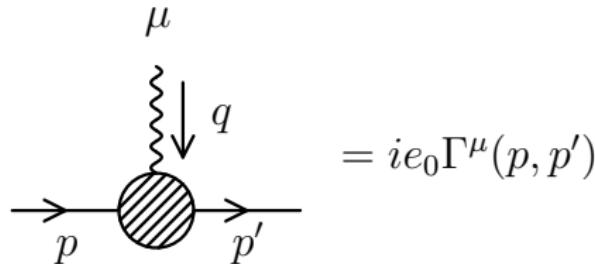
$$\gamma_\psi(\alpha, a) = 2a \frac{\alpha}{4\pi}$$

# Vertex



$$\Gamma^\mu(p, p') = \gamma^\mu + \Lambda^\mu(p, p')$$

# Vertex



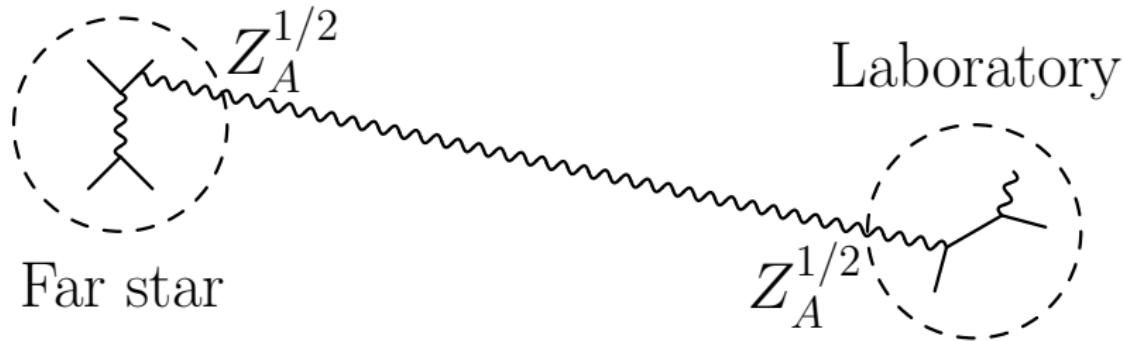
$$\Gamma^\mu(p, p') = \gamma^\mu + \Lambda^\mu(p, p')$$

