

# Two-loop corrections to the pole masses of heavy quarks in the MSSM

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$$\frac{\Delta m_{t,b}}{m_{t,b}} \equiv \frac{M_{t,b}^{pole} - m_{t,b}^{\overline{\text{DR}}}(\bar{\mu})}{m_{t,b}^{\overline{\text{DR}}}(\bar{\mu})}$$

- Motivation
- Calculation method
- Results
- Conclusion

- Estimates of SUSY particles masses are necessary.
- One needs to know values of (running) parameters of MSSM Lagrangian at the “low-energy” ( $EW \approx 100 \text{ GeV}$ , or SUSY breaking  $\approx 1000 \text{ GeV}$ ) scale.
- For large  $\tan \beta$  and/or  $m_0$  predicted values of masses are very sensitive to Yukawa couplings of heavy quarks. This dependence is particularly strong for Higgs boson masses.
- Thus, it is desirable to have determination of the heavy quarks Yukawa couplings as precise as possible.

In order to evaluate Yukawa couplings from

- $M_Z^{pole}, M_W^{pole}$
- $M_t^{pole}$
- $\alpha_s^{(5)}(M_Z), \alpha_{em}$
- $m_b(m_b)$ , etc.

we need to calculate relation between the pole and running masses:

$$\frac{\Delta m_{t,b}}{m_{t,b}} \equiv \frac{M_{t,b}^{pole} - m_{t,b}^{\overline{\text{DR}}}(\bar{\mu})}{m_{t,b}^{\overline{\text{DR}}}(\bar{\mu})}$$

- Pierce, Matchev, Zhang, hep-ph/9606211:

$$\frac{\Delta m_{t,b}}{m_{t,b}} = \mathcal{O}(\alpha_s) + \mathcal{O}(y_t^2) + \mathcal{O}(y_b^2) + \mathcal{O}(g^2) + \mathcal{O}(g'^2)$$

For the  $b$  quark  $\mathcal{O}(\alpha_s)$  terms are reduced significantly by  $\mathcal{O}(y_t^2 + y_b^2)$  ones.

- Bednyakov, Onishchenko, Velizhanin, Veretin, hep-ph/0210258:  
 $\mathcal{O}(\alpha_s^2)$ , obtained by asymptotic expansion in  $m_{t,b}/M_{SUSY}$ , leading term only. For the  $b$  quark this correction is comparable with one-loop one.

- $\mathcal{O}(\alpha_s y_t^2 + \alpha_s y_b^2 + y_t^4 + y_b^4)$  corrections for  $b$ -quark, since these may be of the same order of magnitude as  $\mathcal{O}(\alpha_s^2)$  ones.
- second-order terms of asymptotic expansion for  $t$ -quark, since  $m_t/M_{SUSY} \approx 0.2$ .

full connected propagator of a quark:

$$i(\hat{p} - m - \Sigma(\hat{p}, m_i))^{-1}$$

self-energy of the quark:

$$\Sigma(\hat{p}, m_i) = \hat{p}\Sigma_V(p^2, m_i^2) + \hat{p}\gamma_5\Sigma_A(p^2, m_i^2) + m\Sigma_S(p^2, m_i)$$

pole mass  $M_p$ :

$$\left( (1 + \Sigma_V(M_p^2, m_i^2))^2 - \Sigma_A^2(M_p^2, m_i^2) \right) M_p^2 - m^2 (1 - \Sigma_S(M_p^2, m_i^2))^2 = 0$$

perturbative expansion of the pole mass:

$$\begin{aligned} \frac{M_p - m}{m} &= \alpha M^{(1)} + \alpha^2 M^{(2)}, \\ M^{(1)} &= \Sigma_V^{(1)}(m^2, m_i^2) + \Sigma_S^{(1)}(m^2, m_i), \\ M^{(2)} &= \Sigma_V^{(2)}(m^2, m_i^2) + \Sigma_S^{(2)}(m^2, m_i) + \frac{1}{2} \Sigma_A^{(1)2}(m^2, m_i^2) \\ &+ M^{(1)} \left( \Sigma_V^{(1)}(m^2, m_i^2) \right. \\ &\quad \left. + 2m^2 \frac{\partial}{\partial p^2} \left( \Sigma_V^{(1)}(p^2, m_i^2) + \Sigma_S^{(1)}(p^2, m_i) \right) \right)_{p^2=m^2} \end{aligned}$$



