



An Introduction to the Thermal Model.

I. Phenomenology

II. Theory

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Dense Matter, 29 June - 11 July 2015
Dubna, JINR



Outline

I. Phenomenology

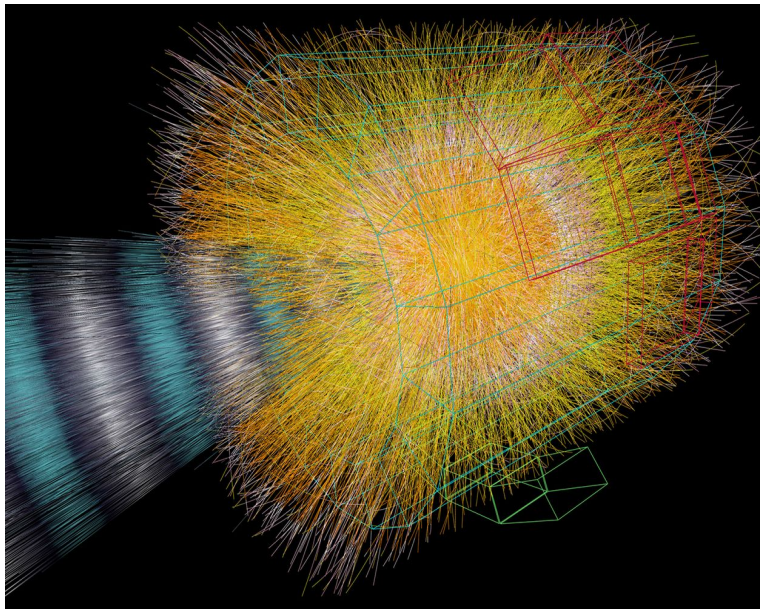
Heavy Ion Collisions at the LHC

Thermal Model

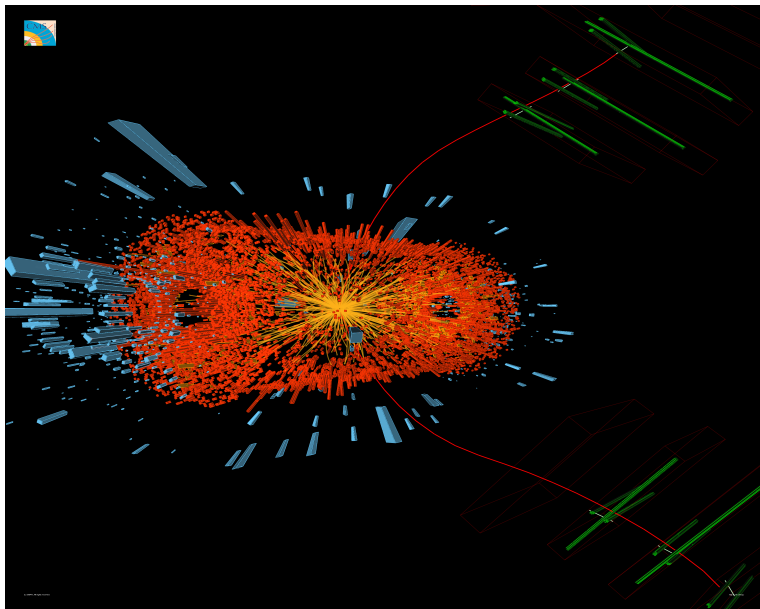
Hagedorn temperature

Heavy Ion Collisions at NICA/FAIR

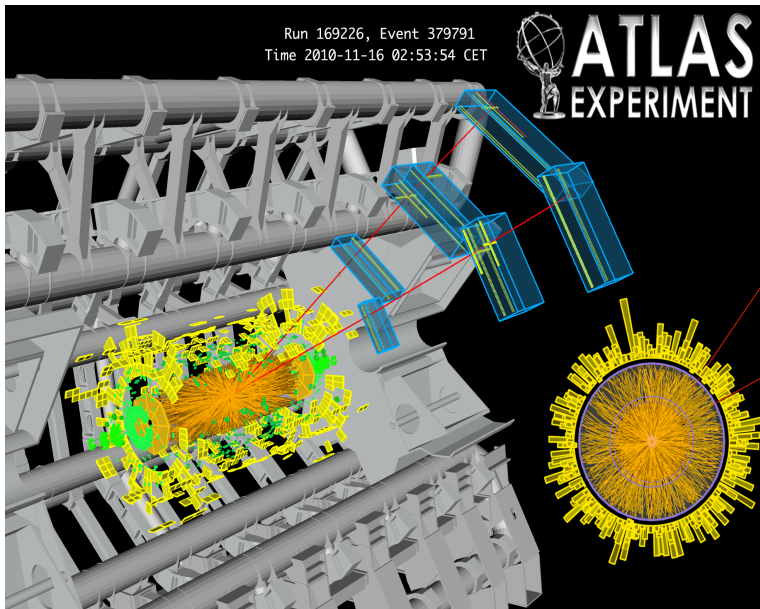
Heavy Ion Collisions in ALICE



Heavy Ion Collisions in CMS



Heavy Ion Collisions in ATLAS



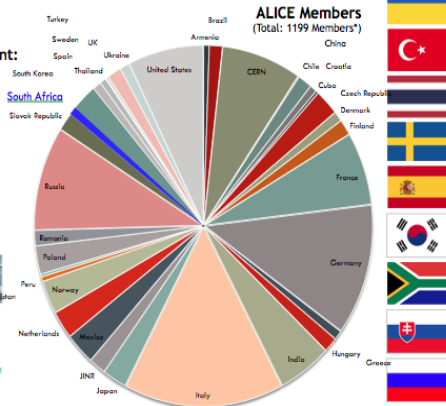
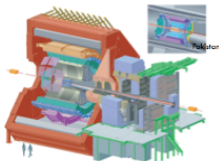
South Africa

