

Joint Institute for Nuclear Research

**VERY HIGH MULTIPLICITY PHYSICS**

*Dubna, April 7 – 9, 2001*

*Proceedings of the II International Workshop*

**Calculation of the Correlators  
for Very High Multiplicity Events  
at LHC**

**J. Budagov, G. Chelkov,  
Y. Kulchitsky, J. Manjavidze,  
A. Olchevsky & A. Sissakian**

*JINR, Dubna*

## 1 Correlation functions

- $K_1(E, n) = \langle \varepsilon \rangle = \int \frac{d^3 p}{(2\pi)^3 2\varepsilon(p)} \varepsilon(p) \frac{d^3 \sigma_n}{dp^3} = \int d\varepsilon \varepsilon \frac{dN}{d\varepsilon}$

If  $d\varepsilon \rightarrow 0$

than  $K_1(E, n) = \langle \varepsilon \rangle = \frac{1}{N} \sum_{i=1}^N E_i$

where  $E_i$  is i-particle energy.

- $K_2(E, n) = \langle (\varepsilon_1 - \langle \varepsilon \rangle) \cdot (\varepsilon_2 - \langle \varepsilon \rangle) \rangle = \langle \varepsilon^2 \rangle - \langle \varepsilon \rangle^2$

- $K_3(E, n) = \langle (\varepsilon_1 - \langle \varepsilon \rangle) \cdot (\varepsilon_2 - \langle \varepsilon \rangle) \cdot (\varepsilon_3 - \langle \varepsilon \rangle) \rangle = \langle \varepsilon^3 \rangle - 2\langle \varepsilon^2 \rangle \langle \varepsilon \rangle + 3\langle \varepsilon \rangle^3$

- Sissakian-Manjavidze prediction for VHM  $R = \frac{K_3^{2/3}}{K_2} \ll 1$

- If  $R = \frac{K_3^{2/3}}{K_2} \ll 1$

then the particle energy spectrum is Gaussian, with

— "temperature":  $K_1^{-1}(E, n)$

— dispersion:  $\sqrt{K_2(E, n)}$

## 2 PYTHIA Sub-processes:

- $q_i q_j \rightarrow q_i q_j$
- $q_i \bar{q}_i \rightarrow q_j \bar{q}_j$
- $q_i \bar{q}_i \rightarrow g g$
- $q_i g \rightarrow q_i g$
- $g g \rightarrow q_i \bar{q}_i$
- $g g \rightarrow g g$

## 2.1 PYTHIA prediction

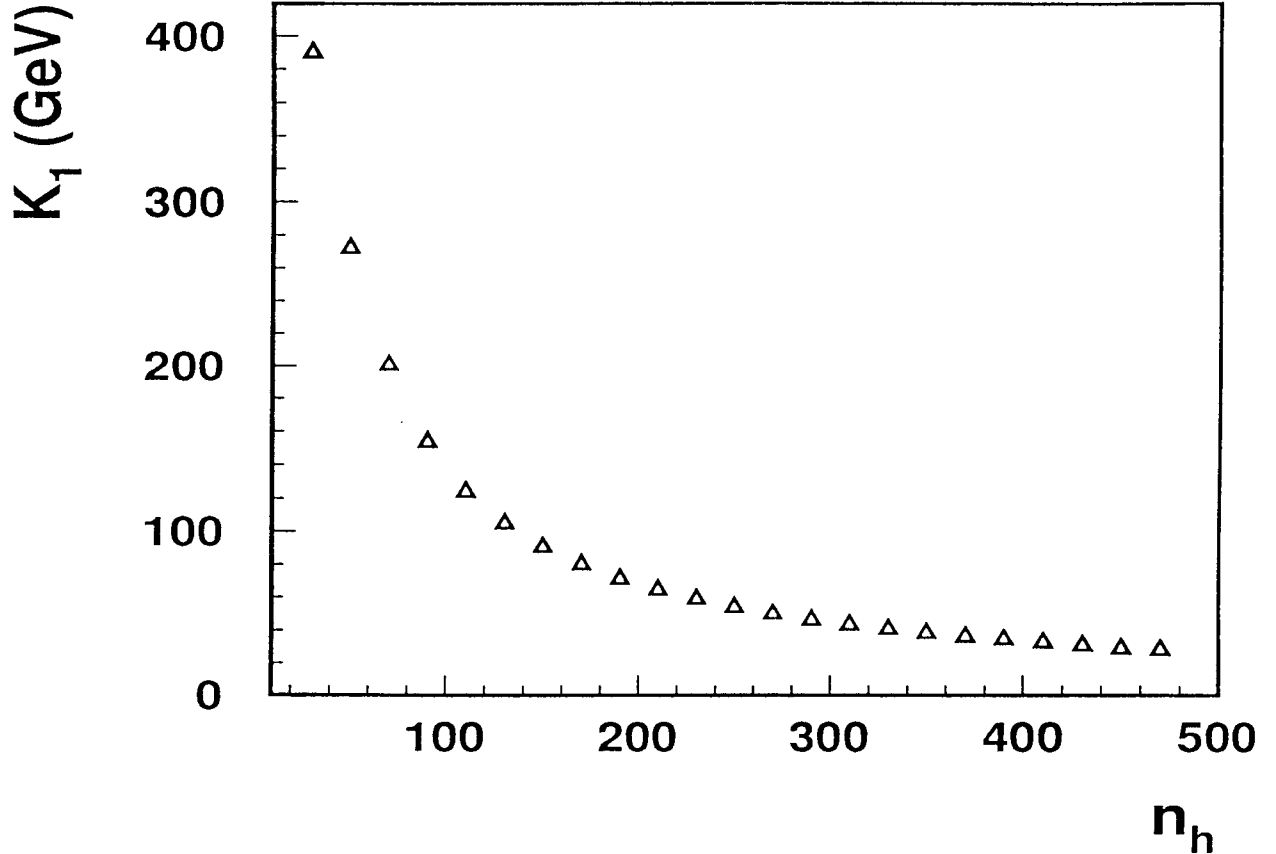


Figure 1: Dependent of  $K_1(\varepsilon, n_h)$  from number of hadrons without cut on the  $p_t$

$$K_1(\varepsilon, n_h) = \langle \varepsilon; n \rangle = \frac{1}{N_n} \int \varepsilon d\varepsilon \frac{dN_n(\varepsilon)}{d\varepsilon} = \frac{1}{N_n} \sum_{i=1}^{N_n} E_i$$

- $n$  - number of particles (hadrons)
- $\varepsilon$  - particles energy
- $N_n$  - number of events with multiplicity  $n_h$
- $dN_n(\varepsilon)/d\varepsilon$  - number of events with multiplicity  $n$  and particle with energy  $\varepsilon$

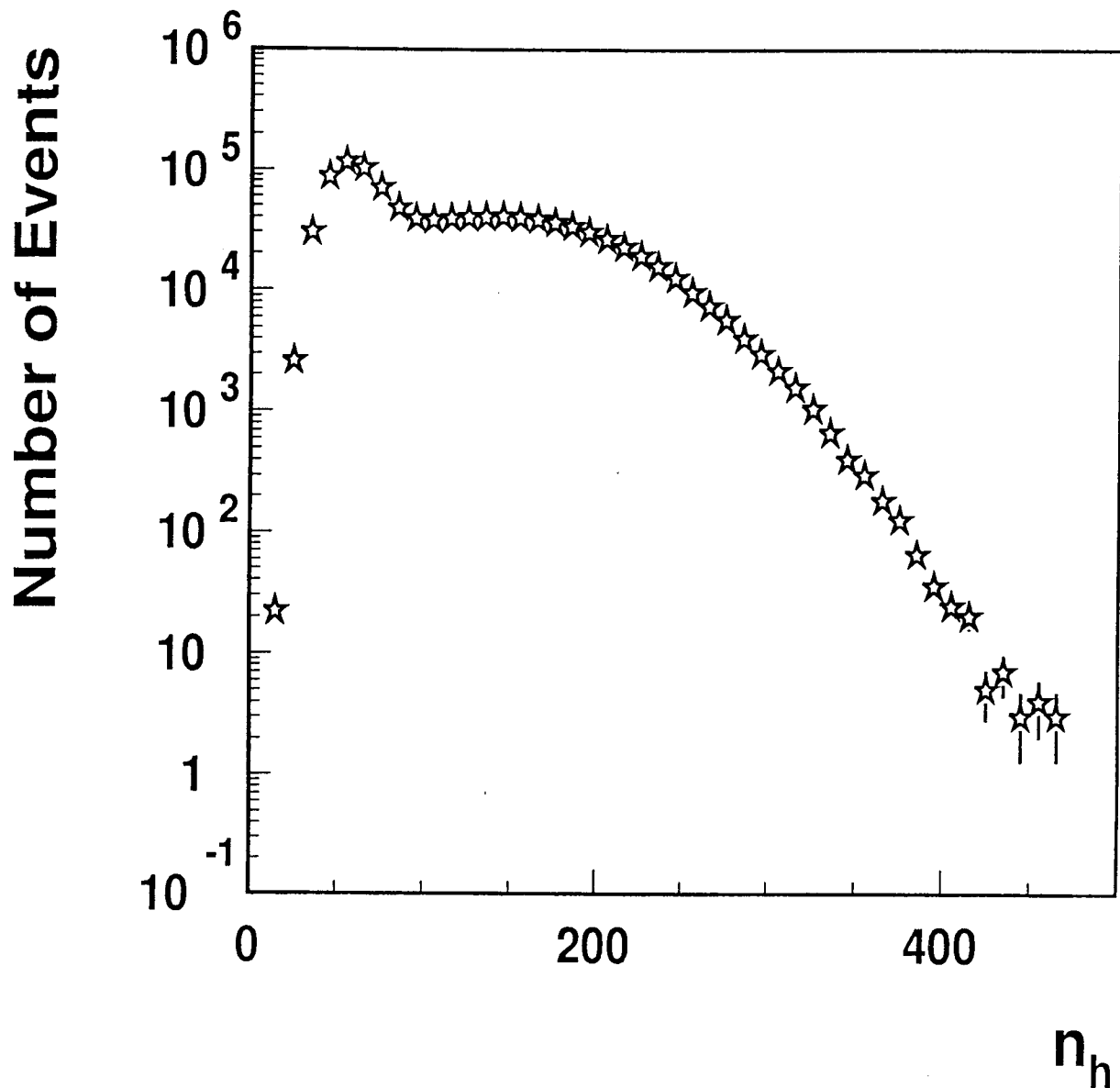


Figure 2: Multiplicity distribution. Without colored partons transverse momentum cutoff  $p_t$

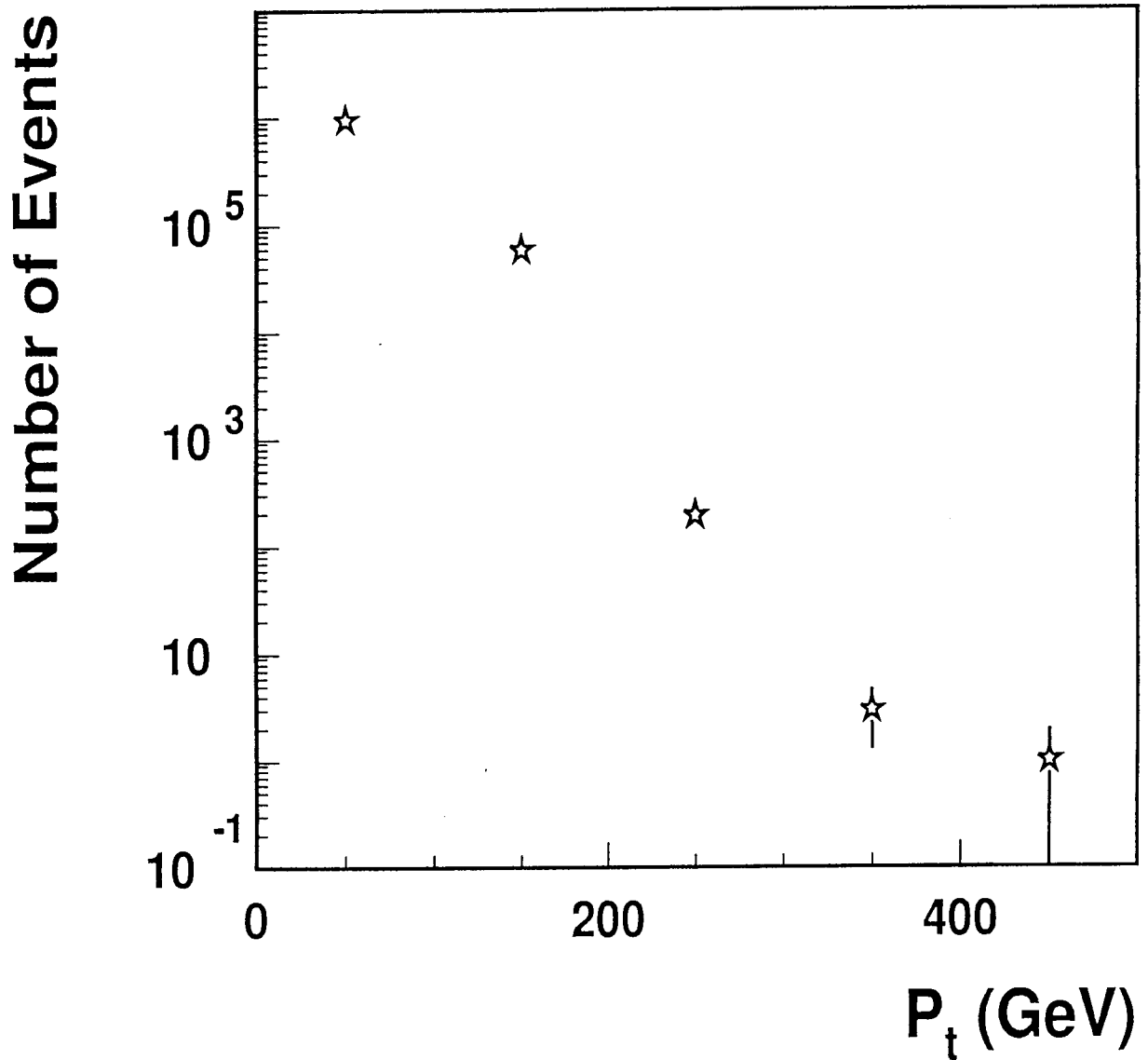


Figure 3:  $p_t$  distribution. Without colored partons transverse momentum cutoff  $p_t$