Asymptotic expansion (F.V. Tkachov):

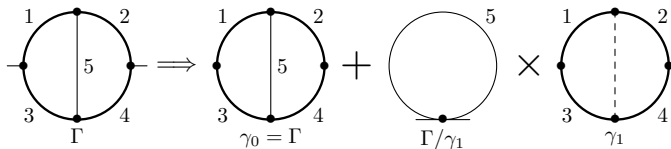
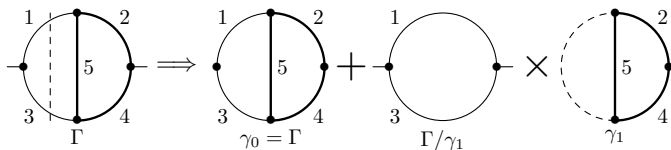
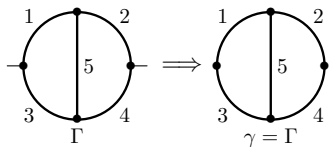
$$F_{\Gamma}(p, m, M) = \sum_{\gamma} F_{\Gamma/\gamma}(p, m) \mathcal{T}_{\gamma}(p_{\gamma}, m) F_{\gamma}(p_{\gamma}, m, M)$$

- $\gamma$ : asymptotically irreducible subgraphs of  $\Gamma$
- $\mathcal{T}_{\gamma}(p_{\gamma}, m)$ : Taylor expansion in small masses and external (with respect to  $\gamma$ ) momenta

about 160 two-loop propagator-type diagrams

$$m_t \ll m_{\tilde{q}}, m_{\tilde{g}}$$

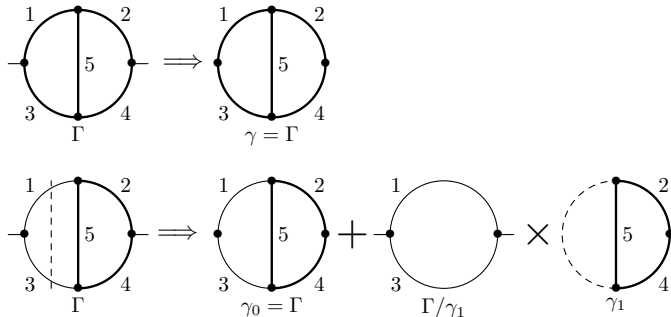
$$\frac{\Delta m_t}{m_t} \approx 1 + \alpha_s \sum_{n=-1}^2 m_t^n \sigma_1^{(n)} + \alpha_s^2 \sum_{n=-1}^2 m_t^n \sigma_2^{(n)}$$



about 1200 two-loop propagator-type diagrams

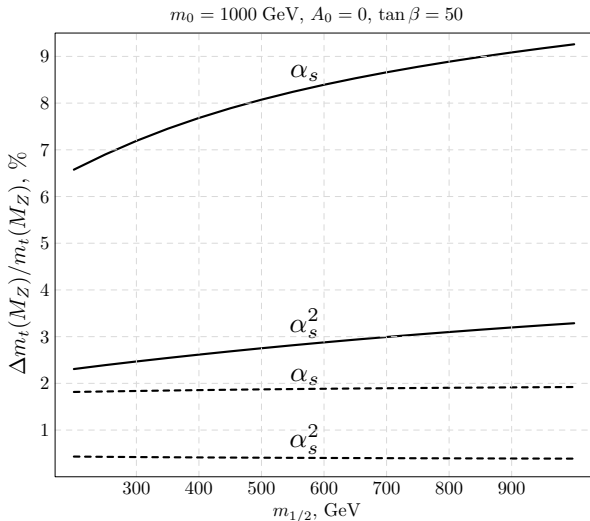
$$m_b \ll m_t, m_{h_0}, m_{H_0}, m_{H^+}, m_{\tilde{q}}, m_{\tilde{\chi}^0}, m_{\tilde{\chi}^+}, m_{G_0}, m_{G^+}$$

$$\frac{\Delta m_b}{m_b} \approx 1 + \alpha \sum_{n=-1}^0 m_b^n \sigma_1^{(n)} + \alpha^2 \sum_{n=-1}^0 m_b^n \sigma_2^{(n)}$$