The ALICE Collaboration

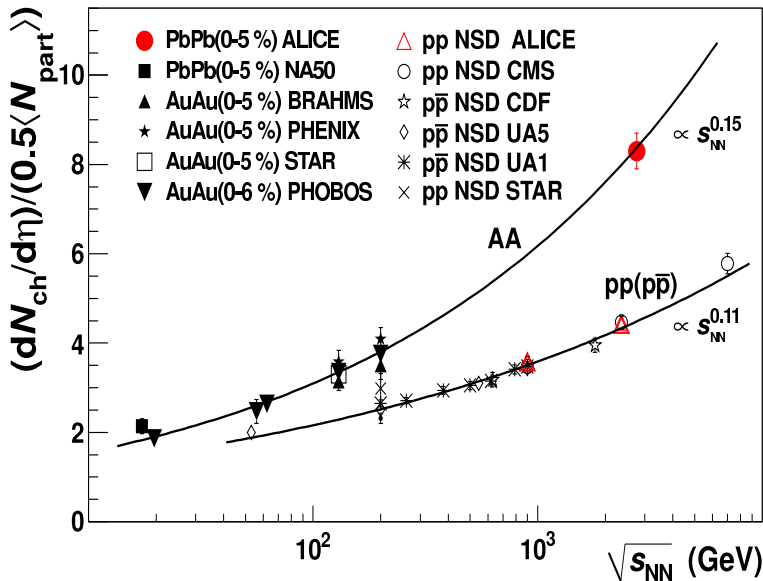
35 Countries - 124 Institutes - 158 MCHF capital cost

History of the ALICE Experiment:

1990-1996 Design
 1992-2002 R&D
 2000-2010 Construction
 2002-2007 Installation
 2008 -> Commissioning
 4 TP addenda along the way:
 1996 Muon spectrometer
 1999 TRD
 2006 EMCAL
 2010 DCAL



Particle Multiplicity in Heavy Ion Collisions



Particle Multiplicity in Heavy Ion Collisions

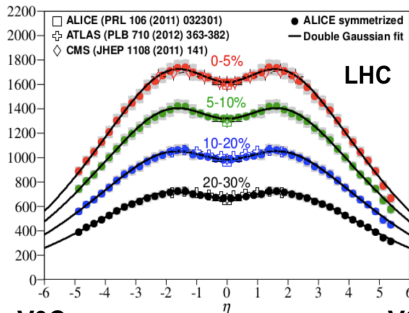


Acceptance for charged particles



ALICE

η coverages
for $z_{\text{vtx}}=0$
(shown at last
AW)



Now:
(T0 now shown)

V0C



V0A

SPD (outer layer)

FMDC



T0A+ (but not hermetic: 80%)

T0C+ (ext)



V0A+ (ext)

After LS2:

MFT



ITS IB (middle layer)

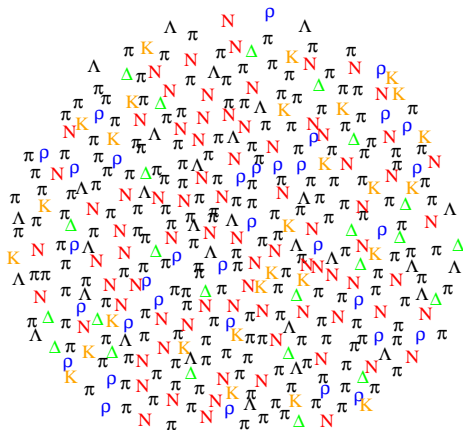
This is (-3.6,-2.5), i.e. the MFT+MUON acc.

Particle Multiplicity in Heavy Ion Collisions

About 16 000 particles are produced in a heavy ion collision.

Hence: Use Concepts from Statistical Mechanics to analyze the final state
e.g. use Energy Density, Particle Density, Pressure, Temperature, Chemical Composition, ...

Hadronic Gas before Chemical Freeze-Out



J.C. and H. Satz, Z. fuer Physik C57, 135, 1993.

Thermal Equilibrium

In thermal equilibrium

$$Z = \text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}}$$

$$\langle N \rangle = \frac{\text{Tr} N e^{-\frac{H}{T} + \frac{\mu N}{T}}}{\text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}}}$$

$$\langle E \rangle = \frac{\text{Tr} E e^{-\frac{H}{T} + \frac{\mu N}{T}}}{\text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}}}$$

Thermal Equilibrium

Particle Number

$$\begin{aligned}
 \langle N \rangle &= \frac{\text{Tr} N e^{-\frac{H}{T} + \frac{\mu N}{T}}}{\text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}}} \\
 &= \frac{T}{Z} \frac{\partial}{\partial \mu} \text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}} \\
 &= T \frac{1}{Z} \frac{\partial Z}{\partial \mu} \\
 &= T \frac{\partial}{\partial \mu} \ln Z
 \end{aligned}$$

Thermal Equilibrium

Average Energy

$$\begin{aligned}
 \langle E \rangle &= \frac{\text{Tr } H e^{\frac{-H}{T} + \frac{\mu N}{T}}}{\text{Tr } e^{\frac{-H}{T} + \frac{\mu N}{T}}} \\
 &= -\frac{1}{Z} \frac{\partial Z}{\partial \beta} + \mu \langle N \rangle \\
 &= T^2 \frac{\partial}{\partial T} \ln Z + \mu \langle N \rangle
 \end{aligned}$$

Thermal Equilibrium

$$\begin{aligned}
 N_i &= g_i V \int \frac{d^3 p}{(2\pi)^3} \exp\left(-\frac{E}{T}\right) e^{\frac{\mu_i}{T}} \\
 &= g_i V \frac{4\pi}{(2\pi)^3} \int p^2 dp \exp\left(-\frac{\sqrt{p^2 + m_i^2}}{T}\right) e^{\frac{\mu_i}{T}} \\
 &= g_i V \frac{4\pi}{(2\pi)^3} T^3 \int x^2 dx \exp\left(-\sqrt{x^2 + m_i^2/T^2}\right) e^{\frac{\mu_i}{T}} \\
 &= g_i V \frac{1}{2\pi^2} T m_i^2 K_2\left(\frac{m_i}{T}\right) e^{\frac{\mu_i}{T}}
 \end{aligned}$$

Full Hydrodynamic Flow

Bjorken scaling + Transverse expansion

After integration over m_T

$$\frac{dN_i/dy}{dN_j/dy} = \frac{N_i^0}{N_j^0}$$

where N_i^0 is the particle yield
as calculated in a fireball **AT REST!**

Effects of hydrodynamic flow cancel out in ratios.