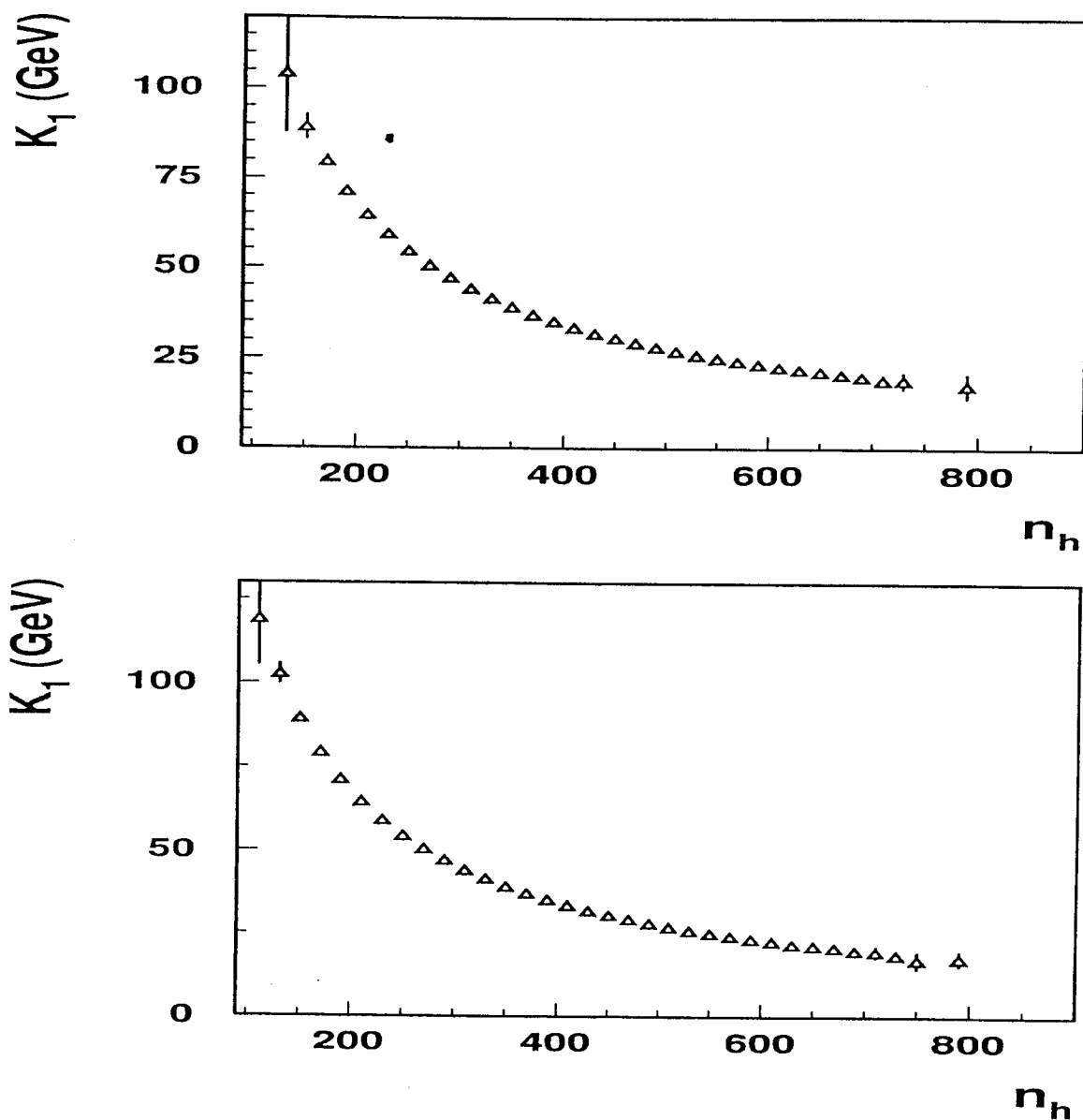


Figure 4: Dependent of  $K_1(\varepsilon, n_h)$  from number of hadrons. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)



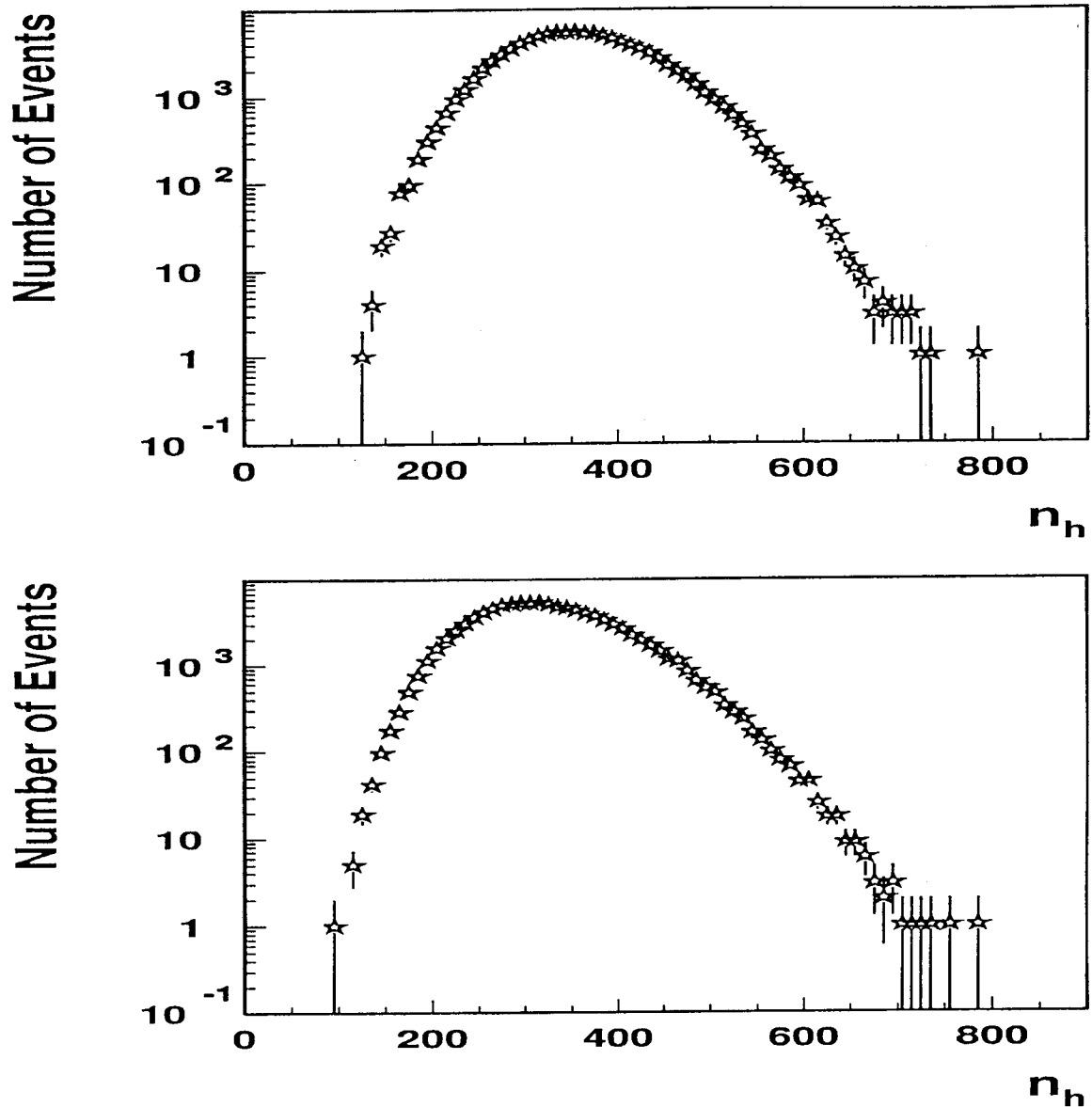


Figure 5: Multiplicity distribution. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

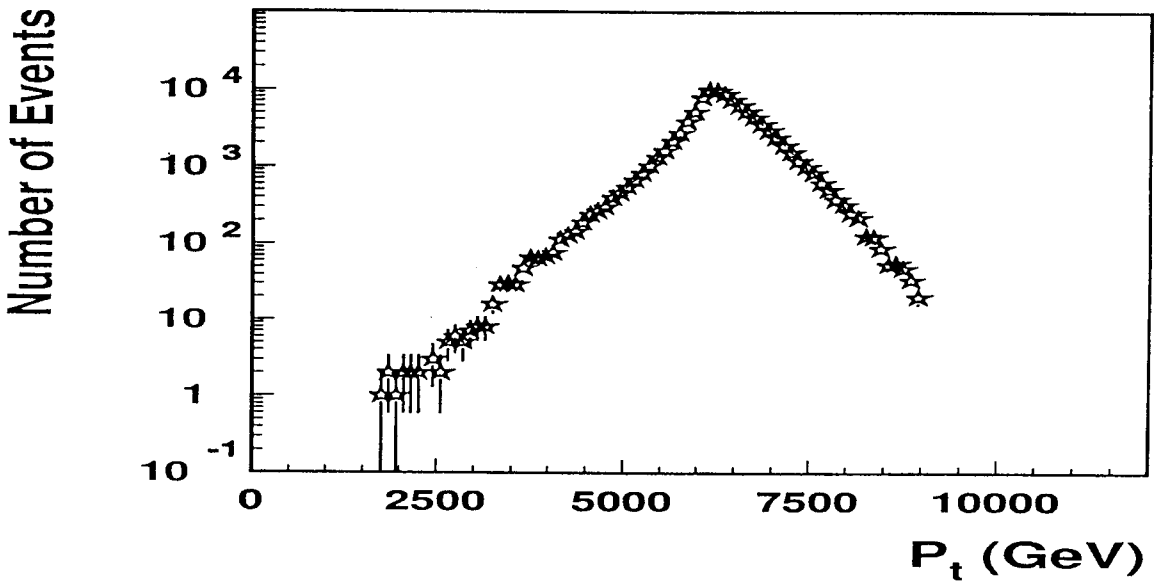
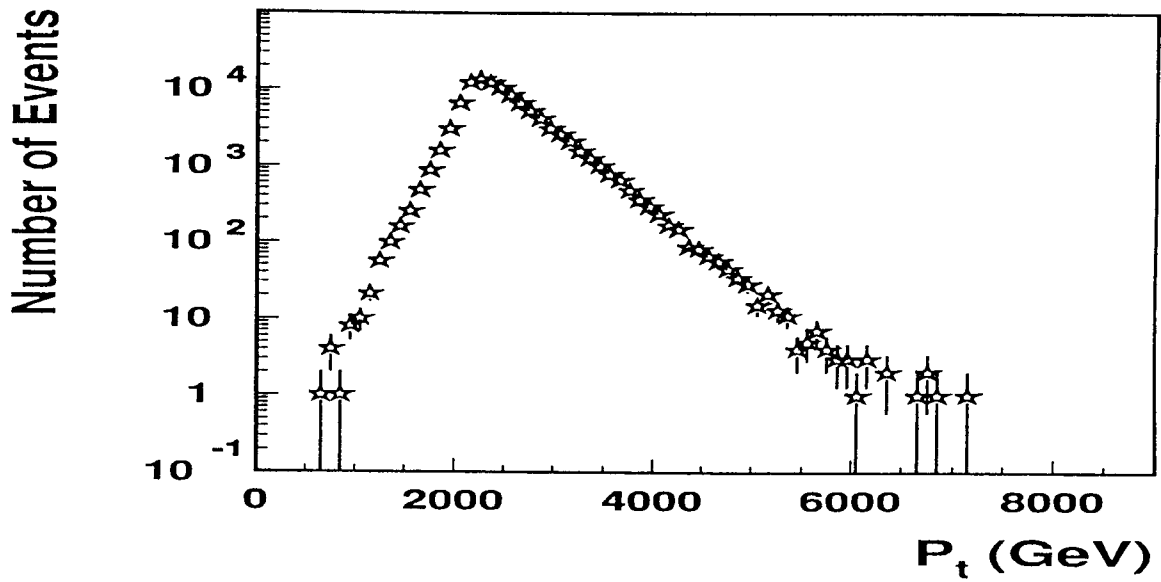


Figure 6:  $p_t$  distribution. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

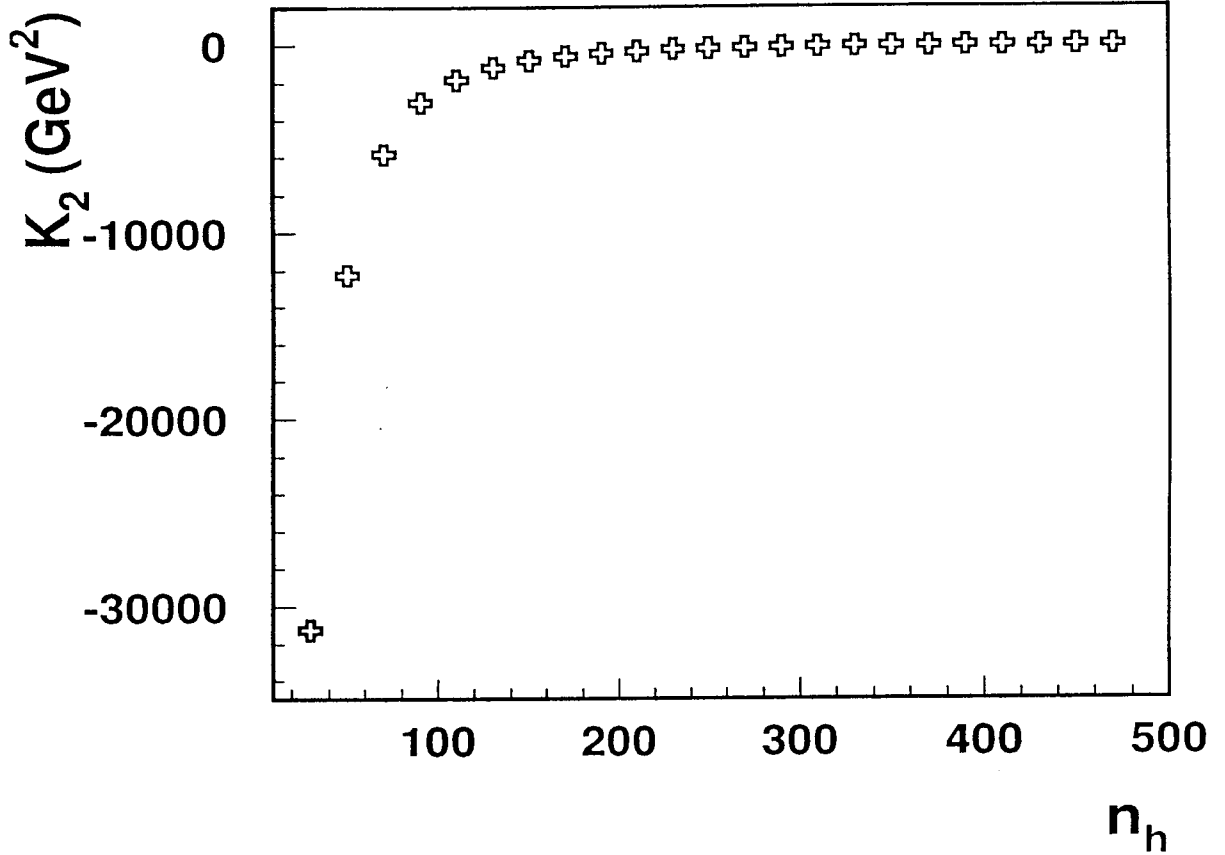


Figure 7: Dependent of  $K_2(\varepsilon, n_h)$  from number of hadrons without cut on the  $p_t$

$$K_2(\varepsilon_1, \varepsilon_2; n_h) = \langle [(\varepsilon_1; n_h) - \langle \varepsilon; n_h \rangle][(\varepsilon_2; n_h) - \langle \varepsilon; n_h \rangle] \rangle$$

$$= \langle \varepsilon^2; n_h \rangle - \langle \varepsilon; n_h \rangle^2$$

$$\langle \varepsilon^2; n_h \rangle = \frac{1}{N_n} \int \varepsilon_1 d\varepsilon_1 \varepsilon_2 d\varepsilon_2 \frac{d^2 N_n}{d\varepsilon_1 d\varepsilon_2}$$

—  $d^2 N_n / d\varepsilon_1 d\varepsilon_2$  - number of events with multiplicity  $n_h$  and particles with energy  $\varepsilon_1$  and  $\varepsilon_2$