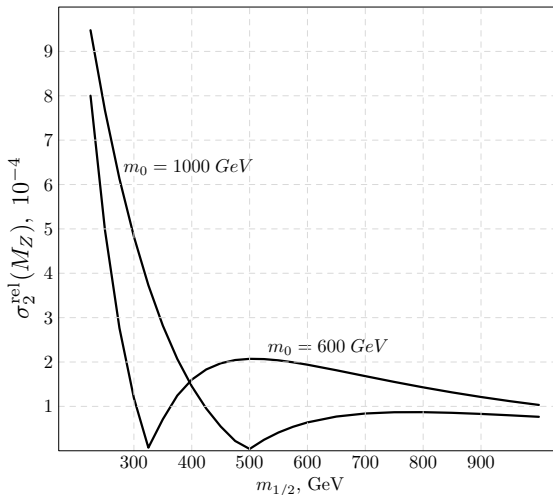
other subgraphs give  $\mathcal{O}(m_b^2)$  contribution.

- Diagram generation: (slightly modified version of) *FeynArts*.
- Analytical and numeric calculations: *GiNaC* C++ library.
- Asymptotic expansion: *prop2exp* C++ library (based on GiNaC).
- Recursive reduction of two-loop vacuum integrals to a by integration by parts method, analytic and numeric evaluation of the master-integral: *bubblesii* C++ library (based on GiNaC too).
- Evaluation of running MSSM parameters and SUSY partner masses: *ffmssmsc* C++ library (forked from ~~*SOFTSUSY*~~ written from scratch).



$\mathcal{O}(\alpha_s^2)$  corrections  
are  $\approx 30\%$  of  
one-loop ones

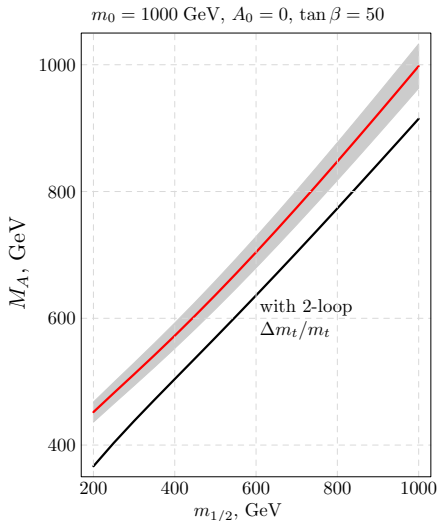
$A_0 = 0, \tan\beta = 50$



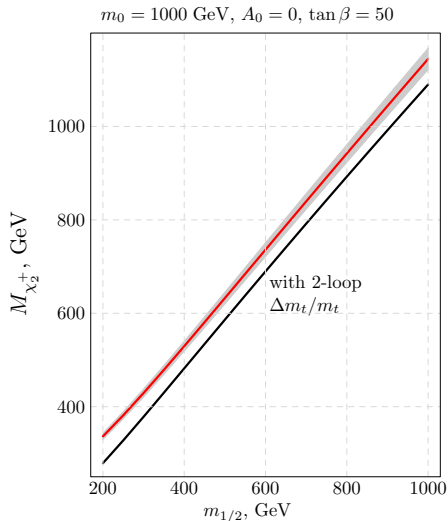
$$\sigma_2^{\text{rel}} \equiv \left| \frac{\Delta^{(2)} m_t/m_t}{\Delta^{(0)} m_t/m_t} \right|$$

second-order terms of asymptotic expansion in  $m_t/M_{SUSY}$  give negligible contribution