Thermal Equilibrium

$$n_i = g_i \frac{1}{2\pi^2} T m_i^2 K_2 \left(\frac{m_i}{T} \right) e^{\frac{\mu_i}{T}}$$

$$\epsilon_i = g_i \frac{1}{2\pi^2} T m_i^3 \left[K_1 \left(\frac{m_i}{T} \right) + 3 \frac{T}{m} K_2 \left(\frac{m_i}{T} \right) \right] e^{\frac{\mu_i}{T}}$$

$$s_i = g_i \frac{1}{2\pi^2} m_i^3 \left[K_1 \left(\frac{m_i}{T} \right) + \frac{4T}{m} K_2 \left(\frac{m_i}{T} \right) - \frac{\mu_i}{m} K_2 \left(\frac{m_i}{T} \right) \right] e^{\frac{\mu_i}{T}}$$

$$P_i = g_i \frac{1}{2\pi^2} T^2 m_i^2 K_2 \left(\frac{m_i}{T} \right) e^{\frac{\mu_i}{T}}$$

Chemical Equilibrium

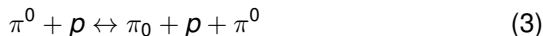
In equilibrium

$$E_1 + E_2 + \dots = E_3 + E_4 + E_5 + \dots \quad (1)$$

for the chemical potentials

$$\mu_1 + \mu_2 + \dots = \mu_3 + \mu_4 + \mu_5 + \dots \quad (2)$$

As an example



leads to

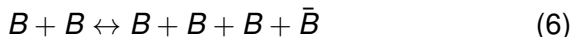
$$\mu_{\pi^0} + \mu_p = \mu_{\pi^0} + \mu_p + \mu_{\pi^0} \quad (4)$$

which leads to

$$\mu_{\pi^0} = 0 \quad (5)$$

Chemical Equilibrium

In equilibrium



$$dE = -pdV + TdS + \mu_B dN_B + \mu_{\bar{B}} dN_{\bar{B}}$$

Due to baryon number conservation one has

$$N_B - N_{\bar{B}} = \text{constant}$$

and

$$dN_B = dN_{\bar{B}}$$

The energy is a minimum for

$$dE = (\mu_B + \mu_{\bar{B}})dN_B = 0 \quad (7)$$

$$\mu_B = -\mu_{\bar{B}} \quad (8)$$

Chemical Equilibrium

In equilibrium

$$N_B = g V \int \frac{d^3 p}{(2\pi)^3} \exp\left(-\frac{E}{T} + \frac{\mu_B}{T}\right)$$

$$N_{\bar{B}} = g V \int \frac{d^3 p}{(2\pi)^3} \exp\left(-\frac{E}{T} - \frac{\mu_B}{T}\right)$$

$$N_B = N_{\bar{B}} \rightarrow \mu_B = 0$$

$$N_B \geq N_{\bar{B}} \rightarrow \mu_B \geq 0$$

$$N_B \leq N_{\bar{B}} \rightarrow \mu_B \leq 0$$

	Chemical Equilibrium	No Chem. Equil.
π	$\exp\left[-\frac{E_\pi}{T}\right]$	$\exp\left[-\frac{E_\pi}{T} + \frac{\mu_\pi}{T}\right]$
N	$\exp\left[-\frac{E_N}{T} + \frac{\mu_B}{T}\right]$	$\exp\left[-\frac{E_N}{T} + \frac{\mu_N}{T}\right]$
\bar{N}	$\exp\left[-\frac{E_N}{T} - \frac{\mu_B}{T}\right]$	$\exp\left[-\frac{E_N}{T} + \frac{\mu_{\bar{N}}}{T}\right]$
Λ	$\exp\left[-\frac{E_\Lambda}{T} + \frac{\mu_B}{T} - \frac{\mu_S}{T}\right]$	$\exp\left[-\frac{E_\Lambda}{T} + \frac{\mu_\Lambda}{T}\right]$
$\bar{\Lambda}$	$\exp\left[-\frac{E_\Lambda}{T} - \frac{\mu_B}{T} + \frac{\mu_S}{T}\right]$	$\exp\left[-\frac{E_\Lambda}{T} + \frac{\mu_{\bar{\Lambda}}}{T}\right]$
K	$\exp\left[-\frac{E_K}{T} + \frac{\mu_S}{T}\right]$	$\exp\left[-\frac{E_K}{T} + \frac{\mu_K}{T}\right]$
\bar{K}	$\exp\left[-\frac{E_K}{T} - \frac{\mu_S}{T}\right]$	$\exp\left[-\frac{E_K}{T} + \frac{\mu_{\bar{K}}}{T}\right]$

The number of particles of type i is determined by:

$$N_i = V g_i \int \frac{d^3 p}{(2\pi)^3} \exp\left(-\frac{E_i}{T} + \frac{\mu_i}{T}\right)$$

For bosons:

$$N_i = V g_i \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\exp\left(\frac{E_i}{T} - \frac{\mu_i}{T}\right) - 1}$$

For fermions:

$$N_i = V g_i \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\exp\left(\frac{E_i}{T} - \frac{\mu_i}{T}\right) + 1}$$

Chemical Equilibrium

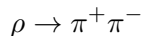
Only conserved quantum numbers matter for chemical equilibrium: In equilibrium

$$\mu_i = B_i \mu_B + Q_i \mu_Q + S_i \mu_S + C_i \mu_C + \dots \quad (9)$$

g_i	m_i	stat	S_i	B_i	Q_i	Particle i
1	0.140	-1	0	0	1.	π^+
1	0.135	-1	0	0	0.	π^0
1	0.140	-1	0	0	-1.	π^-
1	0.547	-1	0	0	0.	η
3	0.770	-1	0	0	1.	ρ^+
3	0.770	-1	0	0	0.	ρ^0
3	0.770	-1	0	0	-1.	ρ^-
3	0.782	-1	0	0	0.	ω
1	0.958	-1	0	0	0.	η'
1	0.980	-1	0	0	0.	f_0
1	0.982	-1	0	0	1.	a_0^+
1	0.982	-1	0	0	0.	a_0^0
1	0.982	-1	0	0	-1.	a_0^-
3	1.019	-1	0	0	0.	ϕ
3	1.170	-1	0	0	0.	
3	1.230	-1	0	0	1.	
3	1.230	-1	0	0	0.	
3	1.230	-1	0	0	-1.	
3	1.229	-1	0	0	1.	
3	1.229	-1	0	0	0.	
3	1.229	-1	0	0	-1.	
5	1.275	-1	0	0	0.	
3	1.282	-1	0	0	0.	
1	1.297	-1	0	0	0.	
1	1.300	-1	0	0	1.	
1	1.300	-1	0	0	0.	