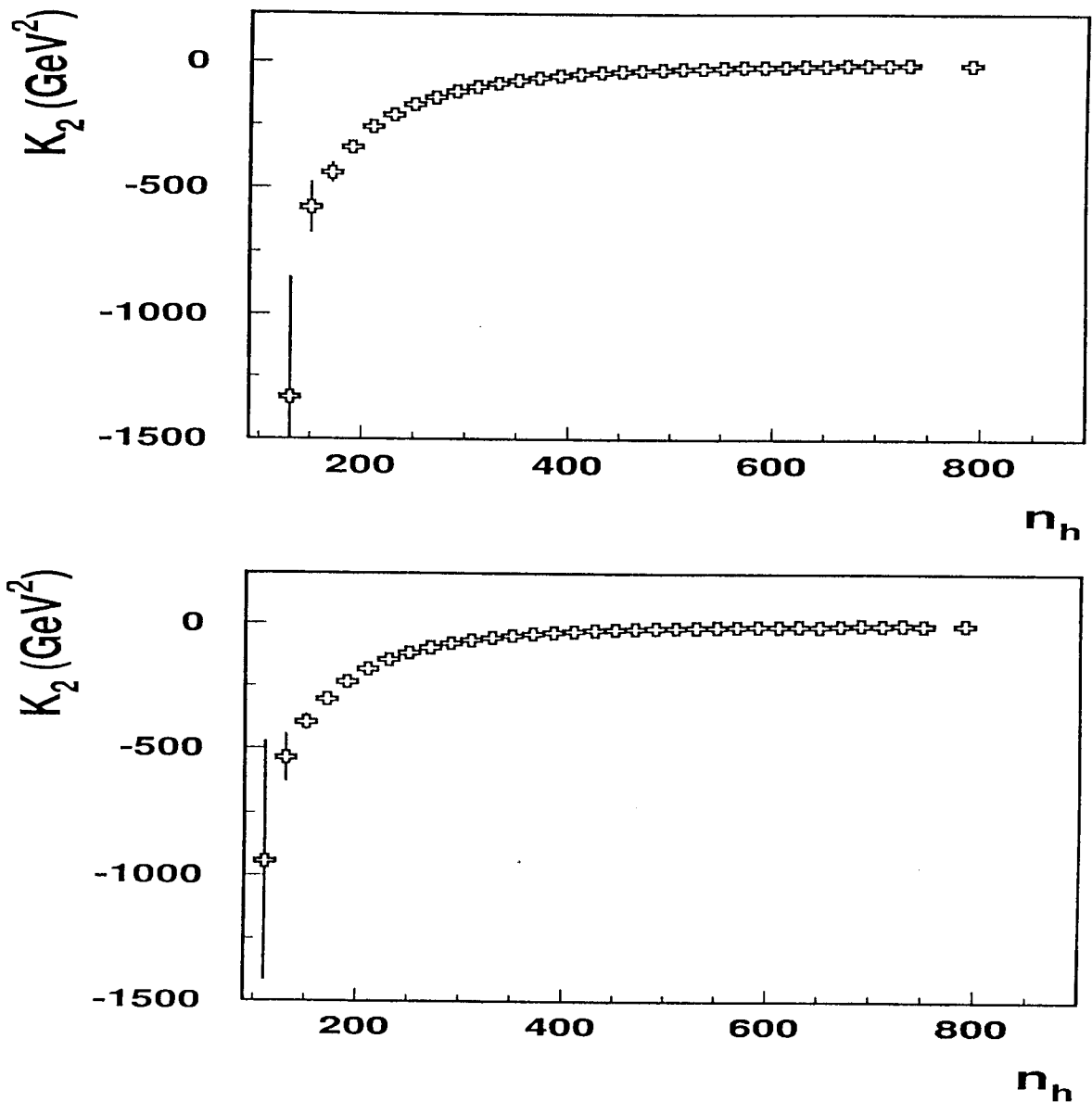


Figure 8: Dependent of  $K_2(\varepsilon, n_h)$  from number of hadrons. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

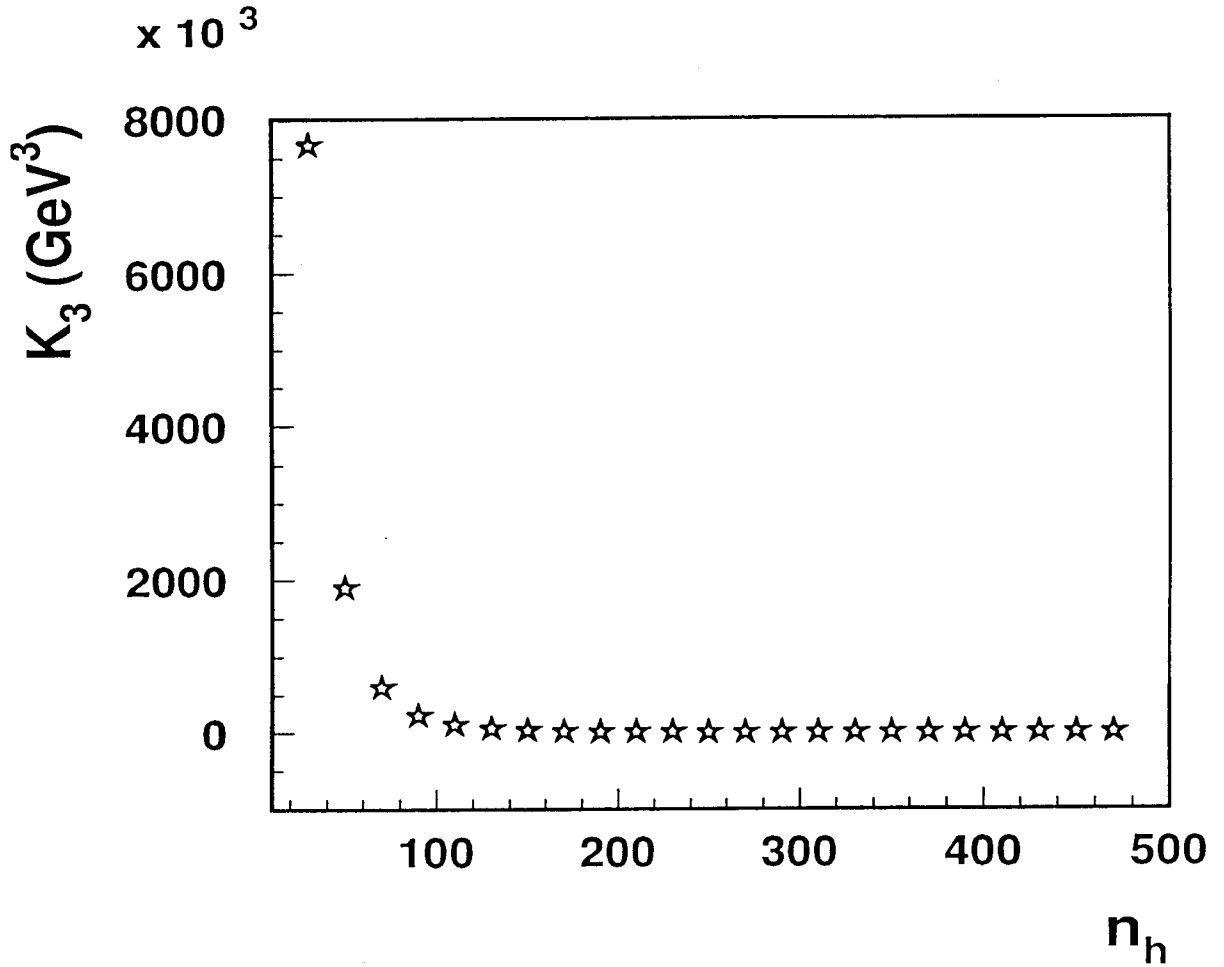


Figure 9: Dependent of  $K_3(\varepsilon, n_h)$  from number of hadrons without cut on the  $p_t$

$$\begin{aligned}
 K_3(\varepsilon_1, \varepsilon_2, \varepsilon_3; n_h) &= \\
 &= \langle [(\varepsilon_1; n_h) - \langle \varepsilon; n_h \rangle][(\varepsilon_2; n_h) - \langle \varepsilon; n_h \rangle][(\varepsilon_3; n_h) - \langle \varepsilon; n_h \rangle] \rangle = \\
 &= \langle \varepsilon^3; n_h \rangle - 2\langle \varepsilon^2; n_h \rangle \langle \varepsilon; n_h \rangle + \langle \varepsilon; n_h \rangle^3
 \end{aligned}$$

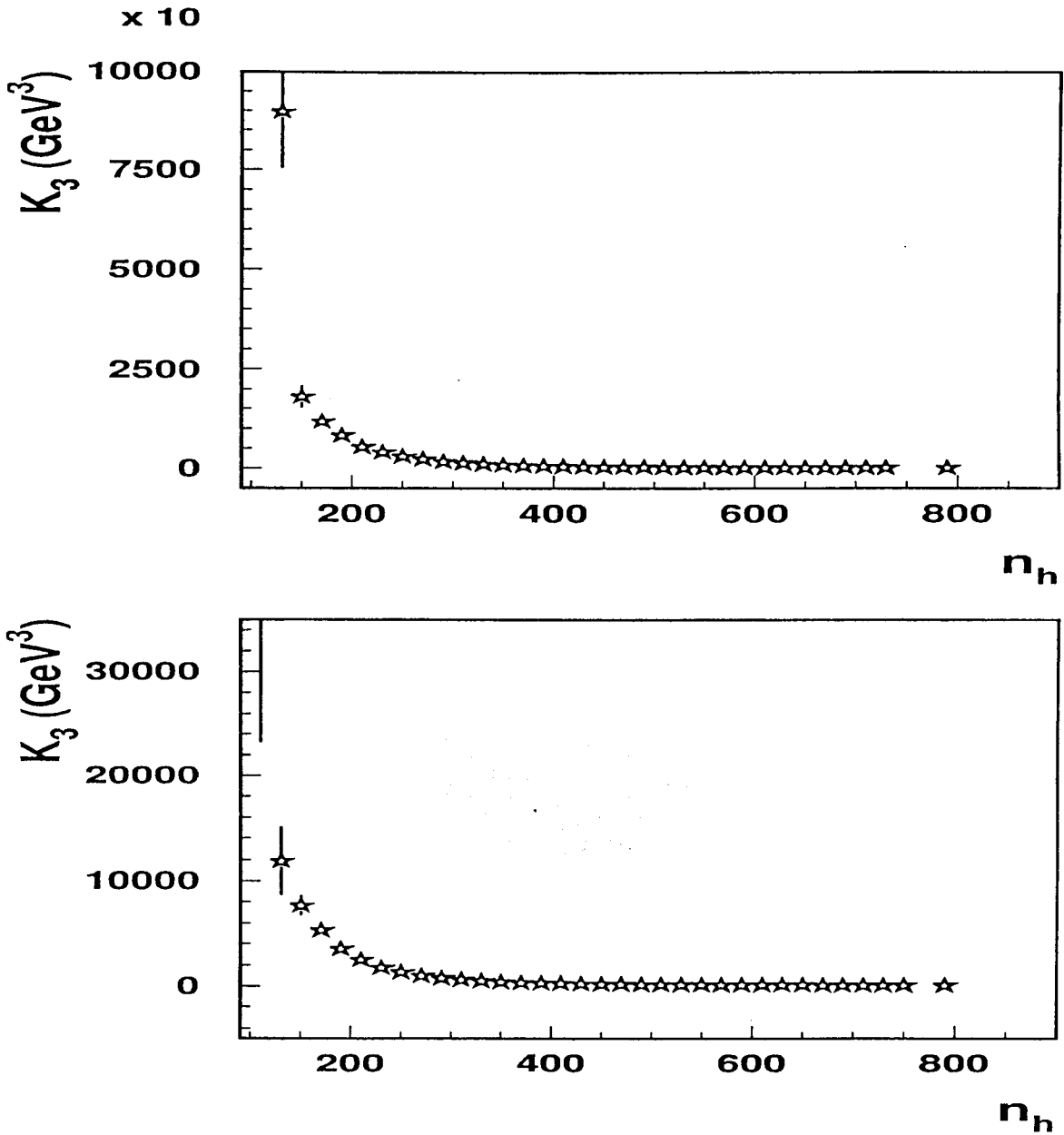


Figure 10: Dependence of  $K_3(\varepsilon, n_h)$  on the number of hadrons. The colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

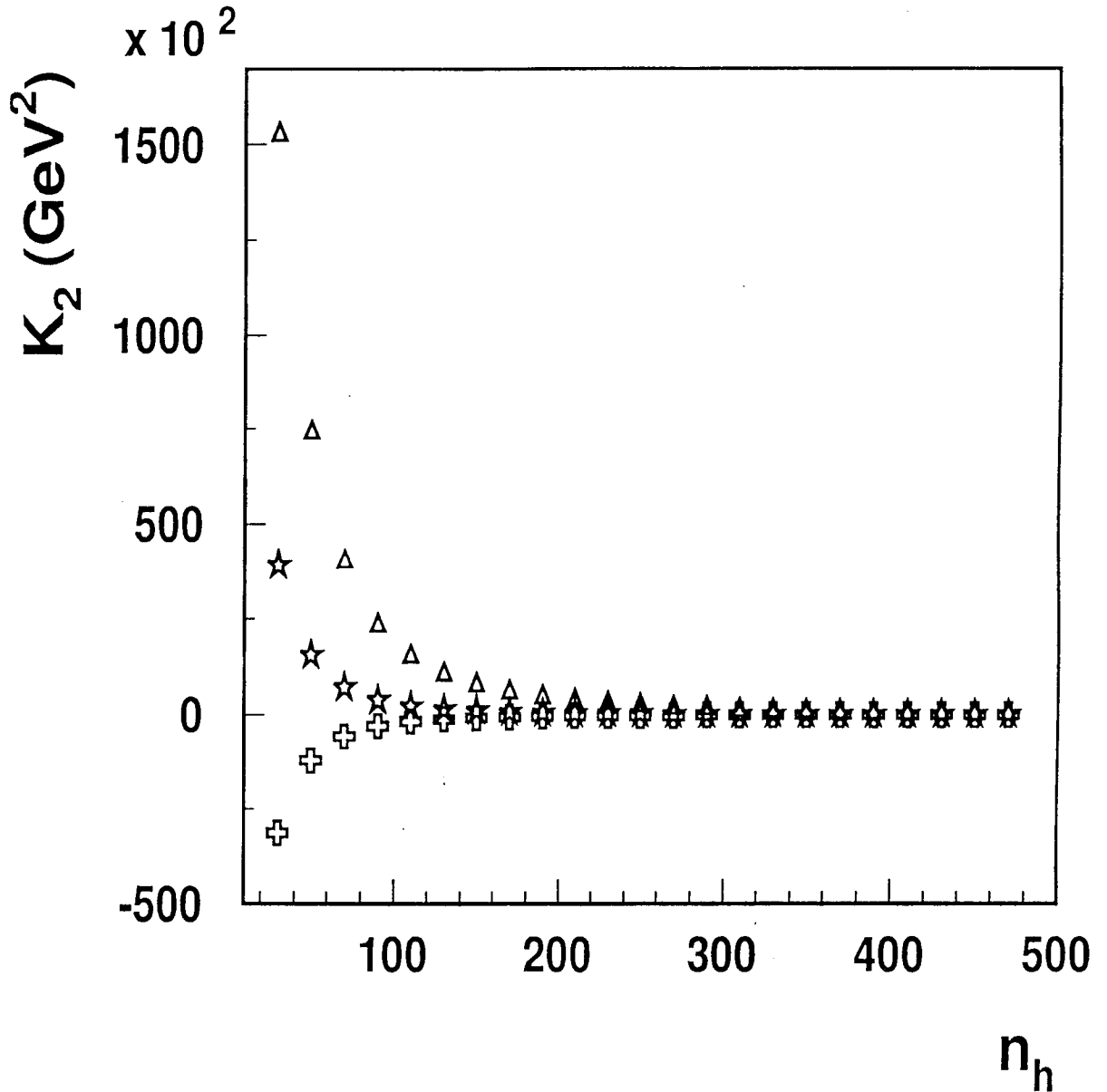


Figure 11: Dependence of  $K_1^2$  (triangles),  $K_2$  (crosses) and  $K_3^{2/3}$  (stars) from number of hadrons without cut on the  $p_t$