When expressed via renormalized quantities,

$$\Gamma^\mu = Z_\Gamma \Gamma_r^\mu$$

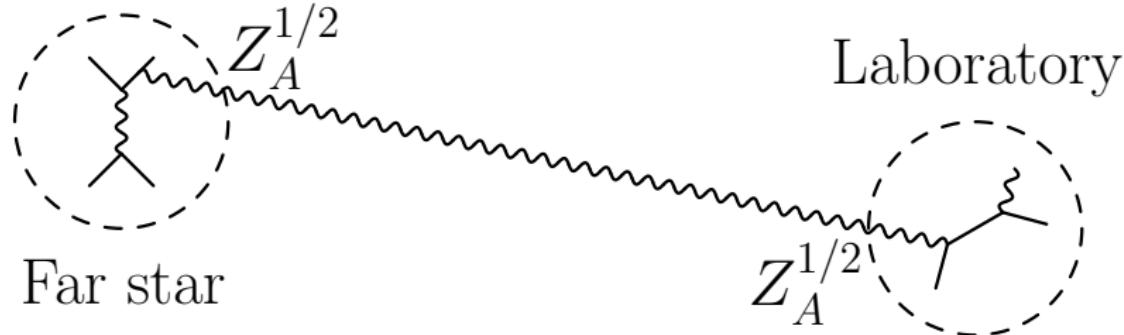
# Matrix element

$S$ -matrix element = vertex  $\times Z_i^{1/2}$  for each  $i$



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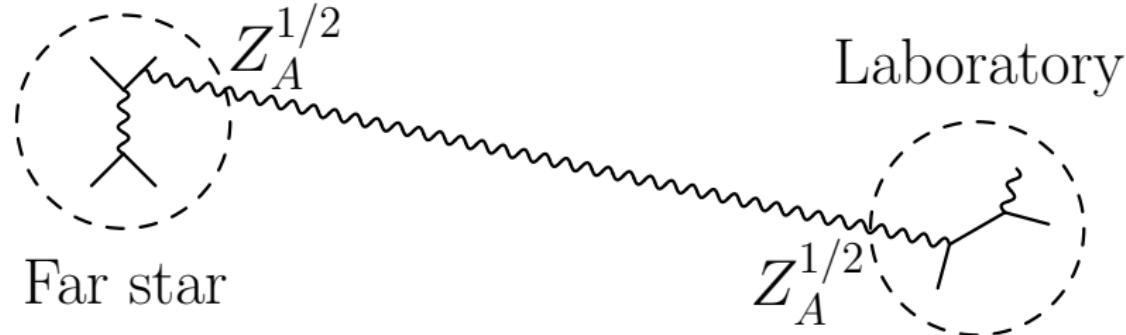
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The physical matrix element  
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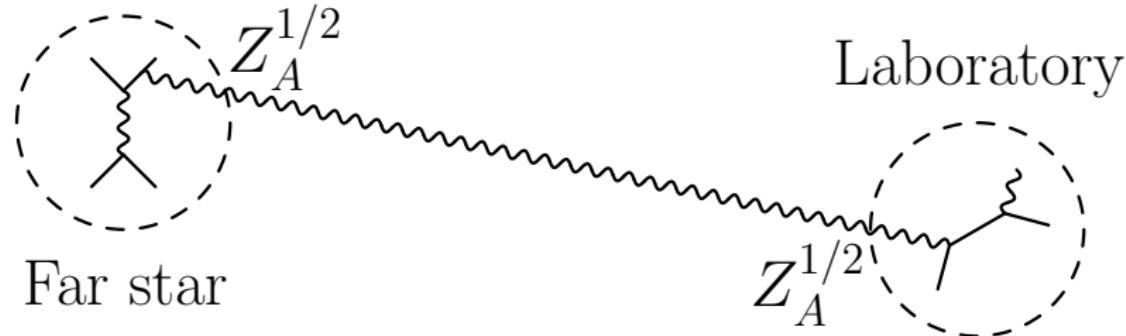
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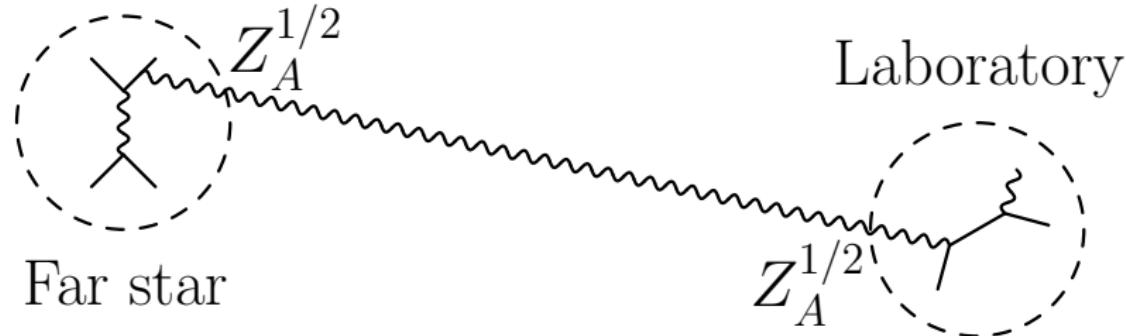
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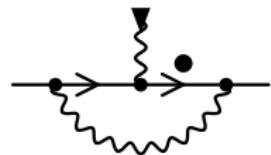
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$$Z_\alpha = (Z_\Gamma Z_\psi)^{-2} Z_A^{-1}$$

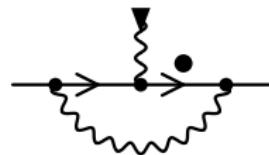
# Ward identity



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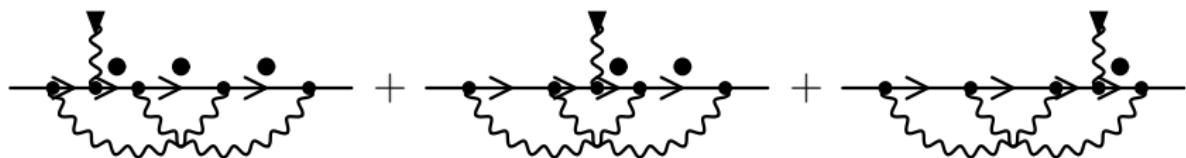
$$= e_0 \left[ \text{Diagram 1} - \text{Diagram 2} \right]$$

The equation shows the Ward identity as a difference between two terms. The first term is labeled  $e_0$  followed by a left bracket containing Diagram 1. The second term is a minus sign followed by a right bracket containing Diagram 2. Diagram 1 is identical to the one above it, but Diagram 2 has a solid line at the second vertex instead of a wavy line.

# Ward identity



# Ward identity



# Ward identity

$$= e_0 \left[ \begin{array}{c} \text{Diagram 1: } \text{---} \text{---} \text{---} \text{---} \\ | \quad | \quad | \quad | \\ \text{---} \text{---} \text{---} \text{---} \end{array} - \begin{array}{c} \text{Diagram 2: } \text{---} \text{---} \text{---} \text{---} \\ | \quad | \quad | \quad | \\ \text{---} \text{---} \text{---} \text{---} \end{array} \\ + \begin{array}{c} \text{Diagram 3: } \text{---} \text{---} \text{---} \text{---} \\ | \quad | \quad | \quad | \\ \text{---} \text{---} \text{---} \text{---} \end{array} - \begin{array}{c} \text{Diagram 4: } \text{---} \text{---} \text{---} \text{---} \\ | \quad | \quad | \quad | \\ \text{---} \text{---} \text{---} \text{---} \end{array} \\ + \begin{array}{c} \text{Diagram 5: } \text{---} \text{---} \text{---} \text{---} \\ | \quad | \quad | \quad | \\ \text{---} \text{---} \text{---} \text{---} \end{array} - \begin{array}{c} \text{Diagram 6: } \text{---} \text{---} \text{---} \text{---} \\ | \quad | \quad | \quad | \\ \text{---} \text{---} \text{---} \text{---} \end{array} \end{array} \right]$$

The equation shows the Ward identity for a four-point vertex. It consists of six terms, each represented by a horizontal line with four vertices. Each vertex has a small circle with an arrow pointing to the right. A wavy line connects the first three vertices. The terms are separated by plus signs (+) and minus signs (-). The first term is preceded by  $e_0$ . The last term is enclosed in square brackets [ ].