The Role of Resonances

Example: ρ 's



Final, observed, number of π^+ is given by

$$N_{\pi^+} = N_{\pi^+}(\text{thermal}) + N_{\pi^+}(\text{resonance decays})$$

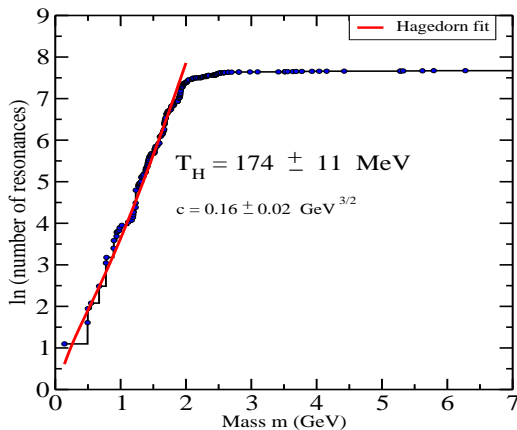
depending on the temperature, over 80% of observed pions are due to resonance decays

The Role of Resonances

Resonances are very important in the thermal model.

Without resonances the thermal model doesn't work.

THE HAGEDORN TEMPERATURE.



Keep on adding the number of hadronic resonances.

J.C. and Dawit Worku, Mod. Phys. Lett. A26 (2011) 1197; arXiv:
1103.1463

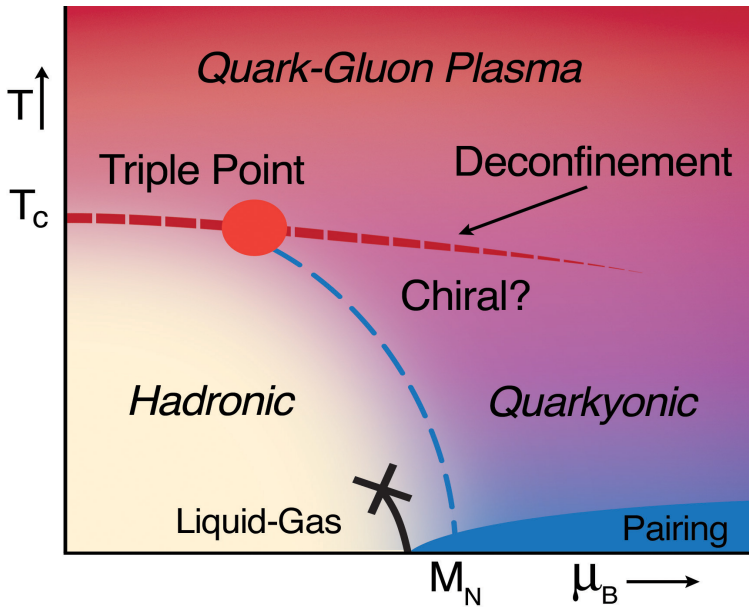
HADRONS DO NOT EXIST ABOVE THE HAGEDORN TEMPERATURE.

Thermodynamic quantities like particle density, energy density, pressure ,... all involve a summation over hadron species:

$$\sum_i \exp -\frac{E_i}{T}$$

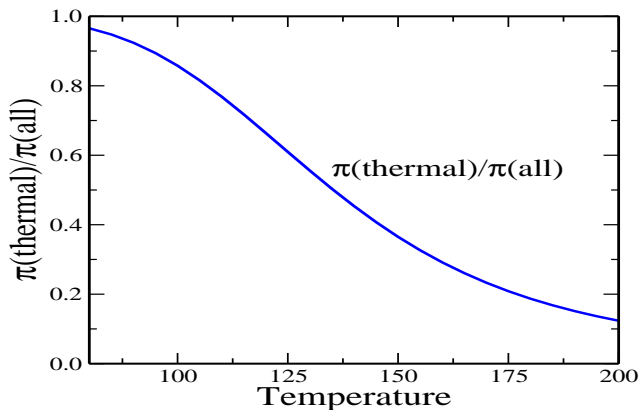
and the sum becomes (too) large due to the number of resonances.

Phase Diagram

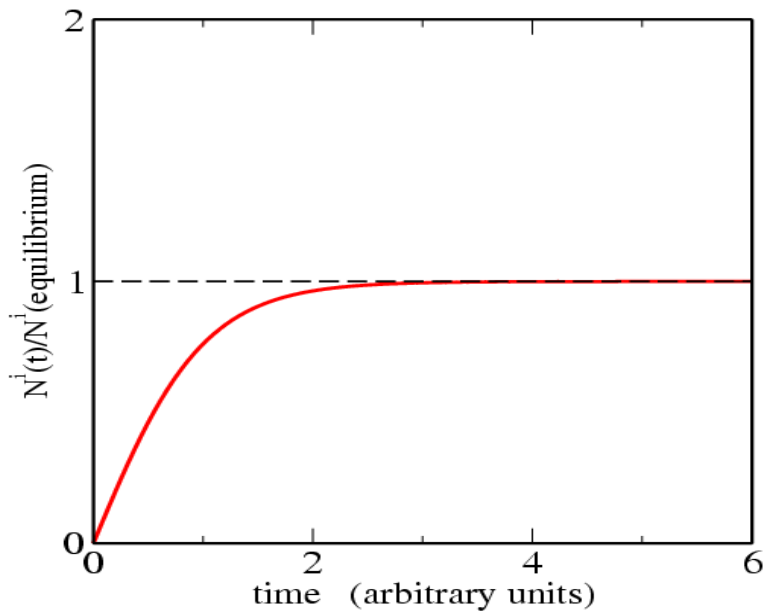


g_i	m_i	stat	S_i	B_i	Q_i	BR $\rightarrow \pi^+$	Particle i
1	0.140	-1	0	0	1.	1.000	π^+
1	0.135	-1	0	0	0.	0.000	π^0
1	0.140	-1	0	0	-1.	0.000	π^-
1	0.547	-1	0	0	0.	0.285	η
3	0.770	-1	0	0	1.	1.000	ρ^+
3	0.770	-1	0	0	0.	1.000	ρ^0
3	0.770	-1	0	0	-1.	0.000	ρ^-
3	0.782	-1	0	0	0.	0.910	ω
1	0.958	-1	0	0	0.	0.965	η'
1	0.980	-1	0	0	0.	0.521	f_0
1	0.982	-1	0	0	1.	1.285	a_0^+
1	0.982	-1	0	0	0.	0.285	a_0^0
1	0.982	-1	0	0	-1.	0.285	a_0^-
3	1.019	-1	0	0	0.	0.155	ϕ
3	1.170	-1	0	0	0.	1.000	h_1
3	1.230	-1	0	0	1.	1.500	
3	1.230	-1	0	0	0.	0.50	
3	1.230	-1	0	0	-1.	0.50	
3	1.229	-1	0	0	1.	1.91	
3	1.229	-1	0	0	0.	0.91	
3	1.229	-1	0	0	-1.	0.91	
5	1.275	-1	0	0	0.	0.69	
3	1.282	-1	0	0	0.	1.00	
1	1.297	-1	0	0	0.	1.11	
1	1.300	-1	0	0	1.	2.00	
1	1.300	-1	0	0	0.	1.50	

Importance of Resonances.