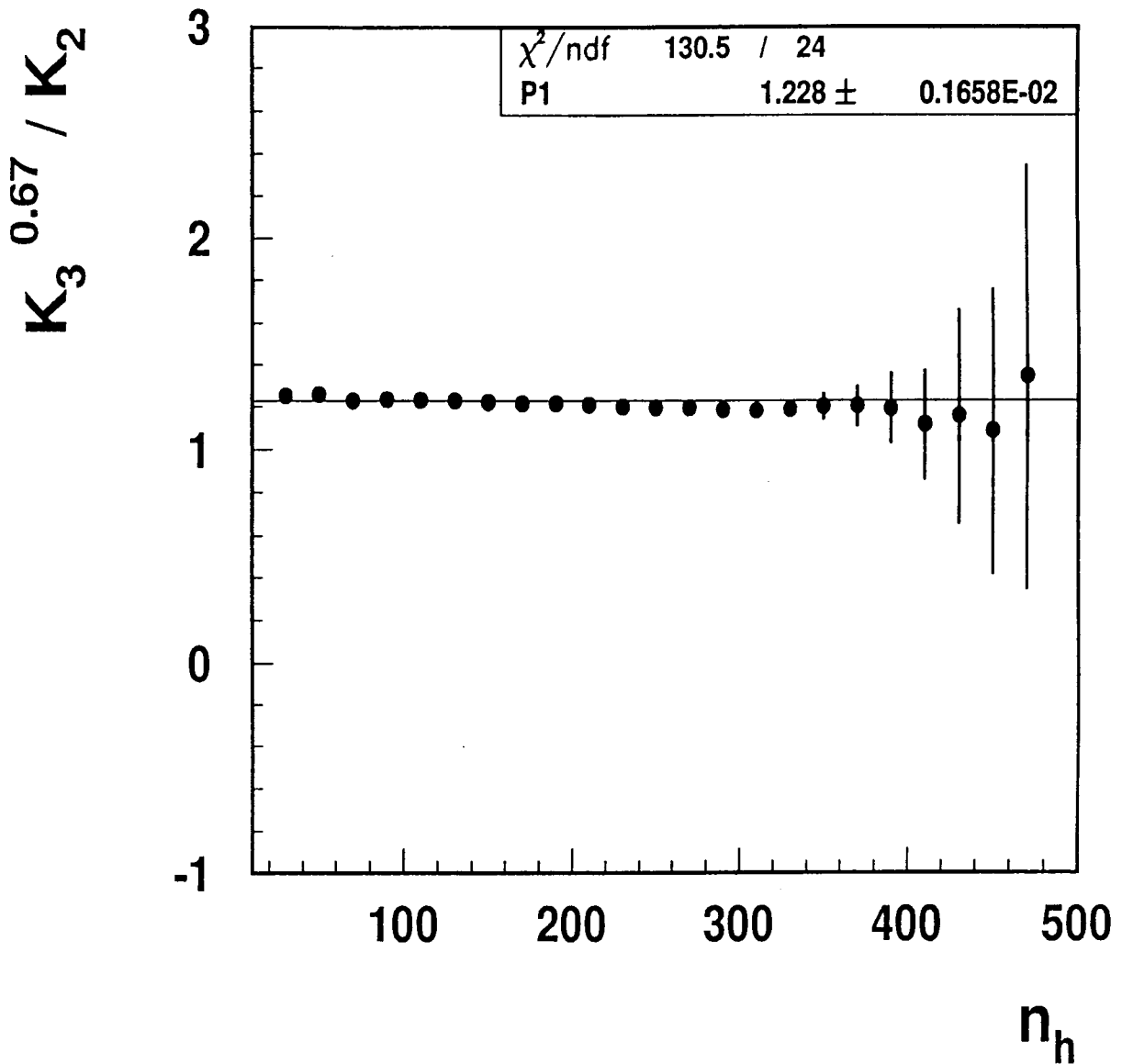


Figure 12: Dependent of Ratio  $K_3^{2/3}/K_2$  from number of hadrons without cut on the  $p_t$



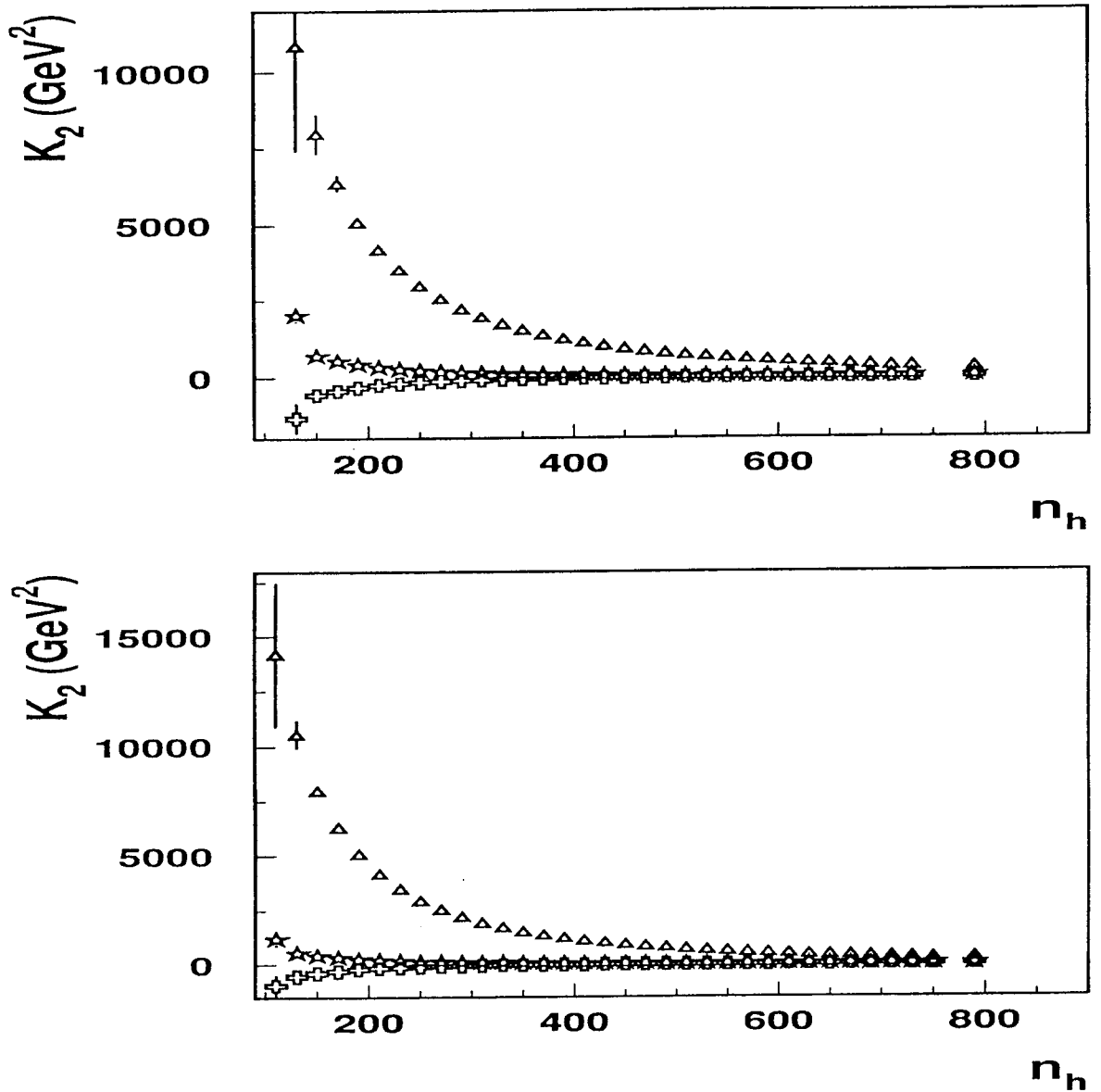


Figure 13: Dependence of  $K_1^2$  (triangles),  $K_2$  (crosses) and  $K_3^{2/3}$  (stars) from number of hadrons. Colored parton transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

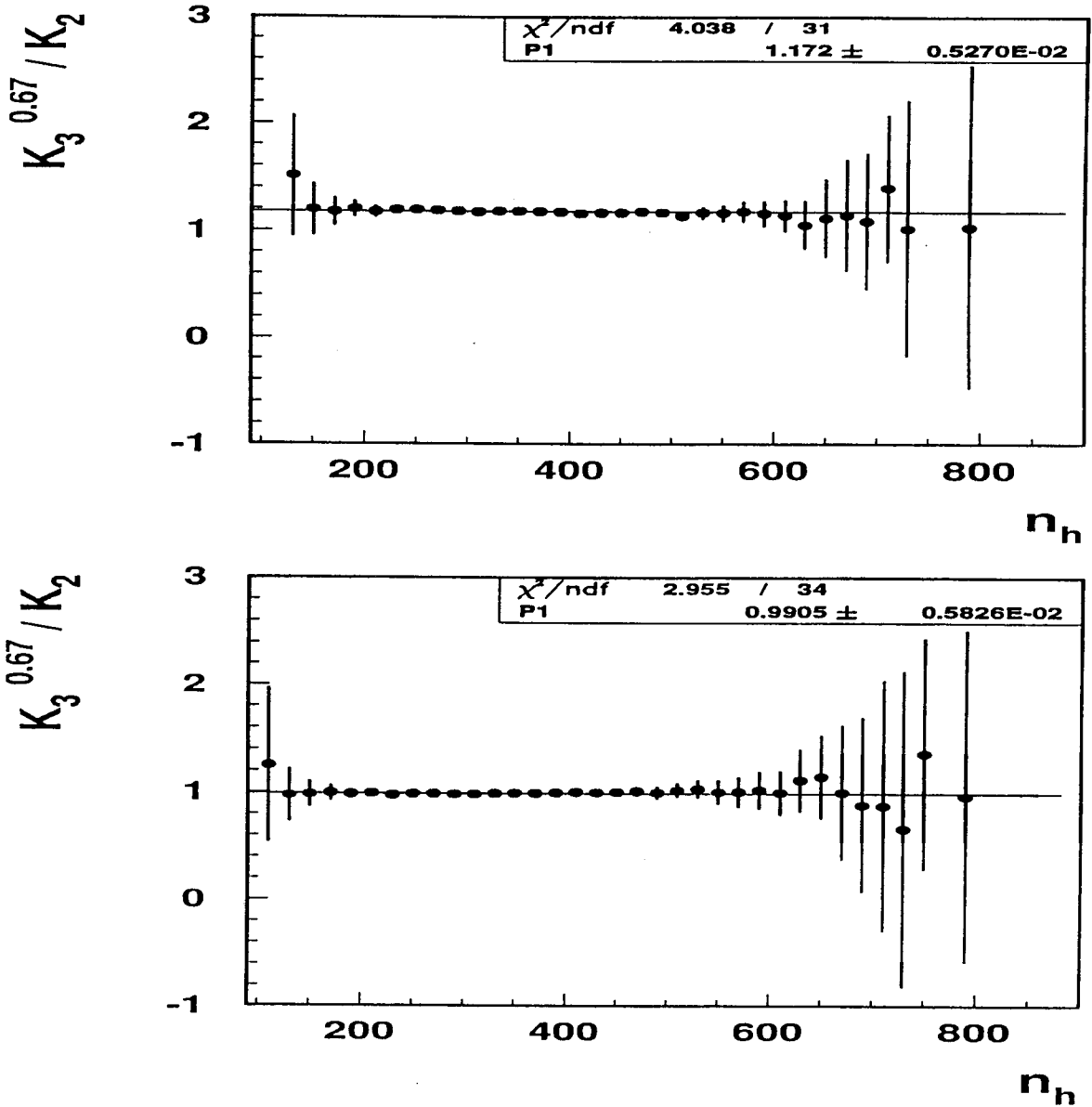


Figure 14: Dependence of ratio  $K_3^{2/3} / K_2$  from number of hadrons. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