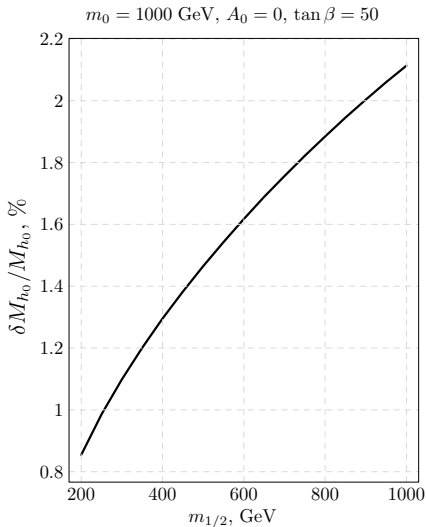




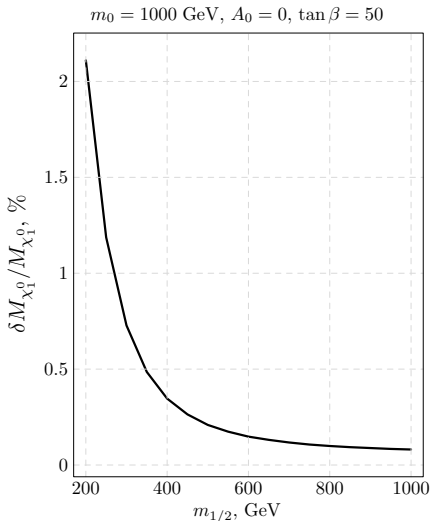
- Two-loop corrections to  $\Delta m_t/m_t$  yields sizable change ( $\approx 15\%$ ) of predicted masses of heavy Higgs bosons and chargino
- This change exceeds discrepancies (indicated by gray region on the plot) between different software for MSSM spectrum calculation



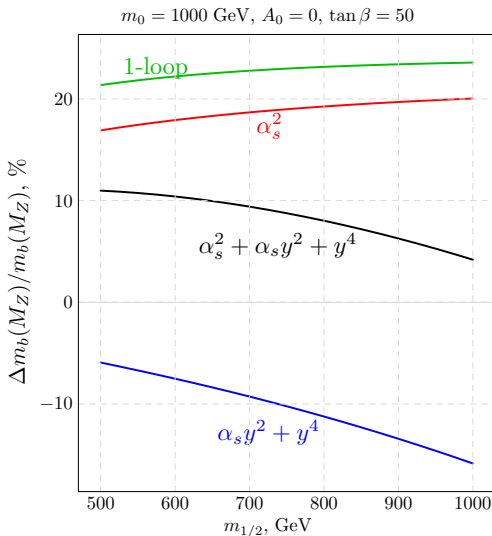
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Masses of squarks, gluino, and relatively light particles (lightest neutralino, lightest Higgs boson) do not obtain any significant changes.



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- $\mathcal{O}(\alpha_s^2)$  correction is positive and comparable to the one-loop MSSM correction
- $\mathcal{O}(\alpha_s y^2 + y^4)$  contribution has the opposite sign and partially compensate  $\mathcal{O}(\alpha_s^2)$  corrections
- $\mathcal{O}(\alpha_s y^2 + y^4 + \alpha_s^2)$  corrections  $\propto 30 - 40\%$  of the one-loop MSSM ones

- $\mathcal{O}(\alpha_s^2)$  MSSM corrections to the relation between pole and running masses of the  $t$  quark were evaluated.
- these corrections are  $\approx 30\%$  of one-loop ones
- second-order terms of asymptotic expansion in  $m_t/M_{SUSY}$  give negligible contribution
- this correction yields  $\gtrsim 15\%$  change of predicted masses of heavy Higgs bosons and chargino

- $\mathcal{O}(\alpha_s y^2 + y^4)$  MSSM corrections to the relation between pole and running masses of the  $b$  quark were evaluated.
- $\mathcal{O}(\alpha_s^2)$  correction is positive and of the same order of magnitude as the one-loop MSSM correction
- $\mathcal{O}(\alpha_s y^2 + y^4)$  contribution has the opposite sign in most “interesting” regions of MSSM parameter space and partially compensate  $\mathcal{O}(\alpha_s^2)$  corrections
- $\mathcal{O}(\alpha_s y^2 + y^4 + \alpha_s^2)$  corrections  $\propto 30 - 40\%$  of the one-loop MSSM ones.

Our results can be used for

- calculation of the MSSM mass spectrum
- renormalization group analysis of the Yukawa coupling unification
- dark matter searches, relic density is sensitive to the masses of heavy quarks for large  $\tan \beta$