Strangeness saturation?



Strangeness saturation?

$$N_i = \boxed{\gamma_s^{|S|}} V g_i \int \frac{d^3 p}{(2\pi)^3} \exp\left(-\frac{E_i}{T} + \frac{\mu_i}{T}\right)$$

with

$\gamma_s < 1$ strangeness under-saturation

$\gamma_s = 1$ strangeness in chemical equilibrium

$\gamma_s > 1$ strangeness over-saturation

SPS data.

	Measurement
Pb–Pb 158A GeV	
$(\pi^+ + \pi^-)/2.$	600 ± 30
K^+	95 ± 10
K^-	50 ± 5
K_S^0	60 ± 12
p	140 ± 12
\bar{p}	10 ± 1.7
ϕ	7.6 ± 1.1
Ξ^-	4.42 ± 0.31
Ξ^-	0.74 ± 0.04
$\bar{\Lambda}/\Lambda$	0.2 ± 0.04

SPS data.

SPS: Chemical Freeze-Out Parameters:

$$T = 156.0 \pm 2.4 \text{ MeV}$$

$$\mu_B = 239 \pm 12 \text{ MeV}$$

$$\gamma_s = 0.862 \pm 0.036$$

F. Becattini, J.C., A. Keränen, E. Suhonen and K. Redlich
Physical Review C64 (2001) 024901.

AGS data.

	Measurement
Au–Au 11.6A GeV	
Participants	363 ± 10
K^+	23.7 ± 2.9
K^-	3.76 ± 0.47
π^+	133.7 ± 9.9
Λ	20.34 ± 2.74
p/π^+	1.234 ± 0.126
\bar{p}	$>0.0185 \pm 0.0018$

AGS data.

AGS: Chemical Freeze-Out Parameters:

$$T = 130.6 \pm 5.5 \text{ MeV}$$

$$\mu_B = 594 \pm 26 \text{ MeV}$$

$$\gamma_s = 0.883 \pm 0.124$$

F. Becattini, J.C., A. Keränen, E. Suhonen and K. Redlich
Physical Review C64 (2001) 024901.

SIS data.

	Measurement
Au–Au 1.7A GeV	
π^+/p	0.052 ± 0.013
K^+/π^+	0.003 ± 0.00075
π^-/π^+	2.05 ± 0.51
η/π^0	0.018 ± 0.007

SIS data.

SIS: Chemical Freeze-Out Parameters:

$$T = 49.7 \pm 1.1 \text{ MeV}$$

$$\mu_B = 818 \pm 15 \text{ MeV}$$

$$\gamma_s = 1 \text{ (fixed)}$$

J. C., H. Oeschler and K. Redlich)
Physical Review C59, (1999) 1663.

RHIC data.

J. C., B. Kämpfer, M. Kaneta, S. Wheaton, N. Xu, Phys. Rev. C71, 0409071 (2005)

Ratio	Experiment	Central	Mid-Central	Peripheral
$\pi_{(2)}^- / \pi_{(2)}^+$	BRAHMS	0.990±0.100		
	PHENIX	0.960±0.177	0.920±0.170	0.933±0.172
	PHOBOS	1.000±0.022		
	STAR	1.000±0.073	1.000±0.073	1.000 ± 0.073
$K_{(2)}^+ / K_{(2)}^-$	PHENIX	1.152±0.240	1.292±0.268	1.322±0.284
	PHOBOS	1.099±0.111		
	STAR	1.109±0.022	1.105±0.036	1.120±0.040
$\bar{p}_{(1)} / p_{(1)}$	PHENIX	0.680±0.149	0.671±0.142	0.717±0.157
$\bar{p}_{(2)} / p_{(2)}$	BRAHMS	0.650±0.092		
	PHOBOS	0.600±0.072		
	STAR	0.714±0.050	0.724±0.050	0.764±0.053
$\bar{\Lambda}_{(1)} / \Lambda_{(1)}$	PHENIX	0.750±0.180	0.798±0.197	0.795±0.197
$\bar{\Lambda}_{(2)} / \Lambda_{(2)}$	STAR	0.719±0.090	0.739±0.092	0.744±0.100
$\Xi_{(2)}^+ / \Xi_{(2)}^-$	STAR	0.840±0.053	0.822±0.114	0.815±0.096
$\bar{\Omega}^+ / \Omega^-$	STAR	1.062±0.410		
$K_{(2)}^- / \pi_{(2)}^-$	PHENIX	0.151±0.030	0.134±0.027	0.116±0.023
	STAR	0.151±0.022	0.147±0.022	0.130±0.019
$K_S^0 / \pi_{(2)}^-$	STAR	0.134±0.022	0.131±0.022	0.108±0.018
$\bar{p}_{(1)} / \pi_{(2)}^-$	PHENIX	0.049±0.010	0.047±0.010	0.045±0.009
$\bar{p}_{(2)} / \pi_{(2)}^-$	STAR	0.069±0.019	0.067±0.019	0.067±0.019
$\Lambda_{(1)} / \pi_{(2)}^-$	STAR	0.043±0.008	0.043±0.008	0.039±0.007
$\Lambda_{(2)} / \pi_{(2)}^-$	PHENIX	0.072±0.017	0.068±0.016	0.074±0.017
$< K^{*0} > / \pi_{(2)}^-$	STAR	0.039±0.011		
$\phi / \pi_{(2)}^-$	STAR	0.022±0.003	0.021±0.004	0.022±0.004
$\Xi_{(2)}^- / \pi_{(2)}^-$	STAR	0.0093±0.0012	0.0072±0.0011	0.0060±0.0008

RHIC data.