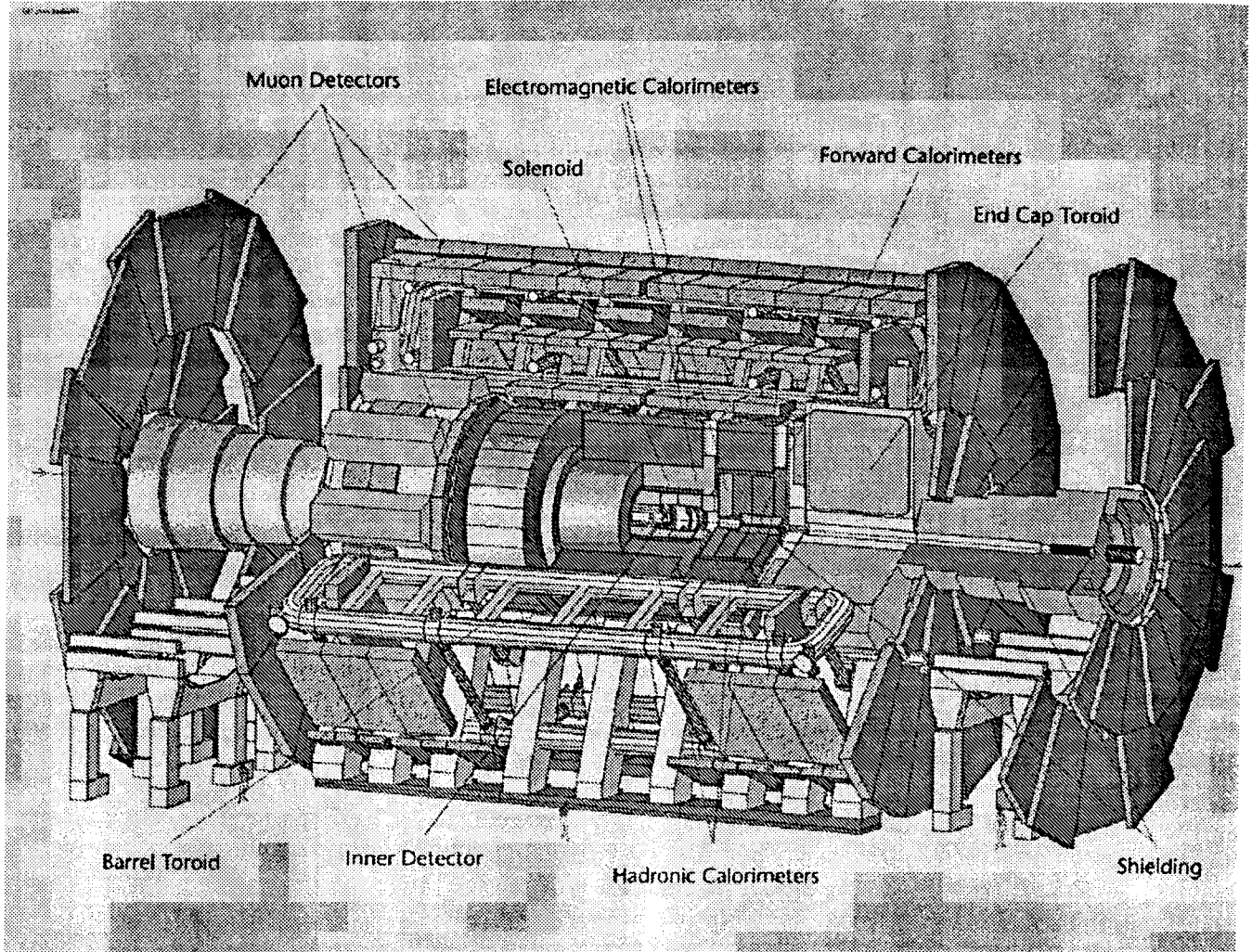


Figure 15: Overall layout of the ATLAS Detector

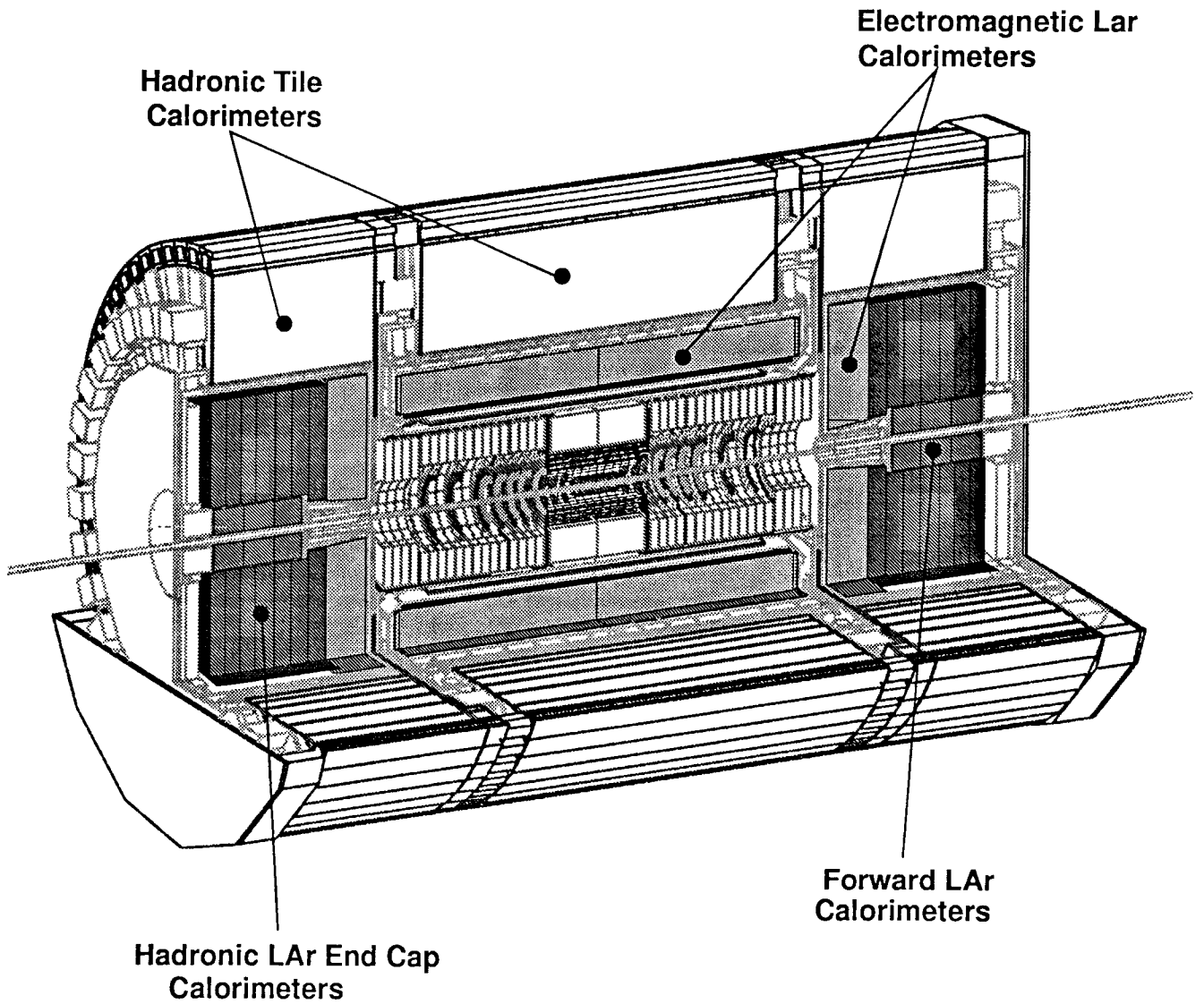


Figure 16: Three-dimensional view of the ATLAS Calorimetry

2.2 ATLFAST prediction for "tower" energy correlators

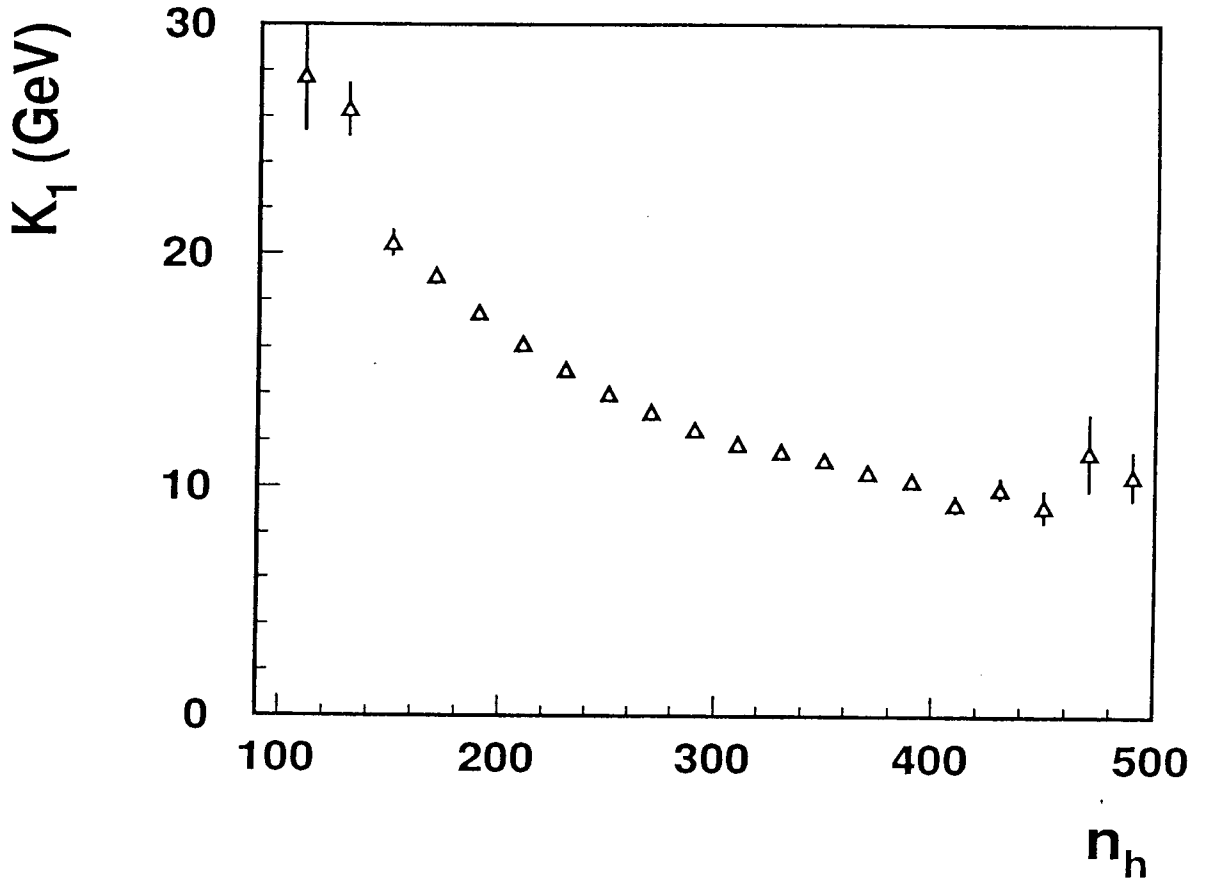


Figure 17: Dependence of  $K_1(\varepsilon, n_h)$  from number of hadrons with cut on the  $p_t = 500$  GeV

- $\varepsilon$  - energy deposited in the "tower"
- $dN_n/d\varepsilon$  - number of events with energy  $\varepsilon$  in the "tower"

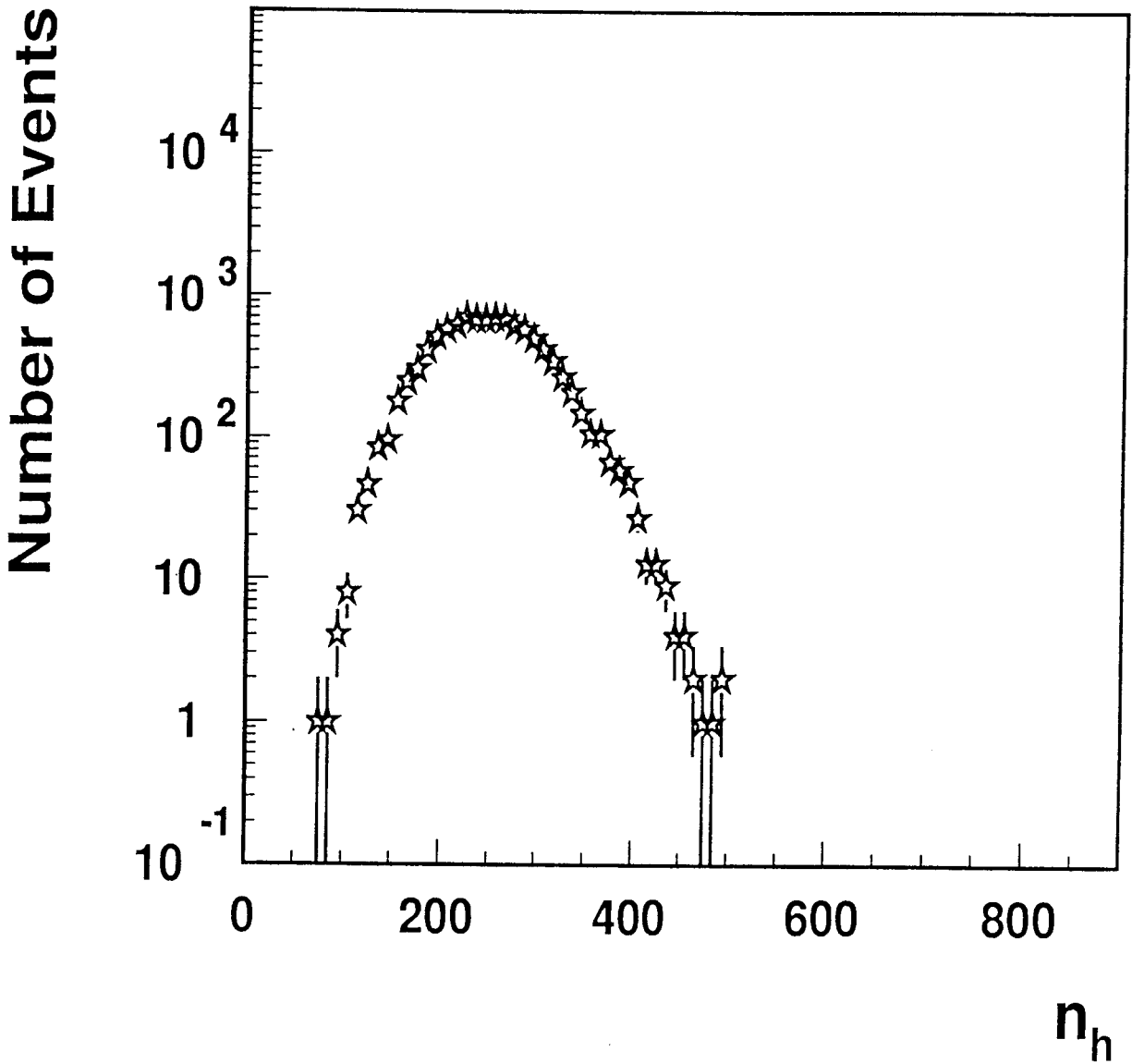


Figure 18: Multiplicity distribution. With colored partons transverse momentum cutoff  $p_t = 500$  GeV

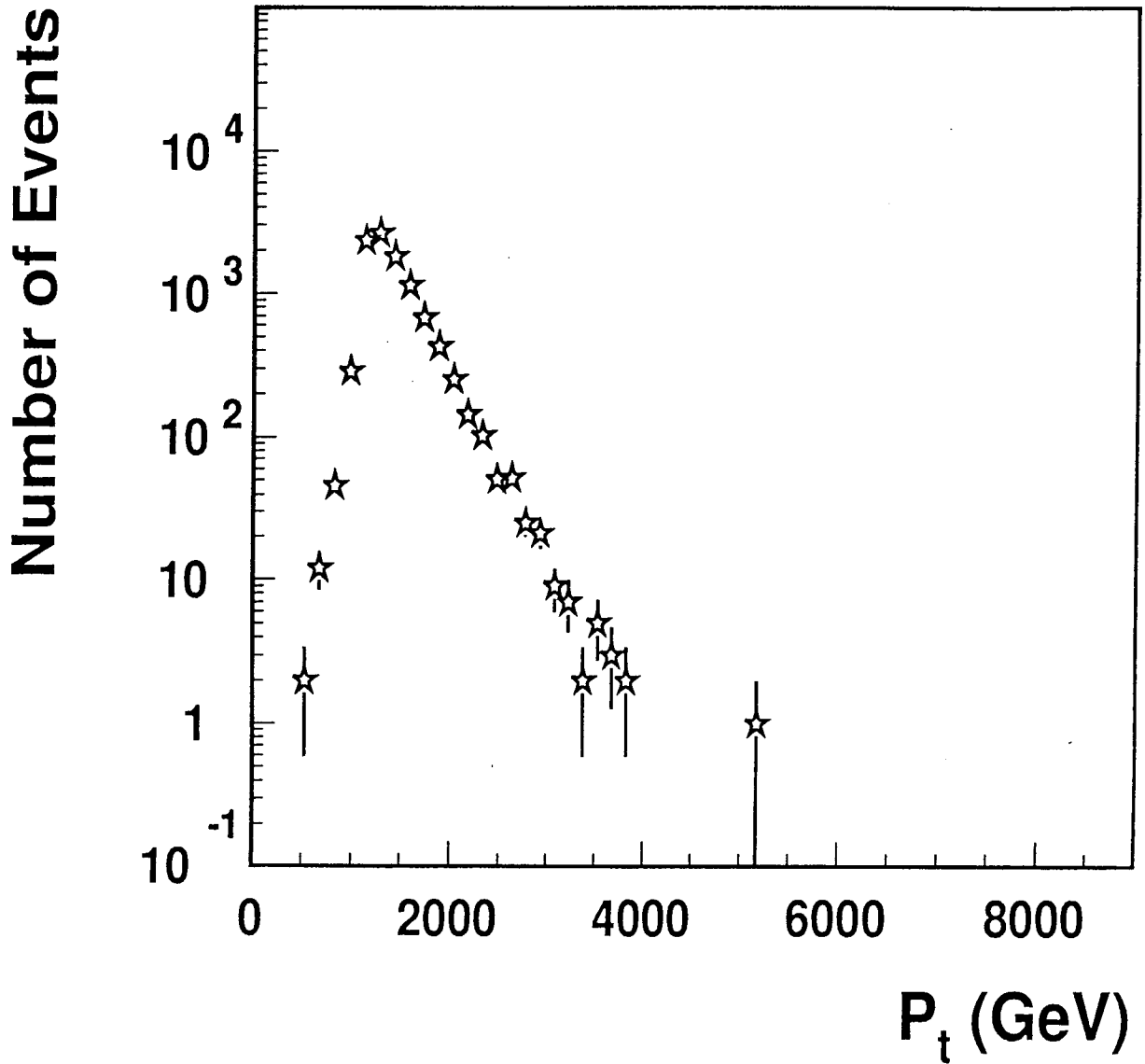


Figure 19:  $p_t$  distribution. With colored partons transverse momentum cutoff  $p_t = 500$  GeV