# Ward identity

$$= e_0 \left[ \begin{array}{c} \text{Diagram 1} \\ \text{Diagram 2} \\ + \text{Diagram 3} \\ + \text{Diagram 4} \end{array} \right] - \left[ \begin{array}{c} \text{Diagram 5} \\ \text{Diagram 6} \end{array} \right]$$

The equation illustrates the Ward identity for a four-point vertex. It shows the sum of six Feynman diagrams, each featuring a horizontal line with four vertices and a wavy line (representing a photon) attached to the first vertex. The diagrams are categorized into two groups by brackets: the top group contains Diagrams 1 through 4, and the bottom group contains Diagrams 5 and 6. The diagrams are labeled with '+' signs between them, indicating they are being summed.

# Ward identity

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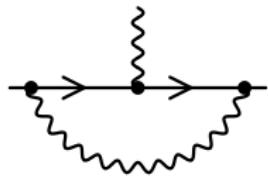
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$$Z_\alpha = Z_A^{-1}$$

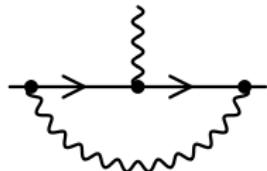
# Direct calculation



$$ie_0 \Lambda^\alpha = \int \frac{d^d k}{(2\pi)^d} ie_0 \gamma^\mu i \frac{\not{k}}{k^2} ie_0 \gamma^\alpha i \frac{\not{k}}{k^2} ie_0 \gamma^\nu$$
$$\times \frac{-i}{k^2} \left( g_{\mu\nu} - \xi \frac{k_\mu k_\nu}{k^2} \right)$$

$$\Lambda^\alpha = -ie_0^2 \int \frac{d^d k}{(2\pi)^d} \frac{\gamma_\mu \not{k} \gamma^\alpha \not{k} \gamma^\mu - \xi k^2 \gamma^\alpha}{(k^2)^2}$$

# Direct calculation



$$ie_0 \Lambda^\alpha = \int \frac{d^d k}{(2\pi)^d} ie_0 \gamma^\mu i \not{k}^2 ie_0 \gamma^\alpha i \not{k}^2 ie_0 \gamma^\nu \\ \times \frac{-i}{k^2} \left( g_{\mu\nu} - \xi \frac{k_\mu k_\nu}{k^2} \right)$$

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Averaging  $\not{k} \gamma^\alpha \not{k} \rightarrow (k^2/d) \gamma_\nu \gamma^\alpha \gamma^\nu$

$$\Lambda^\alpha = -ie_0^2 a_0 \gamma^\alpha \int \frac{d^d k}{(2\pi)^d} \frac{1}{(-k^2)^2}$$

# Direct calculation

$$\Gamma^\alpha = \gamma^\alpha \left[ 1 + a(\mu) \frac{\alpha(\mu)}{4\pi\varepsilon} \right]$$

$$Z_\Gamma = 1 + a \frac{\alpha}{4\pi\varepsilon}$$

agrees with  $Z_\psi$

# Charge renormalization

$e_0^2$  does not depend on  $\mu$ :

$$\frac{d \log \alpha(\mu)}{d \log \mu} = -2\varepsilon - 2\beta(\alpha(\mu))$$

$$\beta(\alpha_s(\mu)) = \frac{1}{2} \frac{d \log Z_\alpha(\alpha_s(\mu))}{d \log \mu}$$

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For a minimal renormalization constant

$$Z_\alpha(\alpha) = 1 + z_1 \frac{\alpha}{4\pi\varepsilon} + \dots$$

we obtain

$$\beta(\alpha) = \beta_0 \frac{\alpha}{4\pi} + \dots = -z_1 \frac{\alpha}{4\pi} + \dots$$

$$Z_\alpha(\alpha) = 1 - \beta_0 \frac{\alpha}{4\pi\varepsilon} + \dots$$

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$$\frac{d}{d \log \mu} \frac{\alpha(\mu)}{4\pi} = -2\beta_0 \left( \frac{\alpha(\mu)}{4\pi} \right)^2$$

$$\frac{d}{d \log \mu} \frac{4\pi}{\alpha(\mu)} = 2\beta_0$$

$$\frac{4\pi}{\alpha(\mu')} - \frac{4\pi}{\alpha(\mu)} = 2\beta_0 \log \frac{\mu'}{\mu}$$

$$\alpha(\mu') = \frac{\alpha(\mu)}{1 + 2\beta_0 \frac{\alpha(\mu)}{4\pi} \log \frac{\mu'}{\mu}}$$