RHIC: Chemical Freeze-Out Parameters:

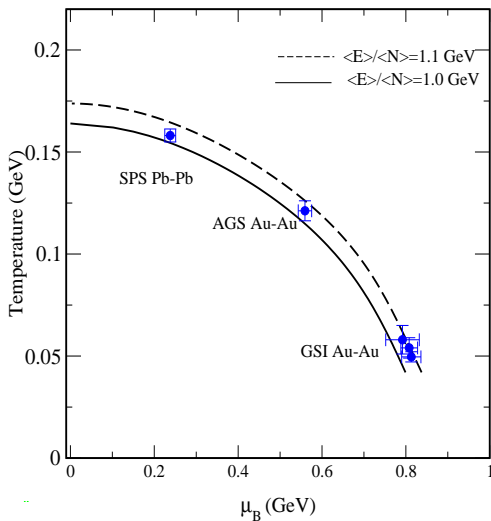
$$T = 169 \pm 4.2 \text{ MeV}$$

$$\mu_B = 39.6 \pm 6 \text{ MeV}$$

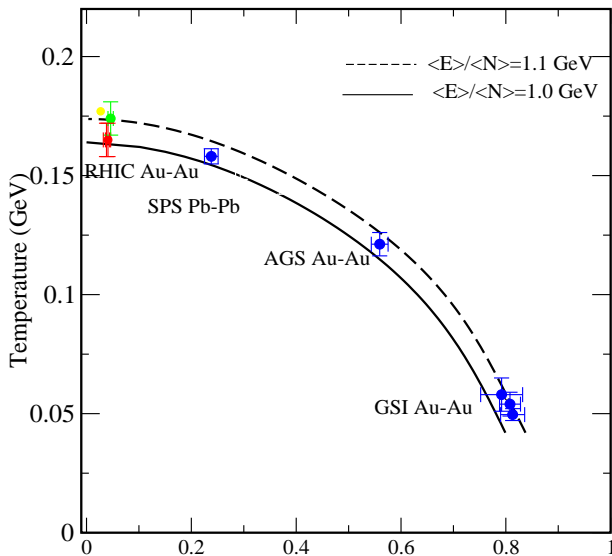
$$\gamma_s = 0.9 \pm 0.1$$

J. C., B. Kämpfer, M. Kaneta, S. Wheaton, N. Xu
Phys. Rev. C71, 0409071 (2005)

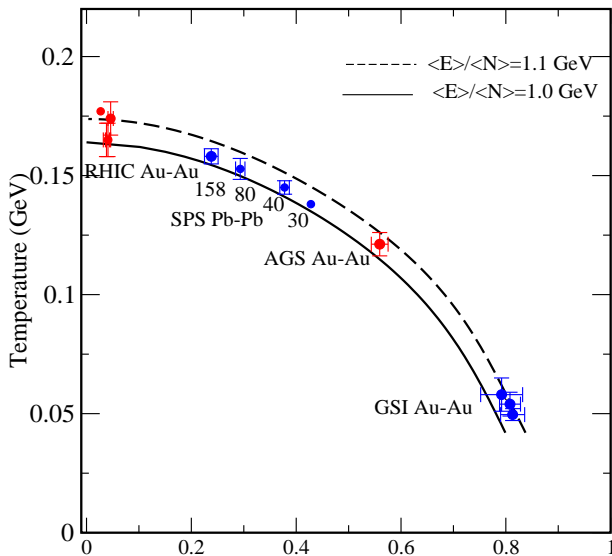
E/N in 1999



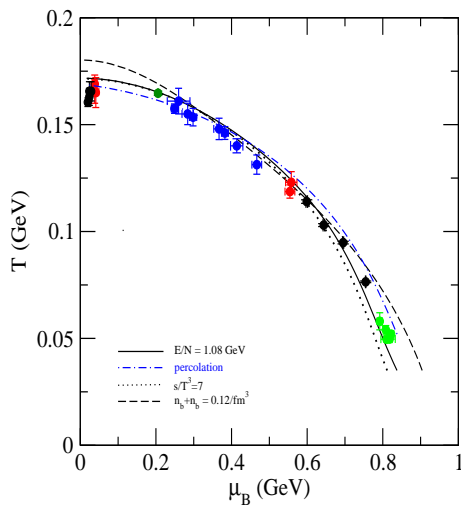
E/N in 2000



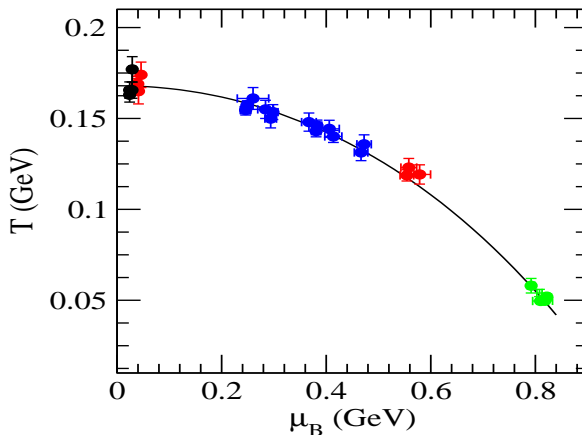
E/N in 2005



Chemical Freeze-Out: Status in 2005

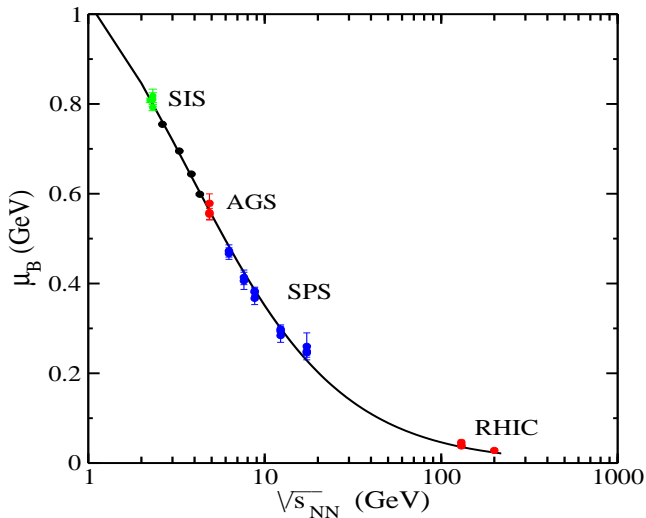


Chemical Freeze-Out: Status in 2005

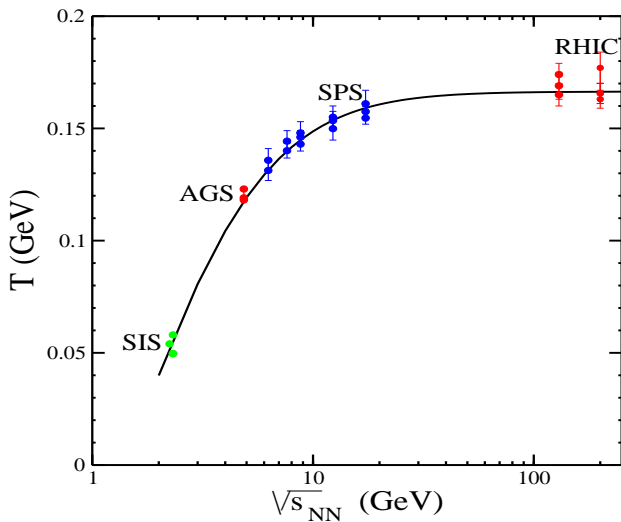


$$T = 0.166 - 0.139 \mu_B^2 - 0.053 \mu_B^4$$

J.C., H. Oeschler, K. Redlich, S. Wheaton,
PR C73, 034905 (2006)