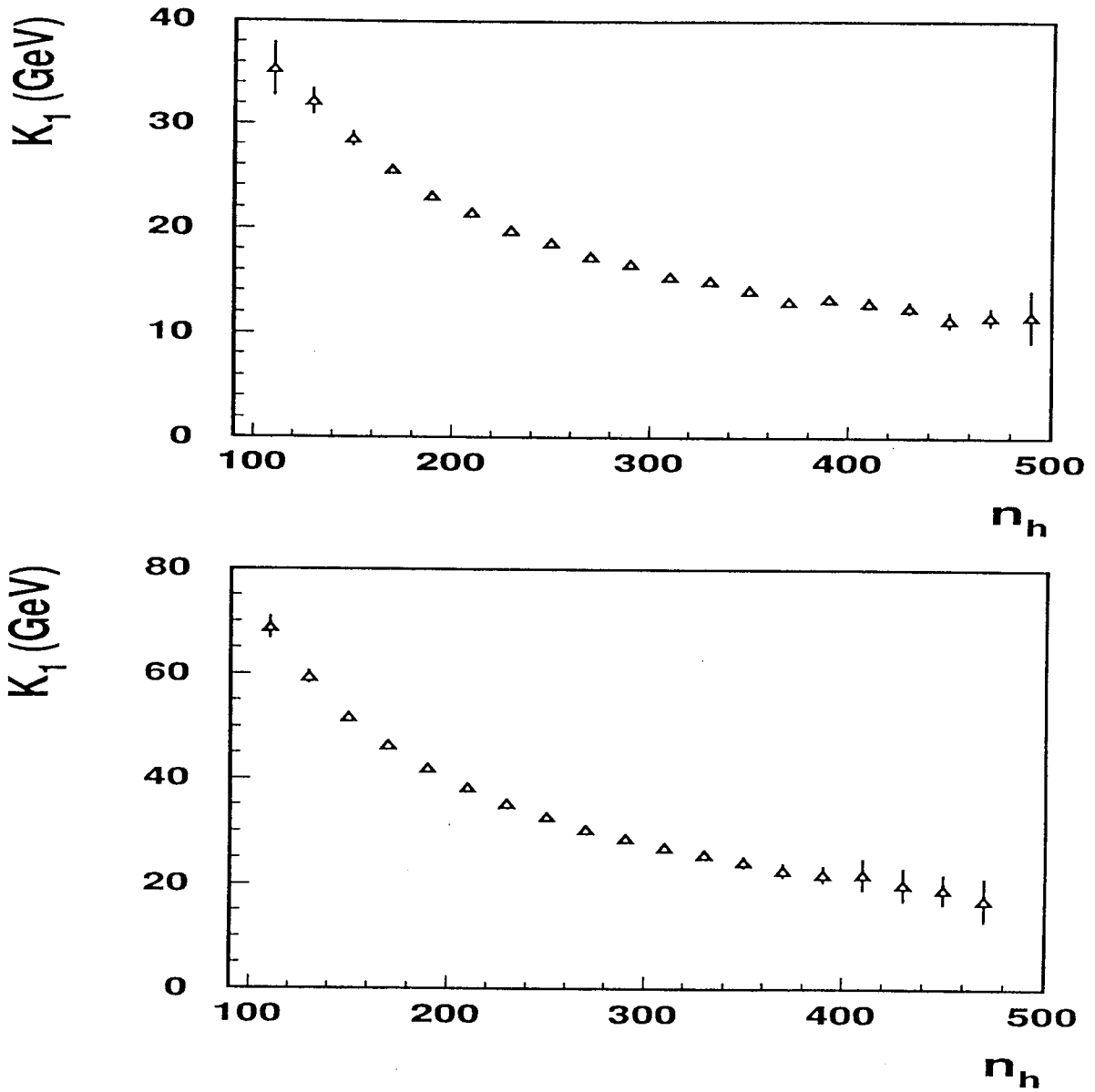


Figure 20: Dependence of  $K_1(\varepsilon, n_h)$  from number of calorimeter towers. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)



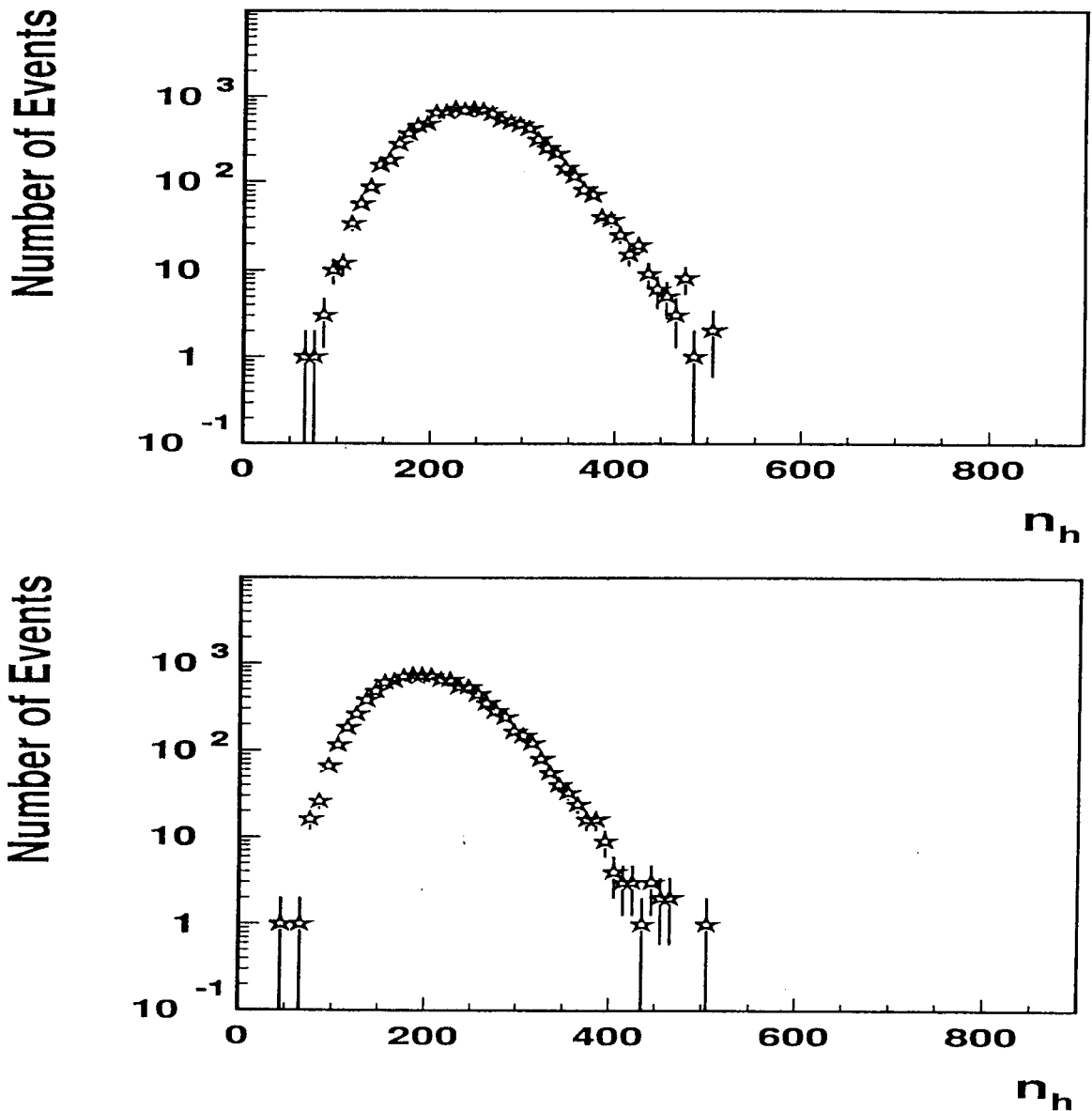


Figure 21: Multiplicity distribution. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

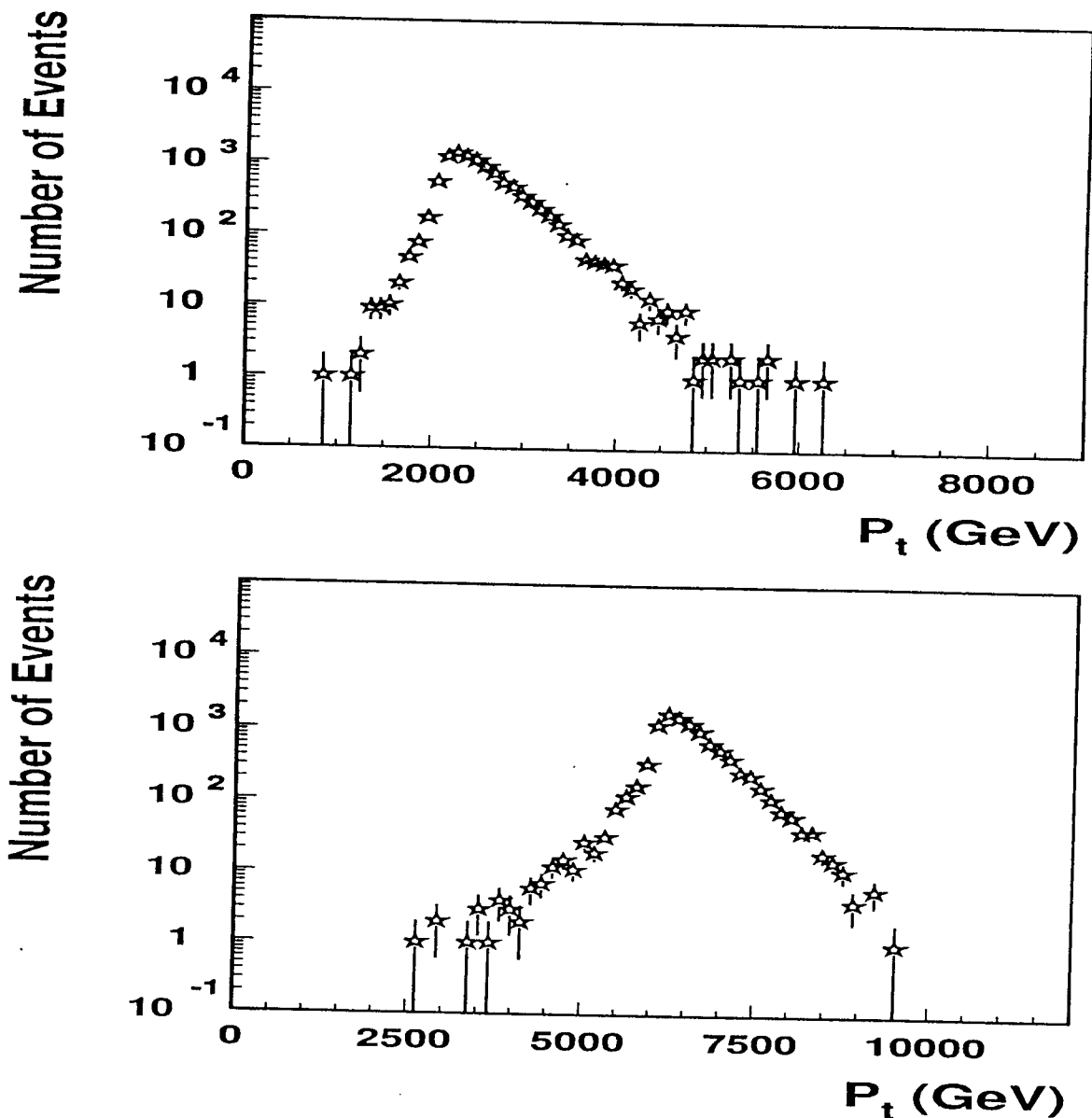


Figure 22:  $p_t$  distribution. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

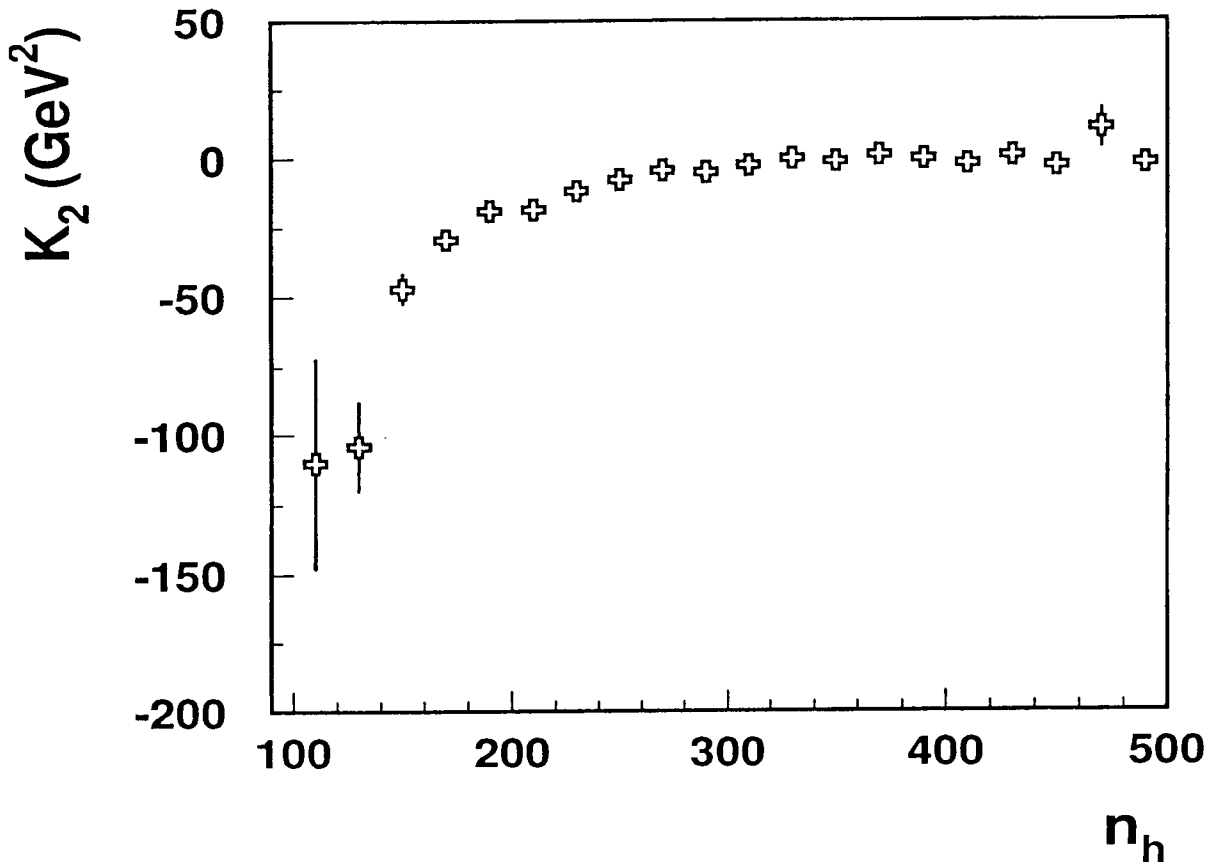


Figure 23: Dependent of  $K_2(\varepsilon, n_h)$  from number of calorimeter towers with cut on the  $p_t = 500$  GeV

$$\begin{aligned}
 K_2(\varepsilon_1, \varepsilon_2; n_h) &= \langle [(\varepsilon_1; n_h) - \langle \varepsilon; n_h \rangle][(\varepsilon_2; n_h) - \langle \varepsilon; n_h \rangle] \rangle \\
 &= \langle \varepsilon^2; n_h \rangle - \langle \varepsilon; n_h \rangle^2
 \end{aligned}$$

$$\langle \varepsilon^2; n_h \rangle = \frac{1}{N_n} \int \varepsilon_1 d\varepsilon_1 \varepsilon_2 d\varepsilon_2 \frac{d^2 N_n}{d\varepsilon_1 d\varepsilon_2}$$

—  $d^2 N_n / d\varepsilon_1 d\varepsilon_2$  - number of events with multiplicity  $n_h$  and particles with energy  $\varepsilon_1$  and  $\varepsilon_2$

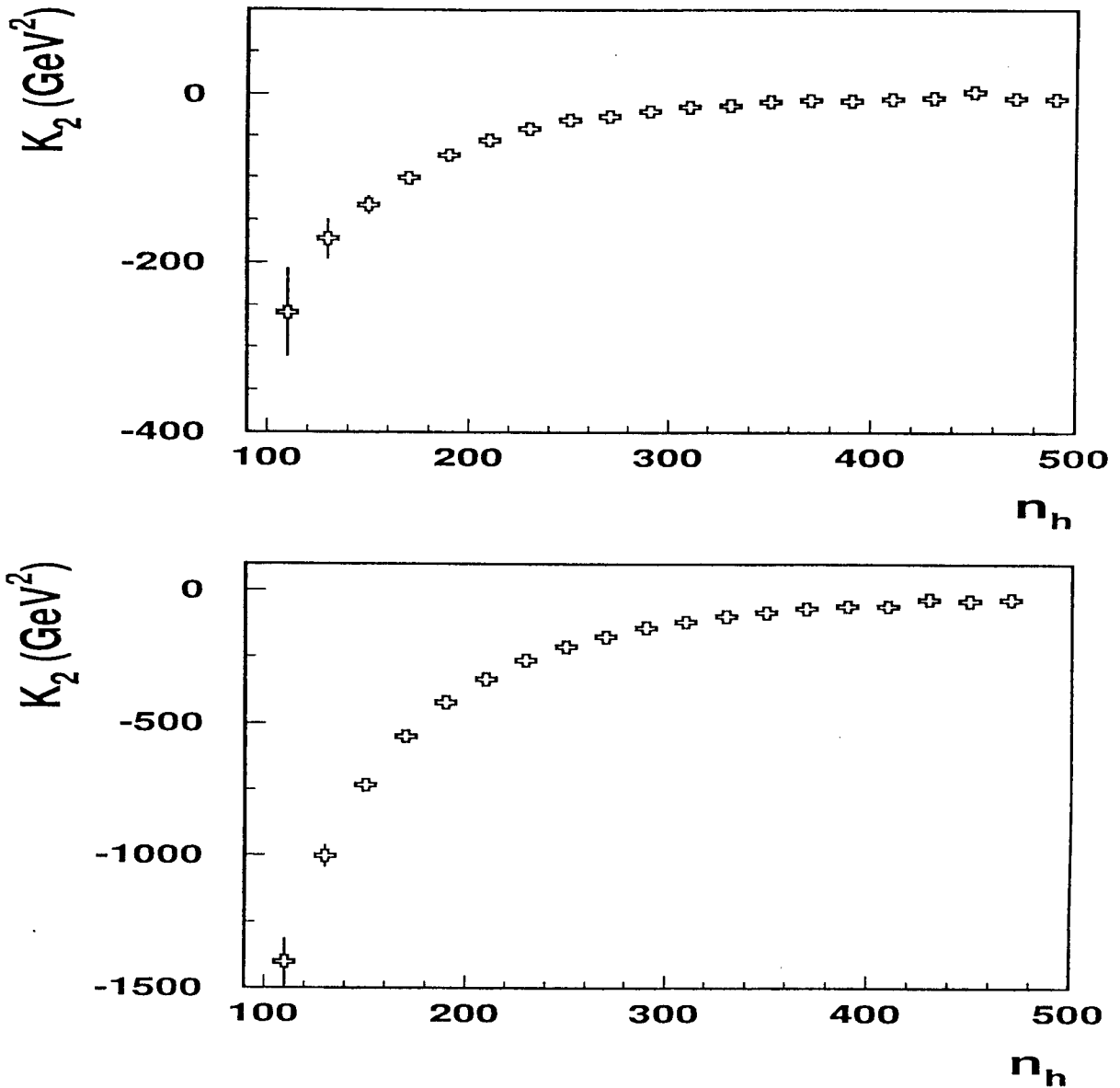


Figure 24: Dependent of  $K_2(\varepsilon, n_h)$  from number of calorimeter towers. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

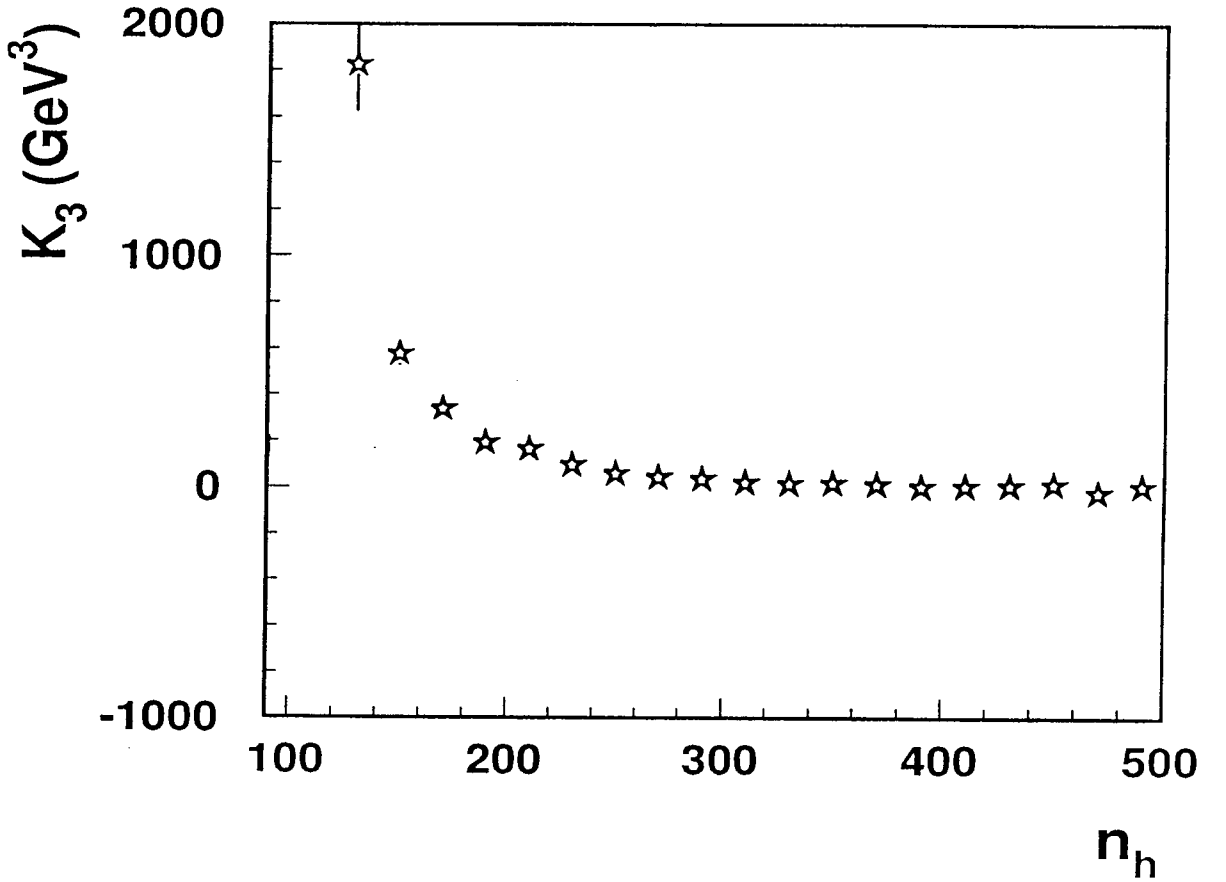


Figure 25: Dependent of  $K_3(\varepsilon, n_h)$  from number of calorimeter towers with cut on the  $p_t = 500$  GeV

$$\begin{aligned}
 K_3(\varepsilon_1, \varepsilon_2, \varepsilon_3; n_h) &= \\
 &= \langle [(\varepsilon_1; n_h) - \langle \varepsilon; n_h \rangle][(\varepsilon_2; n_h) - \langle \varepsilon; n_h \rangle][(\varepsilon_3; n_h) - \langle \varepsilon; n_h \rangle] \rangle = \\
 &= \langle \varepsilon^3; n_h \rangle - 2\langle \varepsilon^2; n_h \rangle \langle \varepsilon; n_h \rangle + \langle \varepsilon; n_h \rangle^3
 \end{aligned}$$

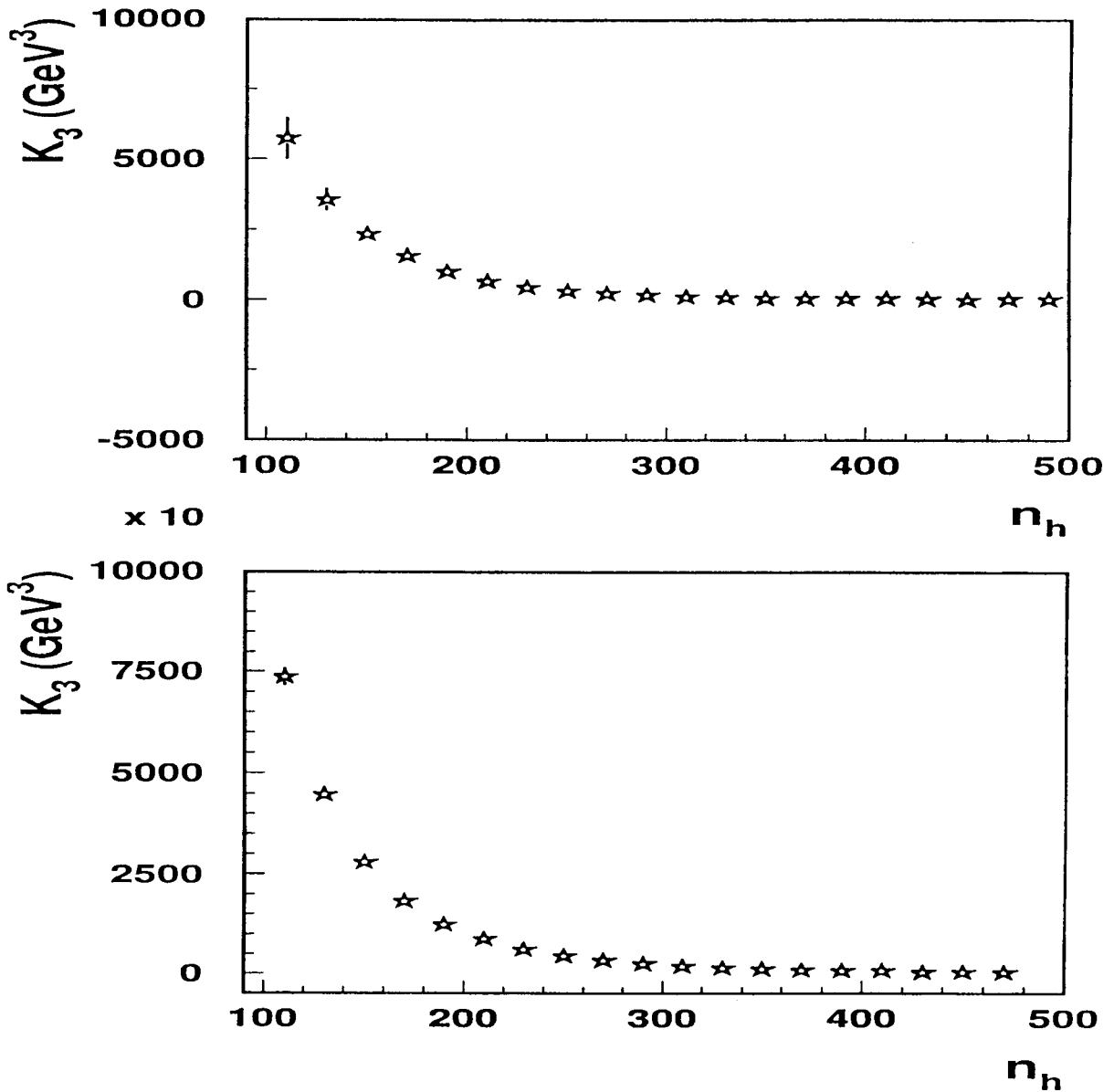


Figure 26: Dependence of  $K_3(\varepsilon, n_h)$  on the number of calorimeter towers. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