μ_B as a function of $\sqrt{s_{NN}}$ 

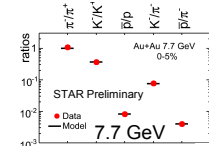
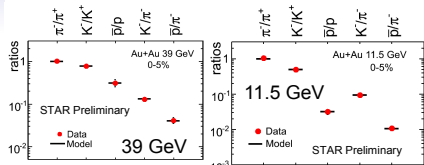
T as a function of $\sqrt{s_{NN}}$



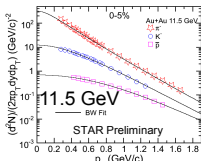
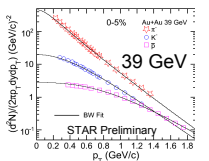


Freeze-out Conditions

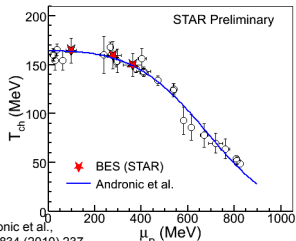
QuickTime™ and a TIFF (Uncompressed) decompressor are needed to see this picture.



Chemical freeze-out:
Particle ratios

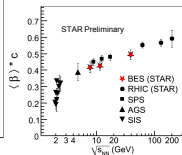
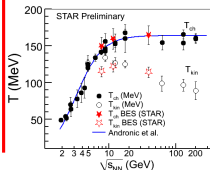


Kinetic freeze-out : Momentum distributions



Andronic et al.,
NPA 834 (2010) 237

QuickTime™ and a TIFF (Uncompressed) decompressor are needed to see this picture.



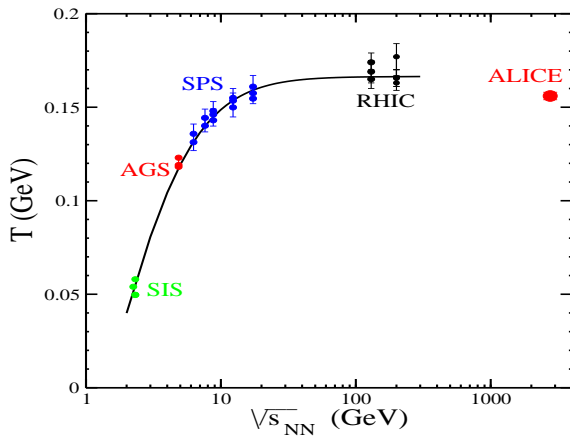
L. Kumar, Energy scan, 27th May

QM2011

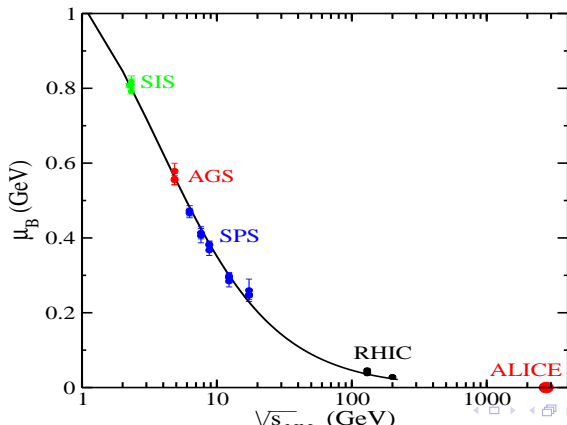
Bedanga Mohanty

4

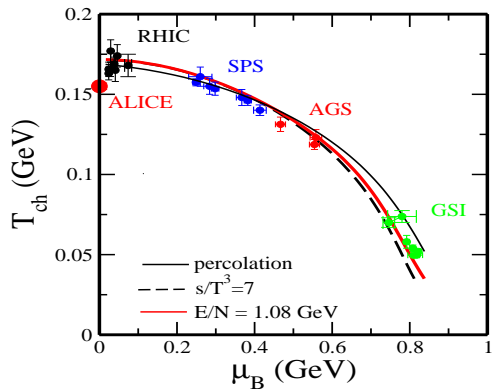
M. Floris: Nucl. Phys. A931, (2014) 103-112



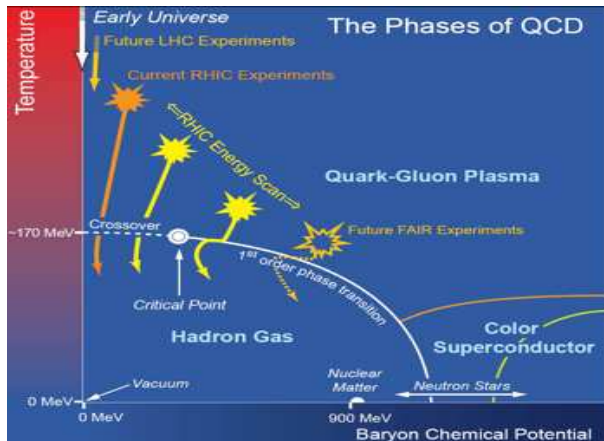
M. Floris: Nucl. Phys. A931, (2014) 103-112