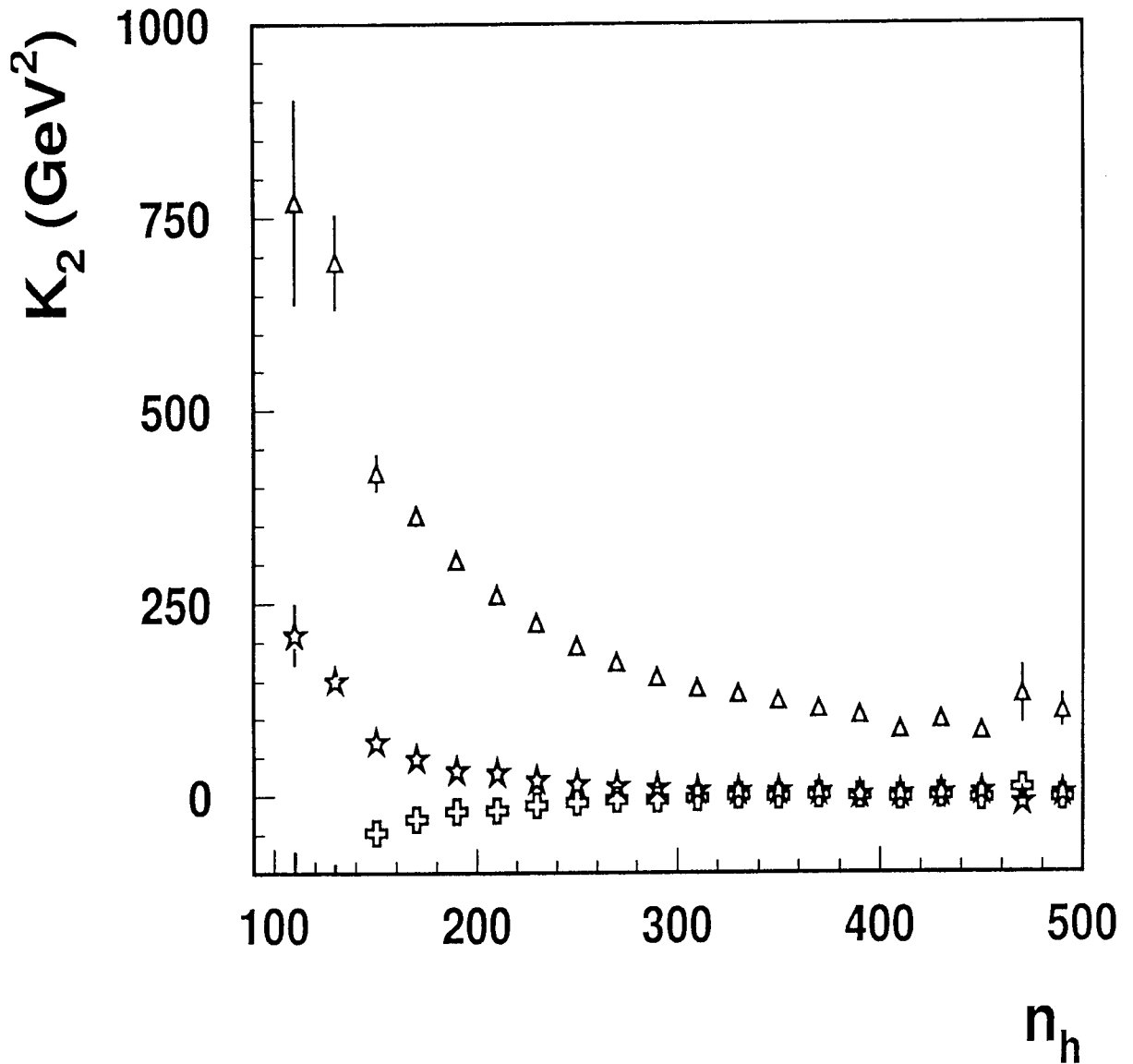


Figure 27: Dependence of  $K_1^2$  (triangles),  $K_2$  (crosses) and  $K_3^{2/3}$  (stars) from number of calorimeter towers with cut on the  $p_t = 500 \text{ GeV}$

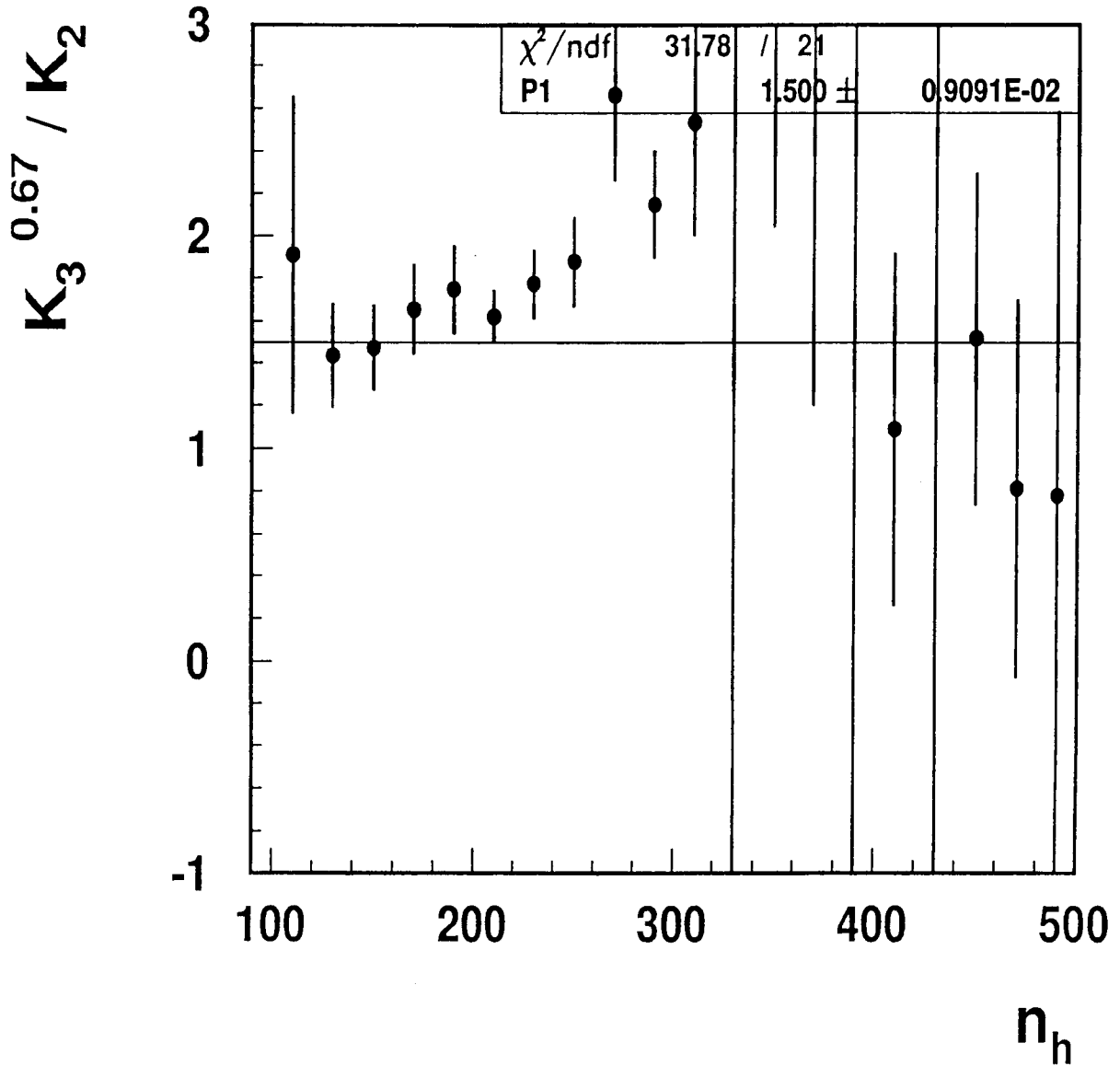


Figure 28: Dependent of Ratio  $K_3^{2/3} / K_2$  from number of calorimeter towers with cut on the  $p_t = 500$  GeV



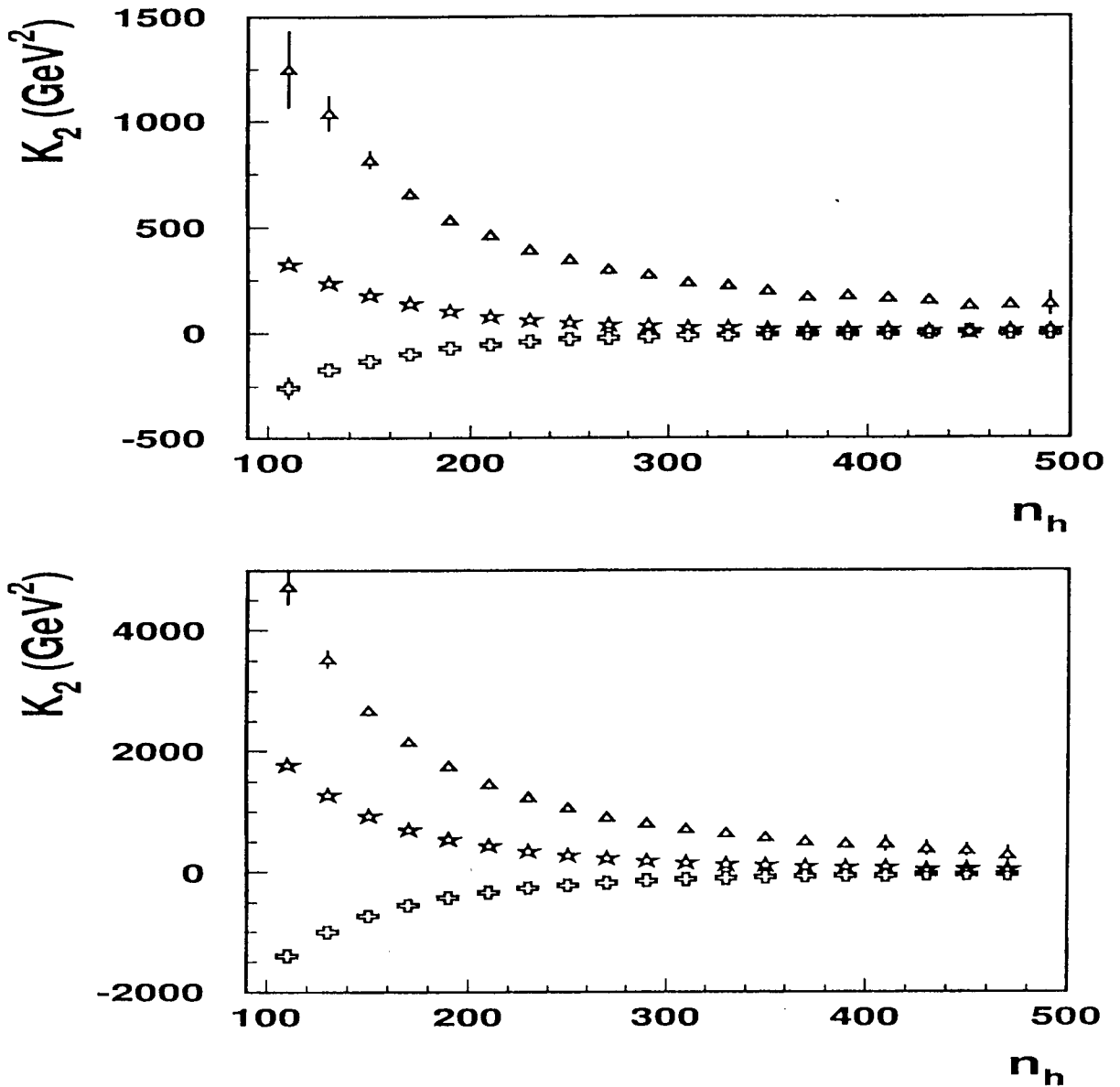


Figure 29: Dependence of  $K_1^2$  (triangles),  $K_2$  (crosses) and  $K_3^{2/3}$  (stars) from number of towers. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

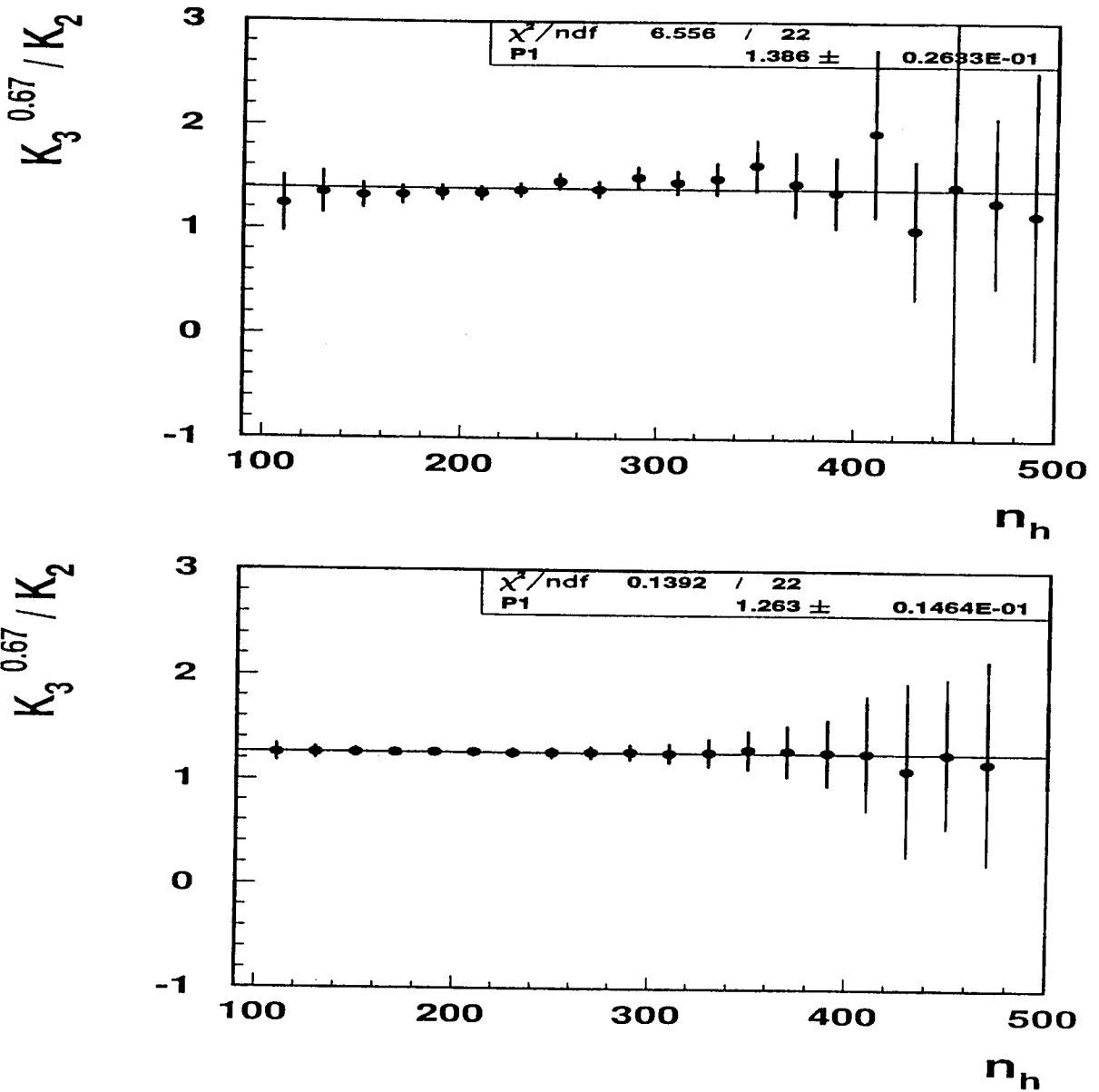


Figure 30: Dependent of ratio  $K_3^{2/3} / K_2$  from number of towers. Colored partons transverse momentum cutoff is  $p_t = 1000$  GeV (top) and  $p_t = 3000$  GeV (bottom)

### 3 Conclusion

- We present results of calculation in PYTHIA correlators  $K_1$ ,  $K_2$ ,  $K_3$  and ratio  $R = \frac{K_3^{2/3}}{K_2}$  versus number of hadrons.
- $R = \frac{K_3^{2/3}}{K_2} \simeq 1$  for all  $p_t$  cuts and does not agree with Sissakian-Manjavidze prediction for VHM  $R = \frac{K_3^{2/3}}{K_2} \ll 1$ . The PYTHIA can not predict the **tendency** to equilibrium.
- Using ATLFAST we calculated correlators  $K_1$ ,  $K_2$ ,  $K_3$  and ratio  $R = \frac{K_3^{2/3}}{K_2}$  versus number of ATLAS calorimeter towers with hit.
- The correlators  $K_1$ ,  $K_2$ ,  $K_3$  have the same dependencies for hadrons and for ATLAS Calorimeter towers.  
 $R_{towers} \approx 1.1 \cdot R_{hadrons}$ .