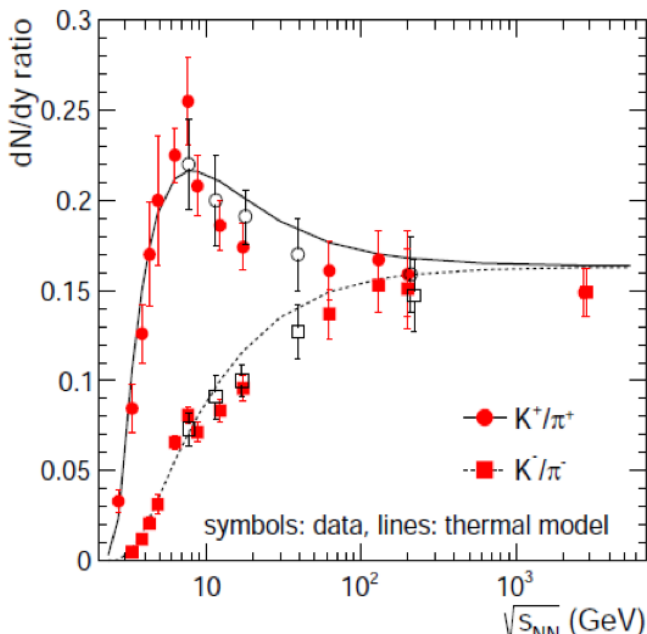


M. Floris: Nucl. Phys. A931, (2014) 103-112



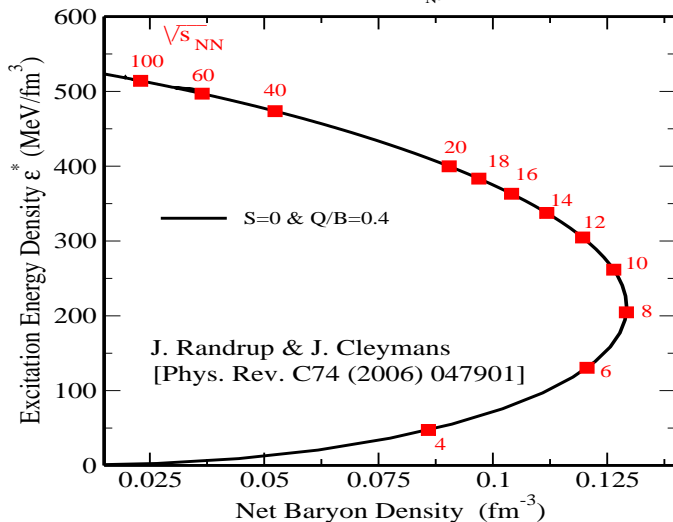
Phase Diagram





Hadronic Freeze-Out

$$\varepsilon^* = \varepsilon - m_N \rho$$



Strangeness in Heavy Ion Collisions vs Strangeness in pp - collisions

Use the Wroblewski factor

$$\lambda_s = \frac{2 \langle s\bar{s} \rangle}{\langle u\bar{u} \rangle + \langle d\bar{d} \rangle}$$

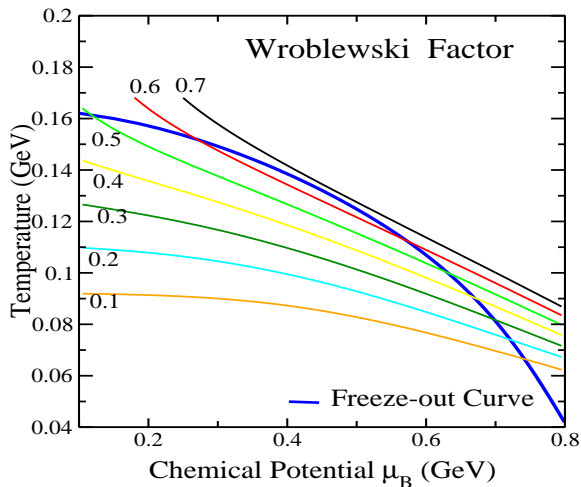
This is determined by the number of **newly** created quark – anti-quark pairs and **before** strong decays, i.e. before ρ 's and Δ 's decay.

Limiting values :

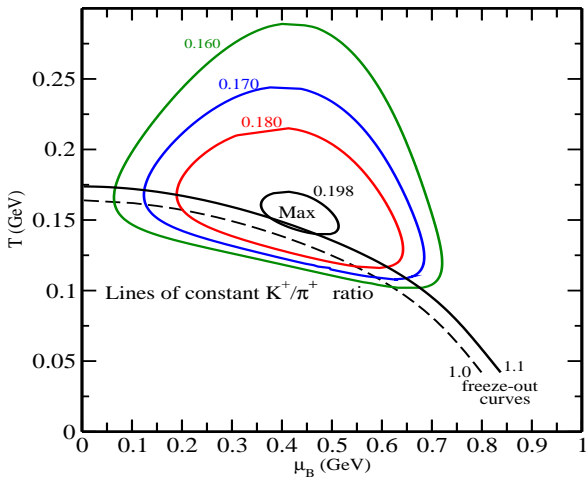
$\lambda_s = 1$ all quark pairs are equally abundant, SU(3) symmetry.

$\lambda_s = 0$ no strange quark pairs.

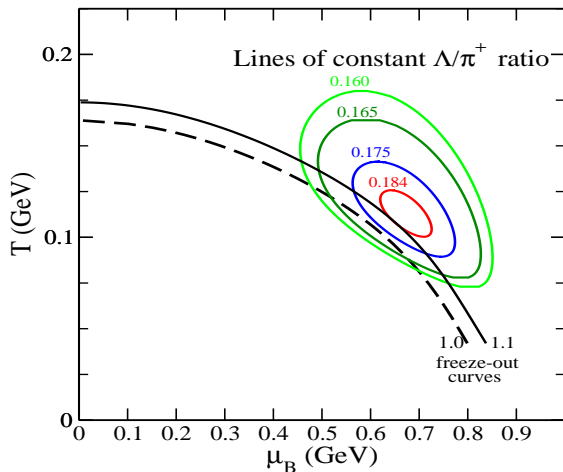
Maxima in particle ratios : Wroblewski factor

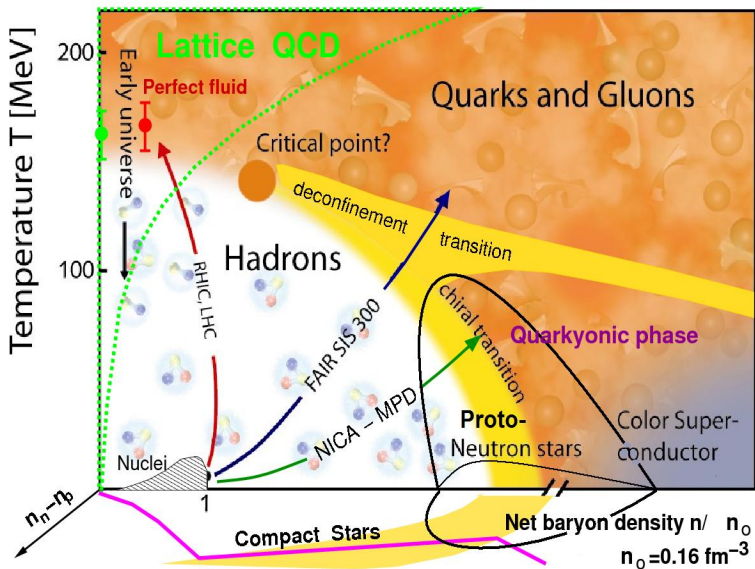


Maxima in particle ratios : K^+/π^+



Maxima in particle ratios : Λ/π^+





Good Luck NICA

